

Università degli Studi di Padova

MODELS OF THEORETICAL PHYSICS

DI

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Introduction

The aim of this document is to collect the notes of *Models of Theoretical Physics* course, held by professors Marco Baiesi and Amos Maritan for the "Physics of data" curriculum in the academic year 2018-2019 (which is the first year of this new curriculum), to have them written in a neater way.

As just told, this document is far from pretending to be perfect and his goal is to help studying in a neater and better way. For this reason, there may be some errors among it.

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I have tried to be as much ruthless as possible in finding and correcting errors and mistakes, and I apologize if some have survived.

I hope that the overall result will anyway be satisfactory.

Padua, October 2018
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Chapter 1

Methods to compute usefull integrals

In this course we will use many kind of particular integrals. For this reason, in this first chapter we will see a brief introduction to all we need to calculate them.

1.1 Gaussian integrals

Let us consider the following integral:

$$Z(A) = \int d^n x \cdot e^{-A_2(\vec{x})} \quad (1.1)$$

where $Z(A)$ is the integral we want to calculate, $A_2(\vec{x})$ is a x quadratic form as the following:

$$A_2(\vec{x}) = \frac{1}{2} \sum_{i,j=1}^n x_i A_{ij} x_j \quad (1.2)$$

Here A is a matrix that can be identified as a metric matrix. In the case we consider A is symmetric with (generally) complex coefficients; furthermore, it has non-negative real parts and non-vanishing eigenvalues a_i .

$$\operatorname{Re}(a_i) \geq 0 \quad \text{and} \quad a_i \neq 0_{\mathbb{C}}$$

(Note: if these conditions are not true the integral would be divergent...).

Let us bring $A \in \mathbb{R}$. In this case it can be digonalized with orthogonal transformation O , for which it holds:

$$\sum_i O_{ij} x_j = x'_j \quad \text{and} \quad |\det(O)| = 1$$

This implies that the Jacobian matrix of the transformation J (which is O itself) has $|\det(J)| = 1$: so when we use it to change coordinates in the integral, we can say that "it has no effect" (we multiply by 1). Furthermore, with that transformation the non-diagonal coefficients of the new matrix become all 0 (the matrix after the transformation is diagonal), so the new coordinates x'_j are independents: this mean that we can evaluate $Z(A)$ by dividing it in the integrals of every single x'_j .

We obtain:

$$Z(A) = (2\pi)^{\frac{n}{2}} \prod_{i=1}^n a_i^{-\frac{1}{2}} = (2\pi)^{\frac{n}{2}} (\det(A))^{-\frac{1}{2}} \quad (1.3)$$

(Note that for a diagonal matrix the product of the eigenvalues is equal to the determinant. Moreover, for Binet's theorem it holds $\det(A \cdot O) = \det(A) \cdot \det(O)$ (true if A and O are square matrix with same dimensions)).

(Note: I think that what just written in the previous note modifies the conditions on $\det(O)$, which I think would be $\det(O) = +1$ and not $|\det(O)| = 1$ otherwise the last formula is meaningless.

This have to be investigated.

However, the important key is that the integral with that tranformation does not become divergent.)

Chapter 2

Path integrals

The most captivating feature of the path-integral technique is that it provides a unified approach to solving problems in different branches of physics. For example in the formalism of the stochastic processes we use: $W(x, t|0, 0) = \sum_{trajectory: x_0 \rightarrow x} e^{-F(trajectory)}$, but this formulation come up from the QM where the probability become the amplitude $K(x, t|0, 0) = \sum_{trajectory: x_0 \rightarrow x} e^{\frac{2\pi i S}{\hbar}(trajectory)}$ where $S = \int L dt$.

This formulation is not mathematically rigorous but the results are correct, since there is always the possibility to formulate it in a proper way but it will take us too much time.

In this kind of sum we use the concept of configuration, which means to assign to every particles of the system, we decide to study, their position and momentum. This will allow us to make averages in this way: $\langle A \rangle = \sum_{config} e^{\beta E(config)} A$

2.1 Diffusion