

Thèse de doctorat:

le long titre de ma thèse

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Abstract

This will be the thesis abstract

Acknowledgments

And here some acknowledgments

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General introduction

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1.1 | Objectifs du travail

1.2 | Organisation de la thèse

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THIS CHAPTER IS BASED ON THE MATERIAL PUBLISHED IN REF. MINOR MODIFICATIONS TO THE TEXT AND FIGURES HAVE BEEN MADE TO FIT THE FORMAT AND NOTATION OF THIS THESIS.

2.1 | Abstract

2.2 | Intro de l'article

2.2.1. sous partie 1

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2.3 | Une autre partie

A general spectral collocation method for computing the dispersion relations of guided acoustic waves in multilayer dissipative structures

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3.1 | abstract

A spectral collocation method is proposed to compute the complex wavenumber-real frequency dispersion relations of guided acoustic waves in multilayer structures involving dissipative materials. The nature of these dissipative materials is initially considered to be arbitrary, i.e., poroelastic, viscoelastic, or viscoacoustic. For a given frequency, the complex wavenumbers as well as the physical fields, which are further used to evaluate the Poynting vectors and analyze the energy flux, are obtained by solving a generalized eigenvalue problem. The latter arises from a set of discretized equations of motion and appropriate boundary (coupling) conditions. These equations of motion and boundary (coupling) conditions are imposed by the nature of the material composing each layer of the structure. A focus is made on poroelastic layers. The dispersion relation of a two-layer elastic-poroelastic structure is analyzed, as well as the energy flows in the structure. The results as calculated with the present spectral collocation method are validated against those obtained with a classical complex root-finding (Müller) method and experiments.

3.2 | Introduction

Dealing with guided acoustic waves in multilayer structures of dissipative (viscoelastic, viscothermal, or poroelastic) material requires the evaluation of complex wavenumber-real frequency dispersion relations. Solving for these dispersion relations consists in finding complex roots of complex characteristic equations, which often arise from complex matrix determinants and can rarely be solved analytically. Various methods [1, 2], among which the Müller algorithm [3], are commonly used to achieve this challenging task. These methods usually rely on preliminary semi-analytical derivations, require multiple initial guesses, and are iterative, so they can get stuck in local minima. Computing the complex wavenumber-real frequency dispersion relation in a reliable and robust way is thus of primary importance. Several numerical methods can be used, such as the Finite Element Method [4], the Semi-Analytical Finite Element method (SAFE) [5–7], or spectral methods [8]. Spectral methods are especially relevant in this framework. They are numerical approximation techniques that seek the solution of a differential equation using a finite series of infinitely differentiable basis functions with *spectral accuracy* [9].

The propagation and dispersion of waves in multilayer systems have long been investigated [10, 11]. Citer [12, 13], plutôt parler de TMM ou présenter logiciels existants (DISPERSE, dispersion box, GEWtool, etc.) ? Spectral methods have been developed to cope with wave radiation problems, but have also emerged in the last decades to solve dispersion relations [14]. Effectively, the numerical scheme boils down to discretizing the geometry along the layer thickness and solving an eigenvalue problem. Here, a collocation method is used to discretize the geometry. A finite-dimension polynomial basis is chosen and the collocation points are distributed according to the roots of this polynomial basis. Leaky modes appear in practice when multilayer structures are surrounded by one or two half-spaces. Accounting for this phenomenon brings us closer to practical considerations regarding these systems, as they would radiate part of the acoustic energy back to the surrounding medium, although it adds hardships to the modeling of the system; the waves modeled in these half-spaces should satisfy some radiation condition. These can be imposed by implementing perfectly matched layers [15], by using a bounded mapping, and applying a Spectral Collocation Method (SCM) to all layers of the extended system [16], or by using a suitable change of variable and solving a non-linear eigenvalue problem [17]. This last method inspires the current work. We propose a SCM suitable for solving complex wavenumber-real frequency dispersion relations in multilayer systems involving dissipative media. As a way of example, we will focus on multilayer poroelastic systems. Mechanical wave propagation in bulk poroelastic materials are well understood and modeled thanks to the seminal contributions of M.A. Biot [18–21] and subsequent contributions, notably by D.L. Johnson *et al.* [22] and J.-F. Allard *et al.* [23, 24] concerning the modeling of viscothermal losses. Although poroelastic materials are bounded in practice and are usually encapsulated in more complex multilayer structures combining layers of materials of different nature in geophysics and acoustics, there are only a few studies on the dispersion relation of guided acoustic waves in poroelastic [25–27], porous [28, 29] or anisotropic poroelastic layers [30]. The SCM has already been used to solve the dispersion relation of guided acoustic waves in elastic layered rings [31, 32] or anisotropic elastic multilayer systems [33], but has never been used to cope with multilayer systems involving dissipative materials of arbitrary nature.

