

# Quantum Theory of Radiation

Marcus Allen Denslow

2026-01-18

# Contents

Chapter 1	Quantum Theory of Radiation	Page 2
1.1	Transverse and Longitudinal Fields	2
1.2	Fourier Decomposition of Radiation Oscillators	3
1.3	The Hamiltonian for the Radiation Field	5
1.4	Canonical Coordinates and Momenta	7

# Chapter 1

# Quantum Theory of Radiation

## 1.1 Transverse and Longitudinal Fields

In non-relativistic Quantum Mechanics, the static Electric field is represented by a scalar potential, magnetic fields by the vector potential, and the radiation field also through the vector potential. It will be convenient to keep this separation between the large static atomic Electric field and the radiation fields, however, the equations we have contain the four-vector  $A_\mu$  with all the fields mixed. When we quantize the field, all E and B fields as well as electromagnetic waves will be made up of photons. It is useful to be able to separate the E fields due to fixed charges from the EM radiation from moving charges. This separation is not Lorentz invariant, but it is still useful. Enrico Fermi showed, in 1930, that  $A_{\parallel}$  together with  $A_0$  give rise to Coulomb interactions between particles, whereas  $A_{\perp}$  gives rise to the EM radiation from moving charges. With this separation, we can maintain the form of our non-relativistic Hamiltonian.

$$\boxed{\mathbf{H} = \sum_j \frac{1}{2m_j} \left( \vec{p} - \frac{e}{c} \vec{A}_{\perp}(\vec{x}_j) \right)^2 + \sum_{i>j} \frac{e_i e_j}{4\pi \|\vec{x}_i - \vec{x}_j\|} + \mathbf{H}_{\text{rad}}} \quad (1.1)$$

Where  $\mathbf{H}_{\text{rad}}$  is purely the Hamiltonian of the radiation (containing only  $\vec{A}_{\perp}$ ) and  $\vec{A}_{\perp}$  is the part of the vector potential which satisfies  $\nabla \cdot \vec{A}_{\perp} = 0$ . Note that  $\vec{A}_{\parallel}$  and  $A_0$  appear nowhere in the Hamiltonian. Instead, we have the Coulomb potential. This separation allows us to continue with our standard Hydrogen solution and just add radiation. We will not derive this result.

In a region in which there are no source terms,

$$j_\mu = 0 \quad (1.2)$$

we can make a gauge transformation which eliminates  $A_0$  by choosing  $\Lambda$  such that

$$\frac{1}{c} \frac{\partial \Lambda}{\partial t} = A_0 \quad (1.3)$$

Since the fourth component of  $A_\mu$  is now eliminated, the Lorentz condition now implies that

$$\vec{\nabla} \cdot \vec{A} = 0 \quad (1.4)$$

Again, making one component of a 4-vector zero is not a Lorentz invariant way of working. We have to redo the gauge transformation if we move to another frame.

If  $j_\mu \neq 0$ , then we cannot eliminate  $A_0$ , since  $\boxed{A}_0 = \frac{j_0}{c}$  and we are only allowed to make gauge transformation for which  $\boxed{\Lambda} = 0$ . In this case we must separate the vector potential into the transverse and longitudinal parts, with

$$\begin{aligned}\vec{A} &= \vec{A}_\perp + \vec{A}_\parallel \\ \vec{\nabla} \cdot \vec{A}_\perp &= 0 \\ \vec{\nabla} \times \vec{A}_\parallel &= 0.\end{aligned}$$

We will now study the radiation field in a region with no sources so that  $\vec{\nabla} \cdot \vec{A} = 0$ . We will use the equations

$$\begin{aligned}\vec{B} &= \vec{\nabla} \times \vec{A} \\ \vec{E} &= -\frac{1}{c} \frac{\partial \vec{A}}{\partial t} \\ \nabla^2 \vec{A} - \frac{1}{c^2} \frac{\partial \vec{A}}{\partial t} &= 0.\end{aligned}$$

## 1.2 Fourier Decomposition of Radiation Oscillators

Our goal is to write the Hamiltonian for the radiation field in terms of a sum of harmonic oscillators Hamiltonians. The first step is to write the radiation field in a simple way as possible, as a sum of harmonic components. We will work in a cubic volume  $V = L^3$  and apply periodic boundary conditions on our electromagnetic waves. We also assume for now that there are no sources inside the region so that we can make a gauge transformation to make  $A_0 = 0$  and hence  $\vec{\nabla} \cdot \vec{A} = 0$ . We decompose the field into its Fourier components at  $t = 0$

$$\vec{A}(\vec{x}, t = 0) = \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \hat{e}^{(\alpha)} \left( c_{k,\alpha}(t = 0) e^{i\vec{k} \cdot \vec{x}} + c_{k,\alpha}^*(t = 0) e^{-i\vec{k} \cdot \vec{x}} \right) \quad (1.5)$$

