



Neutrino Mass and Neutrino Oscillations

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PHYS703 Introduction to Quantum Field Theory

Outline

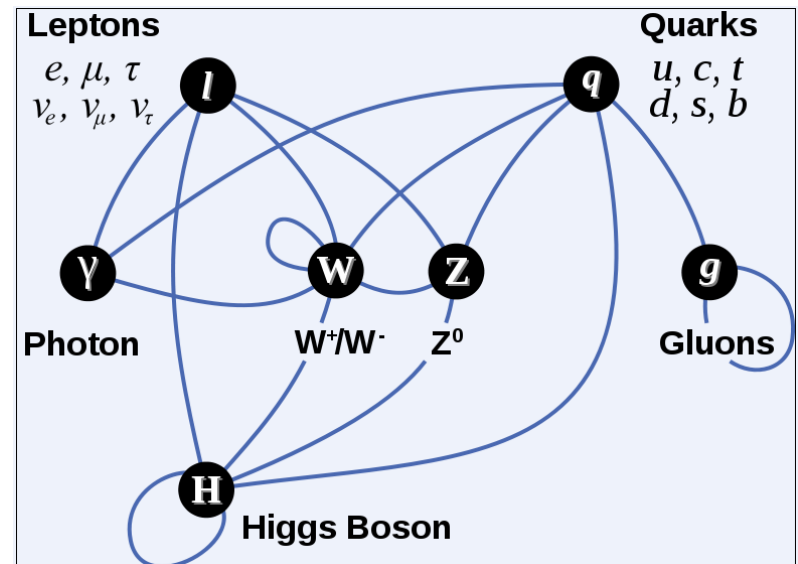
- Neutrinos in the Standard Model
- Phenomenology of neutrino oscillations
- Neutrino oscillation experiments
- Sterile neutrinos
- The see-saw mechanism

Standard Model

- Gauge group: $SU(3)_C \times SU(2)_L \times U(1)_{EM}$
- Gauge bosons: 8 gluons, W^\pm, Z^0, γ
- Quarks and leptons:

1st	u	d	e	ν_e
2nd	c	s	μ	ν_μ
3rd	t	b	τ	ν_τ

- Higgs boson
- 19 parameters



Summary of interactions

About the Neutrino

- Pauli's hypothesis (1934) for beta decay

$$n \rightarrow p + e + \bar{\nu}_e$$

- Experimental discoveries:

ν_e	Cowan, Reines	1956
ν_μ	Lederman, Schwartz, Steinberger	1962
ν_τ	DONUT collaboration	2000

- Properties: $q = 0$, weakly interacting, spin $\frac{1}{2}$, $L = 1$, tiny mass, (left-handed)

Neutrino Mass

- Flavor states: superposition of mass states

$$|\nu_\ell\rangle = \sum_a U_{\ell a} |\nu_a\rangle$$

$$|\nu_\ell(t)\rangle = \sum_a e^{-iE_a t} U_{\ell a} |\nu_a\rangle$$

- Survival or transition probability

$$|\langle \nu_{\ell'} | \nu_\ell(t) \rangle|^2 = \sum_{a,b} |U_{\ell' b} U_{\ell a} U_{\ell' a}^* U_{\ell b}^*| \cos[(E_b - E_a)t - \phi]$$

- Highly relativistic:

$$E \approx p + \frac{m^2}{2p}$$

$$\phi_i = \frac{m_i^2}{2p_i} x$$

- 2-flavor case:

$$|\nu_\ell(t)\rangle = e^{i\phi_1} \cos \theta |\nu_1\rangle + e^{i\phi_2} \sin \theta |\nu_2\rangle$$

Neutrino Oscillations

- Oscillation probability

$$\left| \langle \nu_{\ell'} | \nu_{\ell}(t) \rangle \right|^2 = \sum_{a,b} \left| U_{\ell'b} U_{\ell a} U_{\ell'a}^* U_{\ell b}^* \right| \cos \left(\frac{2\pi x}{L} - \phi \right)$$

$$\Delta m_{ab}^2 = m_a^2 - m_b^2$$

$$L = \frac{4\pi p}{\Delta m_{ab}^2}$$

- Nontrivial effects when $x \neq \kappa L, \kappa \in \mathbb{Z}$
- Same for antineutrino = CPT conservation

PMNS Mixing Matrix

- General neutrino mixing matrix: PMNS

$$U = \begin{pmatrix} c_{12}c_{13} & -s_{12}c_{13} & s_{13} \\ s_{12}c_{13} + c_{12}s_{23}s_{13} & c_{12}c_{23} - s_{12}s_{23}s_{13} & -s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13} & c_{12}s_{23} + s_{12}c_{23}s_{13} & c_{23}c_{13} \end{pmatrix}$$