This work aims at filling this gap and is divided into two sections. First, the SCM is derived for generic multilayer configurations possibly coupled to two surrounding identical fluid half-spaces. A two-layer elastic-poroelastic structure is then considered by way of example. The dispersion relation computed with the SCM method is compared to that obtained with a classical root-finding (Müller) method and that measured. The energy flow is further analyzed for this two-layer configuration, with very little additional computational cost. In the last section, we experimentally analyze the two-layer structure by retrieving its complex dispersion relation from displacement measurements using the SLaTCoW (**Spatial LAPlace Transform for COmplex Wavenumber recovery**). The previously obtained numerical results can thus be correlated with the data of the dispersion relation.

3.3 | Solving the dispersion relations for multilayer structures

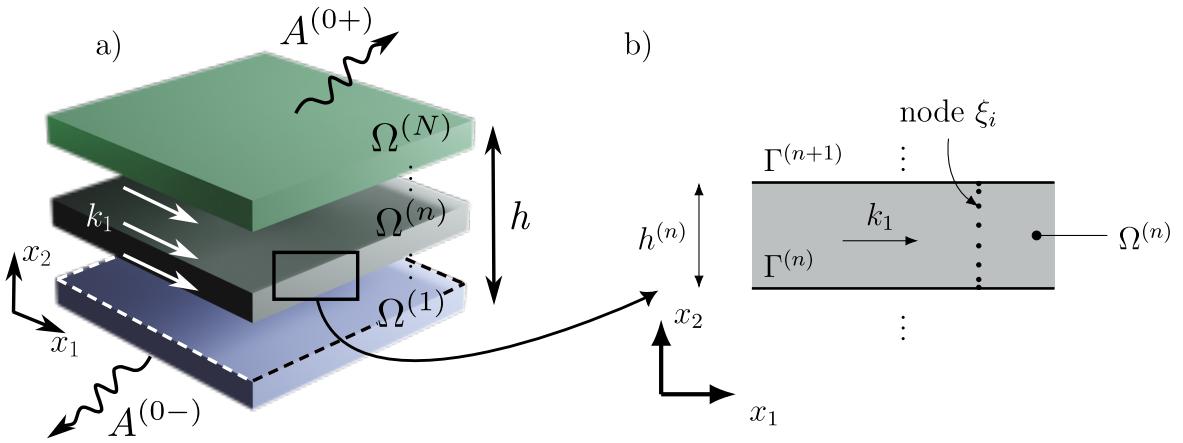


Figure 3.1: a) Sketch of the geometry of the multilayer system and b) zoom on the n -th layer.

3.3.1. Spectral Collocation Method

Figure 3.1a depicts the general geometry of the multilayer structure considered in this article. It is composed of N homogeneous layers of arbitrary nature. The n -th layer has a thickness of $h^{(n)}$ and occupies the domain $\Omega^{(n)}$. The lower and upper interfaces are denoted $\Gamma^{(n-1)}$ and $\Gamma^{(n)}$. The multilayer structure is assumed to be two-dimensional, of infinite extent along x_1 , and bounded along the x_2 -axis between 0 and the total structure thickness $b = \sum_{n=1}^N h^{(n)}$. Outgoing waves radiate above and below the structure in two fluid half-spaces $\Omega^{(0\pm)}$, occupied by the same fluid.

We assume an implicit time dependence $e^{-i\omega t}$. The physical fields associated to the guided waves in the n -th layer propagating along the positive x_1 -axis take the form

$$\Theta^{(n)}(\mathbf{x}) = s^{(n)}(k_1, x_2) e^{ik_1 x_1}, \quad (3.1)$$

where $s^{(n)}(k_1, x_2)$ is the spatial Laplace transform of $\Theta^{(n)}(\mathbf{x})$ and k_1 is the complex valued x_1 -component of the wavenumber. Note that the spatial Laplace transform is used instead of the spatial Fourier transform because k_1 is complex-valued. The x_1 component of the wavenumber is complex-valued. The physical fields that compose $\Theta^{(n)}(\mathbf{x})$ depend on the nature of the material occupying the n -th layer. Introducing $\Theta^{(n)}(\mathbf{x})$ in the corresponding wave equations (of second order in space in this work) reduces to a second-order differential operator $\mathcal{L}^{(n)}$ in x_2 , that depends on the material properties and can further be developed as a second-order polynomial in k_1

$$\mathcal{L}^{(n)}(x_2) s^{(n)}(k_1, x_2) = \left(k_1^2 \mathcal{L}_2^{(n)}(x_2) + k_1 \mathcal{L}_1^{(n)}(x_2) + \mathcal{L}_0^{(n)}(x_2) \right) s^{(n)}(k_1, x_2) = \mathbf{0}, \quad (3.2)$$

where $\mathcal{L}_j^{(n)}$ are the coefficients of the polynomial expansion $j = 0, 1, 2$. Please note that this operator is nothing but the spatial Laplace transform of the equations of motion.

