HW₅

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```
import numpy as np
import matplotlib.pyplot as plt
from math import factorial as fact
```

Exercise 1

Let's develop the functions to compute Stumpff's functions $c_n(z)$, n=1,2,...,5 and their first derivatives $c_n'(z)$, n=1,2,3.

We can re-write the Stumpff's function in a numerical computation way using a recursion method.

From

$$c_n(z) = \sum_{k=0}^{\infty} (-1)^k rac{z^k}{(2k+n)!}$$

we can define

$$t_k(z) = (-1)^k rac{z^k}{(2k+n)!}$$

and by the ration of two adjacent terms we find

$$egin{array}{ccc} rac{t_k}{t_{k-1}} &\longrightarrow & t_k(z) = -rac{z}{(2k+n-1)(2k+n)} \cdot t_{k-1}(z) \end{array}$$

such that

$$c_n(z) = \sum_{k=0}^\infty t_k(z)$$

In the other hand we can write the $c_{n}^{^{\prime}}(z)$ as a linear combination of $c_{n}(z)$

$$c_n'(z) = rac{n}{2} c_{n+2}(z) - rac{1}{2} c_{n+1}(z)$$

For low values of the order n, when the argument is sufficiently far from zero, we can use the Stumpff's functions in term of (hyperbolic) trig functions since they're closely related to the elementary circular function. We use a threshold of 10^{-2} to divide the trigonometric functions into the simple trigonometric functions.

```
# trigonometric functions
pos_func = [lambda x: np.cos(x), lambda x: np.sin(x)/x, lambda x: (1-np.cos(x))/x**2, la
neg_func = [lambda x: np.cosh(x), lambda x: np.sinh(x)/x, lambda x: (np.cosh(x)-1)/x**2,
# recorsive Version
def tk(k, z, n):
```

```
if k == 0:
        return 1/fact(n)
        return - (z)/((2*k + n - 1.)*(2*k + n)) * tk(k-1, z, n)
def uk(k, z, n):
    if k == 0:
        return 1/fact(n)
    else:
        return - (k + 1.)/((2*k + n + 1.)*(2*k + n + 2.)) * tk(k, z, n)
def Stumpffs_Cn(z, n, max_iter = 100, tol = 1e-4):
    c = 0.0
    k = 0.0
    diff = 1.
    if np.abs(z) > 1e-2:
        while (diff > tol) and (k < max_iter):</pre>
            temp = tk(k, z, n)
            diff = np.abs(temp - c)
            c = c + temp
            k += 1
        return c
    elif z>0:
        return pos_func[n]((z)**0.5)
    else:
        return neg_func[n]((-z)**0.5)
def Stumpffs_Cn_deriv_linear_comb(z, n):
    return 0.5*n*Stumpffs_Cn(z, n+2) - 0.5*Stumpffs_Cn(z, n+1)
```

Now let's calculate $c_n(z)$ and $c_n^{\prime}(z)$ for various n=1,2,3,4,5

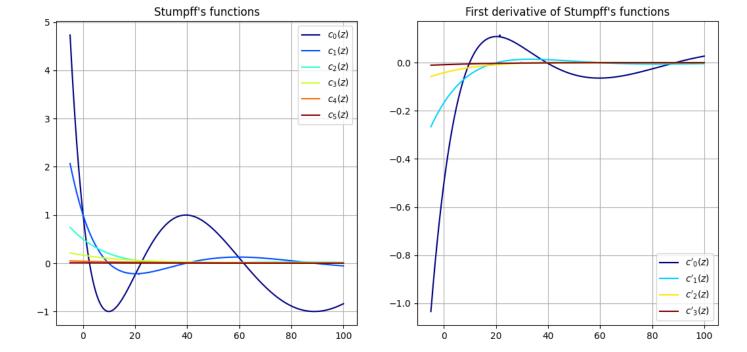
```
In [3]: import matplotlib.cm as cmaps

x = np.linspace(-5, 100, 1000)

fig, ax = plt.subplots(nrows=1, ncols=2, figsize=(13,6))

ax[0].set_title("Stumpff's functions")
for n, col in zip(range(0,6), cmaps.jet(np.linspace(0,1,6))):
    ax[0].plot(x, [Stumpffs_Cn(z, n=n) for z in x], color=col, label =f'$c_{n}(z)$')
    ax[0].legend()
    ax[0].grid()

