

Classical Maxwell Fields

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Chapter 1

Classical Maxwell Fields

1.1 Rationalized Heaviside-Lorentz Units

The SI units are based on a unit of length of the order of human size originally related to the size of the earth, a unit of time approximately equal to the time between heartbeats, and a unit of mass related to the length unit and the mass of water. None of these depend on any even nearly fundamental physical quantities. Therefore many important physical equations end up with extra (needless) constants in them like c . Even with the three basic units defined, we could have chosen the unit of charge correctly to make ϵ_0 and μ_0 unnecessary but instead a very arbitrary choice was made $\mu_0 = 4\pi \times 10^{-7}$ and the Ampere is defined by the current in parallel wires at one meter distance from each other that gives a force of 2×10^{-7} Newtons per meter. The Coulomb is set so that the Ampere is one Coulomb per second. With these choices SI units make Maxwell's equations and our field theory look very messy.

In Rationalized Heaviside-Lorentz units we decrease the field strength by a factor of $\sqrt{4\pi}$ and increase the charges by the same factor, leaving the force unchanged.

$$\begin{aligned}\vec{E} &\rightarrow \frac{\vec{E}}{\sqrt{4\pi}} \\ \vec{B} &\rightarrow \frac{\vec{B}}{\sqrt{4\pi}} \\ \vec{A} &\rightarrow \frac{\vec{A}}{\sqrt{4\pi}} \\ e &\rightarrow e\sqrt{4\pi} \\ \alpha &= \frac{e^2}{\hbar c} \rightarrow \frac{e^2}{4\pi\hbar c} \approx \frac{1}{137}.\end{aligned}$$

It's not a very big change but it would have been nice if Maxwell had started with this set of units. Of course the value of α cannot change, but, the formula for it does because we have redefined the charge e . Maxwell's Equations in CGS units are:

$$\begin{aligned}
\nabla \cdot \vec{B} &= 0 \\
\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} &= 0 \\
\nabla \cdot \vec{E} &= 4\pi\rho \\
\nabla \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} &= \frac{4\pi}{c} \vec{j}.
\end{aligned}$$

The Lorentz Force is

$$\vec{F} = -e \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right) \quad (1.1)$$

When we change to Rationalized Heaviside-Lorentz units, the equations become

$\nabla \cdot \vec{B} = 0$
$\nabla \times \vec{E} + \frac{1}{c} \frac{\partial \vec{B}}{\partial t} = 0$
$\nabla \cdot \vec{E} = \rho$
$\nabla \times \vec{B} - \frac{1}{c} \frac{\partial \vec{E}}{\partial t} = \frac{1}{c} \vec{j}$
$\vec{F} = -e \left(\vec{E} + \frac{1}{c} \vec{v} \times \vec{B} \right)$

That is, the equations remain the same except the factors of 4π in front of the source terms disappear. Of course, it would still be convenient to set $c = 1$ since this has been confusing us about 4D geometry and c is the last unnecessary constant in Maxwell's equations. For our calculations, we can set $c = 1$ any time we want unless we need answer in centimeters.

1.2 The Electromagnetic Field Tensor

The transformation of electric and magnetic fields under a Lorentz boost we established even before Einstein developed the theory of relativity. We know that E-fields can transform into B-fields and vice versa. For example, a point charge at rest gives an Electric field. If we boost to a frame in which the charge is moving, there is an Electric and Magnetic field. This means that the E-field cannot be a Lorentz vector. We need to put the Electric and Magnetic fields together into one (tensor) object to properly handle Lorentz transformations and to write our equations in a covariant way.

The simplest way and the correct way to do this is to make the Electric and Magnetic fields components of a rank 2 (antisymmetric) tensor.

$$F_{\mu\nu} = \begin{pmatrix} 0, B_z, -B_y, -iE_x \\ -B_z, 0, B_x, -iE_y \\ B_y, -B_x, 0, -iE_z \\ iE_x, iE_y, iE_z, 0 \end{pmatrix} \quad (1.2)$$

The fields can simply be written in terms of the vector potential, (which is a Lorentz vector) $A_\mu = (\vec{A}, i\phi)$.

$$F_{\mu\nu} = \frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\mu}{\partial x_\nu} \quad (1.3)$$

Note that this is automatically antisymmetric under the interchange of the indices. As before, the first two (sourceless) Maxwell equations are automatically satisfied for fields derived from a vector potential. We may write the other two Maxwell equations in terms of the 4-vector $j_\mu = (\vec{j}, ic\rho)$.

$$\frac{\partial F_{\mu\nu}}{\partial x_\nu} = \frac{j_\mu}{c} \quad (1.4)$$

Which is why the T-shirt given to every MIT freshman when they take Electricity and Magnetism should say

"...and God said $\frac{\partial}{\partial x_\nu} \left(\frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\mu}{\partial x_\nu} \right) = \frac{j_\mu}{c}$ and there was light":

Of course he or she had not yet quantized the theory in that statement.