The interface (boundary) conditions at $\Gamma^{(n)}$ between the $n+1$ -th and the n -th layers again depend on the nature of these two layers, but can be formally written in the form

$$C^{(n)-} s^{(n)}(k_1) \Big|_{\Gamma^{(n)}} - C^{(n)+} s^{(n+1)}(k_1) \Big|_{\Gamma^{(n)}} = \mathbf{0}, \quad (3.3)$$

where $C^{(n)\pm}$ are interfaces operators, which involve differential operators of maximum second-order in x_2 . Following this procedure and introducing $\mathbf{S} = (s^{(0-)} \dots s^{(n)} \dots s^{(0+)})^\top$, with \top denoting transposition, the

problem can be cast in matrix form, whose coefficients are differential operators in x_2 ,

$$\mathbf{KS}(k_1, x_2) = \begin{pmatrix} \mathbf{C}^{(0^-)} & -\mathbf{C}^{(0^+)} & 0 \\ 0 & \mathcal{L}^{(1)}(x_2) & 0 \\ 0 & \mathbf{C}^{(1^-)} & -\mathbf{C}^{(1^+)} \\ & 0 & \mathcal{L}^{(2)}(x_2) \\ & & \ddots & \ddots & \ddots \\ & & 0 & \mathcal{L}^{(N)}(x_2) & 0 \\ & & 0 & \mathbf{C}^{(N^-)} & -\mathbf{C}^{(N^+)} \end{pmatrix} \mathbf{S}(k_1, x_2) = \mathbf{0}. \quad (3.4)$$

Note that the equations of motion in both upper and lower half-spaces are not required in Eq.(3.4). This system only depends on the values of the two spatial Laplace transforms $s^{(0^\pm)}$ at the upper and lower interfaces. Instead, we specify the nature of the two half-spaces, i.e. a fluid medium, and make use of the Sommerfeld condition, to explicitly give the form of the spatial Laplace transforms $\tilde{p}^{(0^\pm)}(k_1, x_2)$ of the pressure fields $p^{(0^\pm)}(x)$

$$\begin{aligned} \tilde{p}^{(0^-)}(k_1, x_2) &= A^{(0-)} e^{-ik_2^{(0)} x_2}, \in \Omega^{(0^-)}, \\ \tilde{p}^{(0^+)}(k_1, x_2) &= A^{(0+)} e^{ik_2^{(0)} (x_2 - b)}, \in \Omega^{(0^+)}, \end{aligned} \quad (3.5)$$

with $k_2^{(0)} = \sqrt{(k^{(0)})^2 - k_1^2}$, such that $\text{Re}(k_2^{(0)}) \geq 0$, where $k^{(0)} = \omega/c_f$ is the wavenumber in the fluid and c_f the wave velocity. Thus, $s^{(0^\pm)}$ appearing in Eq.(3.4) reduces to $s^{(0^\pm)} = A^{(0\pm)} \cdot \mathbf{C}^{(0^-)}$ and $\mathbf{C}^{(N^+)} = \mathbf{C}^{(N^+)}$. In our fluid case, the interface conditions at $\Gamma^{(0)}$ and $\Gamma^{(N)}$ involve the spatial Laplace transform of the pressure and its first spatial derivative with respect to x_2 whatever the nature of the 1-st and N -th layers. Both matrices $\mathbf{C}^{(j)}$, $j = (0^-), (N^+)$, can thus be formally written as

$$\mathbf{C}^{(j)} = k_1 \mathbf{C}_1^{(j)} + \mathbf{C}_0^{(j)} + k_2^{(0)} \mathbf{C}_{1'}^{(j)}, \quad (3.6)$$

with the subscript denoting the submatrices rearranged according to the power order $k1$. Index $1'$ corresponds to the term in $k_2^{(0)}$. Introducing Eqs.(3.2) and (3.6) in Eq. (3.4) leads to

$$(k_1^2 \mathbf{K}_2 + k_1 \mathbf{K}_1 + \mathbf{K}_0 + ik_2^{(0)} \mathbf{K}_{1'}) \mathbf{S} = \mathbf{0}. \quad (3.7)$$

This eigenvalue problem is non-linear because of the presence of the term $k_2^{(0)}$. Making use of the following changes of variable $k_1 = k^{(0)} (\gamma + \gamma^{-1}) / 2$ and $k_2^{(0)} = k^{(0)} (\gamma - \gamma^{-1}) / 2i$ [34, 35] and companion linearization [36], a generalized eigenvalue problem is formed and solved for γ , from which the dispersion relation can be calculated. More details on this procedure can be found in [17]. Note that the first or the last column of Eq.(3.4) vanishes when radiation is absent on either side of the multilayer system. The situation is different when both fluid half-spaces disappear. The eigenvalue problem Eq.(3.7) is no longer non-linear and the problem can directly be solved after companion linearization. Note also that fluid half-spaces are only considered for the sake of simplicity and that this procedure can be adapted to any kind of material occupying the half-spaces, as described recently [37]. Note finally that this procedure avoids discretization of these two half-spaces.