where  $\hat{e}^{(\alpha)}$  are real unit vectors, and  $c_{k,\alpha}$  is the coefficient of the wave with wave vector  $\vec{k}$  and polarization vector  $\hat{e}^{(\alpha)}$ . Once the wave vector is chose, the two polarization vectors must be picked so that  $\hat{e}^{(1)}$ ,  $\hat{e}^{(2)}$ , and  $\vec{k}$  form a right handed orthogonal system. The components of the wave vector must satisfy

$$k_i = \frac{2\pi n_i}{L} \quad (1.6)$$

due to the periodic boundary conditions. The factor out front is set to normalize the states nicely since

$$\frac{1}{V} \int d^3x e^{i\vec{k} \cdot \vec{x}} e^{-i\vec{k}' \cdot \vec{x}} = \delta_{\vec{k}\vec{k}'} \quad (1.7)$$

and

$$\hat{\epsilon}^{(\alpha)} \cdot \hat{\epsilon}^{(\alpha')} = \delta_{\alpha\alpha'} \quad (1.8)$$

We know the time dependence of the waves from Maxwell's equation,

$$c_{k,\alpha}(t) = c_{k,\alpha}(0) e^{-i\omega t} \quad (1.9)$$

where  $\omega = kc$ . We can now write the vector potential as a function of position and time.

$$\vec{A}(\vec{x}, t) = \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \hat{\epsilon}^{(\alpha)} \left( c_{k,\alpha}(t) e^{i\vec{k} \cdot \vec{x}} + c_{k,\alpha}^*(t) e^{-i\vec{k} \cdot \vec{x}} \right) \quad (1.10)$$

We may need to write this solution in several different ways, and use the best one for the calculation being performed. One nice way to write this is in terms 4-vector  $k_\mu$ , the wave number,

$$k_\mu = \frac{p_\mu}{\hbar} = (k_x, k_y, k_z, ik) = \left( k_x, k_y, k_z, i\frac{\omega}{c} \right) \quad (1.11)$$

so that

$$k_p x_p = k \cdot x = \vec{k} \cdot \vec{x} - \omega t \quad (1.12)$$

We can then write the radiation field in a more covariant way

$$\vec{A}(\vec{x}, t) = \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \hat{\epsilon}^{(\alpha)} \left( c_{k,\alpha}(0) e^{ik_p x_p} + c_{k,\alpha}^*(0) e^{-ik_p x_p} \right) \quad (1.13)$$

A convenient shorthand for calculations is possible by noticing that the second term is just the complex conjugate of the first.

$$\begin{aligned} \vec{A}(\vec{x}, t) &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \hat{\epsilon}^{(\alpha)} \left( c_{k,\alpha}(0) e^{ik_p x_p} + c.c. \right) \\ \vec{A}(\vec{x}, t) &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \hat{\epsilon}^{(\alpha)} c_{k,\alpha}(0) e^{ik_p x_p} + c.c. \end{aligned}$$

Note again that we have made this a transverse field by construction. The unit vectors  $\hat{\epsilon}^{(\alpha)}$  are transverse to the direction of propagation. Also note that we are working in a gauge with  $A_4 = 0$ , so this can also represent

the 4-vector form of the potential. The Fourier Decomposition of the radiation field can be be written very simply.

$$A_\mu = \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} c_{k,\alpha}(0) e^{ik_p x_p} + c.c. \quad (1.14)$$

This choice of gauge makes switching between 4-vector and 3-vector expressions for the potential trivial. Let's verify that this decomposition of the radiation field satisfies the Maxwell equation, just for some practice. It's most convenient to use the covariant form of the equation and field.

$$\begin{aligned} \square A_\mu &= 0 \\ \square \left( \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} c_{k,\alpha}(0) e^{ik_p x_p} + c.c. \right) &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} c_{k,\alpha}(0) \square e^{ik_p x_p} + c.c. \\ &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} c_{k,\alpha}(0) (-k_\nu k_\nu) e^{ik_p x_p} + c.c. = 0 \end{aligned}$$

The result is zero since  $k_\nu k_\nu = k^2 - k^2 = 0$ .

