



UNIVERSIDADE  
FEDERAL  
DE PERNAMBUCO



Departamento  
de Física



INFO

Instituto Nacional de  
Ciência e Tecnologia de Fotônica

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XVII J. A. Swieca School - Curitiba (March, 2023)

# Nonlinear Optics

2nd lecture

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# mini-course - Nonlinear Optics:

1st lecture

Introduction to nonlinear optics (NLO)

2nd lecture

Nonlinear Nanophotonics in Plasmonic media

3<sup>rd</sup> lecture

Stimulated Emission and Random Lasers



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## Research of new materials for NLO

- To increase the efficiency of new frequencies generation
  - Reduce the size of photonic devices
  - Reduce the optical power levels required for NLO
  - Optical processing with low power levels  
*etc, etc, etc.....*
- 

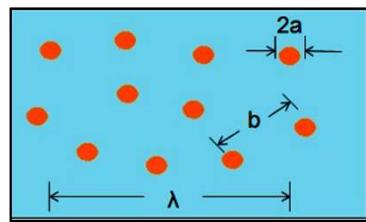
**Materials:** metal NPs, nanoclusters of metal atoms, crystalline nanopowders, glasses doped with NPs, liquid colloids, organic membranes, semiconductor “quantum dots”

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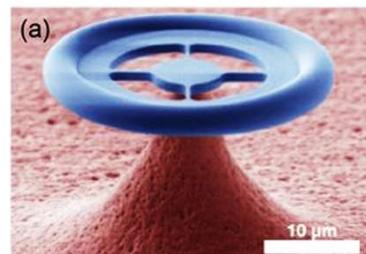
# NANOMATERIALS

# How to increase the NLO response of materials?

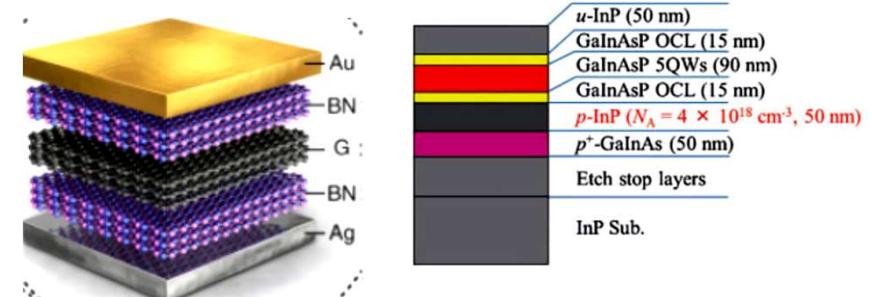
## Examples



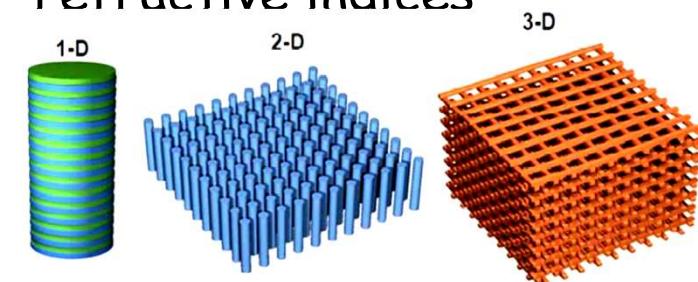
Nanocomposites  
Inclusions with larger NL response  
than the host



Micro-cavities with large quality factor



Hetero-structures with  
components having different  
refractive indices



Photonic crystals

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## When humans started manipulating objects in the nanoscale?

# Roman nanotechnology

Lycurgus cup  
400 yrs A.D.

Gold + silver  
NPs  
embeded in  
the glass



## Gold-Doped Glass: A Maxwell-Garnett Composite



Developmental Glass, Corning Inc.  
Red Glass Caraffe  
Nurenberg, ca. 1700  
Huelsmann Museum, Bielefeld

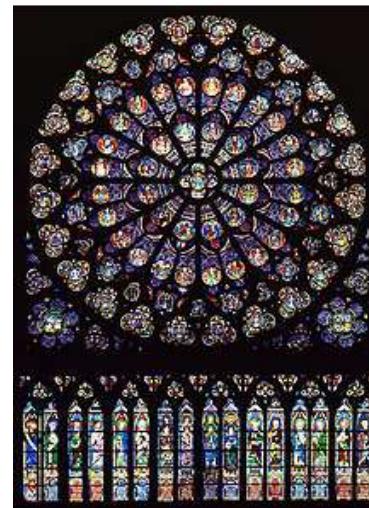


gold volume fraction approximately  $10^{-6}$   
gold particles approximately 10 nm diameter

- Composite materials can possess properties very different from those of their constituents.
- Red color is because the material absorbs very strong in the blue, at the surface plasmon frequency

Slide: R. W. Boyd

# Middle age



Cu, Ag, Au nanoparticles

Michael Faraday's  
Recognition of  
Ruby Gold: the  
Birth of Modern  
Nanotechnology

His 1857 Lecture to the Royal  
Society in London

David Thompson  
DTThompson@adl.com

Figure 1  
Faraday's colloidal ruby gold. Reproduced by Courtesy of the Royal Institution of Great Britain



From the paper in Philosophical Transactions entitled "Experimental relations of gold (and other metals) to light" (1), based on his Bakerian lecture to the Royal Society in London on 5 February 1857, it is clear that Michael Faraday was fascinated by the ruby colour of colloidal gold. The objective of his investigations was to examine the interaction of light with metal particles, but much of this paper focused on various

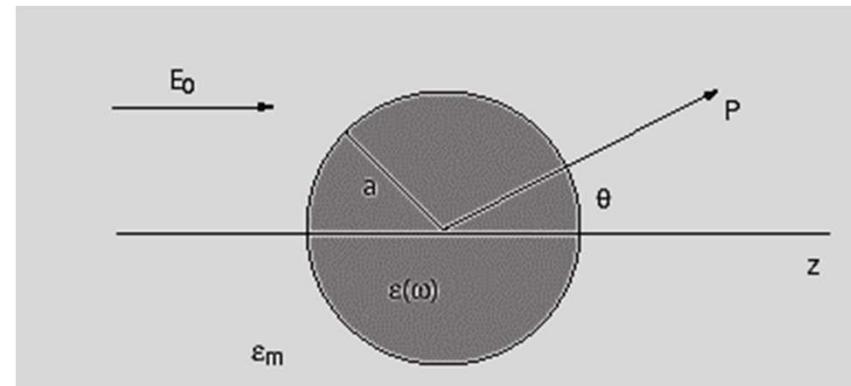
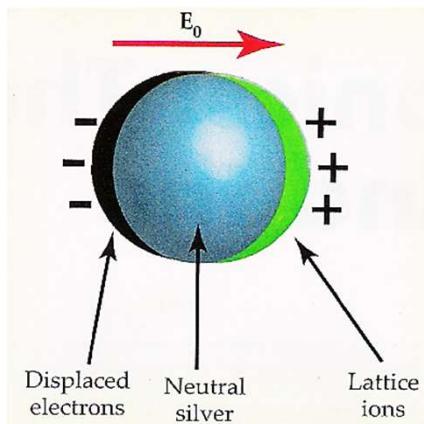
solution such as phosphorus in carbon disulfide in a two phase system. The yellow colour of sodium chloroaurate ( $\text{NaAuCl}_4$ ) changes within minutes to the deep ruby colour of colloidal gold. Faraday concluded that the ruby fluid was gold dispersed in the liquid in a very finely divided metallic form not visible in any of the microscopes available in his day. Nearly 100 years later Turkevich et al (2, 4) used electron microscopic

Experimental Relations of Gold (and Other Metals) to Light,  
M. Faraday, Philos. Trans. R. Soc. London, 147 (1857) 145-181

# Metal nanoparticles for NLO

Dimensions smaller than 100 nm

Quase-static approximation



Solve  
Laplace's  
equation

$$\Phi(r, \theta) = \sum_{l=0}^{\infty} [A_l r^l + B_l r^{-(l+1)}] P_l(\cos \theta)$$

$$\Phi_{\text{in}} = -\frac{3\epsilon_m}{\epsilon + 2\epsilon_m} E_0 r \cos \theta$$

$$\Phi_{\text{out}} = -E_0 r \cos \theta + \frac{\epsilon - \epsilon_m}{\epsilon + 2\epsilon_m} E_0 a^3 \frac{\cos \theta}{r^2}.$$

$$\Phi_{\text{out}} = -E_0 r \cos \theta + \frac{\mathbf{p} \cdot \mathbf{r}}{4\pi \varepsilon_0 \varepsilon_m r^3}$$

$$\mathbf{p} = 4\pi \varepsilon_0 \varepsilon_m a^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m} \mathbf{E}_0.$$

$$\mathbf{p} = \varepsilon_0 \varepsilon_m \alpha \mathbf{E}_0$$

$$\alpha = 4\pi a^3 \frac{\varepsilon - \varepsilon_m}{\varepsilon + 2\varepsilon_m}.$$

**Polarizability**

$$\mathbf{E}_{\text{in}} = \frac{3\varepsilon_m}{\varepsilon + 2\varepsilon_m} \mathbf{E}_0$$

$$\mathbf{E}_{\text{out}} = \mathbf{E}_0 + \frac{3\mathbf{n}(\mathbf{n} \cdot \mathbf{p}) - \mathbf{p}}{4\pi \varepsilon_0 \varepsilon_m} \frac{1}{r^3}.$$

$$\text{Re} [\varepsilon(\omega_{sp}) + 2\varepsilon_m(\omega_{sp})] = 0$$

**Surface plasmon resonance**

# Metal NPs as nanoantennae

$$E_{local} = L E_{light}$$

Local field factor

$$L = \frac{3\epsilon_h(\omega)}{[\epsilon_{NP}(\omega) + 2\epsilon_h(\omega)]}$$

$$\text{Re} [\epsilon_{NP}(\omega_{SP}) + 2\epsilon_h(\omega_{SP})] = 0$$

Localized surface plasmons resonance

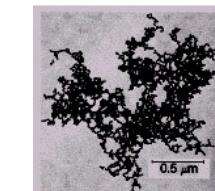
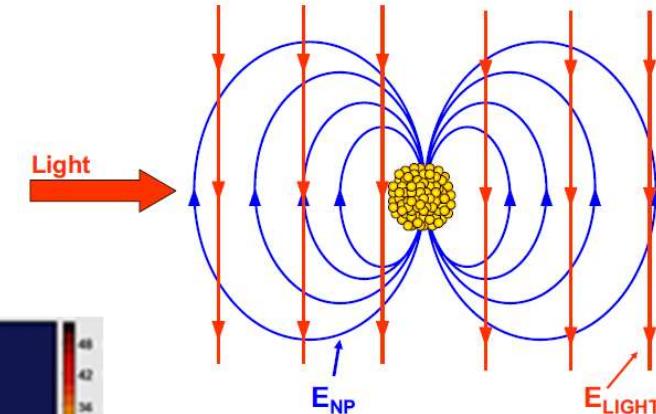
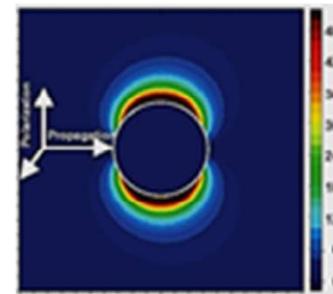
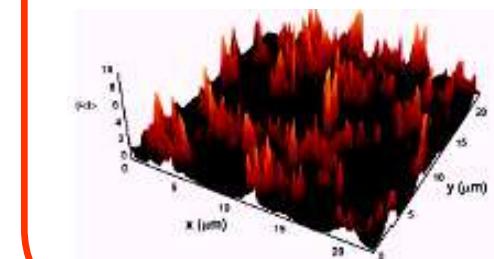


FIG. 1. TEM image of typical gold colloid aggregate. This cluster contains 479 gold particles.

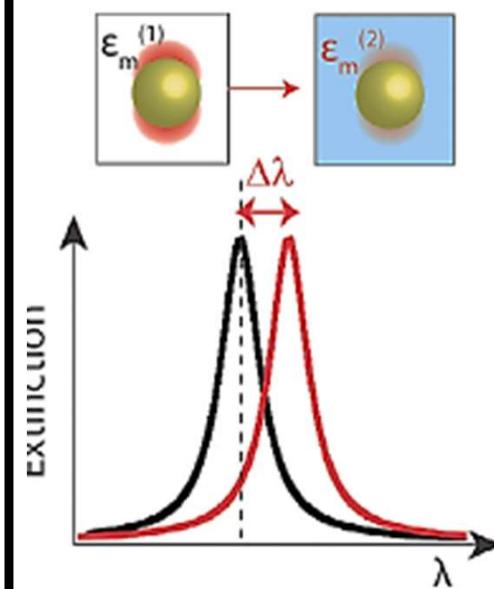
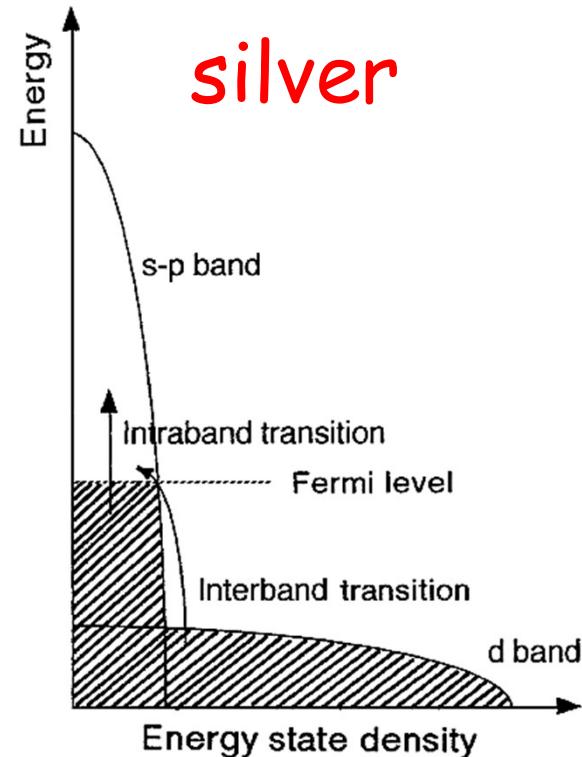
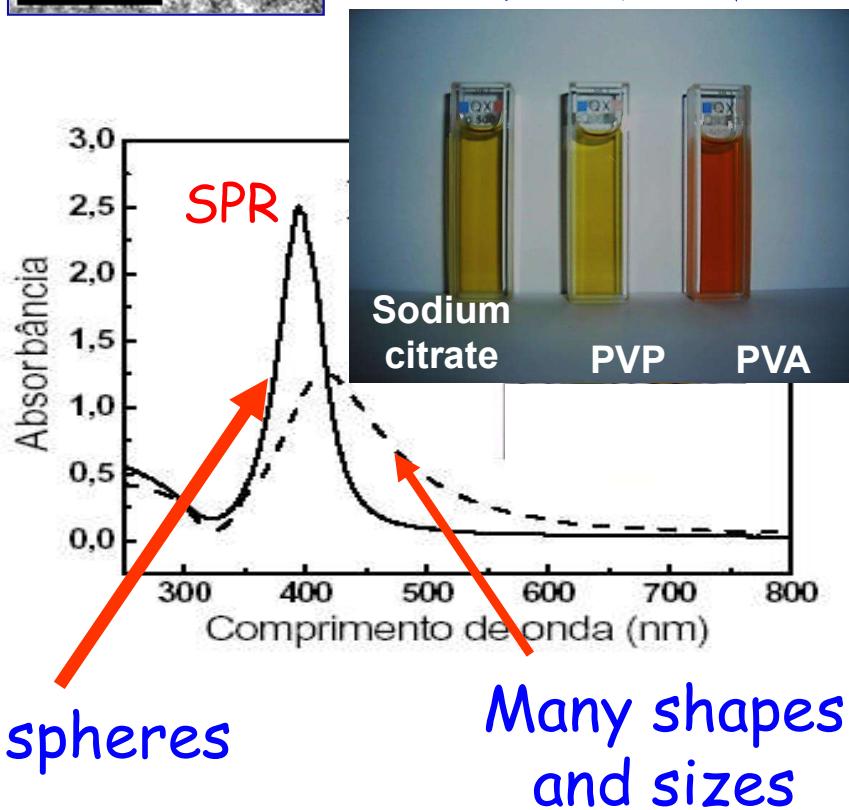
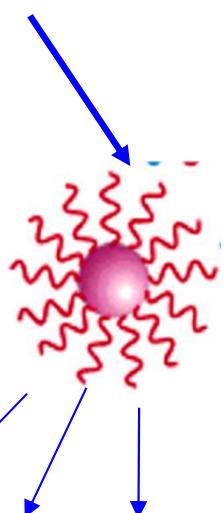
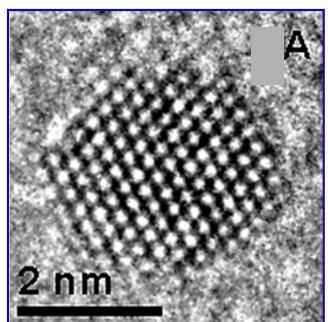
Fractals and "hot-spots"



Optical response increase much when  $\omega \approx \omega_{SP}$

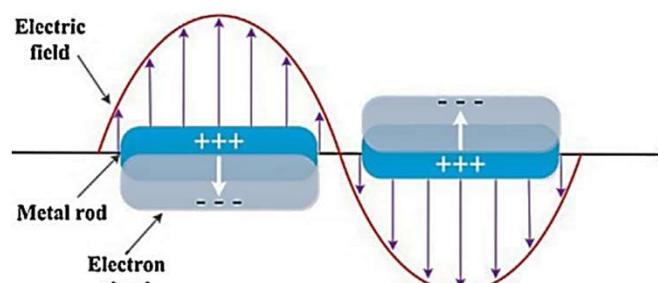
# Colloids with Ag-NPs (spheres)

Stabilizing agents to prevent aggregation

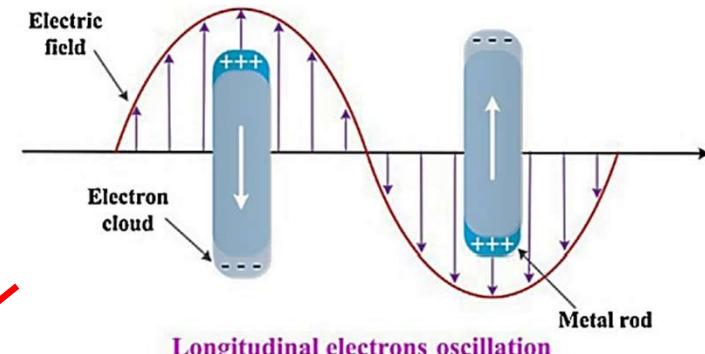


diameter: 4 nm  
≈ 1500 atoms  
≈ 30% on surface

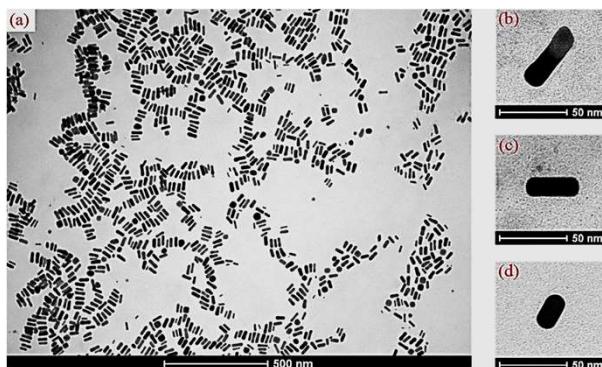
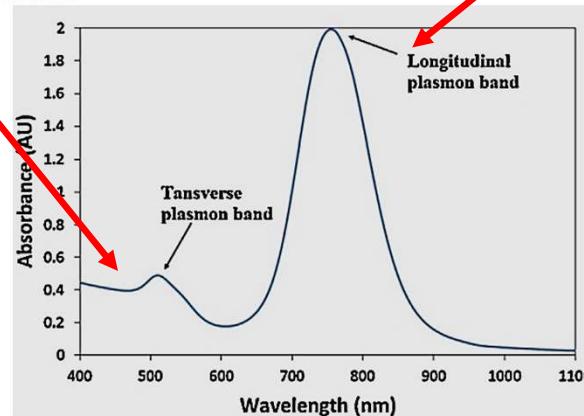
# Gold nanorods



Transverse electrons oscillation

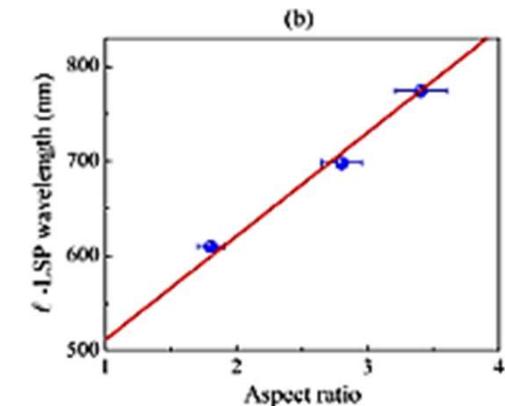
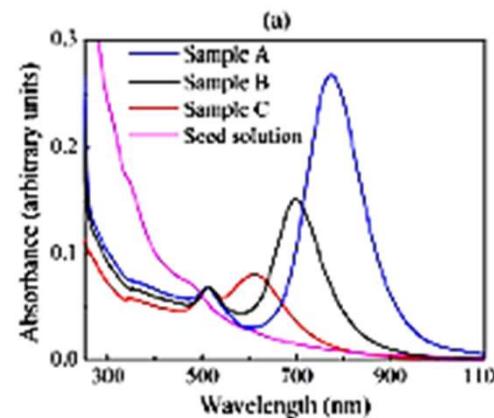