$$c_{ab} = \cos \theta_{ab} \quad s_{ab} = \sin \theta_{ab}$$

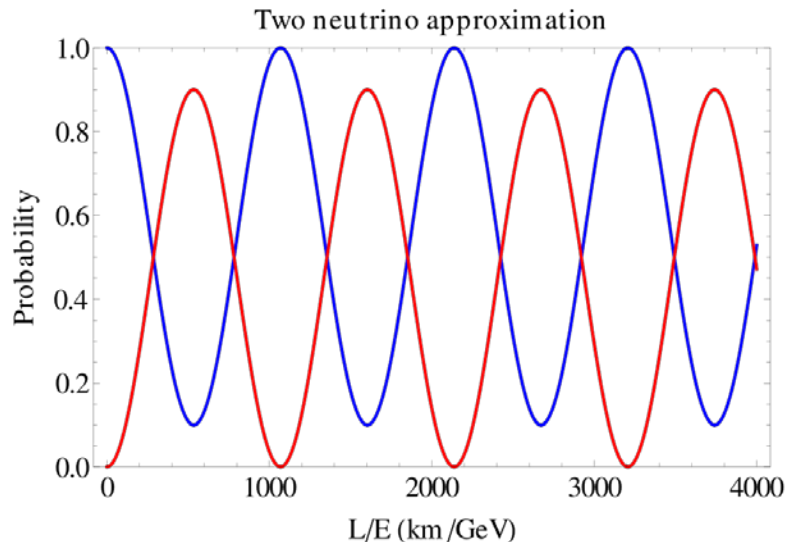
- Similar in form to quark mixing: CKM

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

Two-Flavor Oscillation

- Limiting case: with one small mixing angle

$$\left| \langle \nu_\ell | \nu_\ell(t) \rangle \right|^2 = 1 - \sin^2 2\theta_{ab} \sin^2 \left(\frac{x \Delta m_{ab}^2}{4E} \right)$$



Valid for muon-tau mixing as well as electron- ν_X mixing, where

$$\nu_X = \alpha \nu_\mu + \beta \nu_\tau$$

because one mixing angle is small and two of the mass eigenstates are close relative to the other one.

Mass Term in Lagrangian

- Phenomenology w/ 3 neutrinos:

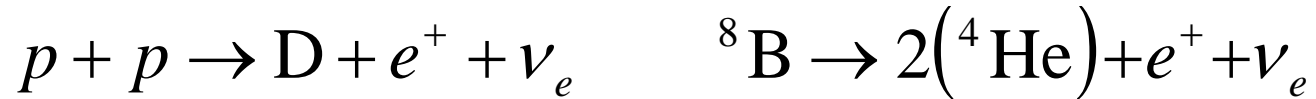
$$L = \frac{1}{2} M_{ij} \bar{\nu}_i (1 + \gamma_5) \nu_j + L_{NC} + L_{CC}$$

- Neutral current term stays the same
- Mixing appears in charged current term

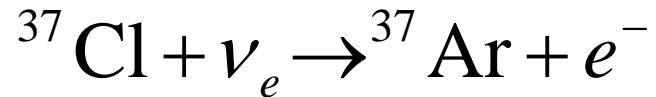
$$L_{cc} = i \frac{g}{\sqrt{2}} U_{ai} W_\mu \left(\bar{\ell}_a \gamma^\mu \gamma_L \nu_i \right) + \text{h.c.}$$

Solar Neutrinos₁₂

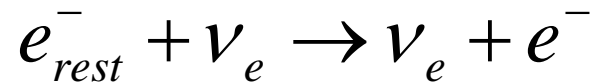
- Nuclear reactions (H-He fusion):



- Detect: inverse beta decay (Davis, 1968)



- Scatter off electron



- Double ratio:

$$R = \frac{(\mu/e)}{(\mu/e)_{MC}}$$

$$(\text{SNO, 2001}) \quad R = 0.35$$

Atmospheric Neutrinos₂₃

- Typical cosmic ray reaction chain:

$$p + X \rightarrow \pi^{\pm} + Y \quad \pi^{\pm} \rightarrow \mu^{\pm} + \nu_{\mu} \quad \mu^{\pm} \rightarrow e^{\pm} + \nu_e + \nu_{\mu}$$

- Best double ratio R :

Super Kamiokande	Sub-GeV	$0.638 \pm 0.16 \pm 0.050$
	Multi-GeV	$0.658 \pm 0.030 \pm 0.078$

- Up-down symmetry shows muon flux deficit

$$\alpha_{\ell} = \left(\frac{U - D}{U + D} \right)_{\ell}$$

$$\alpha_e = 0$$

$$\alpha_{\mu} < 0 \quad (\text{high momenta})$$

Reactor & Accelerator Studies

- Reactors: antineutrinos from beta decay
- CHOOZ: $n + {}^m\text{Gd} \rightarrow {}^{m+1}\text{Gd}^* \rightarrow \text{Gd} + \gamma\text{s}$

$$\Delta m_{13}^2 = 2.5 \times 10^{-5} \text{ eV}^2$$

$$\sin^2 2\theta_{13} = 0.15$$

- Accelerators: Magnetically focused neutrino beam:

$$p \rightarrow \pi, K \rightarrow \mu, e, \nu_{e,\mu}$$

- Muon disappearance, electron appearance

Double Beta Decay

- 35 natural isotopes w/ ground state s.t.