ax[1].set_title("First derivative of Stumpff's functions")
for n, col in zip(range(0,4), cmaps.jet(np.linspace(0,1,4))):
    ax[1].plot(x, [Stumpffs_Cn_deriv_linear_comb(z, n=n) for z in x], color=col, label =f
    ax[1].legend()
    ax[1].grid()
```



Exercise 2

The solution of the universal Kepler equation, for example the determination of the parameter s in the terms of the time of flight $t-t_0$ and the other known parameters μ, r_0, \dot{r}_0 and C of the orbit can be found by applying Newton-Raphson to find the root and after that the derivative of the time of flight with respect to the fictitios time is simply given by the radius vector r(s):

$$egin{aligned} t-t_0 &= r_0 s + r_0 \dot{r}_0 s^2 c_2 \left(-2 C s^2
ight) + \left(\mu + 2 r_0 C
ight) s^3 c_3 \left(-2 C s^2
ight) \ & \\ r(s) &= r_0 c_0 \left(-2 C s^2
ight) + r_0 \dot{r}_0 s c_1 \left(-2 C s^2
ight) + \mu s^2 c_2 \left(-2 C s^2
ight) \end{aligned}$$

```
In [4]:
        from scipy.optimize import newton
        # Newton-Raphson Method
        def Newton_Raphson(func, dfunc, x, tol = 1e-6, max_iter = 1000):
            while (abs(func(x)) > tol) and (i < max_iter):</pre>
                x = x - func(x)/dfunc(x)
                i += 1
             return x
        # Kepler Equation
        def Kepler_eq(t, t0, r0, dr0, mu, C):
             func = lambda s: r0*s + r0*dr0*s**2.*Stumpffs_Cn(z=(-2*C*s**2.), n=2) + (mu + 2*r0*C*)
             return func
        def dev_Kepler_eq(r0, dr0, mu, C):
             func_prime = lambda s: r0 * Stumpffs_Cn(z=(-2*C*s**2.), n=0) + r0*dr0*s*Stumpffs_Cn(
             return func_prime
        # resolve
        def resolve_Kepler_eq(t, t0, r0, dr0, mu, C, s_guess):
             #solution1 = Newton_Raphson(func=Kepler_eq(t, t0, r0, dr0, mu, C), dfunc=dev_Kepler_
             solution2 = newton(func=Kepler_eq(t, t0, r0, dr0, mu, C), x0=s_guess, fprime=dev_Kep
             return solution2
```

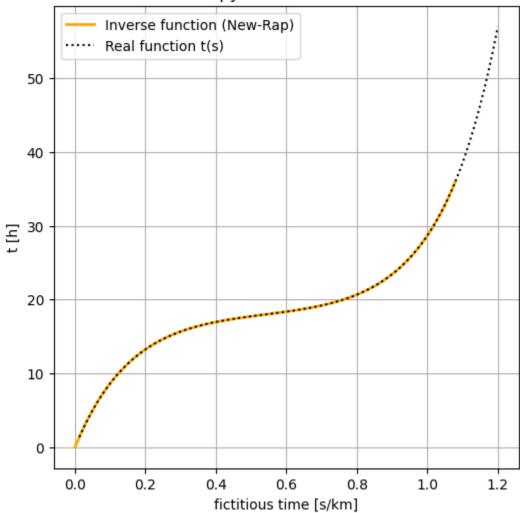
Exercise 3

```
In [5]: # Initial conditions for Venus Flyby
        R0 = np.array([-267733.084163, 199426.194677, 254709.414665]) # Initial position
        R1 = np.array([177071.935393, -334448.764629, -184024.725921]) #Final position
        V0 = np.array([4.168950, -2.598877, -3.925639])
        GM = 324859.2139518842890 # Venus's GM
```

Let's propagate for 36 hours the following Venus-centric dynamical state. To do that we use the data of the Galileo Orbiter