To numerically solve the problem, the spatial Laplace transform of the physical fields in each n -th layer are expanded as

$$\underline{s}^{(n)}(x_2) \approx \underline{s}^{(n)}(x_2) = \sum_{m=0}^M \alpha_m^{(n)} \underline{\psi}_m(\xi^{(n)}), \quad (3.8)$$

where α_m are coefficients of the polynomial expansion, $\underline{\psi}_m$ are the Chebyshev polynomial of order m , and $\xi^{(n)} = 2(x_2 - \sum_{j=1}^{n-1} b^{(j)}) / b^{(n)} - 1$, $\in [-1, 1]$. Underlined variables represent discrete vectors and double-underlined variables discrete matrices. The goal of this approximation is to discretize the differential operators in x_2 appearing in the problem. To do so, the SCM is employed. $\xi^{(n)}$ is discretized on the roots of the Chebyshev polynomials, $\xi_j^{(n)} = \cos\left(\frac{j\pi}{M^{(n)}}\right)$ with $j = 0 \dots M^{(n)}$. Please note that $\xi_j^{(n)}$ discretely runs from 1 to -1 with increasing j , while $\xi^{(n)}$ continuously runs over the same interval but in the opposite direction, i.e. from -1 to 1. These collocation points/nodes form a non-uniform grid along the thickness of each layer.

This discretization is usually preferred because the collocation points are clustered at the interfaces [9]. The locations of these nodes are presented in Fig. 3.1b for the n -th layer with an arbitrary value $M^{(n)} = 8$. The differential operators in the n -th layer are thus represented by differentiation matrices (DMs), such that

$$\begin{aligned} \underline{s}^{(n)}(x_2) &\rightarrow \underline{\underline{I}} \underline{s}^{(n)}(\xi_j), \quad \partial_2 \underline{s}^{(n)}(x_2) \rightarrow \left(\frac{2}{b^{(n)}} \right) \underline{\underline{D}_2} \underline{s}^{(n)}(\xi_j) = \underline{\underline{D}_2}^{(n)} \underline{s}^{(n)}(\xi_j), \\ \partial_{22} \underline{s}^{(n)}(x_2) &\rightarrow \left(\frac{2}{b^{(n)}} \right)^2 \underline{\underline{D}_{22}} \underline{s}^{(n)}(\xi_j) = \underline{\underline{D}_{22}}^{(n)} \underline{s}^{(n)}(\xi_j). \end{aligned} \quad (3.9)$$

where $\underline{\underline{I}}$ is the identity matrix and $\underline{\underline{D}_2}$ and $\underline{\underline{D}_{22}}$ are respectively the first-order and second-order normalized DMs. Contrary to an interpolation with Chebyshev polynomials, DMs interpolate the functions and their spatial derivatives directly and thus do not involve the polynomial expansion coefficients per se. [38]. Note that the normalized DMs are denoted as $\underline{\underline{D}_2}^{(n)}$ and $\underline{\underline{D}_{22}}^{(n)}$ in the following. These matrices are square and have a size of $(M \times M)$. This size also directly depends on the order of the considered Chebyshev polynomials. The formulation prevents the occurrence of roundoff errors, which is of particular importance when SCM is applied to solve eigenvalue problems [39].

The linear operator $\underline{\underline{L}}^{(n)}$ is discretized as a matrix $\underline{\underline{L}}^{(n)}$ with size $(P^{(n)}(M^{(n)} - 1), P^{(n)}(M^{(n)} + 1))$, where $P^{(n)}$ is the number of physical fields required to describe the wave propagation in the layer and again depends on the nature of the material this layer is composed of. The first and the last rows of the DMs are removed because they are used to apply the interface conditions at $\Gamma^{(n+1)}$ and $\Gamma^{(n)}$ respectively. The interface operators $\underline{\underline{C}}^{(n)\pm}$ are effectively expressed in terms of this first row of the DMs of the n -th layer and this last row of the DMs of the $(n + 1)$ -th layer. These rows respectively correspond to the first, i.e. $\xi_{M^{(n)}} = -1$ and $\xi^{(n)} = 1$, and last collocation nodes of the n -th and $(n + 1)$ -th layers, i.e. $\xi_0 = 1$ and $\xi^{(n+1)} = -1$. These vectors are subsequently denoted $\underline{\underline{I}}^{(n+1)-}, \underline{\underline{I}}^{(n)+}, \underline{\underline{D}_2}^{(n+1)-}, \underline{\underline{D}_2}^{(n)+}, \underline{\underline{D}_{22}}^{(n+1)-}$, and $\underline{\underline{D}_{22}}^{(n)+}$. The interface matrix $\underline{\underline{C}}^{(n)-}$ is of size $(P^{(n)} + P^{(n+1)}, P^{(n)}(M^{(n)} + 1))$ and the interface matrix $\underline{\underline{C}}^{(n)+}$ is of size $(P^{(n)} + P^{(n+1)}, P^{(n+1)}(M^{(n+1)} + 1))$.

