# Green Tea: User Guide

(Dated: March 30, 2022)

This guide provides the basic information necessary to setup and run transport simulations with the Green Tea suite of scientific programs. These software tools can solve the self-consistent quantum transport problem by using an *ab-initio* derived Hamiltonian and the Keldysh-Green's functions formalism.

#### I. PREAMBLE

The Green Tea suite of programs consists of two main parts. The first part, located in the directory gt\_h, is devoted to the determination of the Unit-cell Restricted Bloch Functions (URBF) that form the basis set usd for transport simulations Ref. [1], and then to the calculation of the diagonal and off-diagonal blocks of the Hamiltonian matrix in such a URBF basis set. The executable code is named gt\_h.x.

The second part of the Green Tea suite, located in the directory  $\operatorname{gt}_{-}$ t, is devoted to the actual transport simulation and it can handle electron device structures. This part of the programs compute the Green's function elements describing macroscopic quantities, such as the electron charge density and the current. The kinetic equations for the retarded Green's function,  $\mathbf{G}$ , and the lesser(grater)-than Green's function,  $\mathbf{G}^{<(>)}$ , are self-consistently coupled with a non-linear 3D Poisson equation in order to describe the device electrostatics. Phonon scattering is included within the self-consistent Born approximation. This requires the self-consistent solution of the kinetic equations for the retarded and the lesser(greater) Green's functions. The executable is named  $\operatorname{gt}_{-}\mathbf{t}$ .

#### II. DESCRIPTION OF GT\_H

The goal of this program is to extract the Hamiltonian blocks  $[\mathbf{H}_{0,0}]_{\Phi}$  and  $[\mathbf{H}_{0,1}]_{\Phi}$  and the unit-cell restricted Bloch functions  $\{\Phi_{\mathbf{k}}\}$ , which will be used as inputs in the program for transport. An orthorhombic unit cell, with sides  $a_x$ ,  $a_y$  and  $a_z$  is needed to perform transport calculations. The program can compute the Hamiltonian blocks of homogeneous materials, and also the Hamiltonian blocks connecting two different materials (provided that the two materials are lattice matched). This coupling Hamiltonian is required only when

heterostructures along the transport direction are considered (heterostructures along other directions can be included in the unit cell).

The gt\_h program needs as an input the outputs of DFT simulations obtained with Quantum ESPRESSO (QE) for the unit cell of the system at study. The unit cell has to be orthorombic, thus it is typically larger than the primitive unit cell of the system. For monolayer MoS<sub>2</sub> and for most 2D materials, for example, the orthorombic unit cell is twice the primitive (trigonal) unit cell. Once the self-consistent simulation of the orthorhombic unit cell has been obtained by QE, the user has to perform a non self consistent simulation to compute the Bloch functions that will be actually used for the transport problem. Once this is done, from the version 7.0 of QE, it is possible to use a post-processing code named pw2gt.x, which saves relevant information about the plane-wave DFT Hamiltonian and the Bloch functions in the output file named output.dat.

The procedure to generate the Hamiltonian blocks for transport calculations can be summarized as follows:

- 1. Build an orthorhombic unit cell to be used in transport calculations. To this purpose it is advisable to start from the primitive cell first, which is most often reported in the literature, and then extend the primitive cell to an orthorhombic unit cell that will correspondingly reduce the size of the Brillouin zone.
- Perform a "scf" simulation of the orthorhombic unit cell by using QE.
   The present version of Green Tea requires that you use only norm-conserving pseudo-potentials.
- 3. Generate a "bands" input file to produce the QE outputs necessary to calculate the Hamiltonian matrix for the wave-vectors of interest. The program assumes that transport direction is along the x axis. For each transverse wave-vector  $(k_y, k_z)$  one should compute the Bloch functions for the two  $k_x$  wave-vectors  $(0, 1/2) \times 2\pi/a_x$ .

A more accurate reconstruction of the bandstructure in the reduced URBF basis set can be obtained by using the four  $k_x$  values  $(-1/4, 0, 1/4, 1/2) \times 2\pi/a_x$  [1].