Longitudinal electrons oscillation



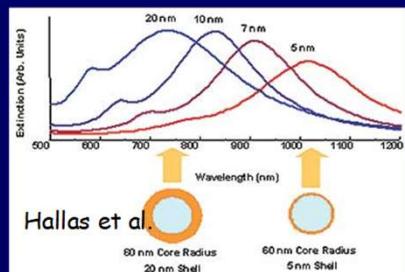
Aspect  
ratio

- A: 3.4  
B: 2.8  
C: 1.8

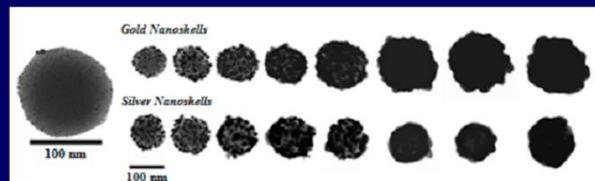


## Metallic nanoshells

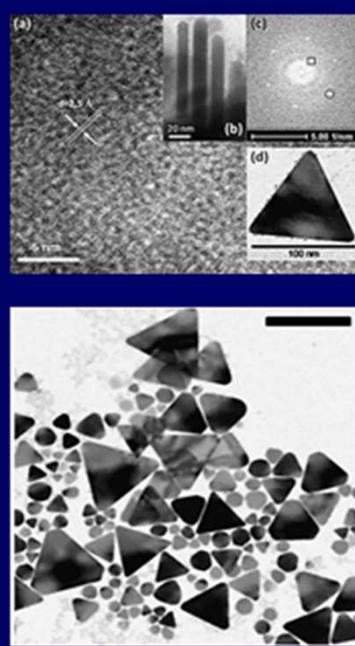
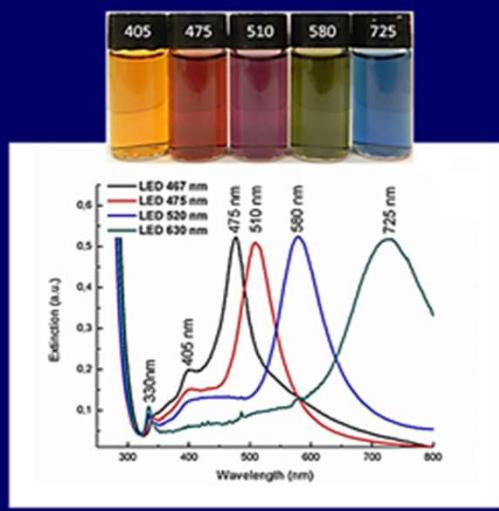
Plasmon frequency depends on the ratio between the shell thickness and the core radius



Langmuir 2013  
29, 4366-4372



## Silver Nanoprisms

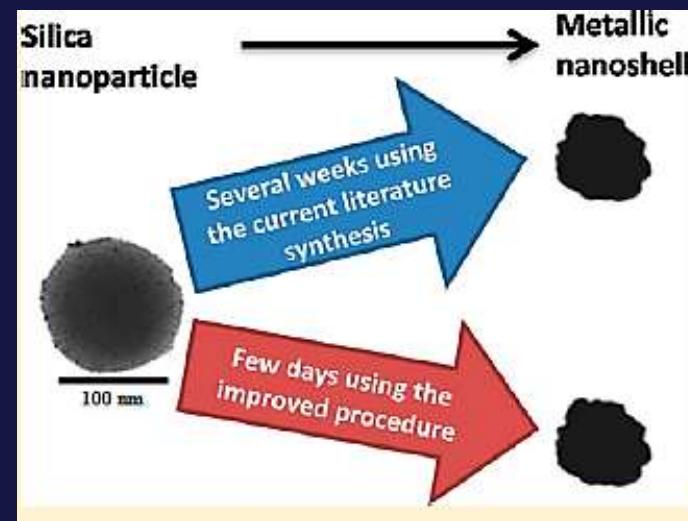


Mater. Chem. Phys. 2014 - to appear

## Improved synthesis of gold and silver nanoshells

Brito-Silva et al.

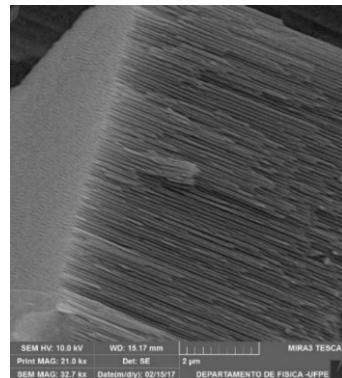
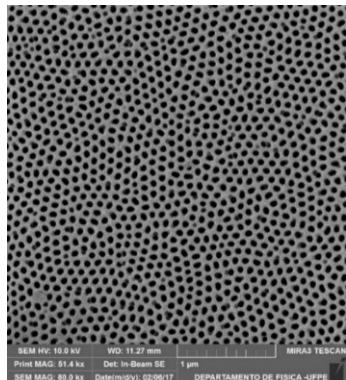
Langmuir 29 (2013) 4366



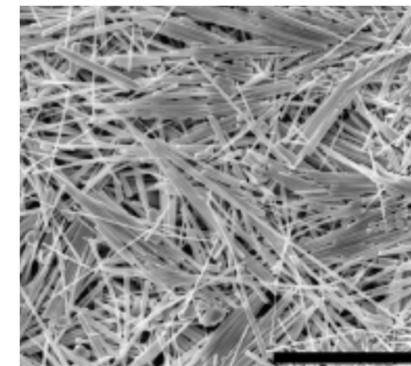
Synthesis of silver nanoprisms:  
A photochemical approach  
using light emission diodes

Saade, de Araújo  
Mater. Chem. Phys. 148 (2014) 1184

# Silver nanowires in alumina membranes electrochemistry



# Silver nanowires (chemical synthesis)

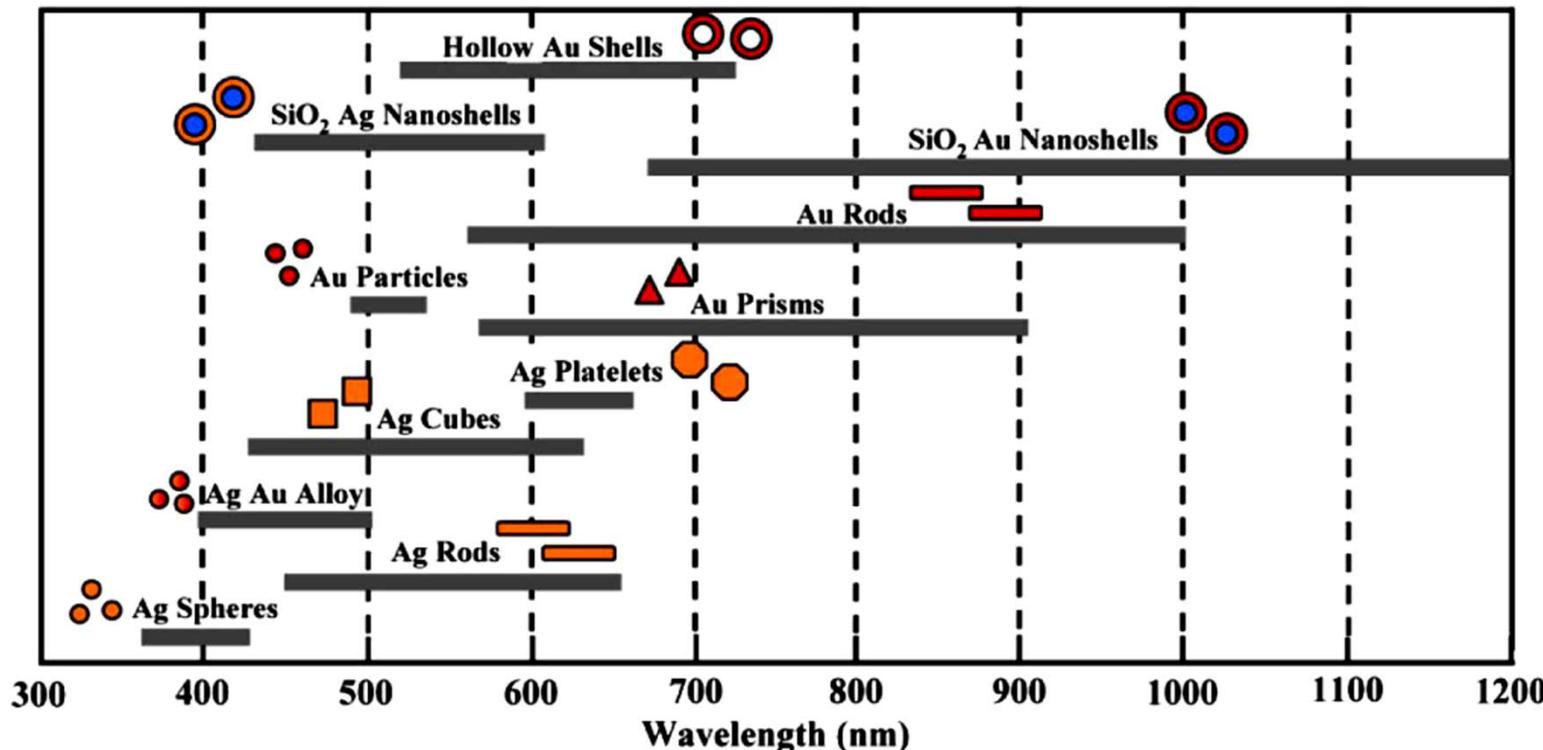


PVP-assisted reaction in EG and employing NaCl

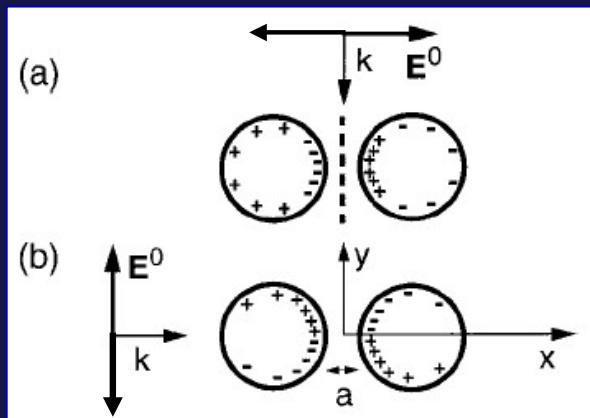
5 microns

Thesis: Manoel L Silva-Neto UFPE 2021

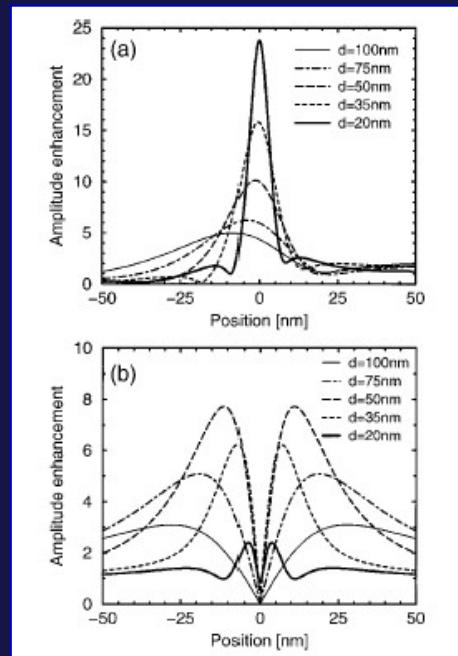
Thesis: Talita Oliveira UFPE 2021



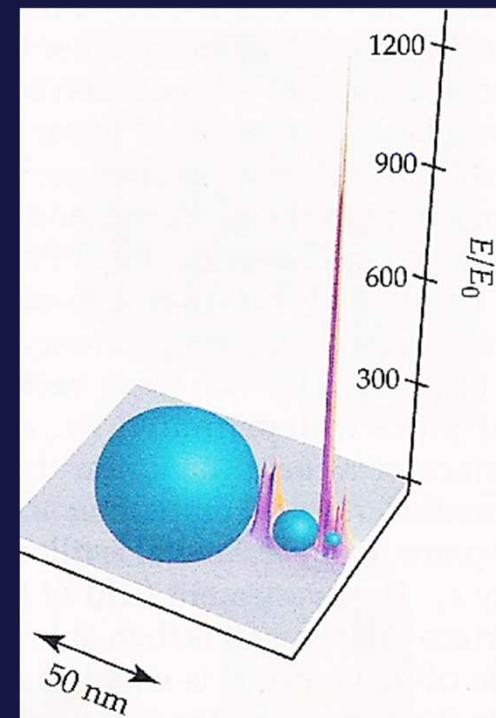
# Coupled nanoparticles



$$d/a = 5$$



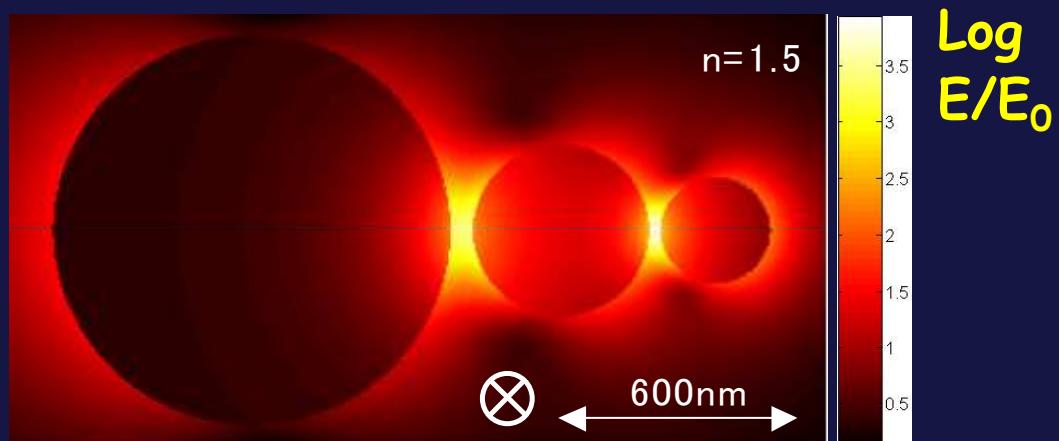
# Plasmonic nanolens



Field increases by 24 for cylinders with diameter  $d=20\text{ nm}$ .

Intensity enhancement  $> 500$

Kottmann, Martin  
Opt. Lett 26 (2001) 1096

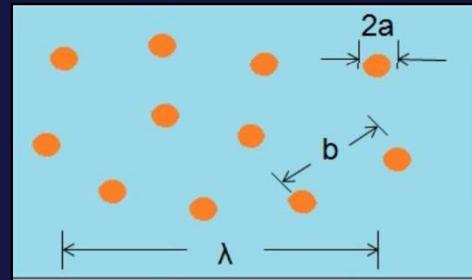


Stockman, Physics Today 2011

# Metal-dielectric nanocomposites

Nanoparticles

Dielectric



$$\lambda > b > 2a$$

- Glasses (bulk and thin films) with metallic NPs
- Polymers with metallic nanostructures
- Liquid colloids with metallic NPs

Optical response is controlled through the volume fraction of the NPs

Colloids with very small  
filling fractions

$$10^{-6} < f < 10^{-4}$$

# Why metal-dielectric nanocomposites?

Wide range of properties can be chosen to meet the needs of various applications

Control of the optical response through the density and shape of nanoparticles

- Can be doped with large amount of RE ions or transition metals
- Large variety of forms and sizes (ex: fibers and waveguides)
- Nonlinearity can be controlled by changing the NPs concentration

## Emission coupling between metal NPs and fluorophores

- SPR does not always enhance the emission properties of the fluorophore.
- Non-radiative decay rate (or quenching) depends on the distance fluorophore - metal particles.
- Emission directionality may be altered

Anger et al. PRL 96 (2006) 113002

Kuhn et al. PRL 97 (2006) 017402

# Glasses and glass-ceramics

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Kassab, Bordon, Reyna, de Araújo

Nanoparticles-based photonic metal-dielectric  
composites: A survey of recent results

Optical Materials X 12 (2021) 100098

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de Araújo, Kassab, da Silva

Optical properties of glasses and glass-ceramics for  
optical amplifiers, photovoltaic devices, color displays,  
optical limiters, and Random Lasers

Optical Materials 131 (2022) 112648

# Nonlinear optics of a nanocomposite

$$P_L + P_{NL} = \epsilon_0 \sum_{N=0}^{\infty} \chi_{eff}^{(2N+1)} E^{(2N+1)}$$

Centro-symmetric media  
 $\chi_{eff}^{(\text{even})} = 0$

Effective 3rd. order  
susceptibility

Local field  
factor

$$\chi_{eff}^{(3)} = f L^2 |L|^2 \chi_{np}^{(3)} + \chi_h^{(3)},$$

$$L = 3\epsilon_h^{(L)} / (\epsilon_{np}^{(L)} + 2\epsilon_h^{(L)})$$

Nonlinear response depends strongly on the laser frequency

$$n_2 \propto \text{Re} \chi_{eff}^{(3)}$$

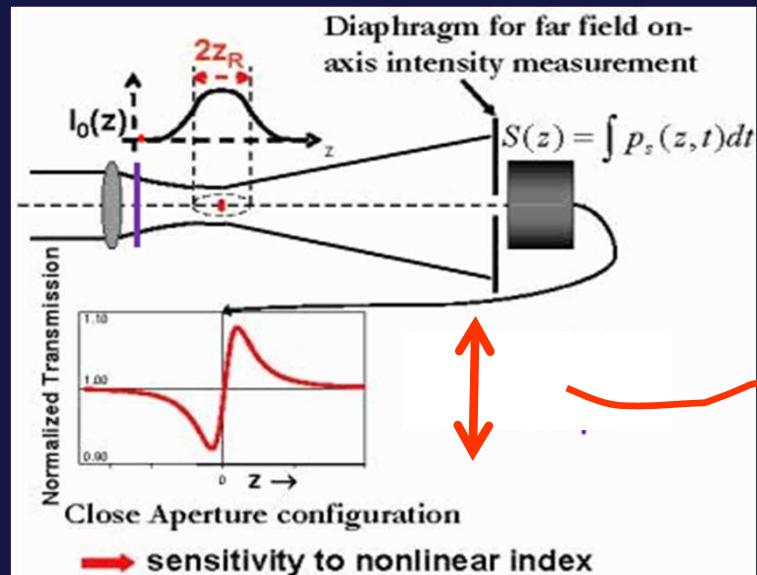
Nonlinear refraction

$$\alpha_2 \propto \text{Im} \chi_{eff}^{(3)}$$

Nonlinear absorption

# When high-order nonlinearities are present:

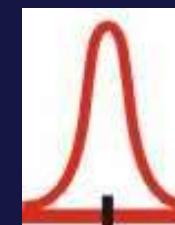
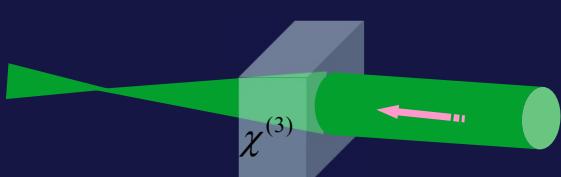
Self - focusing medium