Finally, the full $(\sum_{n=0}^N M^{(n)} P^{(n)} + 2, \sum_{n=0}^N M^{(n)} P^{(n)} + 2)$ -matrix system for the general multilayer system is

$$\underline{\underline{K}} \underline{\underline{S}} = \begin{pmatrix} \underline{\underline{C}}^{(0-)} & -\underline{\underline{C}}^{(0+)} & 0 \\ 0 & \underline{\underline{L}}^{(1)} & 0 \\ 0 & \underline{\underline{C}}^{(1-)} & -\underline{\underline{C}}^{(1+)} \\ & 0 & \underline{\underline{L}}^{(2)} \\ & & \ddots & \ddots & 0 \\ & & & 0 & \underline{\underline{L}}^{(N)} \\ 0 & & & 0 & \underline{\underline{C}}^{(N-)} & -\underline{\underline{C}}^{(N+)} \end{pmatrix} \begin{pmatrix} \underline{\underline{A}}^{(0-)} \\ \underline{\underline{s}}^{(1)} \\ \underline{\underline{s}}^{(2)} \\ \vdots \\ \underline{\underline{s}}^{(N)} \\ \underline{\underline{A}}^{(0+)} \end{pmatrix} = \underline{\underline{0}}. \quad (3.10)$$

This system is cast in the form of a generalized eigenvalue problem that is solved by traditional eigenvalue solvers at each frequency. The dispersion relation for the acoustic guided wave in the dissipative multilayer system is calculated, repeating the procedure for each frequency ω .

3.3.2. Solving the dispersion relation with Müller algorithm

The dispersion relation of multi-layer dissipative structure are more commonly calculated with secant-based algorithm as the Müller algorithm. The latter relies on a semi-analytic description of the fields. These fields are best written in each layer, independently of the nature of the material it is composed of, using the potentials, i.e. decomposing the field $\theta = \nabla\phi + \nabla \times \psi$ into an irrotational component ϕ and a transverse component $\psi = \psi e_3$. The spatial Laplace transform of the potentials in the n -th layer are entirely described by

$$\theta_n(k_1, x_2) = A_X^{(n)+} e^{ik_{2,X}^{(n)}(x_2 - b^{(n)})} + A_X^{(n)-} e^{-ik_{2,X}^{(n)}(x_2 - b^{(n)})}, \quad (3.11)$$

with $A_X^{(n)\pm}$ the up and down-going wave amplitudes, $k_{2,X}^{(n)} = \sqrt{(k_X^{(n)})^2 - k_1^2}$, such that $\text{Re}(k_{2,X}^{(n)}) \geq 0$, and X refers to the type of wave, i.e. shear, compressional or acoustic. If necessary, these Laplace transform representations are complemented by those of the pressure field in the fluid half-spaces provided in Eq.(3.5).

The boundary conditions are applied at each interface and the set of equations is cast into the following matrix form $\mathbf{M}(k_1)\mathbf{q} = 0$, with \mathbf{q} a vector containing the amplitudes of each field/potential. The modes of the system correspond to the wavenumbers k_1 for which the matrix $\mathbf{M}(k_1)$ is singular, i.e.

$$\det(\mathbf{M}(k_1)) = 0. \quad (3.12)$$

These wavenumbers are the complex roots of an implicit complex-valued equation. They are evaluated by running our own implementation of the Müller algorithm [3] from 3 initial guesses. This equation is solved iteratively for each frequency ω , giving the dispersion relation.

In addition, \mathbf{q} can be obtained by computing the nullspace of \mathbf{M} [40] (This null space is computed using singular value decomposition as embedded in the `linalg.null_space` function from the Python package SciPy). The amplitudes of the potentials in the layers are thus evaluated, enabling all the physical fields that depend on these potentials to be calculated. However, the \mathbf{M} matrix becomes increasingly tedious to write as the number of layers increases, and even more so when it comes to calculating the roots of its determinant and nullspaces. Nevertheless, this approach remains a fairly effective way of providing a reference solution for validating numerical results when studying a two-layer system.