As an example, for a 2D material we have  $k_z = 0$  and one can span the entire Brillouin zone by suing  $k_y$  points ranging from 0 to 0.5, with a constant sampling step of  $\Delta k_y =$ 

 $0.1 \times 2\pi/a_y$ . Please notice that only positive  $k_y$  values are needed, therefore a possible choice for the k-points may be

### K\_POINTS crystal

24

-0.25 0.0 0.0 1

 $0.00\ 0.0\ 0.0\ 1$ 

 $0.25\ 0.0\ 0.0\ 1$ 

 $0.50\ 0.0\ 0.0\ 1$ 

-0.25 0.1 0.0 1

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-0.25 0.5 0.0 1

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 $0.50\ 0.5\ 0.0\ 1$ 

**Practical remark.** In order to be sure that QE calculates and outputs the electronic structure only at the selected k points, we should declare **nosym**='.true' in the

namelist &SYSTEM in the input file for the "bands" simulation.

You can also increase the number of computed Bloch functions in order to increase the accuracy by properly setting the variable **nbnd** in the same namelist.

- 4. Run "mpirun -np NP pw.x -in bands.in > bands.out &", being NP the number of your nodes.
- 5. Run the command "pw2gt.x -in pw2gt.in", being pw2gt.in a simple post-processing input of QE. Be careful: pw2gt.x does not run in parallel. This will produce the file output.dat.
- 6. Run "nohup qt\_h.x < input.file &"

It is important to note that  $\operatorname{\mathbf{gt\_h.x}}$  will automatically generate the Hamiltonian block matrices in the reduced basis using all the Bloch functions specified in the "nscf" calculation. Therefore for a given transverse wave-vector  $(k_y, k_z)$ , the number of the basis elements will be nm=nbnd  $\times N_{k_x}$ , being  $N_{k_x}$  the number of k-points sampling the Brillouin zone along the transport direction (typically 2 or 4).

This procedure works safely, in the sense that no spurious states are observed with respect to the electronic structure output by QE, and in fact a close agreement can be achieved with DFT calculations [1]. However this approach tends to produce a fairly large URBF basis set, which that in many practical instances can be further reduced. This is because in transport problems it is typically required to have a proper description of the material bandstructure only in the energy window where the current flow is non negligible, namely for energies close to energy gap in semiconductors or close to the Fermi level in metals.

In gt\_h.x it is therefore possible to refine the basis set by choosing a subset of Bloch functions able to correctly describe the material bandstructure in the energy window relevant for transport problems (typically a few eV or less). This can be done by reducing the sampling of the BZ for bands energetically far from the transport energy window or by eliminating the bands at very low energies. Such an operation can be performed by setting the variable refine='.true.'.

The user should be aware that the refinement procedure can reduce the accuracy of the bandstructure reconstruction, of even . Hence the user should check that the refined basis

works correctly by comparing the bands in the refined basis (output files **Edisp\_k\_m.dat**) with those in the unrefined basis (output files **Evdisp\_k\_m.dat**).

#### A. Input variables for gt\_h.x

Before running **gt\_h.x** you need to modify the input file to setup the simulation to be performed. The variables are described below.

# • &indata\_dimensionality

- ncell (integer). Number of unit cells (either 1 or 2). Use ncell = 1 when a single material is considered. Use ncell = 2, when the Hamiltonian matrix coupling two different materials is needed. In this case the program requires that Hamiltonians and Bloch functions for the two isolated materials were previously calculated.
- $\mathbf{ncy}$  (integer). Number of unit cells along the y axis. If  $\mathbf{ncy} > 1$  the system is periodic along the y axis. This must be consistent with the number of wavevectors in  $\mathbf{output.dat}$ .
- $\mathbf{ncz}$  (integer). Number of unit cells along the z axis. If  $\mathbf{ncz} > 1$  the system is periodic along the z axis. This must be consistent with the number of wavevectors in  $\mathbf{output.dat}$ .
- $\mathbf{nky}$  (integer). Number of  $k_y$  components considered in the simulation ( $\mathbf{nky} \ge 1$ ). Since only  $k_y \ge 0$  are considered (we assume time-reversal symmetry for negative  $k_y$ ),  $\mathbf{nky} \le \mathbf{ncy}/2+1$   $\mathbf{nky}$  must be consistent with the number of wave-vectors in  $\mathbf{output.dat}$ .
- $\mathbf{nkz}$  (integer). Number of  $k_z$  components considered in the simulation ( $\mathbf{nkz} \geq 1$ ). Since only  $k_z \geq 0$  are considered (we assume time-reversal symmetry for negative  $k_z$ ),  $\mathbf{nkz} \leq \mathbf{ncz}/2+1$   $\mathbf{nkz}$  must be consistent with the number of wave-vectors in  $\mathbf{output.dat}$ .