For some peace of mind, let's verify a few terms in the equations. Clearly all the diagonal terms in the field tensor are zero by antisymmetry. Let's take some example off-diagonal terms in the field tensor, checking the (old) definition of the fields in terms of the potential.

$$\vec{B} = \vec{\nabla} \times \vec{A}$$

$$\vec{E} = -\vec{\nabla}\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$F_{12} = \frac{\partial A_2}{\partial x_1} - \frac{\partial A_1}{\partial x_2} = (\vec{\nabla} \times \vec{A})_z = B_z$$

$$F_{13} = \frac{\partial A_3}{\partial x_1} - \frac{\partial A_1}{\partial x_3} = -(\vec{\nabla} \times \vec{A})_y = -B_y$$

$$F_{4i} = \frac{\partial A_i}{\partial x_4} - \frac{\partial A_4}{\partial x_i} = \frac{1}{ic} \frac{\partial A_i}{\partial t} - \frac{\partial(i\phi)}{\partial x_i} = -i \left(\frac{1}{c} \frac{\partial A_i}{\partial t} + \frac{\partial \phi}{\partial x_i} \right) = -i \left(\frac{\partial \phi}{\partial x_i} + \frac{1}{c} \frac{\partial A_i}{\partial t} \right) = iE_i.$$

Let's also check what the Maxwell equations says for the last row in the tensor.

$$\frac{\partial F_{4v}}{\partial x_v} = \frac{j_4}{c}$$

$$\frac{\partial F_{4i}}{\partial x_i} = \frac{ic\rho}{c}$$

$$\frac{\partial(iE_i)}{\partial x_i} = i\rho$$

$$\frac{\partial E_i}{\partial x_i} = \rho$$

$$\vec{\nabla} \cdot \vec{E} = \rho.$$

We will not bother to check the Lorentz transformation of the fields here. It is right.

1.3 The Lagrangian for Electromagnetic Fields

There are not many ways to make a scalar Lagrangian from the field tensor. We already know that

$$\frac{\partial F_{\mu\nu}}{\partial x_\nu} = \frac{j_\mu}{c} \quad (1.5)$$

and we need to make our Lagrangian out of the fields, not just the current. Again x_μ cannot appear explicitly because that violates symmetries of nature. Also we want a linear equation and so higher powers of the field should not occur. A term of the form $m A_\mu A_\mu$ is a mass term and would cause fields to fall off faster than $\frac{1}{r}$. So, the only reasonable choice is

$$F_{\mu\nu} F_{\mu\nu} = 2(B^2 - E^2) \quad (1.6)$$

One might consider

$$e_{\mu\nu\lambda\sigma} F_{\mu\nu} F_{\lambda\sigma} = \vec{B} \cdot \vec{E} \quad (1.7)$$

but that is a pseudo-scalar, not a scalar. That is, it changes sign under a parity transformation. The EM interaction is known to conserve parity so this is not a real option. As with the scalar field, we need to add an interaction with a source term. Of course, we know electromagnetism well, so finding the right Lagrangian is not really guess work. The source of the field is the vector j_μ , so the simple scalar we can write is $j_\mu A_\mu$.

The Lgrangian for Classical Electricity and Magnetism we will try is.

$$\mathcal{L}_{eM} = -\frac{1}{4} F_{\mu\nu} F_{\mu\nu} + \frac{1}{c} j_\mu A_\mu \quad (1.8)$$

In working with this Lagrangian, we will treat each component of A as an independent field. THe next step is to check what the Euler-Lagrange equation gives us.

$$\begin{aligned}
& \frac{\partial}{\partial x_\nu} \left(\frac{\partial \mathcal{L}}{\partial (\partial A_\mu / \partial x_\nu)} \right) - \frac{\partial \mathcal{L}}{\partial A_\mu} = 0 \\
& \mathcal{L} = -\frac{1}{4} F_{\mu\nu} F_{\mu\nu} + \frac{1}{c} j_\mu A_\mu = -\frac{1}{4} \left(\frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\mu}{\partial x_\nu} \right) \left(\frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\mu}{\partial x_\nu} \right) + \frac{1}{c} j_\mu A_\mu \\
& \frac{\partial \mathcal{L}}{\partial (\partial A_\mu / x_\nu)} = -\frac{1}{4} \frac{\partial}{\partial (\partial A_\mu / \partial x_\nu)} \left(\frac{\partial A_\sigma}{\partial x_\lambda} \frac{\partial A_\lambda}{\partial x_\sigma} \right) \left(\frac{\partial A_\sigma}{\partial x_\lambda} - \frac{\partial A_\lambda}{\partial x_\sigma} \right) \\
& = -\frac{1}{4} \frac{\partial}{\partial (\partial A_\mu / \partial x_\nu)} \left(2 \frac{\partial A_\sigma}{\partial x_\lambda} \frac{\partial A_\sigma}{\partial x_\lambda} - 2 \frac{\partial A_\sigma}{\partial x_\lambda} \frac{\partial A_\lambda}{\partial x_\sigma} \right) \\
& = -\frac{1}{4} 4 \left(\frac{\partial A_\mu}{\partial x_\nu} - \frac{\partial A_\nu}{\partial x_\mu} \right) \\
& = -F_{\nu\mu} = F_{\mu\nu} \\
& \frac{\partial}{\partial x_\nu} F_{\mu\nu} - \frac{\partial \mathcal{L}}{\partial A_\mu} = 0 \\
& \frac{\partial}{\partial x_\nu} F_{\mu\nu} - \frac{j_\mu}{c} = 0 \\
& \frac{\partial}{\partial x_\nu} F_{\mu\nu} = \frac{j_\mu}{c}.
\end{aligned}$$