3.4

Application to a two-layer poroelastic-elastic structure

In this section, the calculation of the dispersion relation for guided waves in a two-layer structure, i.e. $N = 2$, is used as an example of the application of the present SCM. This two-layer structure consists of a $h^{(1)}$ -thick poroelastic layer, coated on one side by a $h^{(2)}$ -thick aluminum plate, that is radiating in a fluid half-space, and rigidly backed on the other side as depicted in Fig. 3.2a. The SCM results are analyzed and compared to those obtained with the Müller root-finding method.

3.4.1. Numerical scheme

The pressure field $p^{(1)}$ and the two components of the frame displacement, $u_{1,s}^{(1)}$ and $u_{2,s}^{(1)}$, i.e. $P^{(1)} = 3$, as well as the two-components of the elastic displacement, $u_1^{(2)}$ and $u_2^{(2)}$, i.e. $P^{(2)} = 2$, are used to model the propagation in the poroelastic and in the elastic layers respectively. The $\{\mathbf{u}_s, p\}$ formulation is used [41] for the poroelastic medium. The more usual displacement formulations ($\{\mathbf{u}_s, \mathbf{u}_f\}$ [18] or $\{\mathbf{u}_s, \mathbf{w}\}$ [20]) cannot be used in the present case, because the number of physical fields they rely on, i.e. 4, is different from that of the potentials in the layer, i.e. 3.

In total, 5 physical fields inside the two-layer structure and an additional amplitude for the acoustic wave radiated towards $x_2 \rightarrow \infty$ are needed to solve for the dispersion relation. Each domain is discretized on $\mathcal{M}^{(j)}$, $j = 1, 2$, collocation points.

In more details, the $\{\mathbf{u}_s, p\}$ formulation of the Biot theory [41] is employed to model the propagation in the homogeneous isotropic poroelastic layer. The coupled equations of motion read as

$$\begin{cases} \nabla \cdot \hat{\sigma}_s^{(1)} + \hat{\rho}\omega^2 \mathbf{u}_s^{(1)} + \beta \nabla p^{(1)} = 0, \\ \frac{\nabla^2 p^{(1)}}{\rho_{22}\omega^2} - \frac{\beta}{\phi^2} \nabla \cdot \mathbf{u}_s^{(1)} + \frac{1}{R} p^{(1)} = 0, \end{cases} \quad (3.13)$$

where $\hat{\sigma}_s^{(1)} = \hat{A} \nabla \cdot \mathbf{u}_s^{(1)} \mathbf{I} + 2N_s \epsilon_s$ is the *in vacuo* solid stress tensor, with ϵ_s , the solid-phase strain tensor. The parameters are generally expressed in terms of the elastic coefficients P , Q and R and the effective densities ρ_{11} , ρ_{12} and ρ_{22} initially introduced by Biot [18, 42] as

$$\hat{A} = P - 2N_s - \frac{Q^2}{R}, \quad \hat{\rho} = \rho_{11} - \frac{\rho_{12}^2}{\rho_{22}}, \quad \beta = \phi \left(\frac{\rho_{12}}{\rho_{22}} - \frac{Q}{R} \right). \quad (3.14)$$

These parameters depend on the frame properties and on the effective density $\rho_{eq}(\omega)$ and bulk modulus $K_{eq}(\omega)$ of the fluid phase as reminded in 3.6.A. In short, a poroelastic material is entirely described by the

porosity ϕ , the shear modulus of the frame N_s , the Poisson ratio ν , the density of the frame ρ_1 , the viscous and thermal characteristic lengths Λ and Λ' , tortuosity α_∞ , and the flow resistivity R_f when saturated by a light fluid [18, 22, 23]. All these parameters are real valued except for the shear modulus, which incorporates a viscoelastic damping factor.

Introducing $s^{(1)} = (\pi_1^s, \pi_2^s, p)^\top$, the spatial Laplace transform of Eq. (3.13) is expanded in the form of Eq. (3.2) with

$$\begin{aligned}\underline{\mathcal{L}}_0^{(1)} &= \begin{pmatrix} N_s \partial_{22} + \hat{\rho} \omega^2 & 0 & 0 \\ 0 & \hat{P} \partial_{22} + \hat{\rho} \omega^2 & \beta \partial_2 \\ 0 & -\beta/\phi^2 \partial_2 & 1/R + \partial_{22}/(\rho_{22} \omega^2) \end{pmatrix}, \\ \underline{\mathcal{L}}_1^{(1)} &= \begin{pmatrix} 0 & i\hat{P} \partial_2 & i\beta \\ i(\hat{A} + N_s) \partial_2 & 0 & 0 \\ -i\beta/\phi^2 & 0 & 0 \end{pmatrix}, \text{ and } \underline{\mathcal{L}}_2^{(1)} = \begin{pmatrix} \hat{P} & 0 & 0 \\ 0 & N_s & 0 \\ 0 & 0 & 1/(\rho_{22} \omega^2) \end{pmatrix}.\end{aligned}\quad (3.15)$$