### &indata\_basis

nband\_v (integer). The number of valence bands of the system. This information
 can be extracted from the output of the "scf" calculations by QE.

- **ni** (*integer*). Initial band number of the selected Bloch function set.
- **nf** (*integer*). Final band number of the selected Bloch function set.
- $\mathbf{nkx}$  (integer). Number of  $k_x$  values used for the URBF basis set, with band index  $\mathbf{ni} \leq n \leq \mathbf{nf}$ . Allowed values: 2, 4, 8.
- refine (logical). Default: '.false.'. It is necessary to run at first a simulation with refine='.false.'. If refine='.true.', it is possible to search for a possible reduction of the selected Bloch functions set (this second simulation will be much faster).
- Mplus (integer). Default: 0. If Mplus > 0, it allows you to add a second set of Bloch functions in your reduced basis. Their band index will run from nf+1 to nf+Mplus.
- **nkplus** (integer). Number of  $k_x$  values considered for the Bloch functions with band index  $\mathbf{nf}+1 \leq n \leq \mathbf{nf}+\mathbf{Mplus}$ . Allowed values: 2, 4, 8.
- **Ecut** (*double precision*). The cutoff energy [Ry]. It is recommend to use the same value used in the "nscf" calculation.

### • &indata\_basis2

This namelist has to be specified only if  $\mathbf{ncell} = 2$ , namely if one wishes to compute the coupling Hamiltonian between two materials forming a longitudinal heterojunction. It assumes that you have already run the simulation of the second material, which is supposed to be on the left side of the heterostructure.

- nm2 (integer). The order of the Hamiltonian matrices of the second materials evaluated for the unrefined basis. Please note that the number of elements of the matrix stored in HH00\_nkyz\_0\_nmat\_0.dat must be nm2×nm2.
- mm2 (integer). The order of the Hamiltonian matrices of the second materials evaluated on the refined basis. mm2 ≤ nm2. Please note that the number of elements of the matrix stored in H00\_nkyz\_0\_nmat\_0.dat must be mm2×mm2.
- indir2 (character). The directory where the output files of the second material were saved.

### • &indata\_cell

- **ac1** (double precision). This is  $a_x$ , the side of the orthorhombic unit cell along the x axis [cm].
- **ac2** (double precision). This is  $a_y$ , the side of the orthorhombic unit cell along the y axis [cm].
- **ac3** (double precision). This is  $a_z$ , the side of the orthorhombic unit cell along the z axis [cm].

### • &indata\_inout

- **outdir** (*character*). The directory where the output files will be saved.
- input\_file\_DFT (character). The output file of the "nscf" simulation.

#### B. List of output files of gt\_h.x

We here recall that the Hamiltonian matrix in the URBF reduced basis can be expressed by the block tridiagonal form [1]

$$[\mathbf{H}] = \begin{bmatrix} \mathbf{H}_{0,0} & \mathbf{H}_{0,1} & 0 & 0 & \cdots & \mathbf{H}_{0,1}^{\dagger} \\ \mathbf{H}_{0,1}^{\dagger} & \mathbf{H}_{0,0} & \mathbf{H}_{0,1} & 0 & \cdots & 0 \\ \cdots & \cdots & \ddots & \vdots \\ \mathbf{H}_{0,1} & 0 & \cdots & 0 & \mathbf{H}_{0,1}^{\dagger} & \mathbf{H}_{0,0} \end{bmatrix}$$
(1)

The main output files of **gt\_h.x** are listed below.