$$(Z, A) \rightarrow (Z + 2, A) + 2e + 2\nu_e$$

- Giulio Racah(1937), Wendell Furry(1939):

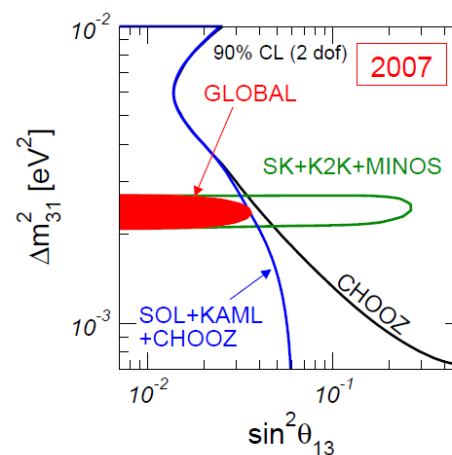
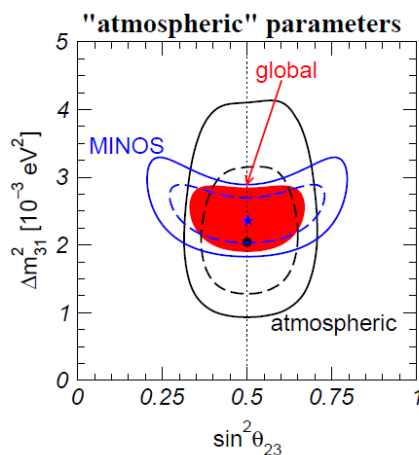
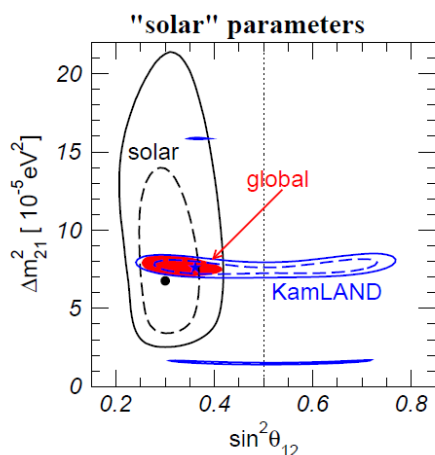
$$(Z, A) \rightarrow (Z + 1, A) + e + \bar{\nu}_e$$

$$(Z + 1, A) + \nu_e \rightarrow (Z + 2, A) + e$$

- Double beta w/o neutrinos = Majorana
- Mass limit (Ge to Se) : $m_\nu \approx 0.35\text{eV}$

Summary of Experiments

- Global analysis of oscillation parameters



Parameter	Best fit	2σ	3σ
Δm_{12}^2 (10^{-5} eV 2)	7.6	7.3 – 8.1	7.1 – 8.3
$ \Delta m_{13}^2 $ (10^{-3} eV 2)	2.4	2.1 – 2.7	2.0 – 2.8
$\sin^2 \theta_{12}$	0.32	0.28 – 0.37	0.26 – 0.40
$\sin^2 \theta_{23}$	0.50	0.38 – 0.63	0.34 – 0.67
$\sin^2 \theta_{13}$	0.007	≤ 0.033	≤ 0.050

3 ———	2 ———
	1 ———
2 ———	
1 ———	3 ———
normal	inverted

Sterile Neutrinos

- Add N sterile neutrinos to 3 active ones

$$M = \begin{pmatrix} m & \mu \\ \mu^T & M \end{pmatrix} \quad \text{dimension (3+N)}$$

- Conserved lepton number: Dirac neutrinos

$$\psi_a = \begin{pmatrix} \nu_a \\ s_a \end{pmatrix} \quad U' = \begin{pmatrix} c_{12}c_{13} & -s_{12}c_{13} & s_{13}e^{-i\delta} \\ s_{12}c_{13} + c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & -s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & c_{12}s_{23} + s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

- Majorana neutrinos: KU , $K = \text{diag}(e^{i\alpha_1}, e^{i\alpha_2}, e^{i\alpha_3})$

See-saw Mechanism

- Single generation of sterile neutrino: $N = 1$

$$M = \begin{pmatrix} 0 & \mu \\ \mu & M \end{pmatrix}$$

$$m^{(\pm)} = \frac{1}{2} \sqrt{M^2 + 4\mu^2} \mp M$$

$$OMO^T = K^2 m$$

- Suppose $M \gg \mu$. Then

$$m^{(-)} \approx \frac{\mu^2}{M} \quad m^{(+)} \approx M$$

$$\begin{array}{l} m^{(+)} \uparrow, m^{(-)} \downarrow \\ m^{(+)} \downarrow, m^{(-)} \uparrow \end{array}$$

Conclusions

- Neutrino oscillations imply neutrino masses can't all be equal, in particular they can't all be zero (but one can be massless).
- The conversion probability is

$$\Pr[\nu_\ell \rightarrow \nu_{\ell'}] \propto \sin^2 2\theta_{ab}, \Delta m_{ab}^2$$

- Sources of neutrino mass: see-saw mechanism, expanded Higgs, GUTs, supersymmetry