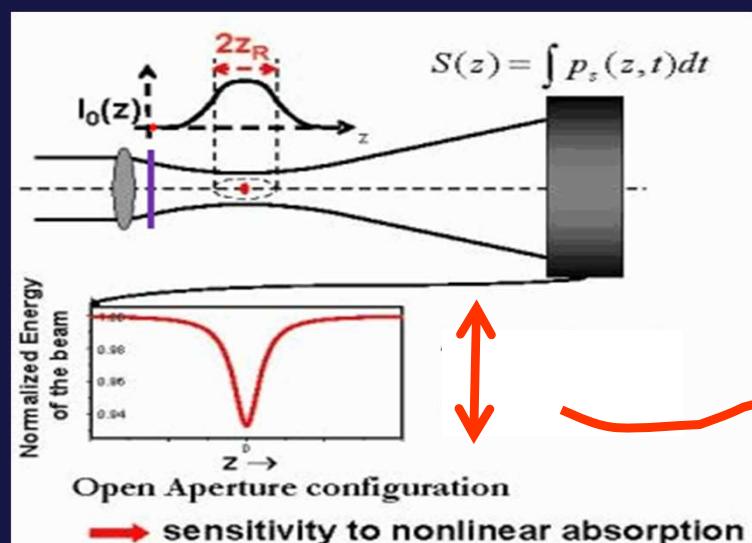
NL refraction

"Closed-aperture" Z scan

$$\Delta T \propto n_2 I + n_4 I^2 + n_6 I^3 + \dots$$



$I(r)$



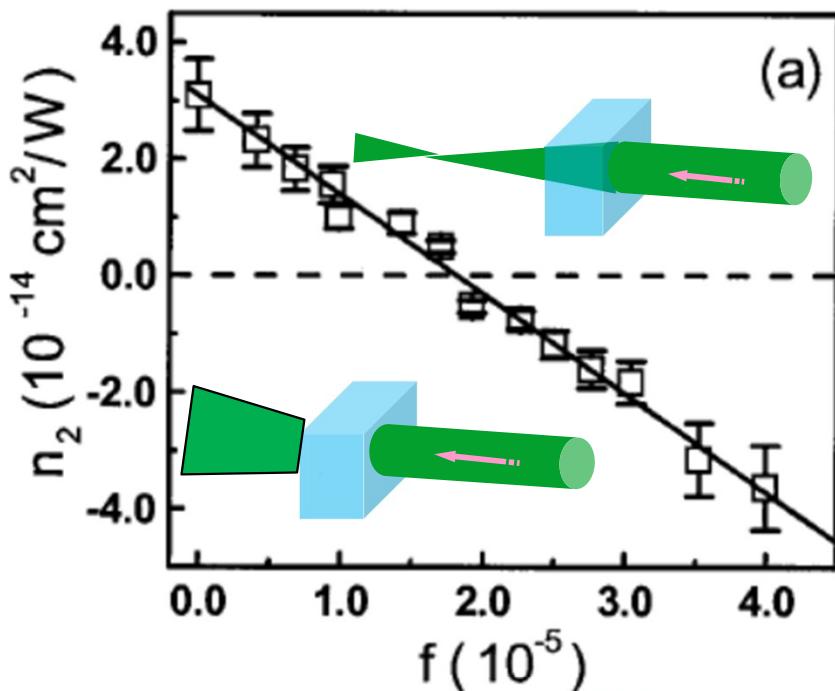
NL absorption

"Open-aperture" Z scan

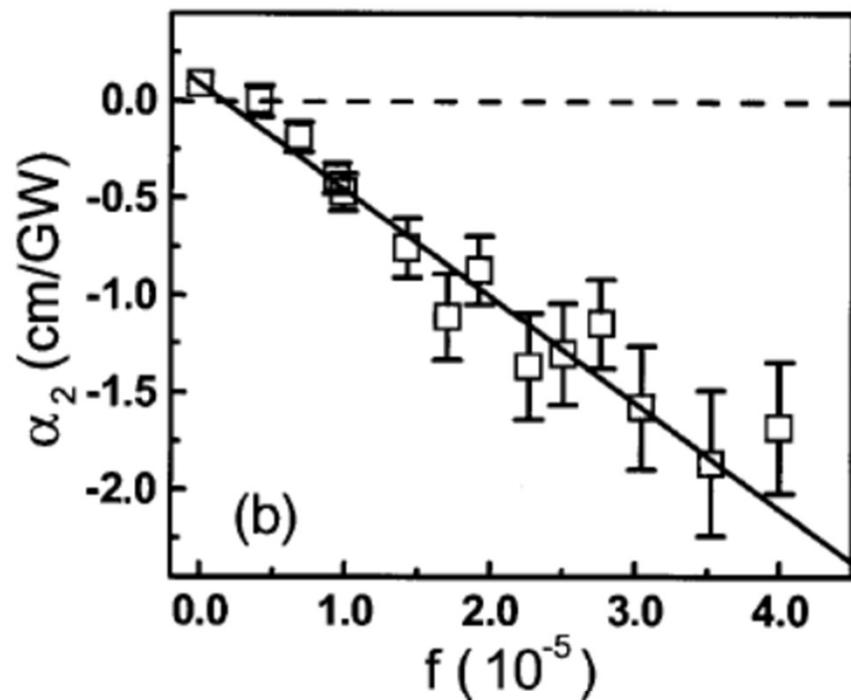
$$\Delta T \propto \alpha_2 I + \alpha_4 I^2 + \alpha_6 I^3 + \dots$$

# Ag-NPs in $CS_2$

532 nm; 80 ps; 7 Hz; 1 GW/cm<sup>2</sup>



$$n_2 \propto Re \chi_{eff}^{(3)}$$



$$\alpha_2 \propto Im \chi_{eff}^{(3)}$$

$$\chi_{eff}^{(3)} = \chi_h^{(3)} + f L^2 |L|^2 \chi_{np}^{(3)}$$

$$L = 3\epsilon_h / (\epsilon_{np} + 2\epsilon_h)$$

$$\chi_{np}^{(3)} = -(6.3 - i1.9) \times 10^{-16} (m/V)^2$$

# Large optical nonlinearity and fast response

Fast response is due to the  
induced dipole relaxation

Surface plasmon optical dephasing,  $T_2$

Measured using the "Persistent Hole-Burning Technique"

$$T_2 < 3fs$$

$T_2$  is influenced by the environment

Falcão-Filho, de Araújo, Rodrigues, Jr.

High-order nonlinearities of aqueous colloids  
containing silver nanoparticles

J. Opt. Soc. Am. B 24 (2007) 2948-2956

$\chi^{(3)}, \chi^{(5)}, \chi^{(7)}, \chi^{(9)}$

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Reyna, de Araújo

Nonlinearity management of photonic composites and  
observation of spatial-modulation instability due to quintic  
nonlinearity

Phys. Rev. A 89 (2014) 063803

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Changing the NPs concentration we may control  
the NL response of the composites

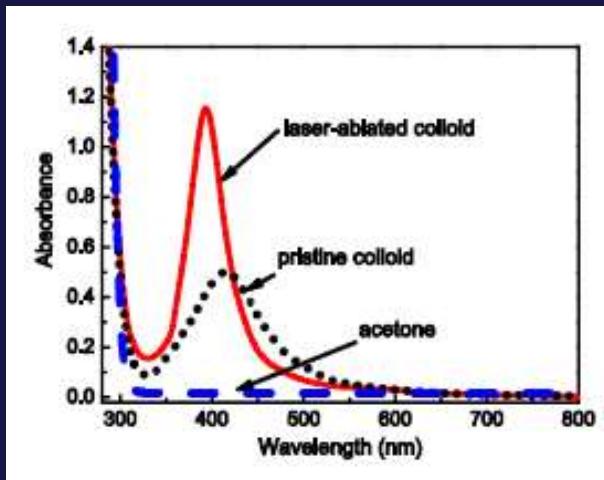
$\chi_{eff}^{(3)}$

$\chi_{eff}^{(5)}$

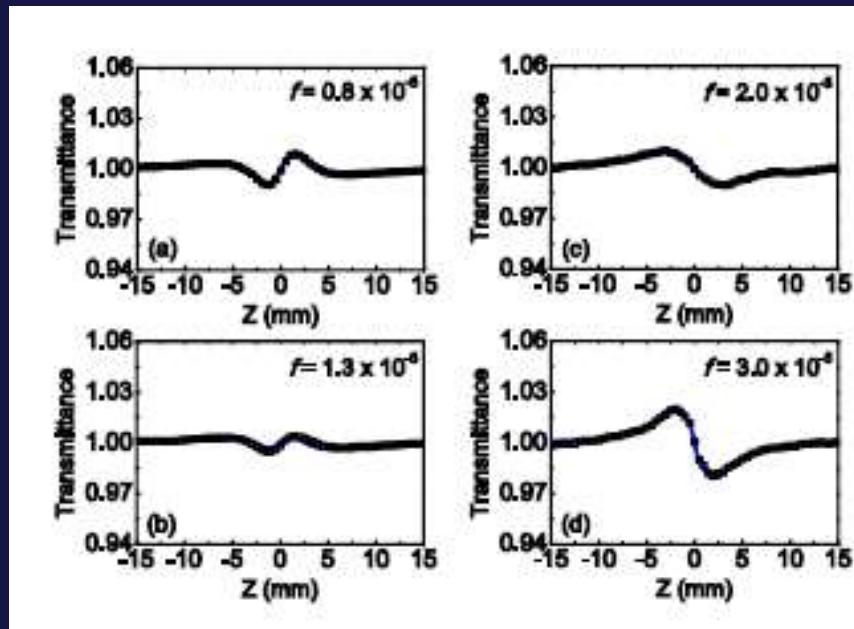
$\chi_{eff}^{(7)}$  ... ....

# Observation of fifth-order refraction in a metal-colloid

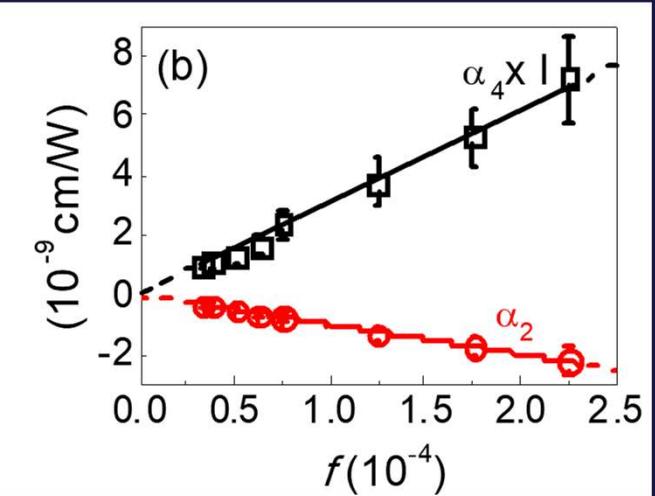
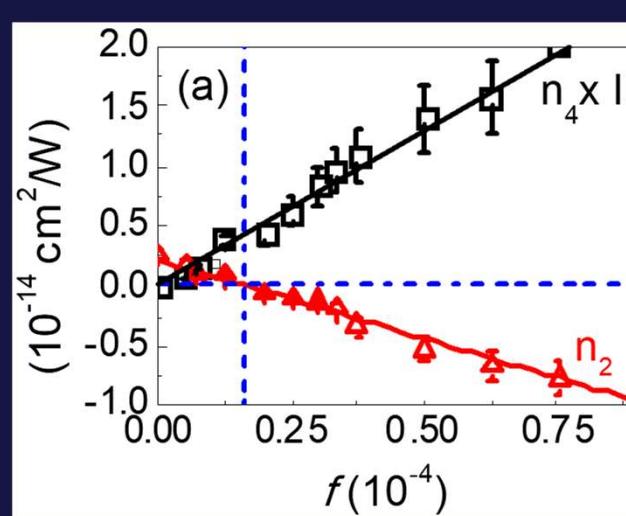
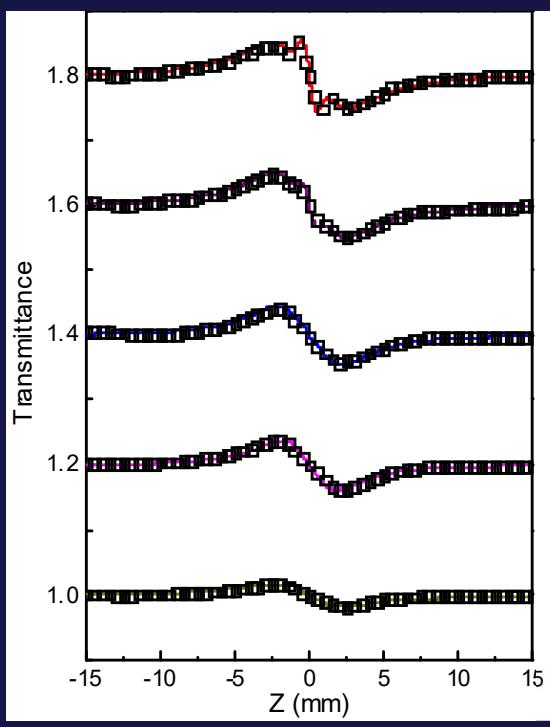
Silver NPs in acetone



$$f = 5.0 \times 10^{-5}$$

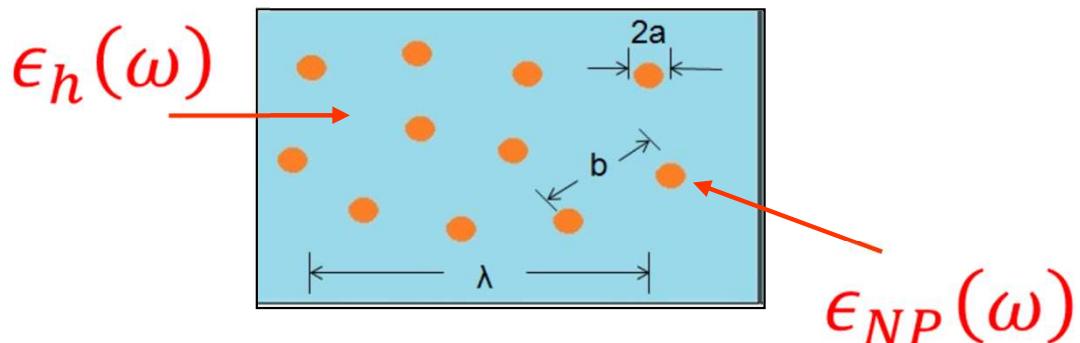


Z-scan  
532 nm  
Single pulses  
 $5 \text{ GW/cm}^2$



$9 \text{ GW/cm}^2$

# Generalized Maxwell-Garnet model



Ag NPs in acetone  
 $\chi_{host}^{(5)}$  - negligible

$a < b < \lambda$   
 Maxwell-Garnet model

$$\chi_{eff}^{(3)} = f L^2 |L|^2 \chi_{np}^{(3)} + \chi_h^{(3)}, \quad L = 3\epsilon_h^{(L)} / (\epsilon_{np}^{(L)} + 2\epsilon_h^{(L)})$$

$$\chi_{eff}^{(5)} = f L^2 |L|^4 \chi_{np}^{(5)} - \frac{6}{10} f L^3 |L|^4 (\chi_{np}^{(3)})^2 - \frac{3}{10} f L |L|^6 |\chi_{np}^{(3)}|^2,$$

$$\begin{aligned} \chi_{eff}^{(7)} = & f L^2 |L|^5 \chi_{np}^{(7)} + \frac{12}{35} f L^4 |L|^6 (\chi_{np}^{(3)})^3 + \frac{3}{35} f |L|^8 \left[ 4L^2 \chi_{np}^{(3)} + |L|^2 (\chi_{np}^{(3)})^* \right] |\chi_{np}^{(3)}|^2 \\ & - \frac{4}{7} f L |L|^6 \left[ 2L^2 \chi_{np}^{(3)} + |L|^2 (\chi_{np}^{(3)})^* \right] \chi_{np}^{(5)}. \end{aligned}$$

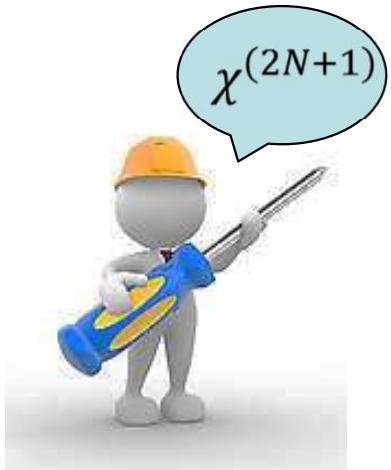
# Nonlinearity Management

For nanocomposites it is possible to vary the material's composition and incident light intensity in order to manage its nonlinearity in such way that the material may present HON on demand.

It means that we may, for instance, to adjust the material's parameters to suppress one particular NL susceptibility and enhance the others

For example, one may suppress the third-order refractive index such that the refraction can be dominate by the fifth-order susceptibility

# Nonlinearity Management

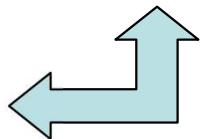


A procedure to obtain exotic metal-dielectric composites

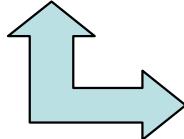
It is possible to suppress one specific order of nonlinearity and enhance another one

Example:  $n_2=0$  and  $n_4 \neq 0$

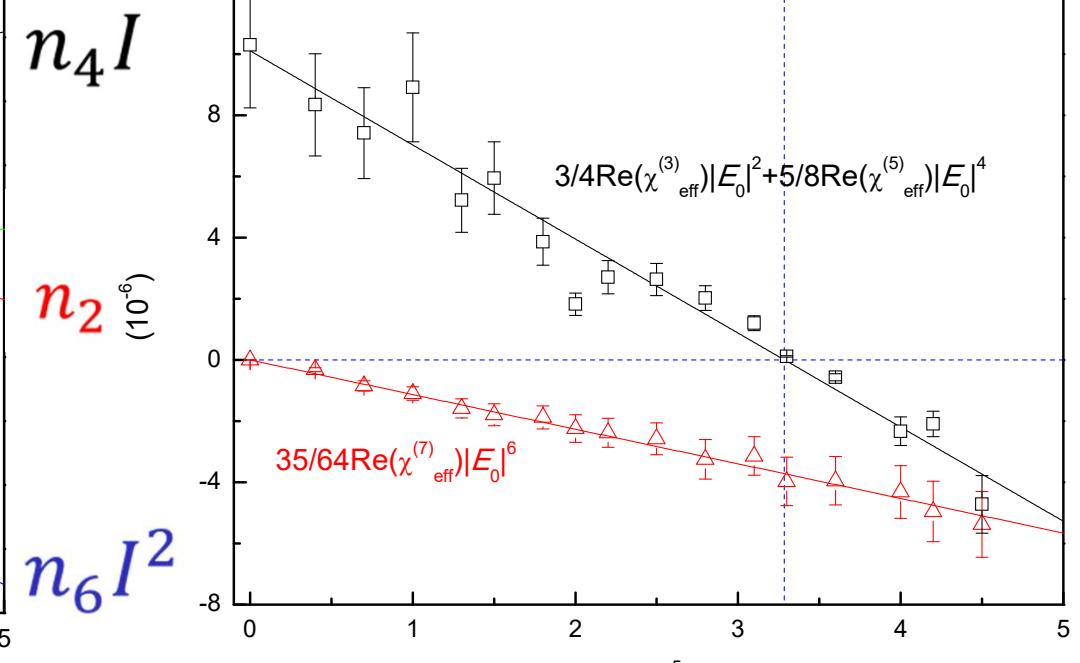
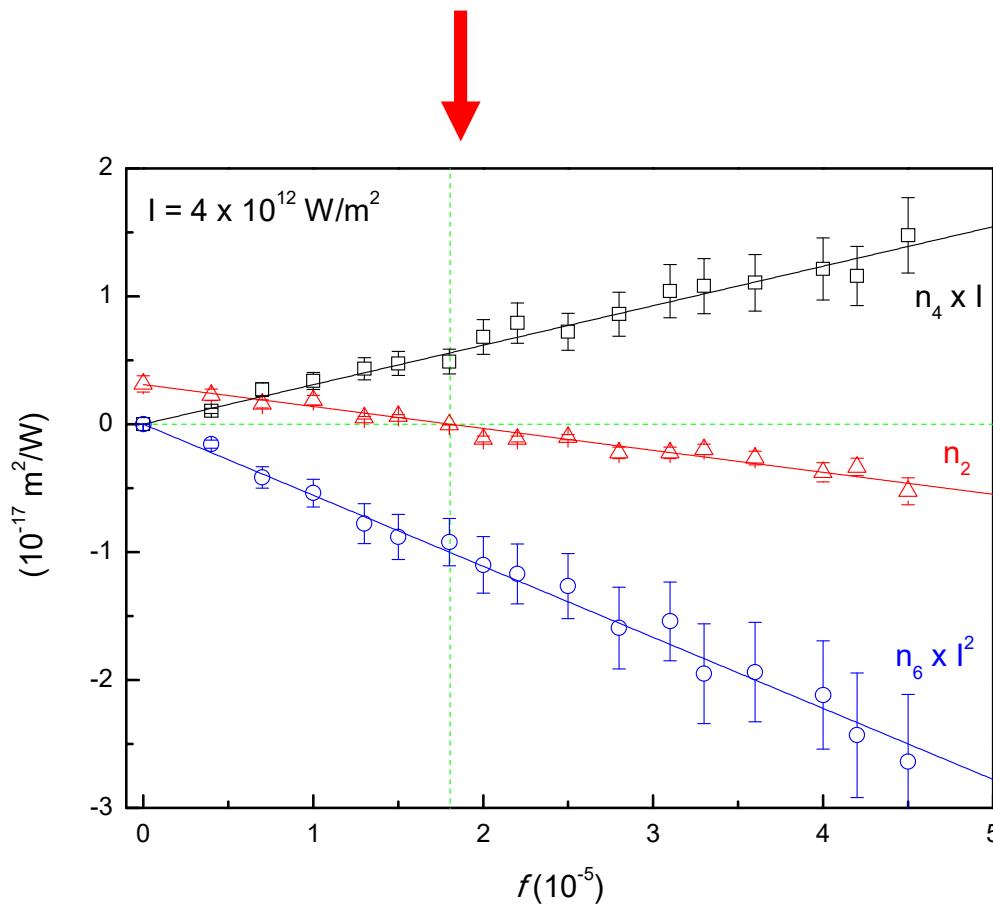
$$n_2 \propto \operatorname{Re} \chi_{eff}^{(3)}$$



