

Università degli Studi di Padova

MODELS OF THEORETICAL PHYSICS

DI

TOMMASO TABARELLI

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Introduction

The aim of this document is to collect the notes of *Models of Theoretical Physics* course, held by professors Marco Baiesi and Amos Maritan for the "Physics of data" curriculum in the academic year 2018-2019 (which is the first year of this new curriculum), to have them written in a neater way.

As just told, this document is far from pretending to be perfect and his goal is to help studying in a neater and better way. For this reason, there may be some errors among it.

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I have tried to be as much ruthless as possible in finding and correcting errors and mistakes, and I apologize if some have survived.

I hope that the overall result will anyway be satisfactory.

Padua, October 2018
Tommaso Tabarelli

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Chapter 1

Methods to compute usefull integrals

(02/10/2018)

In this course we will use many kind of particular integrals. For this reason, in this first chapter we will see a brief introduction to all we need to calculate them.

1.1 Gaussian integrals

Let us consider the following integral:

$$Z(A) = \int d^n x \cdot e^{-A_2(\vec{x})} \quad (1.1)$$

where $Z(A)$ is the integral we want to calculate, $A_2(\vec{x})$ is a x quadratic form as the following:

$$A_2(\vec{x}) = \frac{1}{2} \sum_{i,j=1}^n x_i A_{ij} x_j \quad (1.2)$$

Here A is a matrix that can be identified as a metric matrix. In the case we consider A is symmetric with (generally) complex coefficients; furthermore, it has non-negative real parts and non-vanishing eigenvalues a_i .

$$\operatorname{Re}(a_i) \geq 0 \quad \text{and} \quad a_i \neq 0_{\mathbb{C}}$$

(Note: if these conditions are not true the integral would be divergent...).

Let us bring $A \in \mathbb{R}$. In this case it can be digonalized with orthogonal transformation O , for which it holds:

$$\sum_i O_{ij} x_j = x'_j \quad \text{and} \quad |\det(O)| = 1$$

This implies that the Jacobian matrix of the transormation J (which is O itself) has $|\det(J)| = 1$: so when we use it to change coordinates in the integral, we can say that "it has no effect" (we multiply by 1). Furthermore, with that transformation the non-diagonal coefficients of the new matrix become all 0 (the matrix after the transformation is diagonal), so the new coordinates x'_j are independents: this mean that we can evaluate $Z(A)$ by dividing it in the integrals of every single x'_j .

We obtain:

$$Z(A) = (2\pi)^{\frac{n}{2}} \prod_{i=1}^n a_i^{-\frac{1}{2}} = (2\pi)^{\frac{n}{2}} (\det(A))^{-\frac{1}{2}} \quad (1.3)$$

(Note that for a diagonal matrix the product of the eigenvalues is equal to the determinant. Moreover, for Binet's theorem it holds $\det(A \cdot O) = \det(A) \cdot \det(O)$ (true if A and O are square matrix with same dimensions)).

(Note: I think that what just written in the previous note modifies the conditions on $\det(O)$, which I think would be $\det(O) = +1$ and not $|\det(O)| = 1$ otherwise the last formula is meaningless.

This have to be investigated.

However, the important key is that the integral with that tranformation does not become divergent.)

Since (1.1) and $\det(A)$ are analytic functions of A 's coefficients, we can extend (1.3) to \mathbb{C} .

Example:

$$\begin{aligned} A &= \begin{pmatrix} 3 & -1 \\ -1 & 3 \end{pmatrix} \Rightarrow (3 - \lambda)^2 - 1 = 0 \\ \rightarrow \lambda_{1,2} &= \frac{3 + \sqrt{9-8}}{2} \Rightarrow \begin{matrix} \lambda_1 = 2 \\ \lambda_2 = 4 \end{matrix} \\ Z(A) &= 2\pi \frac{1}{\sqrt{4 \cdot 2}} = \sqrt{\frac{\pi}{2}} \end{aligned}$$

1.2 More general integrals

Let us try to generalize the previous arguments for a more general kind of integral:

$$Z(A) = \int d^n x \cdot e^{-A_2(\vec{x}) + \vec{b} \cdot \vec{x}} \quad (1.4)$$

With $\vec{b} \cdot \vec{x} = \sum_{i=1}^n b_i x_i$. We want to have the integral in the form of (1.1) so we expand the exponent near his maximum x^* .

To find x^* (signs are already changed):

$$\frac{\partial}{\partial x_i} (A_2(x) - \vec{b} \cdot \vec{x}) = 0 \Rightarrow \sum_j A_{ij} x_j = b_i$$

(Note: $A_2 = A/2$: in the derivative the factor 1/2 goes away because we are derivating a quadratic form (if one does the calculation, it can be seen easy).)