- H00\_nkyz\_k\_nmat\_m.dat. The diagonal block of the Hamiltonian matrix in the refined basis set for the k-th lateral wave-vector and the m-th material (see Eq.1). These files will be used by gt\_t.x to perform the self-consistent transport calculations.
- H01\_nkyz\_k\_nmat\_m.dat. The off-diagonal block of the Hamiltonian matrix in the refined basis set for the k-th lateral wave-vector and the m-th material (see Eq.1). These files will be used by gt\_t.x to perform the self-consistent transport calculations.
- H01\_nkyz\_k\_nhet\_h.dat. The Hamiltonian block matrix connecting two different materials in the refined basis set for the k-th lateral wave-vector and the h-th heterojunction. Computed only if ncell=2. These files are used by gt\_t.x to perform the self-consistent transport calculations only for hetero-structures.

- Psi\_Bloch\_nkyz\_k\_nmat\_m.dat. The ortho-normalized Bloch functions in the refined basis set for the k-th lateral wave-vector and the m-th material. These files will be used by gt\_t.x to perform the self-consistent transport calculations.
- HH00\_nkyz\_k\_nmat\_m.dat. The diagonal block of the Hamiltonian matrix in the unrefined basis set for the k-th lateral wave-vector and the m-th material. These files will be used to refine the basis set. Once this operation is finished they can be deleted.
- HH01\_nkyz\_k\_nmat\_m.dat. The off-diagonal block of the Hamiltonian matrix in the unrefined basis set for the k-th lateral wave-vector and the m-th material. These files will be used to refine the basis set. Once this operation is finished they can be deleted.
- **PPsi\_Bloch\_nkyz\_k\_nmat\_m.dat**. The ortho-normalized Bloch functions in the the unrefined basis set for the k-th lateral wave-vector and the m-th material. These files will be used to refine the basis set. Once this operation is finished they can be deleted.
- Evdisp\_ $k_m$ .dat. The energy dispersion [eV] vs  $k_x$  [ $2\pi/a_x$ ] of the n-th (sub)band for the k-th lateral wave-vector computed in the unrefined basis set.
- Edisp\_k\_m.dat. The energy dispersion [eV] vs  $k_x$  [ $2\pi/a_x$ ] of the n-th (sub)band for the k-th lateral wave-vector computed in the refined basis set.

The user should compare  $\mathbf{Edisp}_{-k}$ -m. $\mathbf{dat}$  and  $\mathbf{Evdisp}_{-k}$ -m. $\mathbf{dat}$  to check if the refined basis works well. In particular it is important to double check that no spurious states have been produced by the basis reduction, and that the agreement with the electronic structure output by QE is satisfactorily good, which may also depend on the application.

### III. DESCRIPTION OF GT\_T

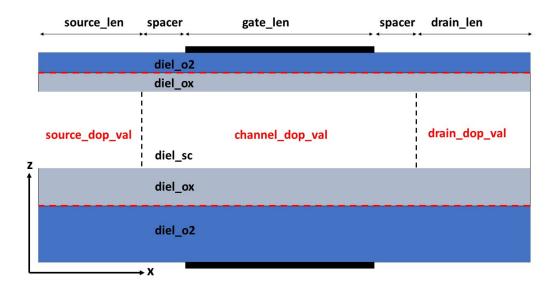


FIG. 1. 2D sketch of the simulated device along the x-z plane.

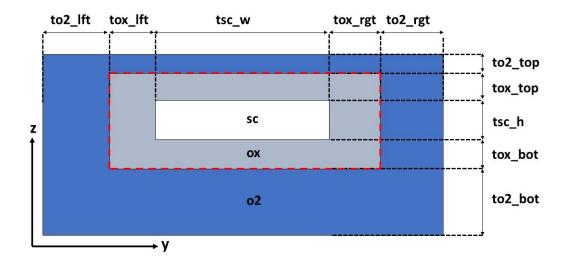


FIG. 2. 2D sketch of the simulated device along the y-z plane. The transport unit cell is delimited by the red dashed lines.

The program self-consistently solves the kinetic equations for the retarded and lesser-than Green's functions and the 3-D nonlinear Poisson equation (see Ref. [1] for more details). This program uses as inputs the outputs of the program **gt\_h**, which have been previ-

ously saved in the files named H00\_nkyz\_k\_nmat\_m.dat, H01\_nkyz\_k\_nmat\_m.dat and Bloch\_nkyz\_k\_nmat\_m.dat.