Finally, the discretized equations of motion read as

$$\underline{\underline{L}}^{(1)} \underline{s}^{(1)} = \begin{pmatrix} (\omega^2 \hat{\rho} - k_1^2 \hat{P}) \underline{\underline{I}} + N_s \underline{\underline{D}}_{22} & i k_1 (\hat{P} - N_s) \underline{\underline{D}}_{22} & -i k_1 \beta \underline{\underline{I}} \\ i k_1 (\hat{P} - N_s) \underline{\underline{D}}_2 & \hat{P} \underline{\underline{D}}_{22} + (\omega^2 \hat{\rho} - k_1^2 N_s) \underline{\underline{I}} & \beta \underline{\underline{D}}_2 \\ -i \beta \omega^2 \underline{\underline{I}} & -\beta \omega^2 \underline{\underline{D}}_2 & \frac{k_1^2}{\rho_{22}} \underline{\underline{I}} + \frac{\omega^2}{R} \underline{\underline{I}} \end{pmatrix} \underline{s}^{(1)} = 0. \quad (3.16)$$

The equation of motion in the elastic medium is derived from the fundamental elasticity equations [43] and reads as,

$$\rho \omega^2 \mathbf{u}^{(2)} + (\lambda + \mu) \nabla (\nabla \cdot \mathbf{u}^{(2)}) + \mu \nabla^2 \mathbf{u}^{(2)} = 0, \quad (3.17)$$

where ρ is the density and λ and μ are the Lamé coefficients. The corresponding discretized equation of motion is

$$\underline{\underline{L}}^{(2)} \underline{s}^{(2)} = \left(\frac{\rho^{(2)}}{\mu} \omega^2 \underline{\underline{I}} + \begin{pmatrix} -k_1^2 (\lambda/\mu + 2) \underline{\underline{I}} + \underline{\underline{D}}_{22} & i k_1 (\lambda/\mu + 1) \underline{\underline{D}}_2 \\ i k_1 (\lambda/\mu + 1) \underline{\underline{D}}_2 & (\lambda/\mu + 2) \underline{\underline{D}}_{22} - k_1^2 \underline{\underline{I}} \end{pmatrix} \right) \underline{s}^{(2)} = 0. \quad (3.18)$$

The interface and boundary conditions depend on the vectors $\underline{\underline{I}}^\pm$, $\underline{\underline{D}}_2^\pm$ and $\underline{\underline{D}}_{22}^\pm$ which corresponds to the last (subscript +) and first row (subscript -) of the DMs. From the bottom to the top of the two-layer system, they are

- Rigid boundary condition at $\Gamma^{(0)}$ which leads to

$$\mathbf{u}^{(1)}(x_2 = 0) = 0; \quad \mathbf{u}_f^{(1)}(x_2 = 0) \cdot \mathbf{n} - \mathbf{u}_s^{(1)}(x_2 = 0) \cdot \mathbf{n} = 0, \quad (3.19)$$

the discretized form of which is,

$$\underline{\underline{C}}^{(0)} \underline{s}^{(1)} = \begin{pmatrix} 0 & \underline{\underline{I}}^+ & 0 \\ 0 & 0 & \underline{\underline{I}}^+ \\ \frac{\phi}{\rho_{22}} \underline{\underline{D}}_2^+ & 0 & -\left(1 + \frac{\rho_{12}}{\rho_{22}}\right) \omega^2 \underline{\underline{I}}^+ \end{pmatrix} \begin{pmatrix} \underline{p}^{(1)} \\ \underline{\underline{u}}_1^{(1)} \\ \underline{\underline{u}}_2^{(1)} \end{pmatrix} = 0. \quad (3.20)$$

- Continuity between the poroelastic and the elastic media [44] at $\Gamma^{(1)}$ which leads to

$$\begin{aligned}\sigma^t(1)(x_2 = b^{(1)}) - \sigma^t(2)(x_2 = b^{(1)}) &= 0; \quad \mathbf{u}^{(2)}(x_2 = b^{(1)}) = \mathbf{u}_s^{(1)}(x_2 = b^{(1)}); \\ \mathbf{u}_f^{(1)}(x_2 = b^{(1)}) \cdot \mathbf{n} - \mathbf{u}_s^{(1)}(x_2 = b^{(1)}) \cdot \mathbf{n} &= 0,\end{aligned}\quad (3.21)$$

where $\sigma^t(1) = \sigma_s - \phi(1 + Q/R)$ is the total stress tensor in the poroelastic layer. The corresponding discretized interface matrices, arising from