So the solution is:

$$x^* = \sum_i (A^{-1})_{ij} b_j$$

Now we change variable $\vec{x} \Rightarrow \vec{y}$ in order to have the integral in the form of (1.1). \vec{y} is the deviation from the min of the exponential:

$$\vec{y} = \vec{x} - \vec{x}^* \Rightarrow x_i = x_i^* + y_i$$

$$\Rightarrow -A_2(\vec{x}) + \vec{b} \cdot \vec{x} = w_2(\vec{b}) - A_2(\vec{y})$$

Notice that $w_2(\vec{b})$ is a costant, so it can be eventually take outside the integral.

$$w_2(\vec{b}) = \frac{1}{2} \sum_{i,j=0}^n b_i (A^{-1})_{ij} b_j = \frac{1}{2} \vec{b} \cdot \vec{x}^*$$

With those changes, we can write:

$$Z(A, \vec{b}) = e^{w_2(\vec{b})} \int d^n y \cdot e^{-A_2(\vec{y})} = e^{w_2(\vec{b})} (2\pi)^{\frac{n}{2}} (\det(A))^{-\frac{1}{2}} \quad (1.5)$$

Example (as before)

$$A = \begin{pmatrix} 3 & -1 \\ -1 & 3 \end{pmatrix} \quad \text{and} \quad b = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$Z(A) = \int \int e^{-\frac{1}{2}(x,y)A\begin{pmatrix} x \\ y \end{pmatrix} + \begin{pmatrix} x \\ 0 \end{pmatrix}} dx dy$$

$$A^{-1} = \frac{1}{\det(A)} \begin{pmatrix} A_{22} & -A_{12} \\ -A_{21} & A_{11} \end{pmatrix} = \frac{1}{8} \begin{pmatrix} 3 & 1 \\ 1 & 3 \end{pmatrix}$$

$$w_2(\vec{b}) = \frac{1}{2}(1,0) \frac{1}{8} \begin{pmatrix} 3 & 1 \\ 1 & 3 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{16}(1,0) \begin{pmatrix} 3 \\ 1 \end{pmatrix} = \frac{3}{16}$$

$$\Rightarrow Z(A, \vec{b}) = e^{\frac{3}{16}} Z(A) = e^{\frac{3}{16}} \frac{\pi}{\sqrt{2}}$$

1.3 Gaussian expectation values

We want now to explore briefly how the expectation values of polynomials with a gaussian distribution works.

$$\langle x_{k_1}, x_{k_2}, \dots, x_{k_m} \rangle = \underbrace{Z^{-1}(A, \vec{0})}_{\text{Normalization factor}} \int d^n x x_{k_1}, x_{k_2}, \dots, x_{k_m} e^{-A_2(\vec{x})} \quad (1.6)$$

Where the "normalization factor" is that for which $\langle 1 \rangle = 1$.

Note: since here we use angular brackets to denote an average made with gaussian weights.

From (1.4) we have:

$$\frac{\partial}{\partial b_k} Z(A, \vec{b}) = \int d^n x x_k e^{-A_2(\vec{x}) + \vec{b} \cdot \vec{x}}$$

In this sense we say that \vec{b} is "completed" to \vec{x} and that Z is the generating function. In the case of only 1 variable average note that if $\vec{b} = 0$ the result is 0, as one expect for a gaussian variable.

Generalizing (1.6) using what just shown, we obtain:

$$\langle x_{k_1}, x_{k_2}, \dots, x_{k_l} \rangle = (2\pi)^{-\frac{n}{2}} (\det A)^{\frac{1}{2}} \left[\frac{\partial}{\partial b_{k_1}} \frac{\partial}{\partial b_{k_2}} \dots \frac{\partial}{\partial b_{k_m}} Z(A, \vec{b}) \right]_{b=0} \quad (1.7)$$

NOT SURE ARGUMENT

We can also notice that if we want to evaluate the expectation value of a power series of x $\langle F(x) \rangle$ we can write the following, thanks to the linearity of the integral (assuming that x is not a vector):

$$\langle F(x) \rangle = \dots \left[F \left(\frac{\partial}{\partial b} \right) e^{w_2(b)} \right]_{b=0}$$

The "..." indicates that every power of x in $F(x)$ has its own normalization factor.

From (1.7) we can obtain one form of Wick's theorem.