```
We start by choosing an guess value os s=0
In [6]: r0 = np.linalg.norm(R0)
         dr0 = np.dot(R0, V0)/r0
                                          #radial component of the velocity
         C = 0.5*np.dot(V0, V0) - GM/r0 #Constant of energy
        T = np.linspace(0, 36*3600, 100)
        Venus_dynamical_states = np.zeros_like(T)
         for i, time in enumerate(T):
             Venus_dynamical_states[i] = resolve_Kepler_eq(t=time, t0=0.0, r0=r0, dr0=dr0, mu=GM,
In [7]: def t_(s, t0, r0, dr0, C, mu):
           return r0*s + r0*dr0*s**2.*Stumpffs_Cn(z=(-2*C*s**2.), n=2) + (mu + 2*r0*C)*s**3. * St
        # True function
         S = np.linspace(0, 1.2, 100)
         t_{int} = np.array([t_{int}(s, 0, r0, dr0, C, GM) for s in S])
         udm = 3600 #Plot the times in hours
         fig, ax= plt.subplots(figsize=(6, 6))
         ax.set_title('Scipy Newton method')
         ax.plot(Venus_dynamical_states, T/udm, color='orange', lw=2, label='Inverse function (Ne
         ax.plot(S, t_list/udm, ls='dotted', color='k', label='Real function t(s)')
         ax.set_xlabel('fictitious time [s/km]')
         ax.set_ylabel('t [h]')
         ax.legend()
         ax.grid()
         plt.show()
         plt.tight_layout()
        C:\Users\bosca\AppData\Local\Temp\ipykernel_23344\1070099785.py:3: RuntimeWarning: inval
        id value encountered in scalar divide
          neg_func = [lambda x: np.cosh(x), lambda x: np.sinh(x)/x, lambda x: (np.cosh(x)-1)/x**
        2, lambda x: (np.sinh(x)-x)/x^{*3}, lambda x: (np.cosh(x)-1-x^{*2}/2)/x^{*4}, lambda x: (np.sinh(x)-x)/x^{*3}
        nh(x)-x-x**3/6)/x**5
```

Scipy Newton method



<Figure size 640x480 with 0 Axes>

We can see from the plot that our root-finding algorithm follows pretty well the teoretical function

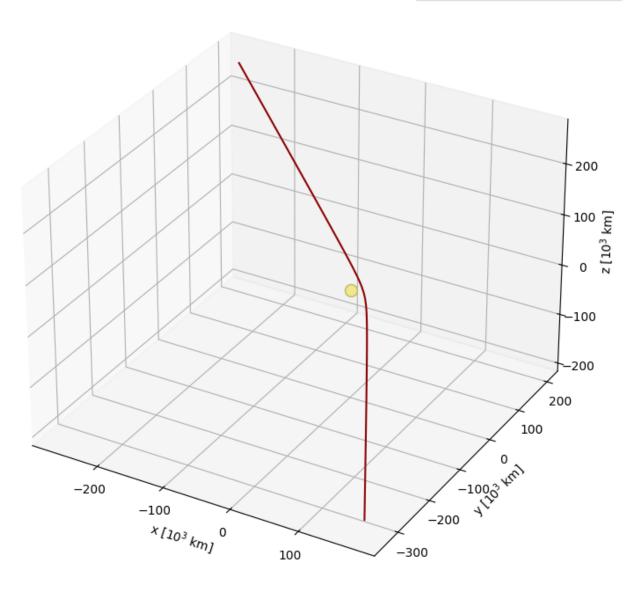
The position and velocity at a given parameter are given by the universal Lagrange coefficients matrix

$$\left[egin{array}{c} r \ \dot{r} \end{array}
ight] = \left[egin{array}{cc} F & G \ \dot{F} & \dot{G} \end{array}
ight] \cdot \left[egin{array}{c} r_0 \ \dot{r}_0 \end{array}
ight]$$

with

$$F = 1 - rac{\mu}{r_0} s^2 c_2 \left(-2C s^2
ight) \ G = t - t_0 - \mu s^3 c_3 \left(-2C s^2
ight) \ \dot{F} = -rac{\mu}{r(s)r_0} ig[s + 2C s^3 c_3 (-2C s^2) ig] \ \dot{G} = 1 - rac{\mu}{r(s)} s^2 c_2 \left(-2C s^2
ight)$$

```
return np.array([F, G, dF, dG])
        '''Radius as a function of s'''
        def r_(s, r0, dr0, C, mu):
          return r0*Stumpffs_Cn(z=(-2*C*s**2.), n=0) + r0*dr0*s*Stumpffs_Cn(z=(-2*C*s**2.), n=1)
        '''Get position and velocity during an orbit'''
        def fly_by(time, mu, R0, V0):
          r0 = np.linalg.norm(R0)
          dr0 = np.dot(R0, V0)/r0
                                         #radial component of the velocity
          C = 0.5*np.dot(V0, V0) - mu/r0 #Constant of energy
          s_new = resolve_Kepler_eq(t=time, t0=0.0, r0=r0, dr0=dr0, mu=GM, C=C, s_guess=1.0) #
          r_s = r_(s=s_new, r0=r0, dr0=dr0, C=C, mu=mu)
          ULM_new = ULM(s=s_new, t=time, t0=0.0, mu=mu, r=r_s, r0=r0, C=C)
          R = np.array([ULM_new[0]*R0 + ULM_new[1]*V0])
          V = np.array([ULM_new[2]*R0 + ULM_new[3]*V0])
          return R, V
In [9]: T = np.linspace(0, 36*3600, 101)
        R_{int} = np.zeros((len(T), 3)) # Replace 3 with the length of the sequence returned by
        V_{list} = np.zeros((len(T), 3)) # Replace 3 with the length of the sequence returned by
        for i, t in enumerate(T):
          R_{i} = fly_{by}(time=t, mu=GM, R0=R0, V0=V0)
        udm = 1e3
        fig = plt.figure(figsize=(8,8))
        ax = fig.add_subplot(projection='3d')
        ax.plot(R_list[:,0]/udm, R_list[:,1]/udm, R_list[:,2]/udm, color='darkred', label='Orbit
        ax.scatter(0,0,0, marker='o', color='khaki', edgecolors='darkkhaki', s=100, label='Plane
        ax.legend()
        ax.set_xlabel('x [$10^3$ km]')
        ax.set_ylabel('y [$10^3$ km]')
        ax.set_zlabel('z [$10^3$ km]')
        ax.set_box_aspect(aspect=None, zoom=0.9)
        plt.tight_layout()
        plt.show()
```