Two different spatial meshes are used: a relatively coarse grid for the Poisson equation and a finer grid for the charges computed by the NEGF solver. The program can simulate 3D, 2D and 1D electron devices. In QE in order to simulate confined systems (either 2D or 1D), it is necessary to insert vacuum regions in the unit cell. These vacuum region should correspond to the internal oxide regions inside Green Tea and they can have assigned an arbitrary dielectric constant. Additional dielectric regions outside the unit cell can be also included in the self-consistent simulation.

The transport direction is assumed to be along the x axis, while transverse directions are along the y-z plane. Two sketches in Figs. 1 and 2 clarify the definition of the geometrical parameters. Non-null values of the oxide thickness either y or z must be used only when the system has a quantum confinement along respectively the y or the z direction.

To run the program use the command "nohup gt\_t.x < input\_file &".

### A. Input variables for gt\_t.x

### • &indata\_mesh

- Ndeltax (integer). Number of the elements inside the unit cell along the x axis in the Poisson mesh. The physical length of the unit cell along the x axis,  $a_x$ , is defined by the variable ac1 in &indata\_cell.
- Ndeltay (integer). Number of the elements inside the unit cell along the y axis in the Poisson mesh. The physical length of the unit cell along the y axis,  $a_y$ , is defined by the variable ac2 in &indata\_cell.
- Ndeltaz (integer). Number of the elements inside the unit cell along the z axis in the Poisson mesh. The physical length of the unit cell along the z axis,  $a_z$ , is defined by the variable ac3 in &indata\_cell.

### • &indata\_lengths

 source\_len (integer). Number of unit cells along the transport direction of the source region (see Fig. 1).

- drain\_len (integer). Number of unit cells along the transport direction of the drain region (see Fig. 1).
- gate\_len (integer). Number of unit cells along the transport direction of the channel region (see Fig. 1).
- spacer (integer). Number of unit cells along the transport direction of the spacer regions (see Fig. 1).

#### • &indata\_oxide

- tox\_top (integer). Number of the elements of oxide in the top region of the unit cell.
- tox\_bot (integer). Number of the elements of oxide in the bottom region of the unit cell.
- tox\_lft (integer). Number of the elements of oxide in the left region layer of the unit cell.
- tox\_rgt (integer). Number of the elements of oxide in the right of the unit cell.
- to2\_top (integer). Number of the elements of additional oxide in the top region outside the unit cell.
- to2\_bot (integer). Number of the elements of additional oxide in the bottom region outside the unit cell.
- to2\_lft (integer). Number of the elements of additional oxide in the left region outside the unit cell.
- to2\_rgt (integer). Number of the elements of additional oxide in the right region outside the unit cell.

### &indata\_channel

- **tsc\_w** (*integer*). Number of the elements of semiconductor inside the unit cell along the y axis. It is assumed that **tsc\_w+tox\_lft+tox\_rgt=Ndeltay**
- **tsc\_h** (*integer*). Number of the elements of semiconductor inside the unit cell along the z axis. It is assumed that **tsc\_h+tox\_bot+tox\_top=Ndeltaz**

### • &indata\_regions

- $\operatorname{num\_mat}$  (integer). Number of materials along the transport direction. If any heterojunction along the transport direction is considered, then  $\operatorname{num\_mat} > 1$ .
- **num\_reg** (*integer*). Number of regions along the transport direction.

### • REGIONS

After the declaration of **num\_reg**, it is necessary to declare **num\_reg** arrays. Each array declares the specific region number, the number of unit cells corresponding to the specific region and the corresponding material number:

### reg\_i nc\_i mat\_i

Be careful: the sum of nc\_i over reg\_i must equal source\_len + drain\_len + gate\_len + 2×spacer.

# • &indata\_heterojunctions

- **num\_het** (*integer*). Number of heterojunctions (HJs) along the transport direction. *Default*: 0.

#### • HETEROJUNCTIONS

If **num\_het** > 0, then it is necessary to declare **num\_het** arrays specifying the triplet given by the number of the HJ, the number of material at the left and the number of the material at the right of the HJ:

het\_i mat\_lft mat\_rgt

# • &indata\_doping

- source\_dop\_val (double precision). Doping concentration in the source region [cm<sup>-3</sup>]. If positive (negative), it assumes donors (acceptors).
- drain\_dop\_val (double precision). Doping concentration in the drain region [cm<sup>-3</sup>]. If positive(negative), it assumes donors (acceptors).
- channel\_dop\_val (double precision). Doping concentration in the channel region
   [cm<sup>-3</sup>]. If positive (negative), it assumes donors (acceptors).