$$n_4 \propto \operatorname{Re} \chi_{eff}^{(5)}$$



# Nonlinearity management: Silver NPs + $CS_2$



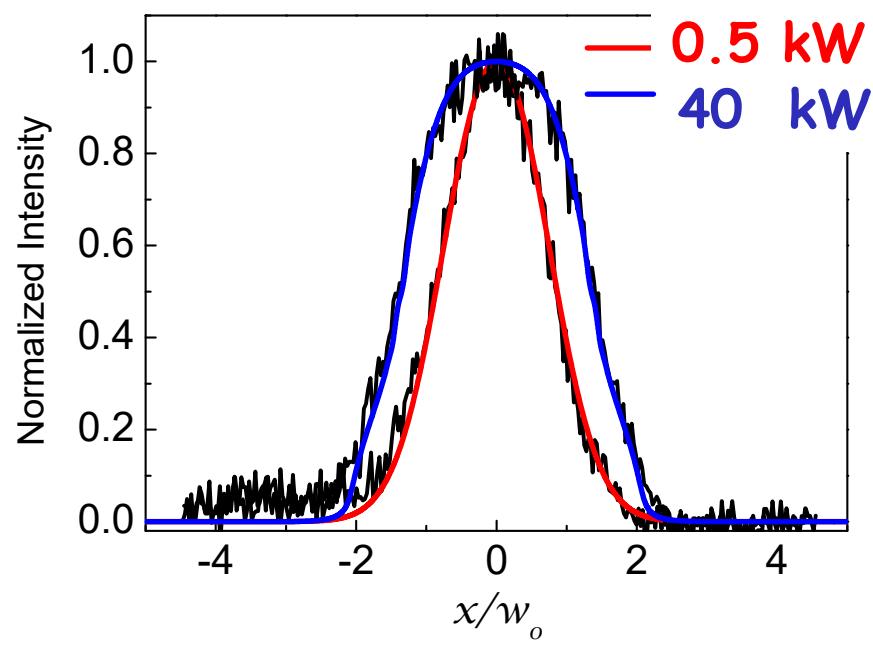
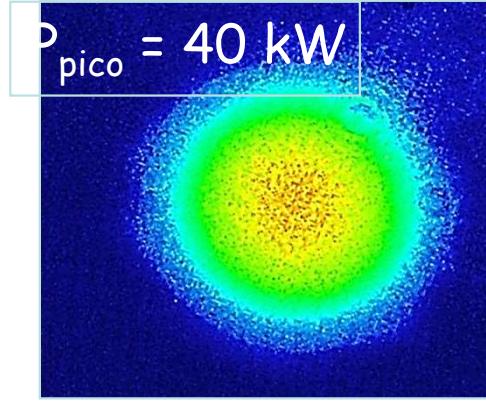
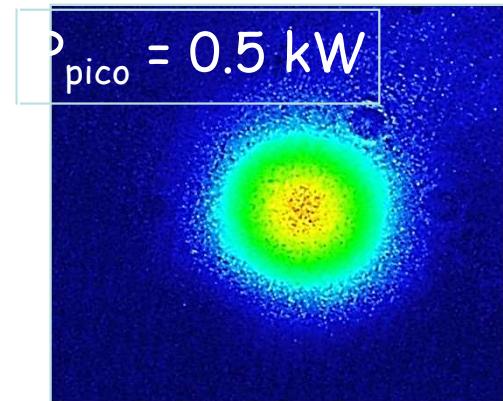
$4 \times 10^8 \text{ W/cm}^2$

Effective  $\chi_{\text{eff}}^{(7)}$

$1 \times 10^8 \text{ W/cm}^2$

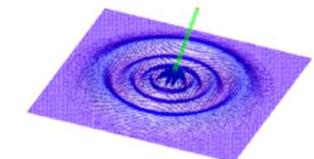
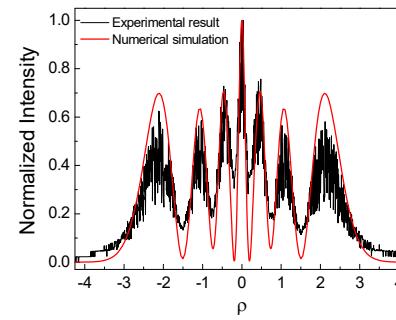
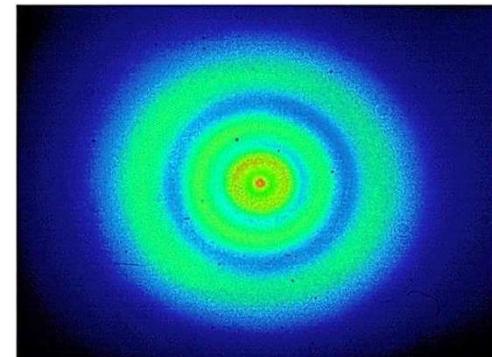
# Self defocusing due to $\chi^{(7)}$

$$f = 3.3 \times 10^{-5}$$



Input plane out of  
the lens focus

$70 \text{ GW/cm}^2$



# Cross-phase modulation with two counter - propagating beams

---

$$-\frac{\partial A_1}{\partial z} - \frac{i}{2k} \left( \frac{\partial^2 A_1}{\partial x^2} + \frac{\partial^2 A_1}{\partial y^2} \right) = \frac{ikn_2}{n_0} \left( |A_1|^2 + 2|A_2|^2 \right) A_1 + \frac{ikn_4}{n_0} \left( |A_1|^4 + 6|A_1|^2 |A_2|^2 + 3|A_2|^4 \right) A_1,$$

$$\frac{\partial A_2}{\partial z} - \frac{i}{2k} \left( \frac{\partial^2 A_2}{\partial x^2} + \frac{\partial^2 A_2}{\partial y^2} \right) = \frac{ikn_2}{n_0} \left( |A_2|^2 + 2|A_1|^2 \right) A_2 + \frac{ikn_4}{n_0} \left( |A_2|^4 + 6|A_1|^2 |A_2|^2 + 3|A_1|^4 \right) A_2,$$

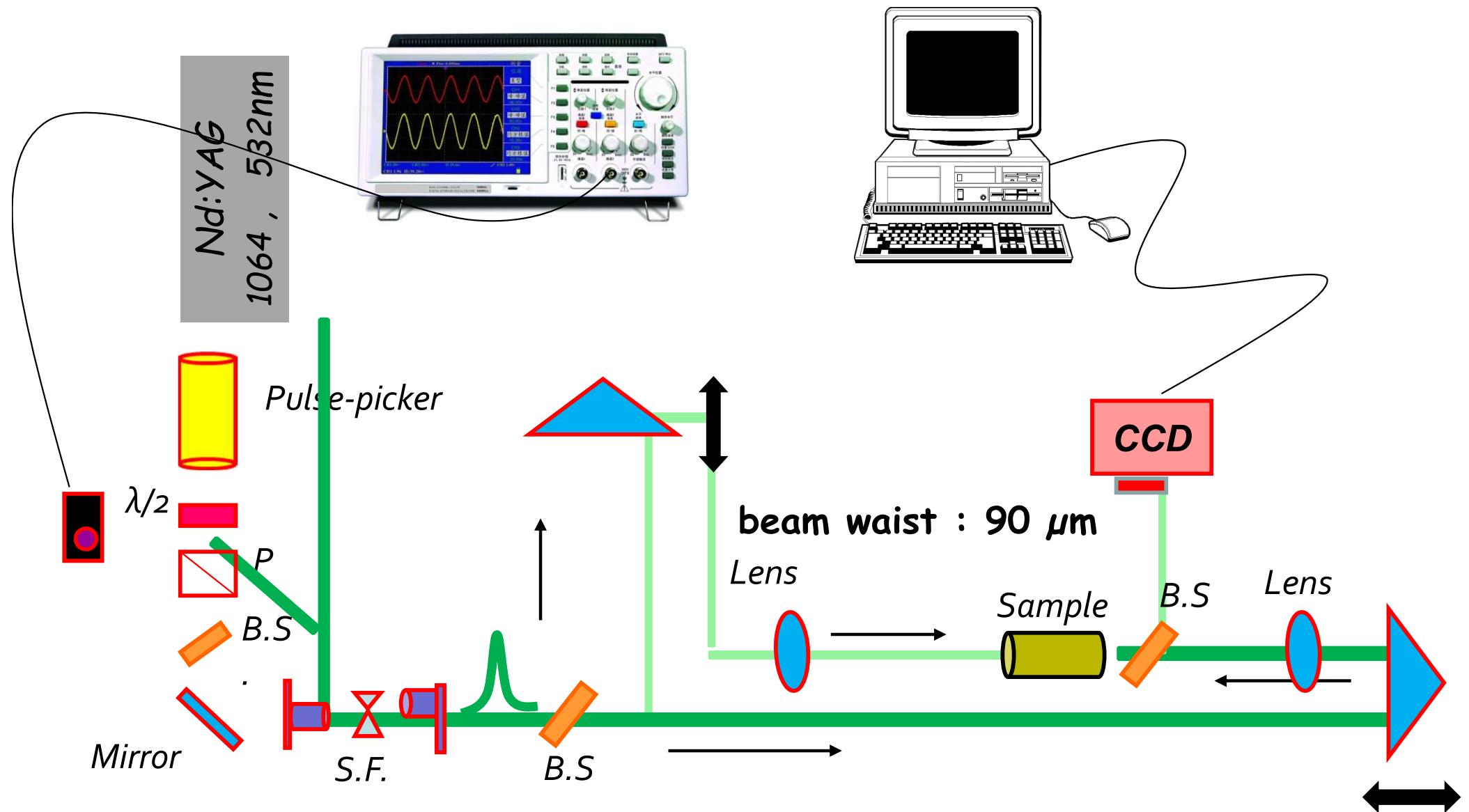


Third-order



Fifth-order

# Spatial Cross-Phase Modulation

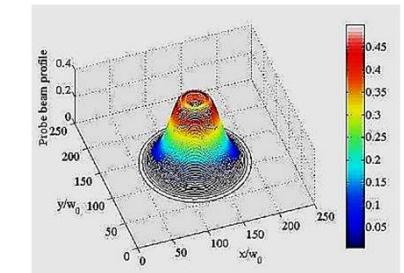
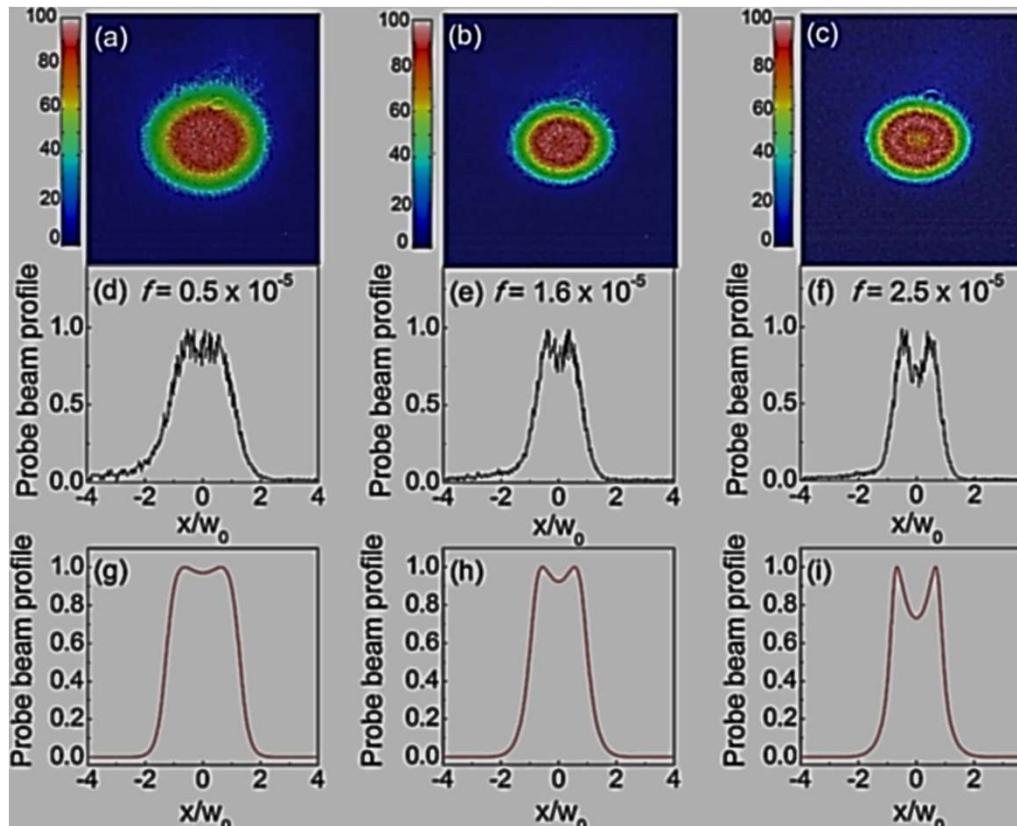
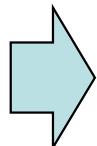


NPs: 9 nm, L=5 cm,  $I_{\text{pump}} = 2 \text{ GW/cm}^2$ ,  $I_{\text{probe}} = 0.1 I_{\text{pump}}^{30}$

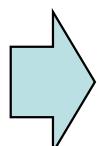
# Counter-propagating beams

Ag NPs + acetone  
First observation of Spatial Modulational Instability  
due to  $\chi_{eff}^{(5)}$

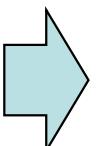
Probe beam profile



Profiles from the images



Theory



$$Re \chi_{eff}^{(3)} = 0$$

$$Re \chi_{eff}^{(5)} > 0$$

$$n_2 = 0 \quad n_4 = +3.2 \times 10^{-25} \text{ cm}^4/W^2$$

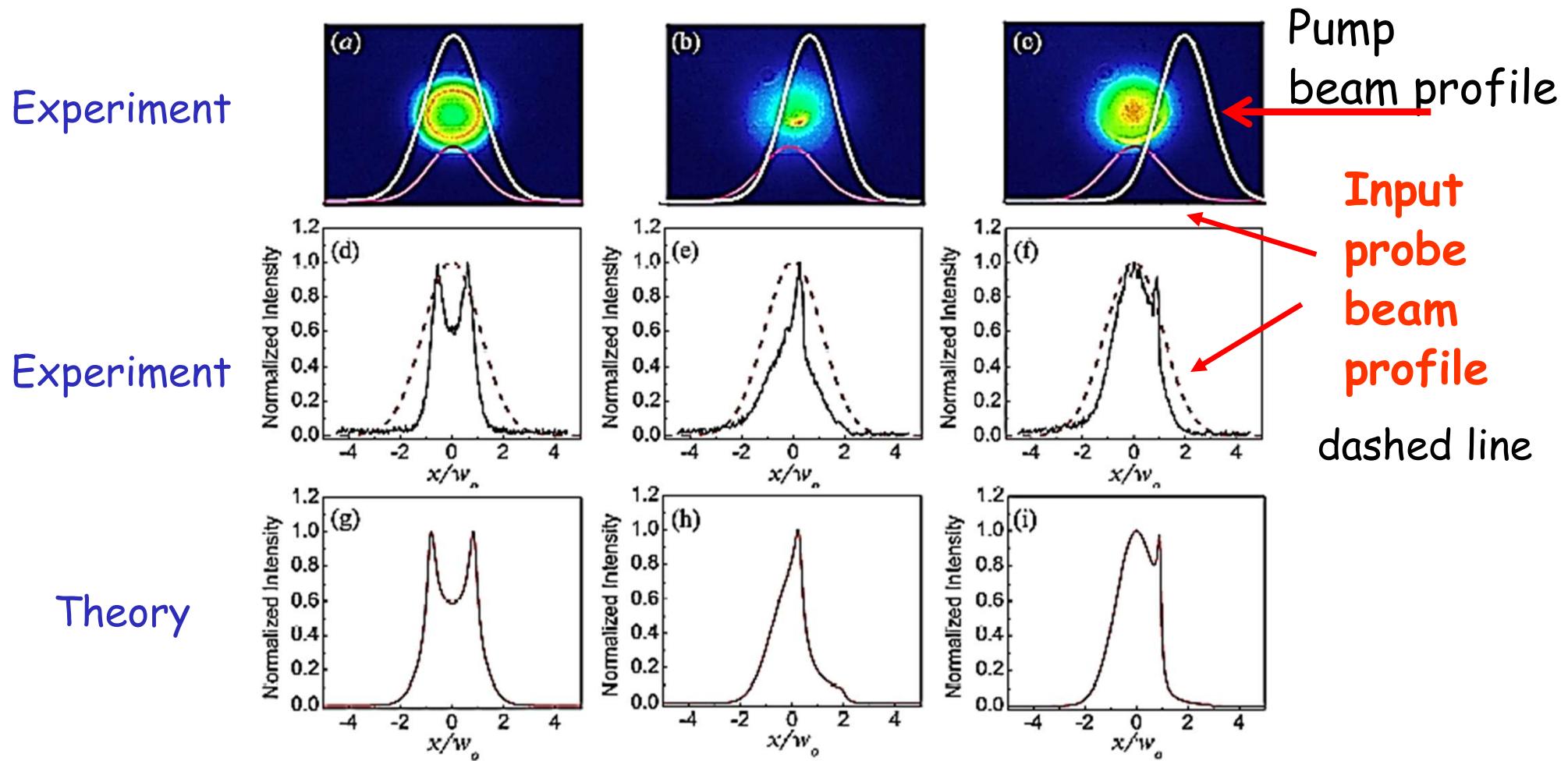
# Cross-phase modulation

## Co-propagating beams

$$\begin{aligned} -2ik \frac{\partial E_1}{\partial z} + \Delta E_1 = & -\frac{\omega^2}{c^2} \left[ 3\chi_{eff}^{(3)} (|E_1|^2 + 2|E_2|^2) E_1 \right. \\ & + 10\chi_{eff}^{(5)} (|E_1|^4 + 6|E_1|^2|E_2|^2 + 3|E_2|^4) E_1 \\ & \left. + 35\chi_{eff}^{(7)} (|E_1|^6 + 18|E_1|^2|E_2|^4 + 12|E_1|^4|E_2|^2 + 4|E_2|^6) E_1 \right] \end{aligned}$$

$$\begin{aligned} 2ik \frac{\partial E_2}{\partial z} + \Delta E_2 = & -\frac{\omega^2}{c^2} \left[ 3\chi_{eff}^{(3)} (2|E_1|^2 + |E_2|^2) E_2 \right. \\ & + 10\chi_{eff}^{(5)} (3|E_1|^4 + 6|E_1|^2|E_2|^2 + |E_2|^4) E_2 \\ & \left. + 35\chi_{eff}^{(7)} (4|E_1|^6 + 12|E_1|^2|E_2|^4 + 18|E_1|^4|E_2|^2 + |E_2|^6) E_2 \right]. \end{aligned}$$

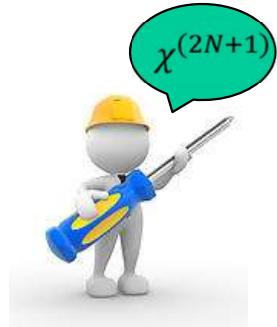
# Induced focusing due to the seventh-order susceptibility



$$n_2 = 0; n_4 = 0; n_6 < 0$$

Analogous to  $\chi^{(3)}$ : Hickmann, Gomes, de Araújo, PRL 68 (1992) 3547

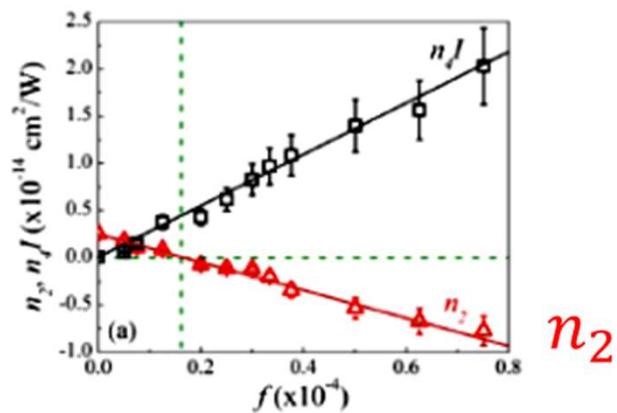
# "Nonlinearity management"