```
In [10]: R_final, V_final = fly_by(time=36*3600, mu=GM, R0=R0, V0=V0)
    print('The final dynamical state is :')
    print('R: ', R_final[0], '[km]')
    print('V: ', V_final[0], '[km/s]')
    dist = np.linalg.norm(R_list[0]-R1)
    print(f'\n The distance from the true value R:{R_final} [km], is {dist} km')

The final dynamical state is :
    R: [ 176956.92970277 -334433.44069936 -184558.85545383] [km]
    V: [ 2.36979447 -5.24495423 -2.53289844] [km/s]
The distance from the true value R:[[ 176956.92970277 -334433.44069936 -184558.8554538
```

Exercise 4

3]] [km], is 821803.883972285 km

The Universal Lambert's equation gives the energy constant needed to orbit from to in a given time.

Similarly in Ex. 2 we use Newton-Raphson method to find the root of the equation that we can consider as the combination of the 3 equations. From the hand-calculated resolution, The derivate of the equation with respect to *z* is the following:

$$rac{A^2}{8 \mu s} + rac{3 s c_3(z)}{2 c_2(z)} \Biggl(rac{A \sqrt{c_2(z)}}{4} - \mu s^2 c_2^{'}(z) \Biggr) + \mu s^3 c_3^{'}(z)$$

and we try to use z=1 as starting point so that our algorithm converge

```
In [11]: def Lambert_eq(R0, R, Dt, mu):
    r0 = np.linalg.norm(R0)
    r = np.linalg.norm(R)
    Df = np.arccos(np.dot(R0, R)/(r0*r)) # true anomaly

A = (np.sqrt(r0*r)*np.sin(Df))/(np.sqrt(1 - np.cos(Df)))
    s2 = lambda Z: (mu*Stumpffs_Cn(z=Z, n=2))**(-1) * (r0 + r - A*(Stumpffs_Cn(z=Z, n=1)))

f = lambda Z: A*np.sqrt(s2(Z)*Stumpffs_Cn(z=Z, n=2)) + mu*(s2(Z))**1.5*Stumpffs_Cn(z=Z, n=1))

f = lambda Z: A**2/(8*mu*s2(Z)**0.5) + 3*s2(Z)*Stumpffs_Cn(z=Z, n=3)/(2*Stumpf)

z_new = newton(func=f, x0=1.0, fprime=f_prime, tol=1e-12)

F = 1 - (GM/r0)*s2(z_new)*Stumpffs_Cn(z=z_new, n=2)
    G = 36*3600 - GM*s2(z_new)**1.5*Stumpffs_Cn(z=z_new, n=3)

V0 = (R- F*R0)/G

    return V0
```

Exercise 5

45383] km

The initial velocity can be found by inverting with $r=F\mathbf{r}_0+G\dot{\mathbf{r}}_0$ with F, G being the Lagrange coefficient in Ex.3

```
V = Lambert_{eq}(R0, R1, 36*3600, GM)
In [12]:
         print(f'The initial velocity found by the Lambert equation is {V} [km/s]')
         print(f'The real value is {V0} [km/s]')
         print(f'The velocity difference is {VO-V} km/s (with magnitude {np.linalg.norm(VO-V)} km
         The initial velocity found by the Lambert equation is [ 3.13466068 -4.3618355 -3.137947
         5 ] [km/s]
         The real value is [ 4.16895 -2.598877 -3.925639] [km/s]
         The velocity difference is [ 1.03428932    1.7629585   -0.7876915 ] km/s (with magnitude 2.
         190487377935409 km/s)
         R_V, _ = fly_by(time=36*3600, mu=GM, R0=R0, V0=V)
In [13]:
         print(f'The final position with this velocity is \{R_V[0]\}\ km,')
         R_V0, _ = fly_by(time=36*3600, mu=GM, R0=R0, V0=V0)
         print(f'while the one found in the propagation is {R_V0[0]} km')
         The final position with this velocity is [ 177071.935393 -334448.764629 -184024.725921]
         while the one found in the propagation is [ 176956.92970277 -334433.44069936 -184558.855
```