If there is one $\frac{\partial}{\partial b_i}$ in (1.7), it pulls down some b_j because $w_2(\vec{b}) = \frac{1}{2} \sum_{i,j=0}^n b_i (A^{-1})_{ij} b_j$. Without another $\frac{\partial}{\partial b_k}$ (also with $k \neq i$), setting $\vec{b} = \vec{0}$ makes all components $b_j \rightarrow 0$ in $\langle \dots \rangle$.

$$\lim_{b_1 \rightarrow 0} \langle b_1, \dots \rangle \rightarrow 0$$

\Rightarrow Another $\frac{\partial}{\partial b_k}$ is needed to ensure the average does not tend to 0.

$\Rightarrow l$ index in (1.7) must be even.

Theorem 1.3.1 (Wick's theorem) *Using arguments just above, for every $k_p k_q$ from the list $\{k_1, k_2, \dots, k_l\}$, with associated $(A^{-1})_{k_p k_q}$, the average of l (multiplied) variables can be calculated as following:*

$$\begin{aligned} \langle x_{k_1}, x_{k_2}, \dots, x_{k_l} \rangle &= \sum_{\text{all possible pairings } P \text{ of } \{k_1, \dots, k_l\}} A_{k_{P_1} k_{P_2}}^{-1} \dots A_{k_{P_{l-1}} k_{P_l}}^{-1} \\ &= \sum_{\text{all possible pairings } P \text{ of } \{k_1, \dots, k_l\}} \langle x_{k_{P_1}} x_{k_{P_2}} \rangle \cdot \dots \cdot \langle x_{k_{P_{l-1}}} x_{k_{P_l}} \rangle \end{aligned}$$

Example:

Let us see an application 1-dim of the theorem: with $x, b \in \mathbb{R}$

$$\begin{aligned} \langle x^2 \rangle &= \frac{1}{Z(A, 0)} \frac{\partial^2}{\partial b^2} \int_{-\infty}^{\infty} x^2 e^{-\frac{A}{2}x + bx} dx \Big|_{b=0} \\ &= \frac{\partial^2}{\partial b^2} e^{\frac{b^2}{2A}} \Big|_{b=0} = \frac{\partial}{\partial b} \left[\frac{b}{A} e^{\frac{b^2}{2A}} \right]_{b=0} = \frac{1}{A} \end{aligned}$$

$$\begin{aligned} \langle x^4 \rangle &= \frac{1}{Z(A, 0)} \frac{\partial^4}{\partial b^4} \int_{-\infty}^{\infty} x^4 e^{-\frac{A}{2}x + bx} dx \Big|_{b=0} = \frac{\partial^4}{\partial b^4} e^{\frac{b^2}{2A}} \Big|_{b=0} = \frac{\partial^3}{\partial b^3} \left[\frac{b}{A} e^{\frac{b^2}{2A}} \right]_{b=0} \\ &= \frac{\partial^2}{\partial b^2} \left[\frac{1}{A} e^{\frac{b^2}{2A}} + \frac{b}{A} \cdot \frac{b}{A} e^{\frac{b^2}{2A}} \right]_{b=0} = \frac{\partial}{\partial b} \left[\frac{1}{A} \cdot \frac{b}{A} e^{\frac{b^2}{2A}} + 2 \frac{b}{A^2} e^{\frac{b^2}{2A}} + \frac{b^2}{A^2} \cdot \frac{b}{A} e^{\frac{b^2}{2A}} \right]_{b=0} \\ &= \left[\frac{1}{A^2} e^{\frac{b^2}{2A}} + \frac{b}{A^2} \cdot \frac{b}{A} e^{\frac{b^2}{2A}} + \frac{2}{A^2} e^{\frac{b^2}{2A}} + \frac{2b}{A^2} \cdot \frac{b}{A} e^{\frac{b^2}{2A}} + \frac{3b^2}{A^3} e^{\frac{b^2}{2A}} + \frac{b^3}{A^3} \cdot \frac{b}{A} e^{\frac{b^2}{2A}} \right]_{b=0} \\ &= \frac{3}{A^2} \end{aligned}$$

Remembering that $\langle x^2 \rangle = \frac{1}{A}$ we can write:

$$\langle x^4 \rangle = 3 \langle x^2 \rangle^2$$

Using Wick's theorem we can rewrite this last result: With $l = 4$ and A as following:

$$A = \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & 1 & \\ & & & 1 \end{pmatrix}$$

we can write:

$$\langle x_1 x_2 x_3 x_4 \rangle = \langle x_1 x_2 \rangle + \langle x_3 x_4 \rangle + \langle x_1 x_3 \rangle \langle x_2 x_4 \rangle + \langle x_1 x_4 \rangle \langle x_2 x_3 \rangle$$

Now come back to the 1-dim case, using $x = x_{1,2,3,4}$, obtaining:

$$\langle x^4 \rangle = \langle x^2 \rangle \cdot \langle x^2 \rangle \cdot 3 = 3 \langle x^2 \rangle^2$$

which is the same result as before.