Z-scan  
Ag-NPs (<10 nm) + CS<sub>2</sub>

$\chi_{eff}^{(3)}$ ,  $\chi_{eff}^{(5)}$ ,  $\chi_{eff}^{(7)}$

Reyna, de Araújo. Opt. Express 22 (2014) 22456-22469

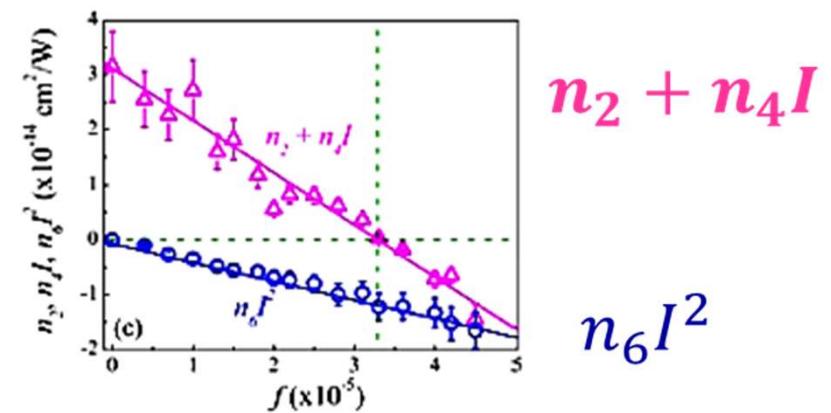


$$f = 1.6 \times 10^{-5}$$

$$4 \times 10^8 \text{ W/cm}^2$$

$$\operatorname{Re} \chi^{(3)} = 0$$

$$\operatorname{Re} \chi^{(5)} \neq 0$$



$$f = 3.3 \times 10^{-5}$$

$$1 \times 10^8 \text{ W/cm}^2$$

$$\operatorname{Re} (\chi^{(3)} + I \times \chi^{(5)}) = 0$$

$$\operatorname{Re} \chi^{(7)} \neq 0$$

# Non linear propagation of light

$$\nabla^2 \vec{E} - \frac{1}{c^2} \frac{\partial^2 \vec{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \vec{P}_L}{\partial t^2} + \mu_0 \frac{\partial^2 \vec{P}_{NL}}{\partial t^2}$$

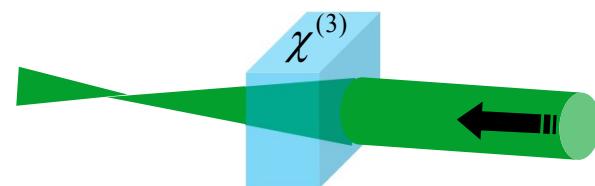
$$\vec{E}(\vec{r}, t) = \vec{A}(\vec{r}) \exp[i(kz - \omega t)]$$

$$i \frac{\partial \vec{A}}{\partial z} = -\frac{1}{2k} \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) \vec{A} - \frac{n_2}{n_0} k |\vec{A}|^2 \vec{A}$$

difração



auto-focalização

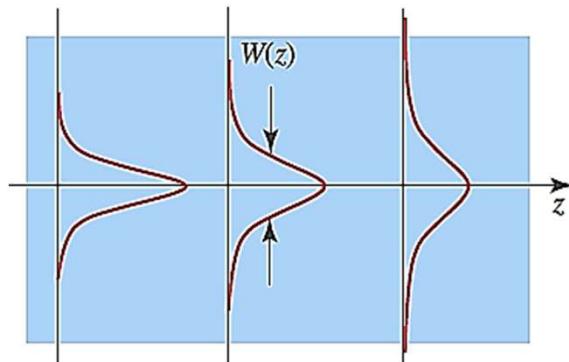


# Optical Spatial Solitons

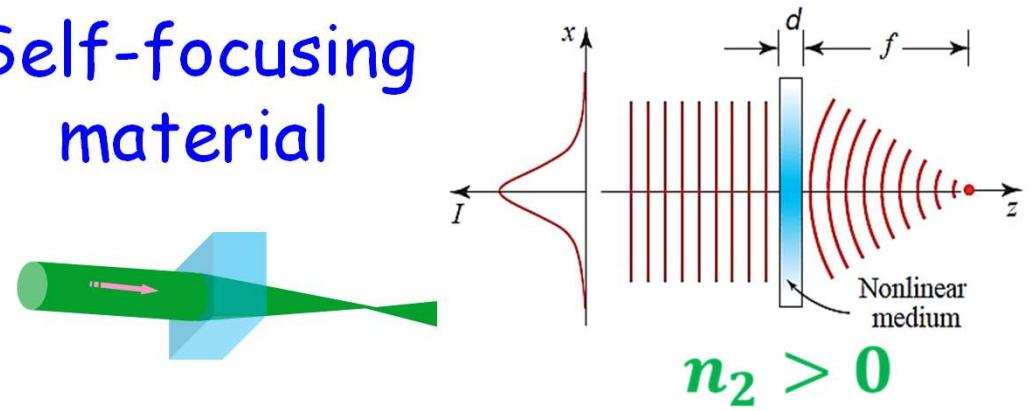
Solutions of the Maxwell's equation including nonlinear terms

## NL Schrödinger Equation

Difraction

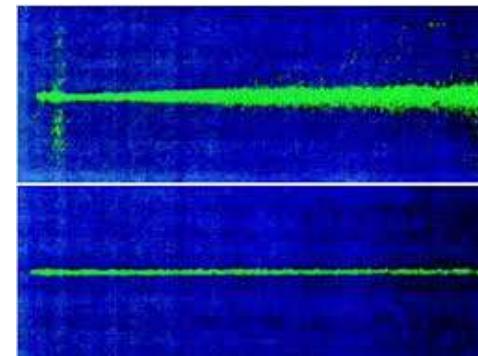
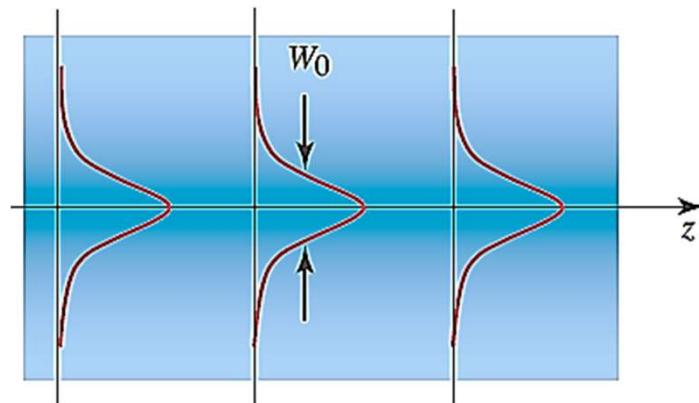


Self-focusing material



$$n = n_0 + n_2 I + \dots$$

Bright  
Spatial  
soliton

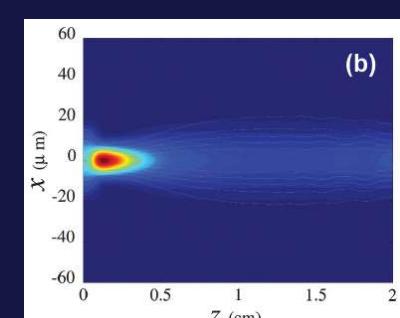
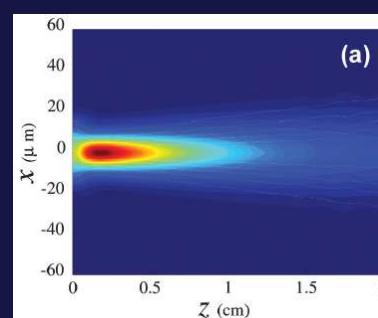
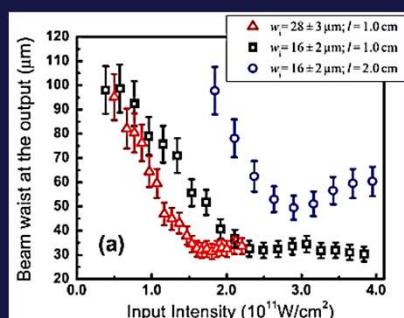
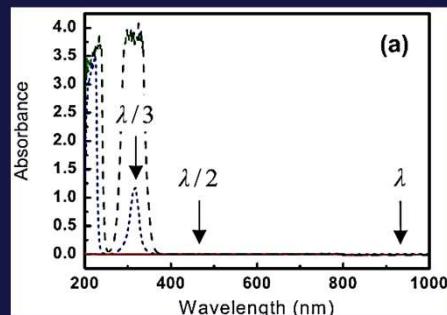


Low intensity  
High intensity  
Soliton

# 1st. lecture

## First demonstration of (2+1)D soliton propagating in a homogeneous medium with local nonlinearity

Falcão-Filho, de Araújo, Boudebs, Leblond, Skarka  
Robust two-dimensional spatial solitons in liquid carbon disulfide  
Phys. Rev. Lett. 110 (2013) 013901.



Low intensity

soliton

Very important: contributions of third and fifth order  
of opposite signs

## $CS_2$ : stable (2+1)D soliton

$$Re \chi^{(3)} > 0$$

$$Re \chi^{(5)} < 0$$

---

Is it possible to observe a stable (2+1)D soliton  
in a system with:

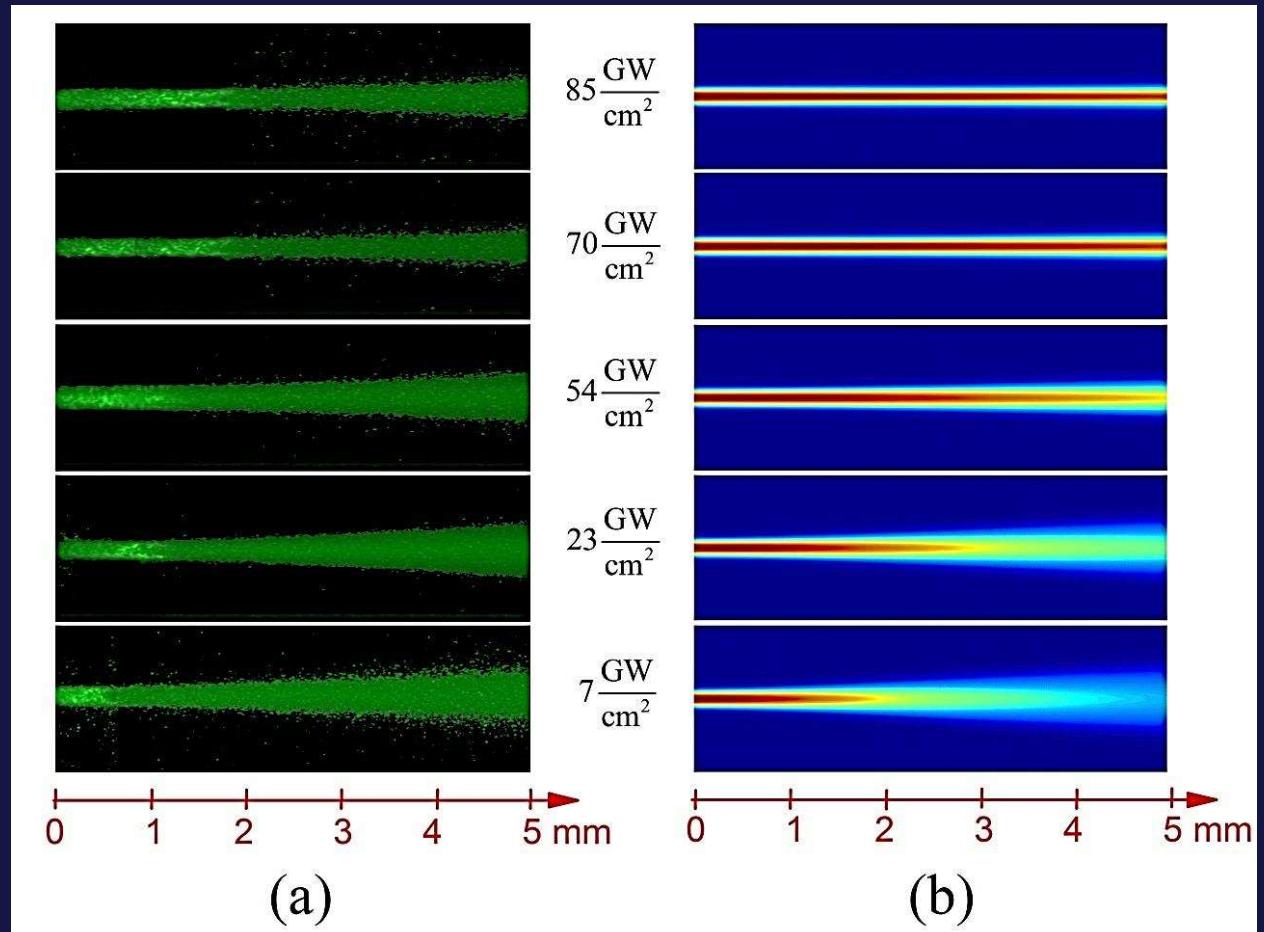
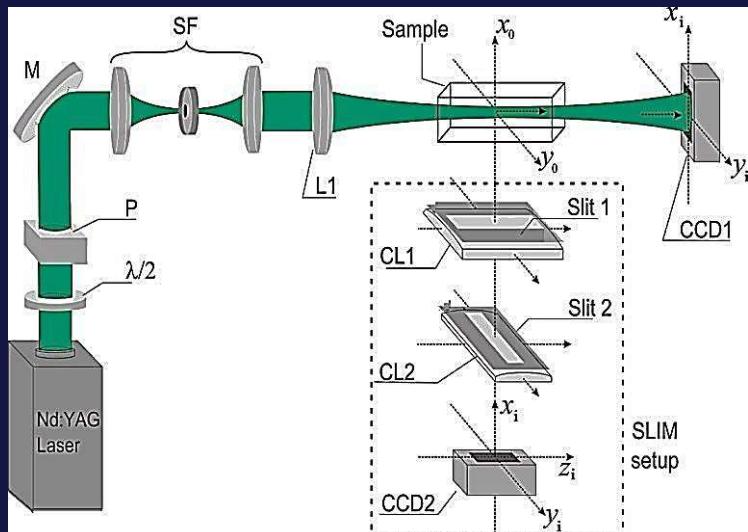
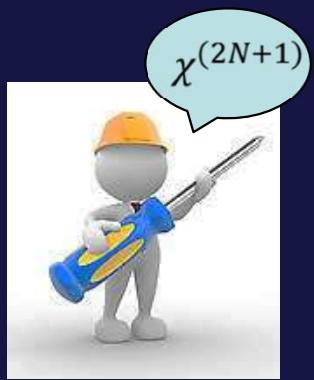
$$Re \chi^{(3)} = 0 , \quad Re \chi^{(5)} > 0 , \quad Re \chi^{(7)} < 0 \quad ?$$

$$2ik \frac{\partial E}{\partial z} + \Delta E = -\frac{\omega^2}{c^2} \left[ 3\chi_{eff}^{(3)} |E|^2 E + 10\chi_{eff}^{(5)} |E|^4 E + 35\chi_{eff}^{(7)} |E|^6 E \right]$$

# First observation of 2D Spatial-Solitons in a quintic-septimal medium

Silver NPs  
in acetone

$$n_2 = 0; \quad n_4 > 0; \quad n_6 < 0$$

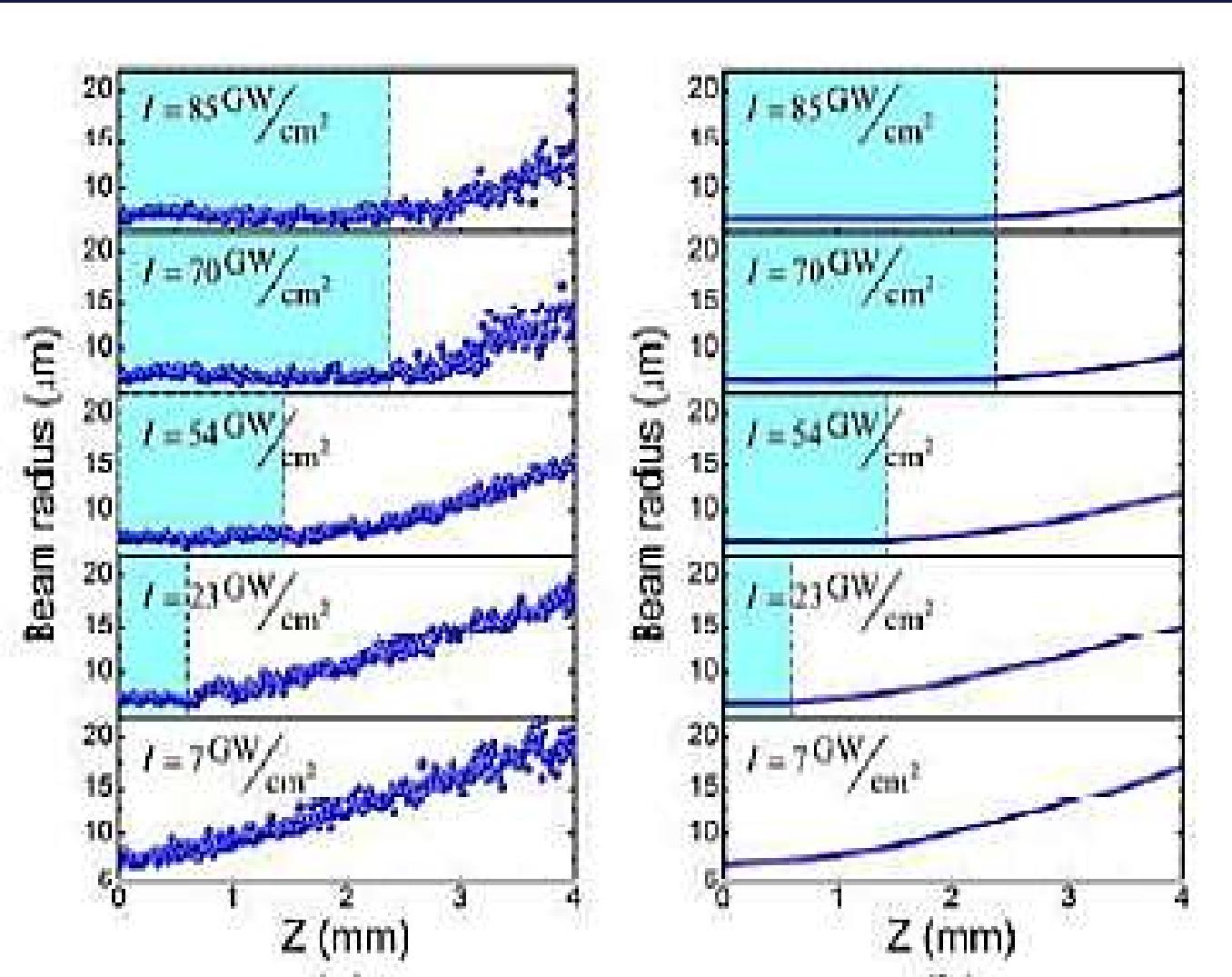


experiment

theory

Reyna, Jorge, de Araújo Phys. Rev. A 90 (2014) 063835  
de Araújo et al. J. Lumin. 169 (2016) 492-496

$$2ik \frac{\partial E}{\partial z} + \Delta E = -\frac{\omega^2}{c^2} \left[ 3\chi_{\text{eff}}^{(3)} |E|^2 E + 10\chi_{\text{eff}}^{(5)} |E|^4 E + 35\chi_{\text{eff}}^{(7)} |E|^6 E \right]$$