Let's also verify that  $\vec{\nabla} \cdot \vec{A} = 0$

$$\begin{aligned} \vec{\nabla} \cdot \left( \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \hat{\epsilon}^{(\alpha)} c_{k,\alpha}(t) e^{i\vec{k} \cdot \vec{x}} + c.c. \right) &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 c_{k,\alpha}(t) \hat{\epsilon}^{(\alpha)} \cdot \vec{\nabla} e^{i\vec{k} \cdot \vec{x}} + c.c. \\ &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 c_{k,\alpha}(t) \hat{\epsilon}^{(\alpha)} \vec{k} e^{i\vec{k} \cdot \vec{x}} + c.c. = 0. \end{aligned}$$

The result here is zero because  $\hat{\epsilon}^{(\alpha)} \cdot \vec{k} = 0$

### 1.3 The Hamiltonian for the Radiation Field

We now wish to compute the Hamiltonian in terms of the coefficients  $c_{k,\alpha}(t)$ . This is an important calculation because we will use the Hamiltonian formalism to do the quantization of the field. We will do the calculation using the covariant notaion (while Sakurai outlines an alternate calculation using 3-vectors). We have already calculated th Hamiltonian density of a classical EM field.

$$\mathcal{H} = F_{\mu 4} \frac{\partial A_\mu}{\partial x_4} + \frac{1}{4} F_{\mu\nu} F_{\mu\nu} \quad (1.15)$$

$$\begin{aligned} \mathcal{H} &= \left( \frac{\partial A_4}{\partial x_\mu} - \frac{\partial A_4}{\partial x_4} \right) \frac{\partial A_\mu}{\partial x_4} + \frac{1}{4} \left( \frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\mu}{\partial x_\nu} \right) \left( \frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\mu}{\partial x_\nu} \right) \\ \mathcal{H} &= - \frac{\partial A_4}{\partial x_\mu} \frac{\partial A_\mu}{\partial x_4} + \frac{1}{2} \left( \frac{\partial A_\nu}{\partial x_4} \frac{\partial A_\nu}{\partial x_4} - \frac{\partial A_\nu}{\partial x_\nu} \frac{\partial A_\mu}{\partial x_\nu} \right). \end{aligned}$$

Now let's compute the basic element of the above formula for our decomposed radiation field.

$$\begin{aligned}
A_\mu &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} \left( c_{k,\alpha}(0) e^{ik_p x_p} + c_{k,\alpha}^*(0) e^{-ik_p x_p} \right) \\
\frac{\partial A_\mu}{\partial x_\nu} &= \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} \left( c_{k,\alpha}(0) (ik_\nu) e^{ik_p x_p} + c_{k,\alpha}^*(0) (-ik_\nu) e^{-ik_p x_p} \right) \\
\frac{\partial A_\mu}{\partial x_\nu} &= i \frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} \frac{\omega}{c} \left( c_{k,\alpha}(0) e^{ik_p x_p} - c_{k,\alpha}^*(0) e^{-ik_p x_p} \right) \\
\frac{\partial A_\mu}{\partial x_4} &= -\frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} \frac{\omega}{c} \left( c_{k,\alpha}(0) e^{ik_p x_p} - c_{k,\alpha}^*(0) e^{-ik_p x_p} \right).
\end{aligned}$$

We have all the elements to finish the calculation of the Hamiltonian. Before pulling this all together in a brute force way, it's good to realize that almost all the terms will give zero. We see that the derivative of  $A_\mu$  is proportional to a 4-vector, say  $k_\nu$  and to a polarization vector, say  $\epsilon_\mu^{(\alpha)}$ . The dot products of the 4-vectors, either  $k$  with itself, or  $k$  with  $\epsilon$  are zero. Going back to our expression for the Hamiltonian density, we can eliminate some terms.

$$\begin{aligned}
\mathcal{H} &= -\frac{\partial A_\mu}{\partial x_4} \frac{\partial A_\mu}{\partial x_4} + \frac{1}{2} \left( \frac{\partial A_\nu}{\partial x_\mu} \frac{\partial A_\nu}{\partial x_\mu} - \frac{\partial A_\nu}{\partial x_\mu} \frac{\partial A_\mu}{\partial x_\nu} \right) \\
\mathcal{H} &= -\frac{\partial A_\nu}{\partial x_4} \frac{\partial A_\mu}{\partial x_4} + \frac{1}{2} (0 - 0) \\
\mathcal{H} &= -\frac{\partial A_\mu}{\partial x_4} \frac{\partial A_\mu}{\partial x_4}.
\end{aligned}$$

The remaining term has a dot product between polarization vectors which will be nonzero if the polarization vectors are the same. (Note that this simplification is possible because we have assumed no sources in the region.) The total Hamiltonian we are aiming at, is the integral of the Hamiltonian density.

$$H = \int d^3x \mathcal{H} \tag{1.16}$$

When we integrate over the volume only products like  $e^{ik_p x_p} e^{-ik_p x_p}$  will give a nonzero result. So when we multiply one sum over  $k$  by another, only the terms with the same  $k$  will contribute to the integral, basically because the waves with different wave number are orthogonal.