## • &indata\_gate

 top\_gate (logical). If '.true.', it considers a gate in the top side of the channel region.

- bot\_gate (logical). If '.true.', it considers a gate in the bottom side of the channel region.
- lft\_gate (logical). If '.true.', it considers a gate in the left side of the channel region.
- rgt\_gate (logical). If '.true.', it considers a gate in the right side of the channel region.

### • &indata\_device

- chtype (character). It can be 'n' for n-type devices (carrier transport in the conduction (sub)bands), 'p' for p-type devices (carrier transport in the valence (sub)bands) and 't' for tunnel devices (carrier transport in both the conduction and valence (sub)bands).
- **DIEL\_SC** (double precision). Relative dielectric constant of the semiconductor.
- DIEL\_OX (double precision). Relative dielectric constant of the oxide inside the unit cell.
- **DIEL\_O2** (double precision). Relative dielectric constant of the oxide outside the unit cell.

### &indata\_dimensionality

- **ncy** (integer). Number of unit cells along the y axis. If **ncy** > 1 the system is periodic along the y axis. This defines the sampling of the wave vector  $k_y$  as  $\Delta k_y = 1/\text{ncy}$  in units of  $2\pi/a_y$ .
- **ncz** (integer). Number of unit cells along the z axis. If **ncz** > 1 the system is periodic along the z axis. This defines the sampling of the wave vector  $k_z$  as  $\Delta k_z = 1/\mathbf{ncz}$  in units of  $2\pi/a_z$ .
- $\mathbf{nky}$  (integer). Number of  $k_y$  components considered in the simulation ( $\mathbf{nky} \ge 1$ ). Since only  $k_y \ge 0$  are considered (we assume time-reversal symmetry for negative  $k_y$ ),  $\mathbf{nky} \le \mathbf{ncy}/2+1$ .
- $\mathbf{nkz}$  (integer). Same as  $\mathbf{nky}$ , but for  $k_z$ .

### • SELECTED\_KYZ

The list of the expected  $(k_y, k_z)$  has to be specified in the next  $\mathbf{nky} \times \mathbf{nkz}$  lines. The values must obey  $0 \le k_y \le 0.5$  and  $0 \le k_z \le 0.5$  in units of  $2\pi/a_y$  and  $2\pi/a_z$ , respectively. The third column is a logical (either T or F), which includes or not the specific  $(k_y, k_z)$  in the calculation.

## • &indata\_basis

- **nsolv** (*integer*). Number of highest valence (sub)band saved in output files.
- **nsolc** (*integer*). Number of lowest conduction (sub)band saved in output files.
- g\_spin (double precision). It can be either 1 (spin-degeneracy off) or 2 (spin-degeneracy on).
- npol (integer). Number of polarizations. Same as in the output of gt\_h.x. It can be either 1 or 2.
- magnetic (logical). Default: 'F'. It must be '.true.' for magnetic materials. If this is the case, the user must specify also the variable updw.
- updw (character). It can be either 'up' or 'dw' depending on the considered polarization.
- Ecut (double precision). Cutoff energy [Ry]. This indirectly determines also the spatial grid of the charges computed by the NEGF solver. It should be the same used in QE.

### BASIS\_INFO

**num\_mat** arrays have to be specified with the information about the Hamiltonian block matrices computed by **gt\_h**. Each line contains the number of the material, the number of the basis elements, the number of valence states, the offset energy [eV] (at least one material has to be have zero offset).

### mat\_i NB\_i NV\_i offset\_i

These numbers must be consistent with the Hamiltonian blocks stored in the input files H00\_nkyz\_k\_nmat\_m.dat and H01\_nkyz\_k\_nmat\_m.dat.