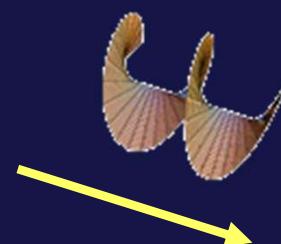


experiment

theory

# Optical vortices

- Beams with phase singularity
- Zero field in the center of the vortex
- Helical wavefront
- Phase



$$\phi(t, z, \theta) = kz + \omega t + m\theta$$

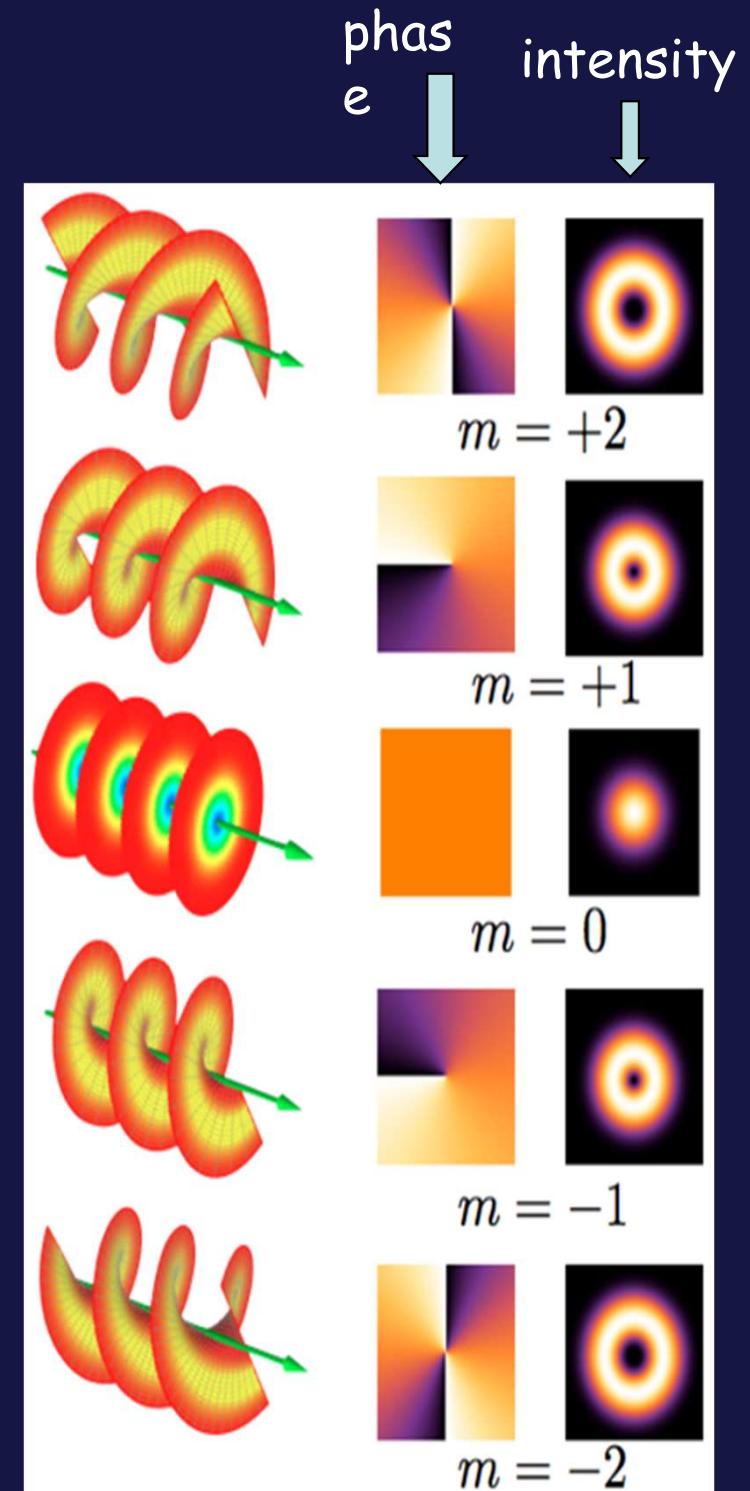
$m = 0$

Plane wave

$m \neq 0$

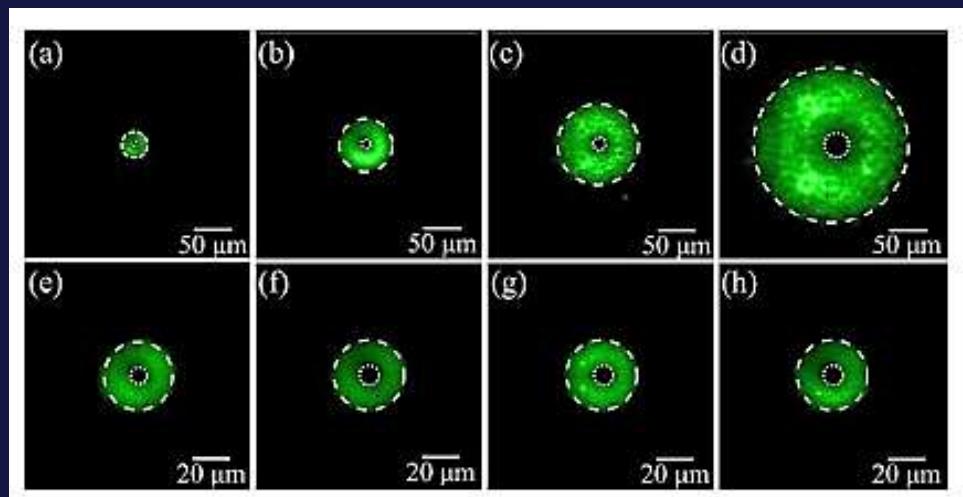
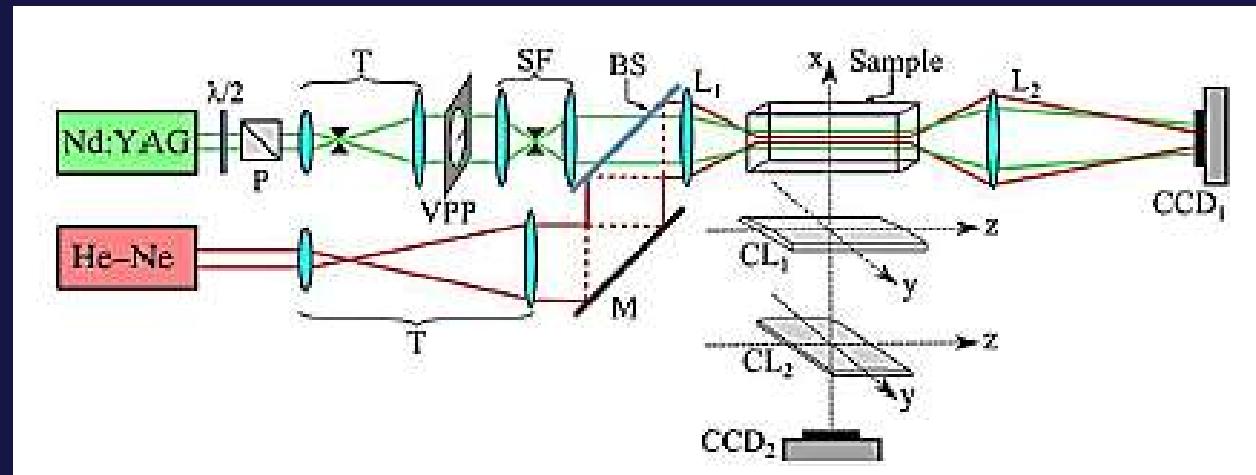
Wave with topological phase

$m$  is the "topological charge"



## Guiding and confinement of light induced by optical vortex solitons in a cubic-quintic medium

ALBERT S. REYNA\* AND CID B. DE ARAÚJO



$Z = 0$

3

5

10 mm

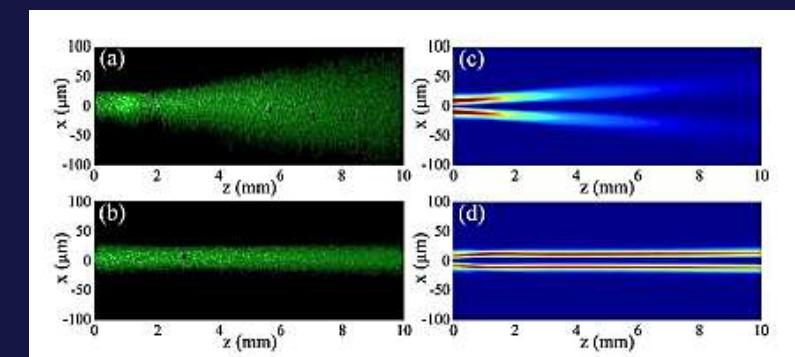
0.1

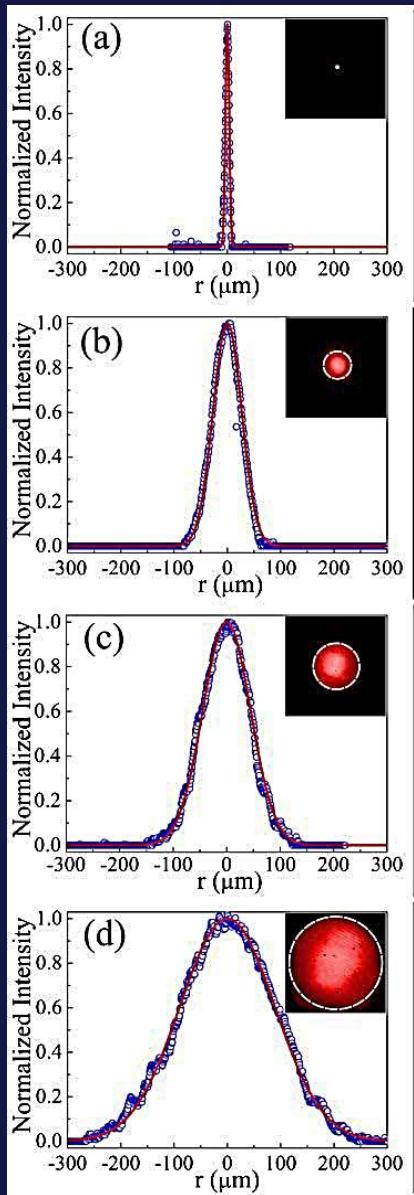
3.0

$GW/cm^2$

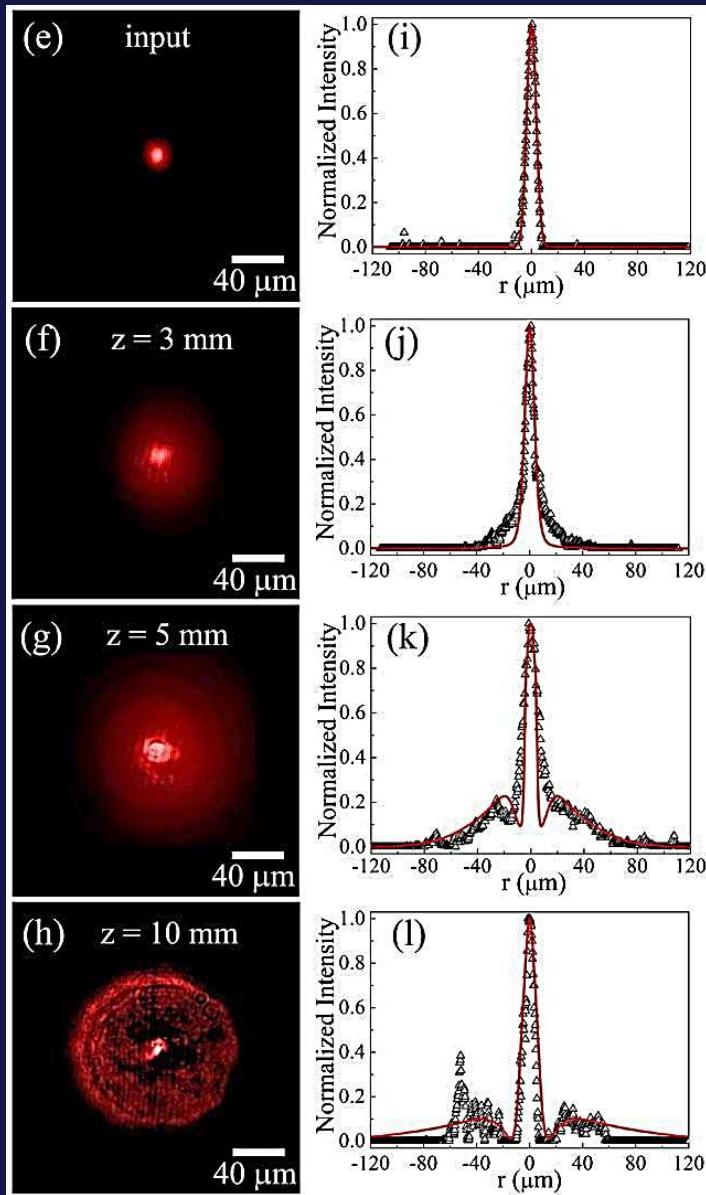
10 mm - 25  $Z_R$

42

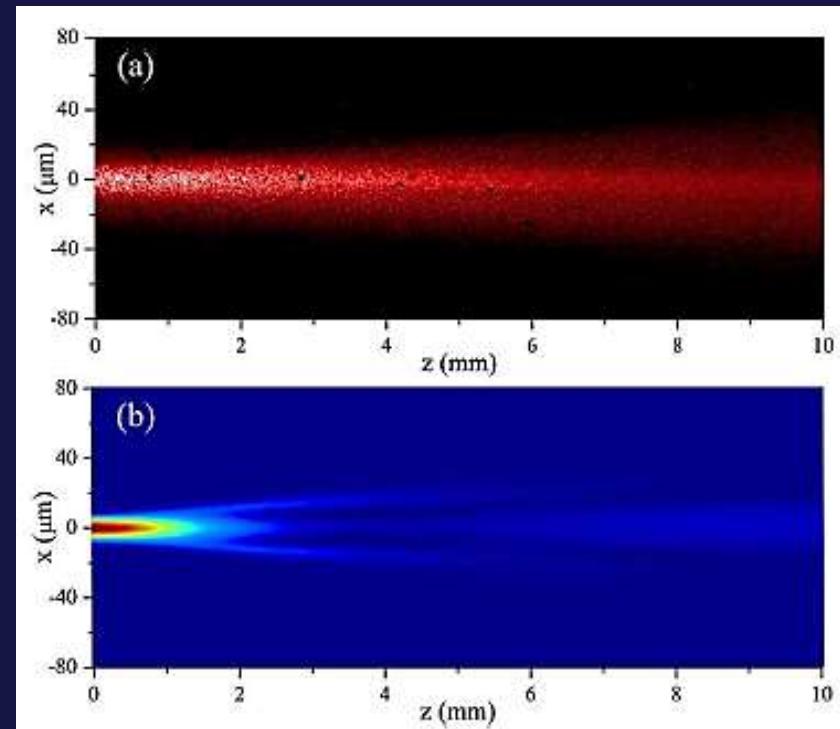




HeNe



Guided HeNe



$$I_{OVS} = 3.0 \text{ GW/cm}^2$$

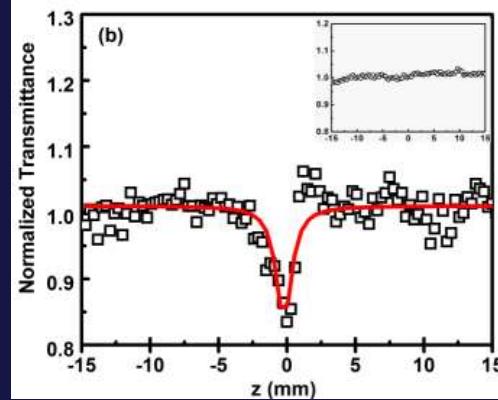
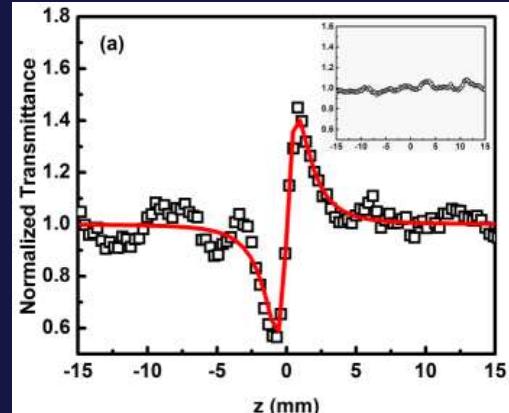
$$I_{\text{HeNe}} = 0.1 \text{ GW/cm}^2$$

How to address the long standing problem of discovering a very good material for all-optical switching?

We need a material with large NL refraction and low NL absorption

In general large NL refraction presents large NL absorption

# PbO- $\text{GeO}_2$ films with gold NPs for all-optical switching



RF sputtering

800 nm

150 fs

Figure-of-merit  
enhanced by two  
orders of magnitude

Germanate film	$n_2/\lambda\alpha_2$
As grown	$8.3 \times 10^{-4}$
With Au NPs	$>2.1 \times 10^{-1}$

# Optimization procedure for the design of all-optical switches based on metal-dielectric nanocomposites

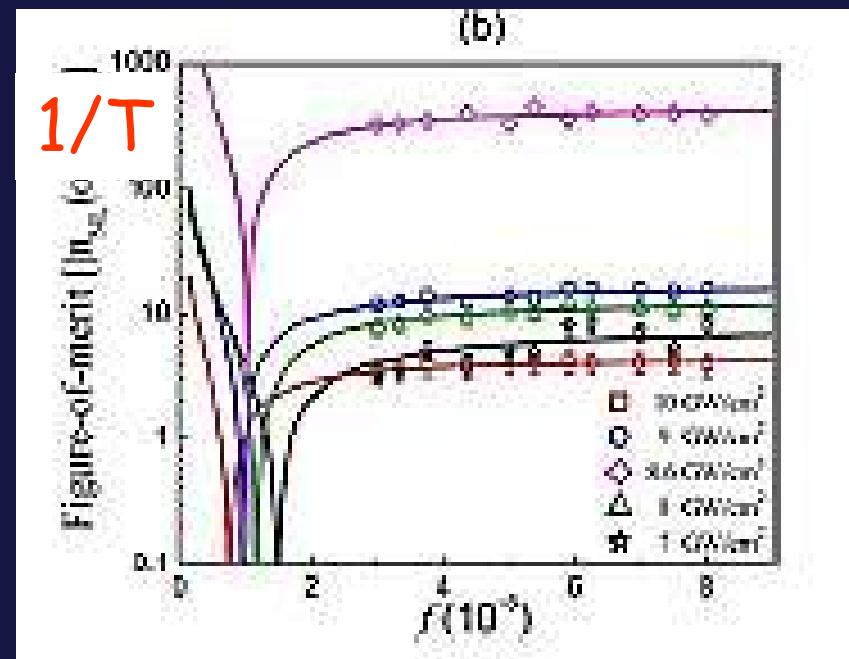
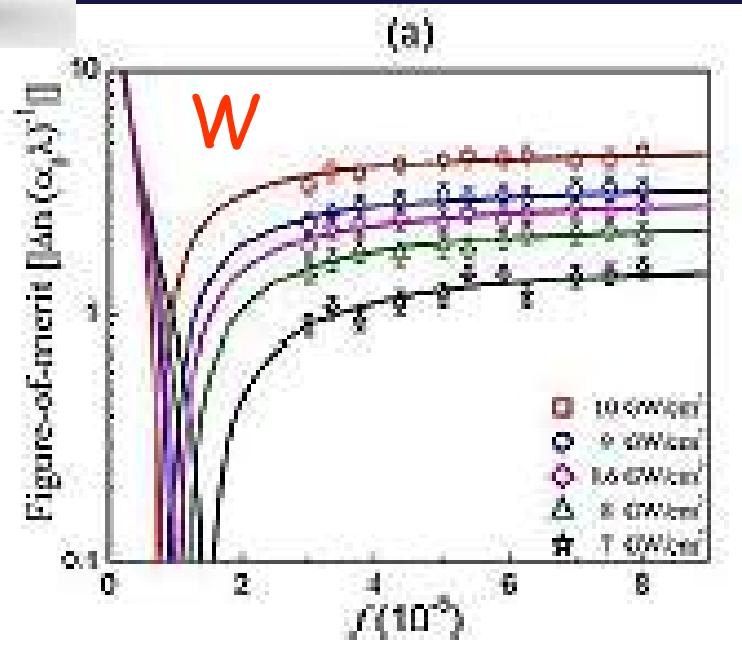
$$\Delta n = n_2 I + n_4 I^2 + n_6 I^3 + \dots = n_{NL} I$$

$$\alpha_{NL} = \alpha_2 + \alpha_4 I + \alpha_6 I^2 + \dots$$



$$W = \frac{\Delta n}{\lambda \alpha_0} > 1$$

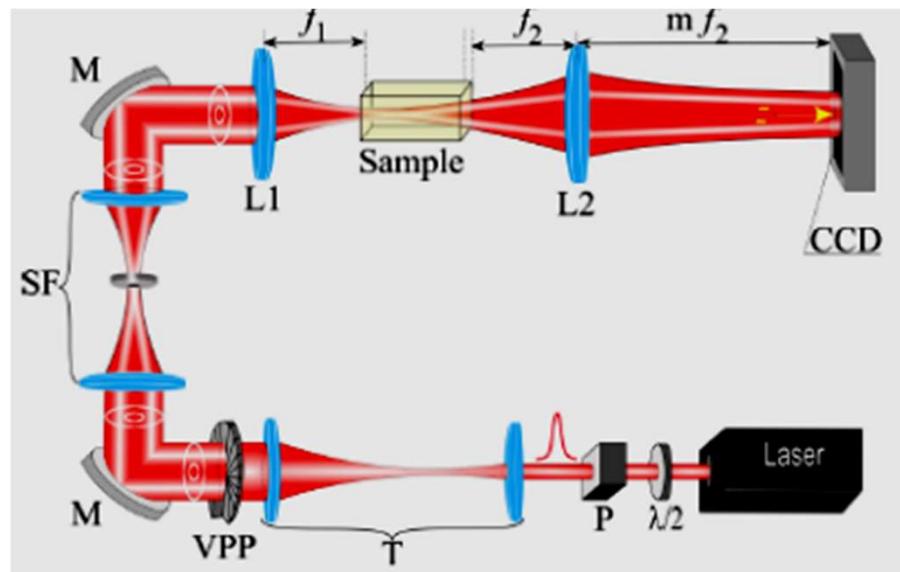
$$T = \frac{\lambda \alpha_{NL}}{n_{NL}} < 1$$



These results show that it is possible to have an efficient all-optical switch if a nanocomposite is made according to the nonlinearity management procedure presented

Challenge for  
materials scientists

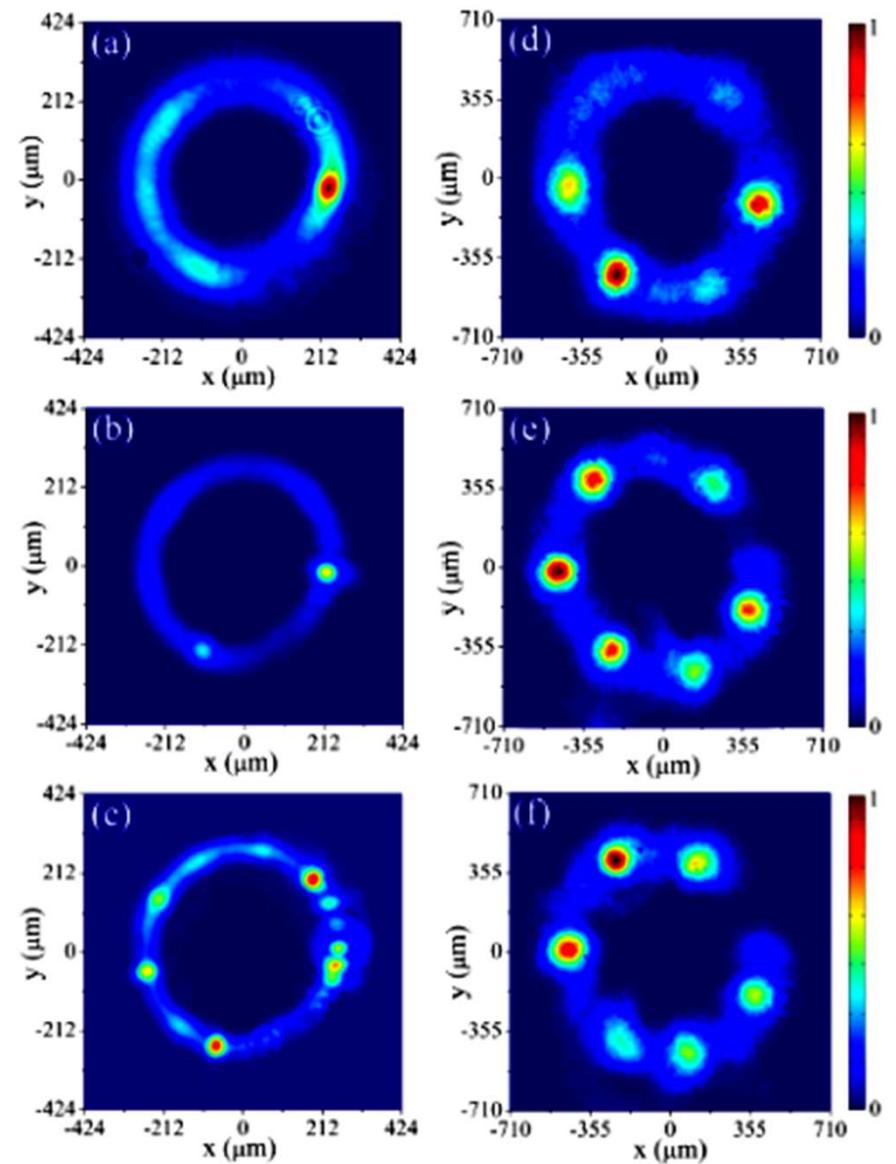
# Soliton clusters



Reyna et al.