$$\frac{1}{V} \int d^3x e^{ik_p x_p} e^{-ik'_p x_p} = \delta_{kk'} \tag{1.17}$$

$$\begin{aligned}
H &= \int d^3x \mathcal{H} \\
\mathcal{H} &= -\frac{\partial A_\mu}{\partial x_4} \frac{\partial A_\mu}{\partial x_4} \\
\frac{\partial A_\mu}{\partial x_4} &= -\frac{1}{\sqrt{V}} \sum_k \sum_{\alpha=1}^2 \epsilon_\mu^{(\alpha)} \left( c_{k,\alpha}(0) \frac{\omega}{c} e^{ik_p x_p} - c_{k,\alpha}^*(0) \frac{\omega}{c} e^{-ik_p x_p} \right) \\
H &= - \int d^3x \frac{\partial A_\mu}{\partial x_4} \frac{\partial A_\mu}{\partial x_4} \\
H &= - \int d^3x \frac{1}{V} \sum_k \sum_{\alpha=1}^2 \left( c_{k,\alpha}(0) \frac{\omega}{c} e^{ik_p x_p} - c_{k,\alpha}^*(0) e^{-ik_p x_p} \right) \\
H &= - \sum_k \sum_{\alpha=1}^2 \left( \frac{\omega}{c} \right)^2 \left[ -c_{k,\alpha}(t) c_{k,\alpha}^*(t) - c_{k,\alpha}^*(t) c_{k,\alpha}(t) \right] \\
H &= \sum_k \sum_{\alpha=1}^2 \left( \frac{\omega}{c} \right)^2 \left[ c_{k,\alpha}(t) c_{k,\alpha}^*(t) + c_{k,\alpha}^*(t) c_{k,\alpha}(t) \right] \\
H &= \sum_{k,\alpha} \left( \frac{\omega}{c} \right)^2 \left[ c_{k,\alpha}(t) c_{k,\alpha}^*(t) + c_{k,\alpha}^*(t) c_{k,\alpha}(t) \right].
\end{aligned}$$

This is the result we will use to quantize the field. We have been careful not to commute  $C$  and  $C^*$  here in anticipation of the fact that they do not commute.

It should not be a surprise that the terms that made up the Lagrangian gave a zero contribution because  $\mathcal{L} = \frac{1}{2} (E^2 - B^2)$  and we know that E and B have the same magnitude in radiation field. (There is one wrinkle we have glossed over; terms with  $\vec{k}' = -\vec{k}$ .)

## 1.4 Canonical Coordinates and Momenta

We now have the Hamiltonian for the radiation field

$$\boxed{H = \sum_{k,\alpha} \left( \frac{\omega}{c} \right)^2 \left[ c_{k,\alpha}(t) c_{k,\alpha}^*(t) + c_{k,\alpha}^*(t) c_{k,\alpha}(t) \right]} \quad (1.18)$$

It was with the Hamiltonian that we first quantized the non-relativistic motion of particles. The position and momentum became operators which did not commute. Lets define  $c_{k,\alpha}$  to be the time dependent Fourier coefficient.

$$\ddot{c}_{k,\alpha} = -\omega^2 c_{k,\alpha} \quad (1.19)$$

We can then simplify our notation a bit

$$H = \sum_{k,\alpha} \left( \frac{\omega}{c} \right)^2 \left[ c_{k,\alpha} c_{k,\alpha}^* + c_{k,\alpha}^* c_{k,\alpha} \right] \quad (1.20)$$

This now clearly looks like the Hamiltonian for a collection of uncoupled oscillators; one oscillator for each wave vector and polarization.

We wish to write the Hamiltonian in terms of a coordinate for each oscillator and the conjugate momenta. The coordinate should be real so it can be represented by a Hermitian operator and have a physical meaning. The simplest choice for a read coordinates is  $c + c^*$ . With a little effort we can identify the coordinate

$$Q_{k,\alpha} = \frac{1}{c} (c_{k,\alpha} + c_{k,\alpha}^*) \quad (1.21)$$

and its conjugate momentum for each oscillator,

$$P_{k,\alpha} = -\frac{i\omega}{c} (c_{k,\alpha} - c_{k,\alpha}^*) \quad (1.22)$$

The Hamiltonian can be written in terms of these

$$\begin{aligned} H &= \frac{1}{2} \sum_{k,\alpha} \left[ P_{k,\alpha}^2 + \omega^2 Q_{k,\alpha}^2 \right] \\ &= \frac{1}{2} \sum_{k,\alpha} \left[ -\left(\frac{\omega}{c}\right)^2 (c_{k,\alpha} - c_{k,\alpha}^*)^2 + \left(\frac{\omega}{c}\right)^2 (c_{k,\alpha} + c_{k,\alpha}^*)^2 \right] \\ &= \frac{1}{2} \sum_{k,\alpha} \left( \frac{\omega}{c} \right)^2 \left[ - (c_{k,\alpha} - c_{k,\alpha}^*)^2 + (c_{k,\alpha} + c_{k,\alpha}^*)^2 \right]. \end{aligned}$$