### • &indata\_cell

- **ac1** (double precision). This is  $a_x$ , the side of the orthorhombic unit cell along the x axis [cm].
- **ac2** (double precision). This is  $a_y$ , the side of the orthorhombic unit cell along the y axis [cm].
- **ac3** (double precision). This is  $a_z$ , the side of the orthorhombic unit cell along the z axis [cm].

## • &indata\_convergence

- **ERROR\_INNER** (double precision). Tolerance of the Poisson solver. Default:  $10^{-10}$ .
- MAX\_ITER\_INNER (integer). Maximal number of iterations in the Poisson solver. Default: 30.
- ERROR\_OUTER (double precision). Tolerance of the self-consistent Poisson-NEGF loop. Default:  $10^{-3}$ .
- MAX\_ITER\_OUTER (integer). Maximal number of iterations of the selfconsistent Poisson-NEGF loop. Default: 30.
- alphapot (double precision). Sub-relaxation factor of the self-consistent Poisson-NEGF loop (0 < alphapot  $\leq$  1). Default: 1.
- Nomp (integer). Number of threads for the OPENMP parallelization. Be careful: by increasing Nomp you increase the memory requirement. Default: 1.

## • &indata\_energy

- **Eop** (double precision). Unit of energy intervals in the NEGF solver [eV].
- Nop<sub>-g</sub> (integer). Number of Eop equal to the phonon energy such that  $\hbar\omega_{op} = \text{Nop}_{-g} \times \text{Eop}$ .
- **Dop**-g (double precision). Multiplication factor in front of the optical phonon self-energy  $\hbar D_{op}^2/(2\rho\omega_{op})$ . Units are [eV<sup>2</sup>/cm<sup>3</sup>]. If non null, then the optical phonon scattering is considered within the self-consistent Born approximation (SCBA). Default: 0.

- **Dac** (double precision). Multiplication factor in front of the acoustic phonon self-energy  $D_{ac}^2 k_B T/(\rho v_s^2)$ . Units are [eV<sup>2</sup>/cm<sup>3</sup>]. If non null, then the acoustic phonon scattering is considered within the SCBA. Default: 0.
- Nsub (integer). Number of energy points computed for each Eop. Default: 3.
- SCBA\_alpha (double precision). Sub-relaxation factor of the SCBA loop. Default: 1.
- SCBA\_tolerance (double precision). Tolerance of the SCBA loop. Default:
  10<sup>-3</sup>. In some cases it could be necessary to decrease this tolerance in order to ensure the conservation of the current flowing though the device.
- SCBA\_max\_iter (integer). Maximal number of iteration in the SCBA loop.
   Default: 50. In some cases it could be necessary to increase this number to achieve the convergence of the SCBA loop.
- eta (double precision). Imaginary term used in the calculation of the contact self-energies [eV]. Default:  $10^{-6}$ .
- TEMP (double precision). Lattice temperature [K]. Default: 300.
- **NKT** (integer). Number of  $K_BT$  layers computed above the top of the barrier for n-type devices or below the bottom of the barrier for p-type devices. Default: 10.

### • &indata\_stimulus

- onlyT (logical). If '.true.', then only the transmission is computed in the flat band regime without computing the charges and with no self-consistent loop. If '.false.', then the self-consistent solution is computed.
- **VGMIN** (double precision). Initial  $V_{GS}$  value [V].
- VGMAX (double precision). Final  $V_{GS}$  value [V].
- **DELTAVG** (double precision). Value of  $\Delta V_{GS}$  [V]. It must be always non null.
- **VDMIN** (double precision). Initial  $V_{DS}$  value [V].
- **VDMAX** (double precision). Final  $V_{DS}$  value [V].
- **DELTAVD** (double precision). Value of  $\Delta V_{DS}$  [V]. It must be always non null.

- workgate (double precision). It can be used to change the work-function of the metal gate [eV]. workgate = 0 means that the metal work-function is aligned to the energy reference defined as the valence band of the semiconductor at the  $\Gamma$  point.

# • &indata\_inout

- **outdir** (*character*). The directory where output files will be saved.
- inputdir (character). The directory where files where the Hamiltonian and the
   Bloch functions files computed by the gt\_h program are located.

### B. List of output files of gt\_t.x

Main output files of gt\_t.x are listed below.