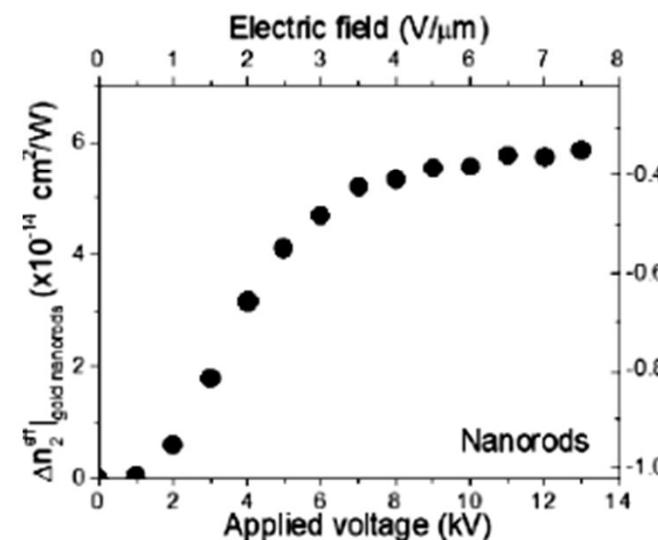
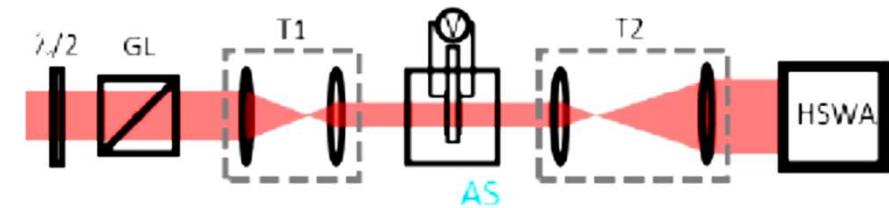
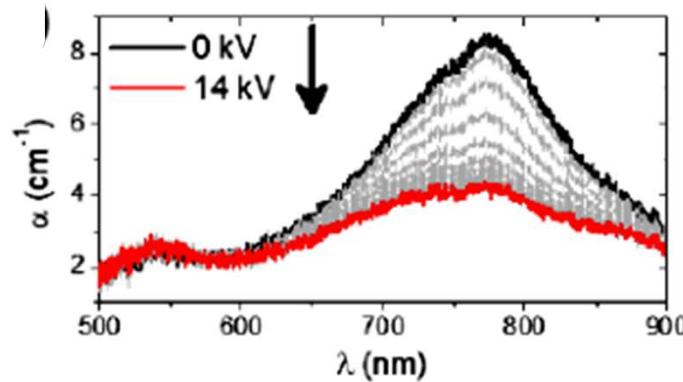
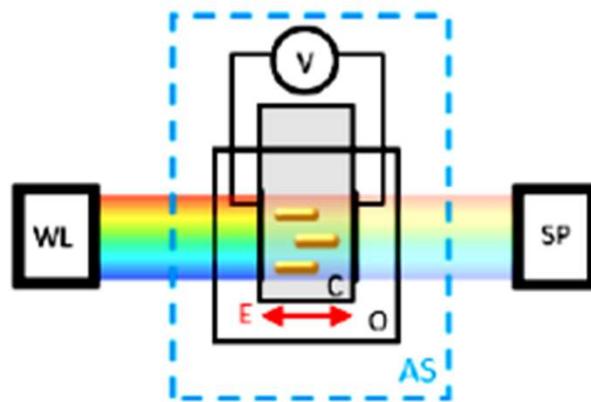
Observation and analysis of creation,  
decay, and regeneration of  
annular soliton clusters in a lossy  
cubic-quintic optical medium,

Phys. Rev. A 102 (2020) 033523



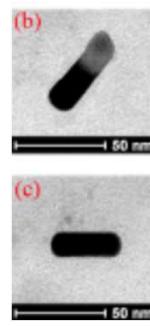
# Nonlinear refractive index of electric field aligned gold nanorods suspended in index matching oil measured with a Hartmann-Shack wavefront aberrometer

M. Maldonado et al.

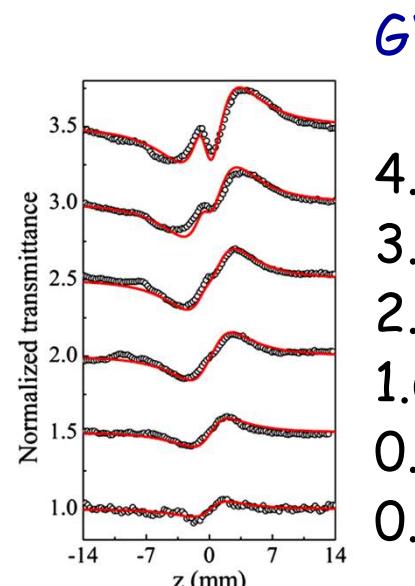


# Light Scattering, Absorption, and Refraction due to High-Order Optical Nonlinearities in Colloidal Gold Nanorods

Oliveira et al. *J. Phys. Chem. C* 2019, 123, 12997–13008

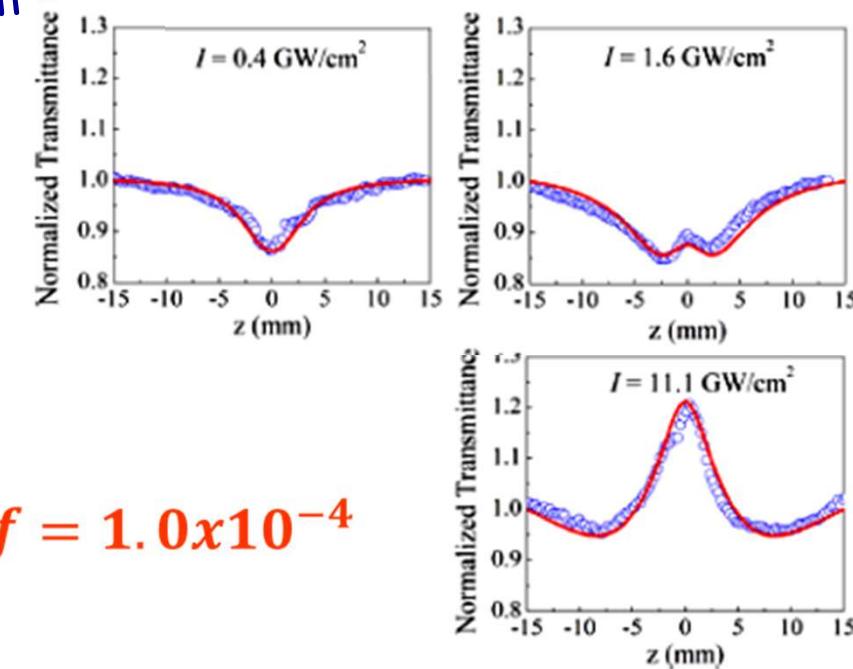


$\chi^{(3)}$

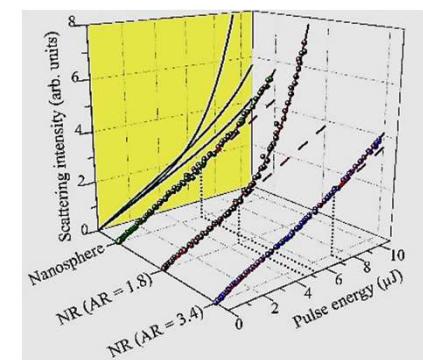
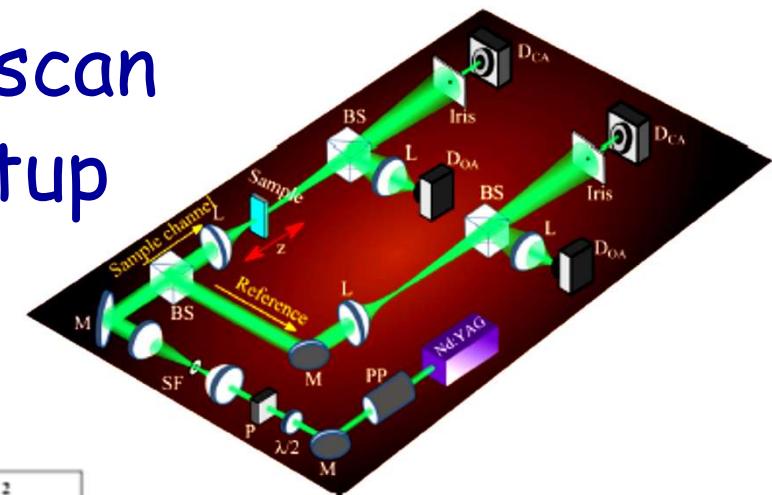


$$f = 1.0 \times 10^{-4}$$

RA: 3.4



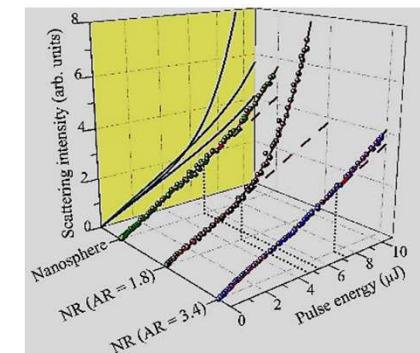
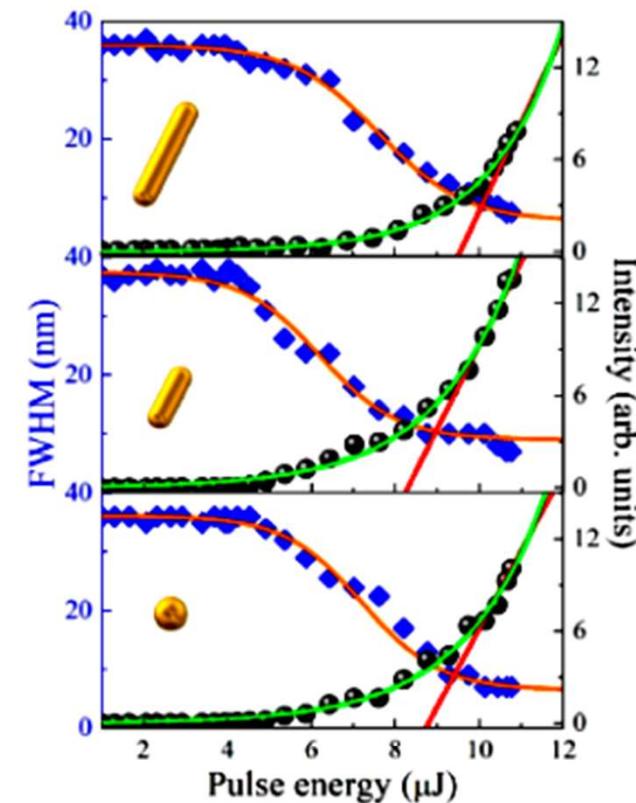
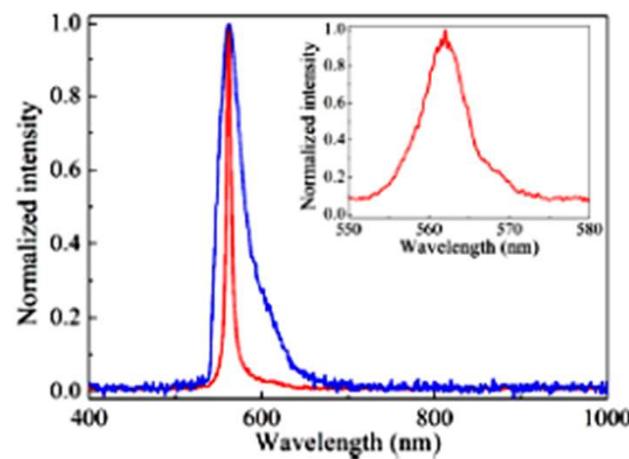
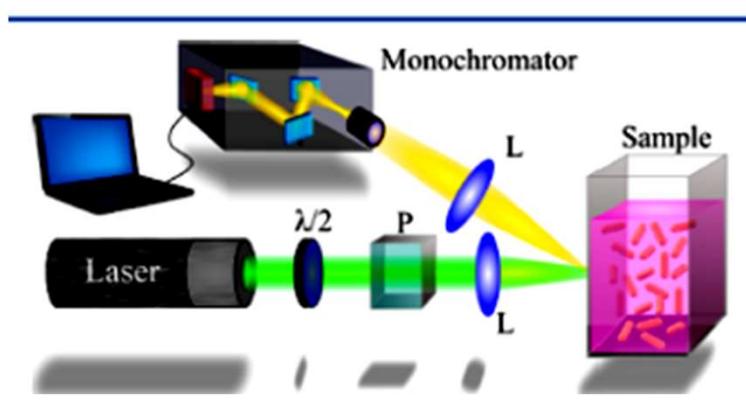
Z-scan  
setup



NL scattering

# Influence of the Fifth-Order Nonlinearity of Gold Nanorods on the Performance of Random Lasers

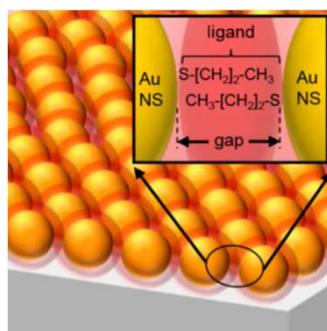
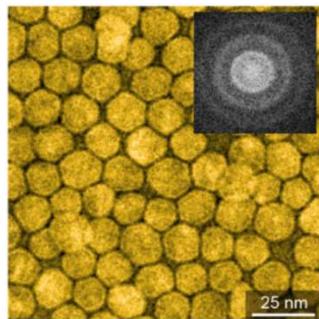
Oliveira et al. *J. Phys. Chem. C* 2020, 124, 10705–10709



Nonlinear scattering: 5th. order

# Gold metasurface

samples: Jake Fontana  
NRL USA



spheres : 15 nm diam.  
gap: 0.6 nm

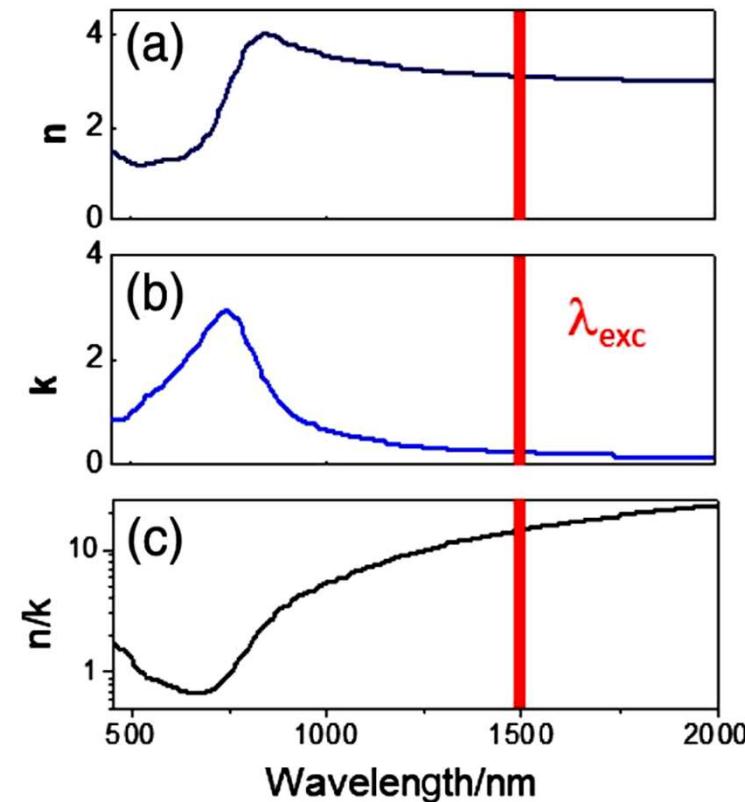
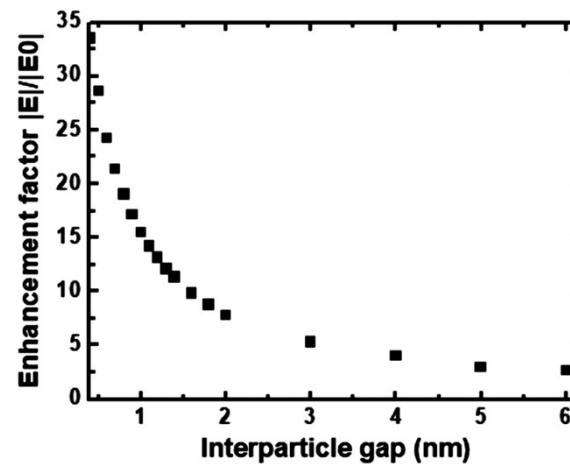
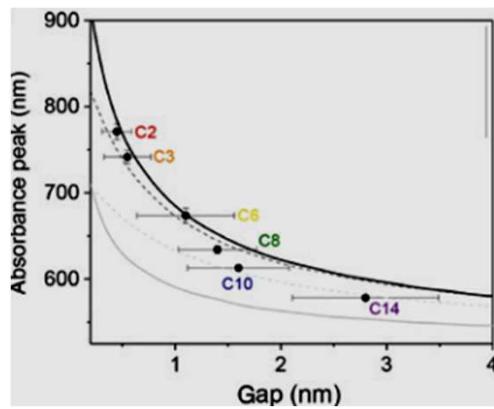
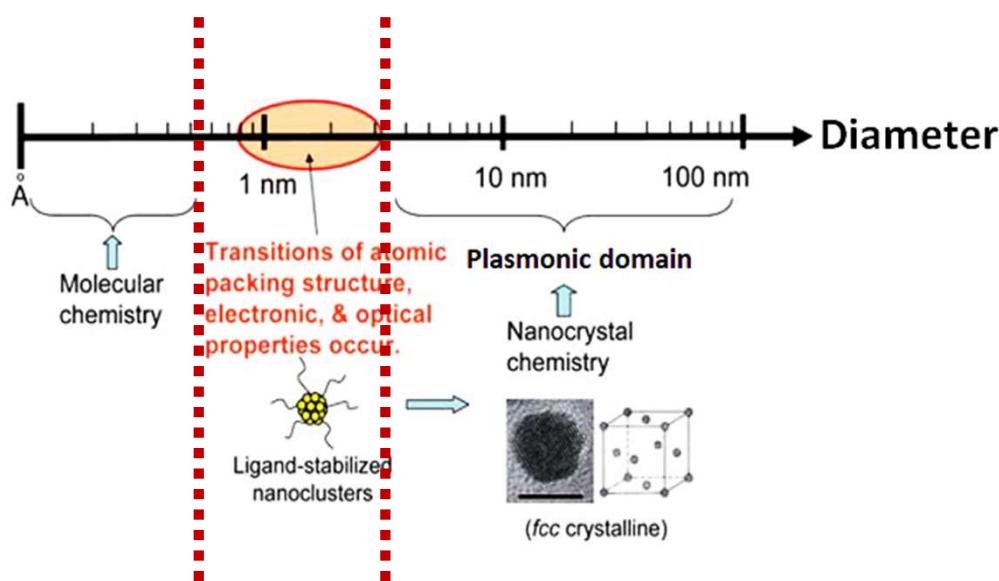


