## Corregidum to "Novel whitening approaches in functional settings"

Marc Vidal\* 1,2 and Ana M. Aguilera\*2

<sup>1</sup>Ghent University, Belgium <sup>2</sup>University of Granada, Spain

1. Due to production errors, Equation 3 in pp. 3 is written as

$$\langle f,g
angle_{\mathbb{M}}=\sum_{j=1}^{\infty}\lambda_{j}^{-1}\left\langle f,\gamma_{j}
ight
angle \left\langle g,\gamma_{j}
ight
angle =\left\langle \Gamma^{\dagger1/2}f,\Gamma^{\dagger1/2}g
ight
angle f,\quad g\in\mathbb{M},$$

while it was originally written as

$$\langle f,g
angle_{\mathbb{M}} = \sum_{j=1}^{\infty} \lambda_j^{-1} \left\langle f,\gamma_j \right
angle \left\langle g,\gamma_j 
ight
angle = \left\langle \Gamma^{\dagger 1/2} f,\Gamma^{\dagger 1/2} g 
ight
angle \quad f,g \in \mathbb{M}.$$

- 2. Definition 1 states: "The whitening operator  $\Psi$  transforms a functional variable X into a new element  $X = \Psi(X \mu)$  with zero mean and covariance operator being exactly the identity inside H." The notion of "inside" might be somewhat ambiguous, but refers to the idea that whitening is an effective mapping in a closed subspace of H.
- 3. Considering that we work with elements that are already in the range space of  $\Gamma$ , a whitening transformation takes an element from this space to  $\mathbb{M}$  (its closure). Thus,  $\Gamma^{\dagger 1/2} : \operatorname{ran}(\Gamma) \to \mathbb{M}$  instead of  $\Gamma^{\dagger 1/2} : \mathbb{M} \to \mathbb{M}$ . In the case of correlation-based whitening operators, we need to assume X is in  $\operatorname{ran}(\operatorname{span}(e_k)_{k\geq 1})$ ; see comment bellow on the operator  $\mathcal{R}$ .
- 4. In §4, the sentence "As  $2\text{tr}(\Gamma_{XX})$  is the only dependence between the original and the whitened variable, the minimization problem can be reduced to the maximization of  $\text{tr}(\Gamma_{XX})$ ." reads also as ".. is the only *dependent term*...".
- 5. Note that, in §4 the term  $\operatorname{tr}(\Gamma_{\mathbb{X}})$  in the quadratic distances diverges (the trace of  $\Gamma_{\mathbb{X}}$  is an infinite sum of ones). However,  $\operatorname{tr}(\Gamma_{\mathbb{X}})$  is not accounted for in the proof. In order for these distances to converge, one has to consider regularization or finite space dependency. Furthermore, the operator  $\Gamma_{X\mathbb{X}}$  coincides with  $\Gamma^{1/2}$  if  $\Psi \equiv \Gamma^{\dagger 1/2}$ . Note we only know that this operator belongs to the class of Hilbert Schmidt operators (from the trace property of the autocovariance operator), but this fact does not necessarily imply  $\Gamma^{1/2}$  has finite trace. Hence, we implicitly assume that  $\operatorname{tr}(\Gamma^{1/2}) < \infty$  is satisfied.
- 6. In the first paragraph of the technical proofs, some discussion is necessary on how the Picard condition applies to functional random variables. The Picard condition is typically given as in Eq. 2 for a deterministic element  $h \in H$ . We can enforce the Picard condition for almost every realization of  $X: \sum_{j=1}^{\infty} \langle X(\omega), \gamma_j \rangle^2 / \lambda_j < \infty$  for almost every  $\omega$ .

<sup>\*</sup>For correspondence: marc.vidalbadia@ugent.be (M.V.), aaguiler@ugr.es (A.M.A)

This ensures that  $\mathbb{X}(\omega) = \Gamma^{\dagger 1/2} X(\omega)$  is well-defined for typical realizations. We know from Proof of proposition 1 that  $E \| \mathbb{X} \|^2 = \infty$ , since  $\operatorname{tr}(\varGamma_{\mathbb{X}})$  diverges (is an infinite sum of ones). Even if the expected norm diverges, individual realizations  $X(\omega)$  may still satisfy the Picard condition. This reflects a weaker notion of existence, where the random variable  $\mathbb{X}$  is not strongly integrable but still acts meaningfully within the RKHS.

We wrote  $\sum_{j=1}^{\infty} \langle X, \gamma_j \rangle^2 / \lambda_j < \infty$ , where  $\langle X, \gamma_j \rangle^2$  is wrongly understood as a random variable. The same occurs with the following sentence which was originally written in Mas paper for deterministic elements.

7. In the technical proof of Proposition 1, it is stated that "The operator  $P_{\overline{\text{ran}}(\Gamma^{1/2})}$  is compact...". However, this characterization of the projection operator as compact appears to be mistaken since projection operators are typically not compact in infinite-dimensional spaces.

Characterization of 
$$\mathcal{R} = \mathcal{V}^{\dagger 1/2} \Gamma \mathcal{V}^{\dagger 1/2}$$

Suppose *X* is expanded as  $X = \sum_{k=1}^{\infty} \langle X, e_k \rangle e_k$  and note the following:

$$\mathcal{V} = \sum_{k=1}^{\infty} \mathcal{P}_{e_k} \Gamma \mathcal{P}_{e_k} = \sum_{k=1}^{\infty} \langle \Gamma e_k, e_k \rangle \mathcal{P}_{e_k} = \sum_{k=1}^{\infty} E(\langle X, e_k \rangle^2) (e_k \otimes e_k) = \sum_{k=1}^{\infty} \eta_k (e_k \otimes e_k). \tag{1}$$

Now, consider the operator

$$\mathcal{R} \equiv E\left(\sum_{j=1}^{\infty} \frac{\langle X, e_j \rangle}{\eta_j^{1/2}} e_j \otimes \frac{\langle X, e_j \rangle}{\eta_j^{1/2}} e_j\right). \tag{2}$$

Observe that  $\mathcal{V}^{\dagger 1/2} \Gamma \mathcal{V}^{\dagger 1/2} = \sum_{j=1}^{\infty} \eta_j^{-1/2} E(\mathcal{P}_{e_j} X \otimes \mathcal{P}_{e_j} X) \eta_j^{-1/2}$ , where  $\mathcal{V}^{\dagger}$  is the Moore-Penrose inverse of  $\mathcal{V} = \sum_{j=1}^{\infty} \eta_j \mathcal{P}_{e_j}$ . This shows that  $\mathcal{V}^{1/2\dagger} \Gamma \mathcal{V}^{1/2\dagger}$  is equivalent to the operator  $\mathcal{R}$ , as defined in 2. Note this operator bears resemblance to the classical correlation matrix in the multivariate setting.

## REFERENCES

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