- **convergence\_LOG.dat**. It saves the evolution of the error when the self-consistent solution is computed.
- Edisp\_n\_nkyz\_k\_nmat\_m.dat. Energy dispersion [eV] vs  $k_x$  [ $2\pi/a_x$ ] of the n-th band, the k-th lateral wavevector and the m-th material.
- doping\_profile.dat. Spatial profile of the doping along the transport direction.
- bal\_current\_vg\_d.dat. Transfer characteristics,  $I_{DS}$  [A] vs  $V_{GS}$  [V], at the d-th  $V_{DS}$  in the absence of phonon scattering. In laterally periodic systems, if you want to know the current density, you have to divide  $I_{DS}$  by the width (area) of the transport unit cell.
- scat\_current\_vg\_d.dat. Transfer characteristics,  $I_{DS}$  [A] vs  $V_{GS}$  [V], at the d-th  $V_{DS}$  in the presence of phonon scattering. In laterally periodic systems, if you want to know the current density, you have to divide  $I_{DS}$  by the width (area) of the transport unit cell.
- bal\_current\_vd\_g.dat. Output characteristics,  $I_{DS}$  [A] vs  $V_{DS}$  [V], at the g-th  $V_{GS}$  in the absence of phonon scattering. In laterally periodic systems, if you want to know the current density, you have to divide  $I_{DS}$  by the width (area) of the transport unit cell.

- scat\_current\_vd\_g.dat. Output characteristics,  $I_{DS}$  [A] vs  $V_{DS}$  [V], at the g-th  $V_{GS}$  in the presence of phonon scattering. In laterally periodic systems, if you want to know the current density, you have to divide  $I_{DS}$  by the width (area) of the transport unit cell.
- bal\_LDOS\_convergence\_vd\_d\_vg\_g.dat. Local density of states vs x [nm] and E [eV] for the g-th gate voltage and the d-th drain voltage in the absence of phonon scattering.
- scat\_LDOS\_convergence\_vd\_d\_vg\_g.dat. Local density of states vs x [nm] and E [eV] for the g-th gate voltage and the d-th drain voltage in the presence of phonon scattering.
- bal\_Jdens\_convergence\_vd\_d\_vg\_g.dat. Spatial distribution of the current spectrum vs x [nm] and E [eV] for the g-th gate voltage and the d-th drain voltage in the absence of phonon scattering.
- scat\_Jdens\_convergence\_vd\_d\_vg\_g.dat. Spatial distribution of the current spectrum vs x [nm] and E [eV] for the g-th gate voltage and the d-th drain voltage in the presence of phonon scattering.
- bal\_Jspectrum\_convergence\_kyz\_k\_vd\_d\_vg\_g.dat. Current spectrum for the k-th lateral wave-vector, the g-th gate voltage and the d-th drain voltage in the absence of phonon scattering.
- scat\_Jspectrum\_convergence\_kyz\_k\_vd\_d\_vg\_g.dat. Current spectrum for the k-th lateral wave-vector, the g-th gate voltage and the d-th drain voltage in the presence of phonon scattering.
- bal\_subc\_convergence\_kyz\_k\_vd\_d\_vg\_g.dat. Lowest conduction (sub)band [eV] along x [nm] for the k-th lateral wave-vector, the g-th gate voltage and the d-th drain voltage in the absence of phonon scattering.
- scat\_subc\_convergence\_kyz\_k\_vd\_d\_vg\_g.dat. Lowest conduction (sub)band [eV] along x [nm] for the k-th lateral wave-vector, the g-th gate voltage and the d-th drain voltage in the presence of phonon scattering.

- bal\_subv\_convergence\_kyz\_k\_vd\_d\_vg\_g.dat. Highest valence (sub)band [eV] along x [nm] for the k-th lateral wave-vector, the g-th gate voltage and the d-th drain voltage in the absence of phonon scattering.
- scat\_subv\_convergence\_kyz\_k\_vd\_d\_vg\_g.dat. Highest valence (sub)band [eV] along x [nm] for the k-th lateral wave-vector, the g-th gate voltage [V] and the d-th drain voltage [V] in the presence of phonon scattering.

 $[1]\,$  M. G. Pala, P. Giannozzi,  $\,$  and D. Esseni, Phys. Rev. B  $\bf 102,\,045410$  (2020).