Table 2. Comparison between the NL Optical Coefficients of the Gold MS at 800 nm and 1500 nm, Corrected for Linear Absorption\*

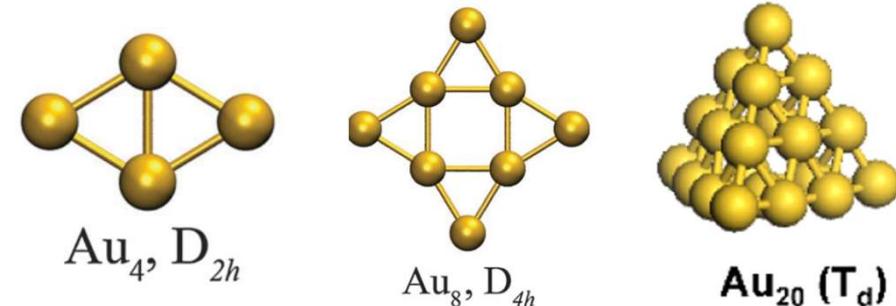
$\lambda$ (nm)	$n_2$ ( $\text{cm}^2 \text{W}^{-1}$ )	$\alpha_2$ ( $\text{cmW}^{-1}$ )	$\text{Re}\chi^{(3)}$ ( $\text{cm}^2 \text{V}^{-2}$ )	$\text{Im}\chi^{(3)}$ ( $\text{cm}^2 \text{V}^{-2}$ )
800	$-7.90 \times 10^{-9}$	$-0.90 \times 10^{-4}$	$-3.37 \times 10^{-10}$	$-1.90 \times 10^{-10}$
1500	$-1.05 \times 10^{-10}$	$3.00 \times 10^{-6}$	$-2.73 \times 10^{-12}$	$8.08 \times 10^{-13}$

Menezes et al. JOSA B 36 (2019) 1485-1491

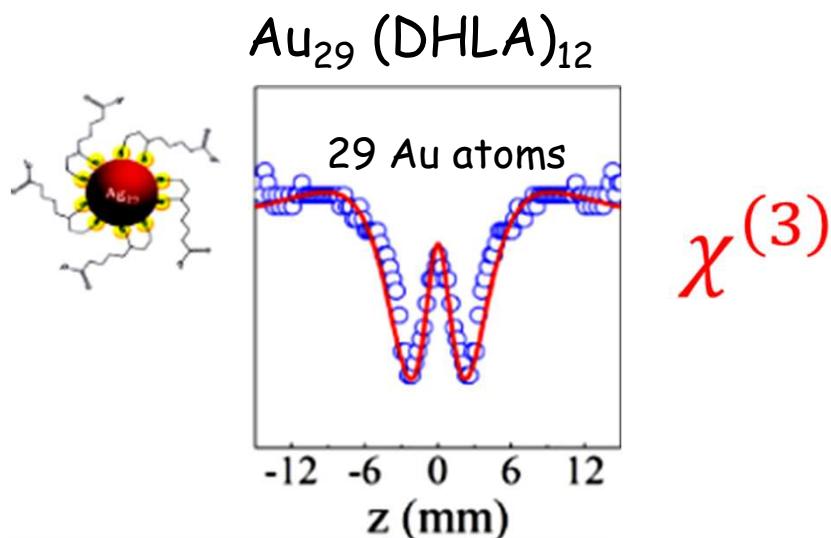
Gomes et al. Nanophotonics 9 (2020) 725-740



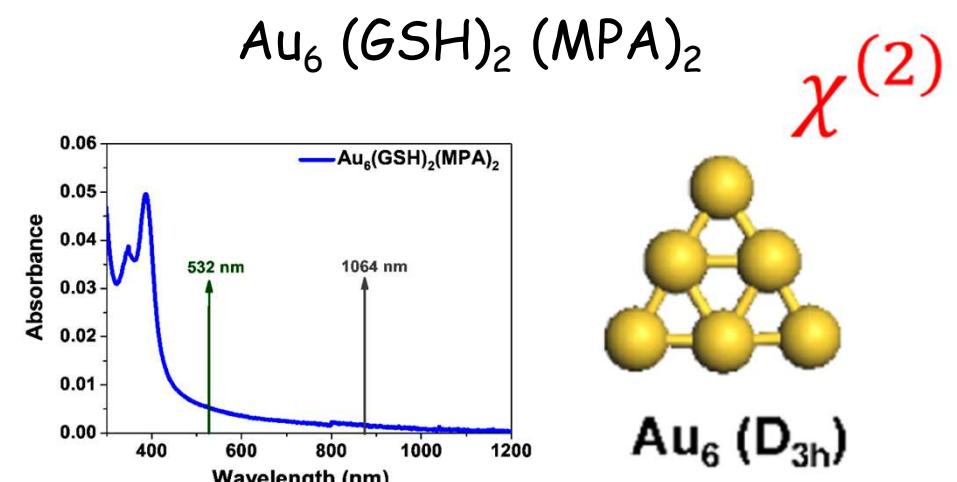
## Gold nanoclusters - Vacuum -



Geometry changes when adding ligands

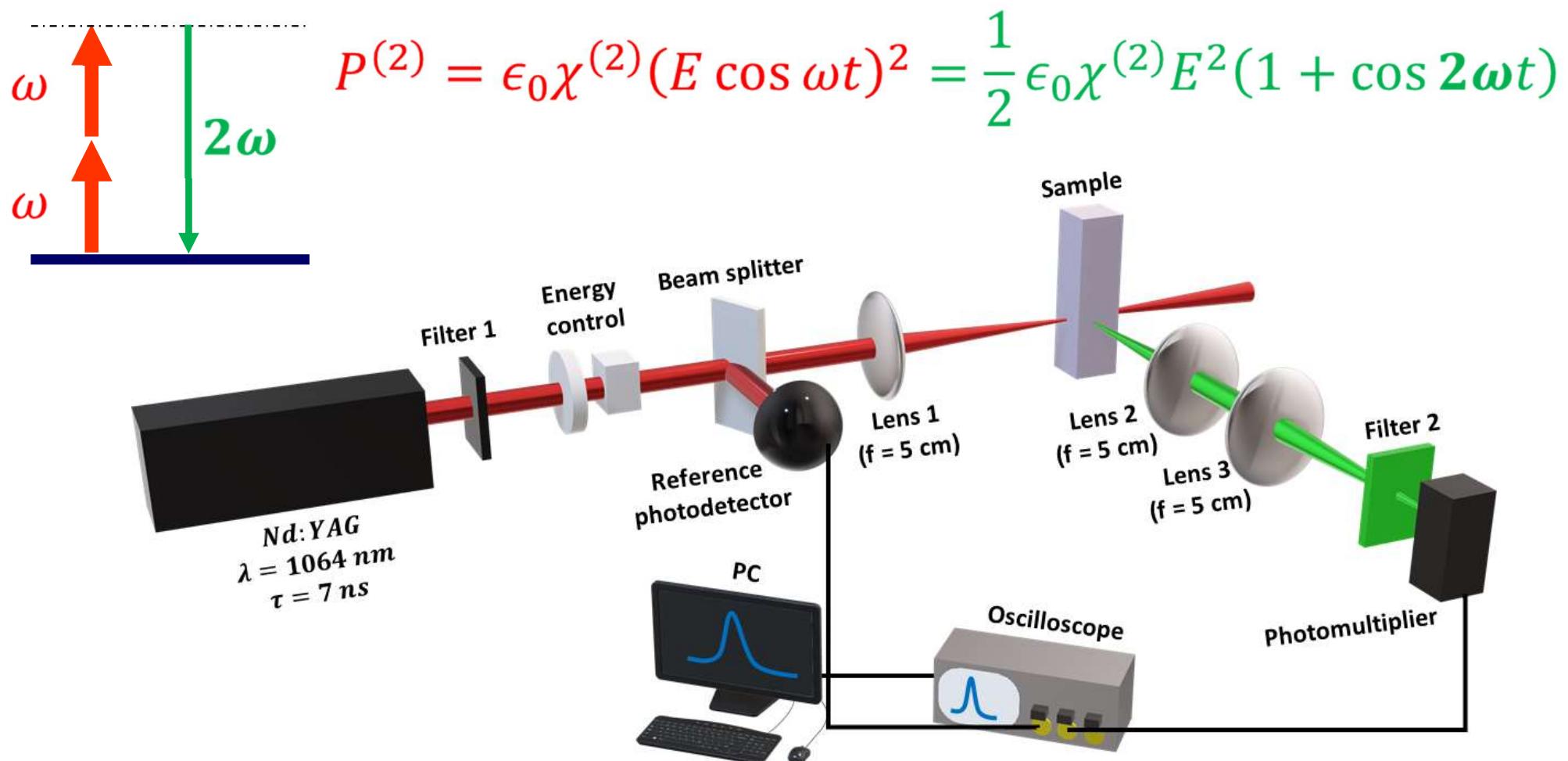


NL refraction and absorption  
Saturation of 2PA



Second harmonic generation by clusters  
of six gold atoms

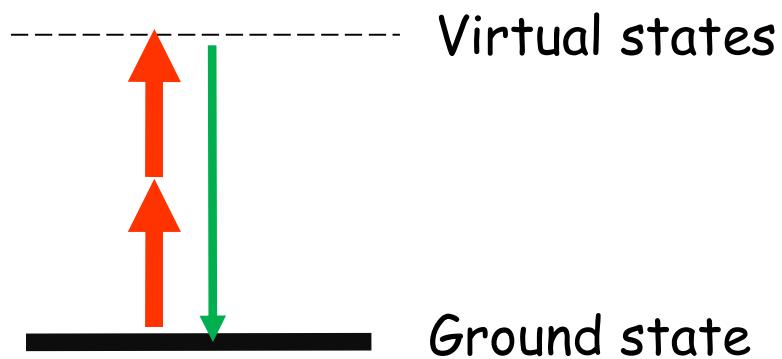
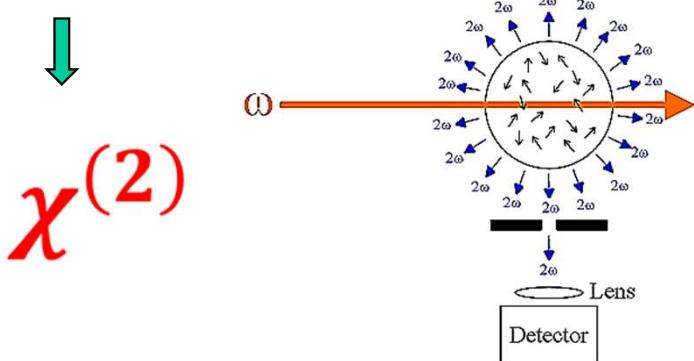
# Incoherent second harmonic generation hyper Rayleigh scattering



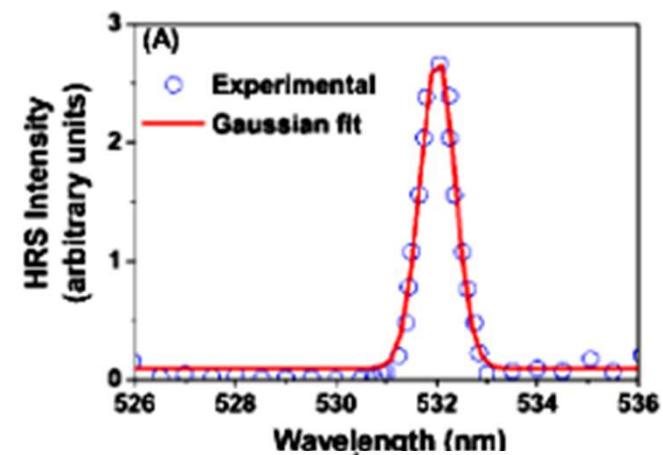
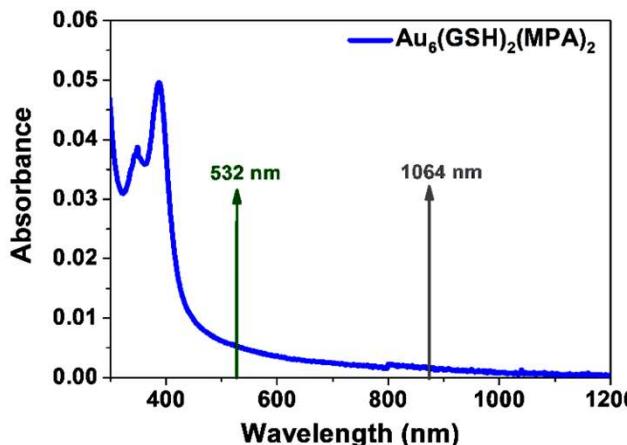
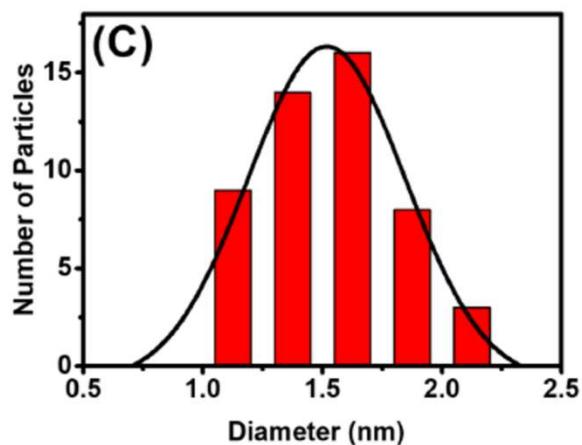
# Second order nonlinearity

Excited states

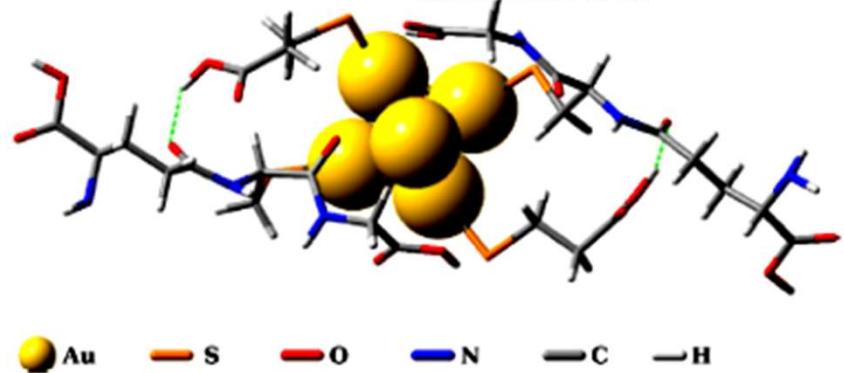
hyper-Rayleigh scattering



$\text{Au}_6(\text{GSH})_2(\text{MPA})_2$  A. Patra - India



Gold nanocluster	Wavelength (nm)	$10^{20} \beta(2\omega)$ (esu)	$10^{20} \beta(2\omega)$ per gold atom (esu)	Reference
Au <sub>6</sub>	1064	760	127	This work
Au <sub>15</sub>	802	509	34	12
Au <sub>25</sub>	802	128	5.1	12



# Summary

Metal composites present large NL susceptibility which depends on the shape and volume fraction of NPs

Metallic NPs can be nucleated inside different media allowing enhancement of:

- luminescence properties (Stokes and anti-Stokes)
- optical gain/amplification in waveguides
- random lasers, DFB lasers
- all-optical switching, etc.

## Nonlinearity Management



The control of NPs volume fraction allows suppression and/or enhancement of nonlinear optical contributions

**Robust two-dimensional spatial solitons in liquid carbon disulfide**  
Phys. Rev. Lett. 110 (2013) 013901.

Two-dimensional solitons in a quintic-septimal medium.  
Phys. Rev. A 90 (2014) 063835.

Nonlinearity management of photonic composites and observation  
of spatial-modulation instability due to quintic nonlinearity.  
Phys. Rev. A 89 (2014) 063803.

Spatial phase modulation due to quintic and septimal nonlinearities in  
metal colloids. Opt. Express 22 (2014) 22456.

An optimization procedure for the design of all-optical switches based  
on metal-dielectric nanocomposites. Opt. Express 23 (2015) 7659 .

Robust self-trapping of optical vortex beams in a saturable  
optical medium. Phys. Rev. A 93 (2016) 013840.

Taming the emerging beams after the split of optical vortex solitons in a  
saturable. Phys. Rev. A 93 (2016) 013843.

Guiding and confinement of light induced by optical vortex solitons in a  
cubic-quintic medium. Opt. Lett. 41 (2016) 191.

Observation and analysis of soliton clusters

Phys. Rev. A 102 (2020) 033523

Control of the nonlinear refractive index by applied voltage

Optics Express 26 (2018) 20298

Investigation of high order nonlinearities of colloidal gold nanorods

J. Phys. Chem. C 123 (2019) 12997-13008

J. Phys. Chem. C 124 (2020) 10705-10709

Nonlinear optical behavior of gold metasurface

Nanophotonics 9 (2020) 725-740

NL refraction, NL absorption, Saturated 2PA of silver nanocluster with 29 atoms

J. Phys. Chem. C 122 (2018) 18682

First hyperpolarizability of gold clusters with six atoms

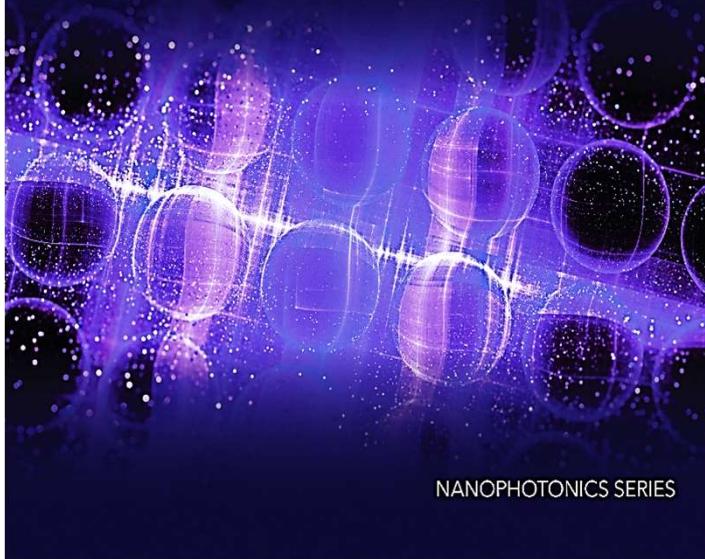
J. Phys. Chem. C 124 (2020) 15440-15447



SERIES EDITOR: DAVID L. ANDREWS

# METAL NANOSTRUCTURES FOR PHOTONICS

LUCIANA REYES PIRES KASSAB  
CID BARTOLOMEU DE ARAUJO



ELSEVIER (2019)

Section I (Influence of metallic nanoparticles on luminescence of ions in solids)

Section II (Near and far-field optical phenomena associated and/or influenced by metallic nanoparticles)

Section III (Photonic materials and devices with improved performance influenced by metallic nanoparticles)

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Reyna, de Araújo

High-order optical nonlinearities  
in plasmonic nanocomposites - a review  
Adv. Opt. Photon. 9 (2017) 720-774

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Reyna, de Araújo. Beyond third-order optical nonlinearities in liquid suspensions of metal-nanoparticles and metal-nanoclusters.

J. Opt. 24 (2022) 104006.

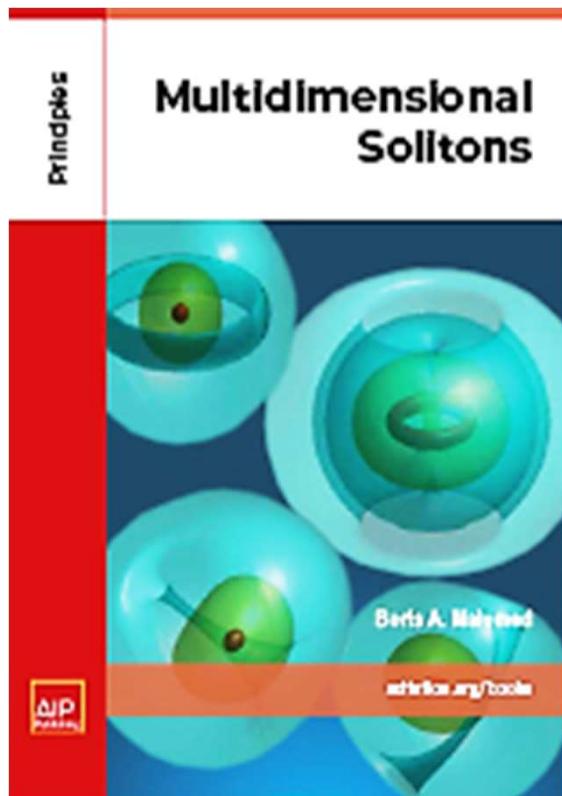
## Review Article

# Frontiers in multidimensional self-trapping of nonlinear fields and matter

Y. V. Kartashov, G. E. Astrakharchik, B. A. Malomed & L. Torner

*Nature Reviews Physics* volume 1, pages 185–197 (2019)

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Author: Boris A. Malomed

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