$$\underline{\underline{C}}^{(1)+} \underline{s}^{(2)} - \underline{\underline{C}}^{(1)-} \underline{s}^{(1)} = \underline{\underline{C}}^{(1)} \begin{pmatrix} \underline{p}^{(1)} & \underline{\underline{u}}_1^{(1)} & \underline{\underline{u}}_2^{(1)} & \underline{\underline{u}}_1^{(2)} & \underline{\underline{u}}_2^{(2)} \end{pmatrix}^\top = 0, \quad (3.22)$$

are cast as

$$\underline{\underline{C}}^{(1)} = \begin{pmatrix} -N_s \underline{D}_2^+ & -ik_1 N_s \underline{I} & 0 & \mu \underline{D}_2^- & ik_1 \mu \underline{I}^- \\ -ik_1 \hat{A} \underline{I}^+ & -\hat{P} \underline{D}_2^+ & \phi \left(1 + \frac{Q}{R} \right) \underline{I}^+ & ik_1 \lambda \underline{I}^- & (\lambda + 2\mu) \underline{D}_2^- \\ 0 & -\underline{I}^+ & 0 & \underline{I}^- & 0 \\ 0 & 0 & -\underline{I}^+ & 0 & \underline{I}^- \\ -\frac{\phi}{\rho_{22}} \underline{D}_2^+ & \left(1 + \frac{\rho_{12}}{\rho_{22}} \right) \omega^2 \underline{I}^+ & 0 & 0 & 0 \end{pmatrix}, \quad (3.23)$$

with the left-hand side corresponding to $\underline{\underline{C}}^{(1)-}$ and the right-hand side to $\underline{\underline{C}}^{(1)+}$.

- Continuity between the elastic and the fluid media at $\Gamma^{(2)}$ which leads to

$$\begin{aligned} \sigma_{12}^{(2)}(x_2 = b) &= 0; & \sigma_{22}^{(2)}(x_2 = b) &= -p^{(0+)}; \\ u_2^{(2)}(x_2 = b) &= u_2^{(0+)}(x_2 = b). \end{aligned} \quad (3.24)$$

The corresponding discretized interface matrix $\underline{\underline{C}}^{(2)} \begin{pmatrix} \underline{u}_1^{(2)} & \underline{u}_2^{(2)} & A^{(0+)} \end{pmatrix}^T = 0$ is

$$\underline{\underline{C}}^{(2)} = \begin{pmatrix} \underline{D}_2^+ & ik_1 \underline{I}^+ & 0 \\ ik_1 \lambda \underline{I}^+ & (\lambda + 2\mu) \underline{D}_2^+ & 1 \\ 0 & \omega^2 \underline{I}^+ & \frac{ik_2^{(0)}}{\rho_f} \end{pmatrix}. \quad (3.25)$$

Combining the discretized interface conditions, Eqs.(3.20), (3.23) and, (3.25), with the discretized equations of motion, Eqs.(3.16) and (3.18), leads to the following system

$$\begin{pmatrix} \underline{\underline{C}}^{(0)} & 0 & 0 \\ \underline{\underline{L}}^{(1)} & 0 & 0 \\ \underline{\underline{C}}^{(1)-} & -\underline{\underline{C}}^{(1)+} & 0 \\ 0 & \underline{\underline{L}}^{(2)} & 0 \\ 0 & \underline{\underline{C}}^{(2)-} & -\underline{\underline{C}}^{(2)+} \end{pmatrix} \begin{pmatrix} \underline{\underline{S}}^{(1)} \\ \underline{\underline{S}}^{(2)} \\ A^{(0+)} \end{pmatrix} = \underline{0}. \quad (3.26)$$

This system is further expanded in the form of Eq.(3.7) and solved for a set of eigenvalues γ and associated eigenvectors \mathbf{S} at each frequency ω . Wavenumbers k_1 are then obtained from γ . The dispersion curve is thus calculated over a given frequency range. However, the condition $\text{Re}(k_2^0) \geq 0$ should be imposed to sort out the solution that do not meet it.

Although usual and simple, this two-layer structure is challenging from a numerical point of view because it involves materials with very large impedance contrasts and layers whose thicknesses differ by several orders of magnitude. These problems are partially solved by normalizing each line and rearranging the terms of the matrix Eq. (3.7), leading to a better numerical conditioning of the system.

3.4.2. Dispersion relation

The material properties of the $b^{(1)} = 52$ mm-thick poroelastic (melamine) layer are listed in Table 3.1, while the properties of the $b^{(2)} = 1$ mm-thick aluminium plate are the density $\rho = 2700 \text{ kg.m}^{-3}$ and the Lamé coefficients $\lambda = 60.75 \text{ GPa}$ and $\mu = 26.03 \text{ GPa}$. The saturating fluid and that occupying the half-space is air with density $\rho_f = 1.213 \text{ kg.m}^{-3}$, the heat capacity ratio $\kappa = 1.4$, kinematic compressibility $\mu_f = 1.839 \times 10^{-5} \text{ Pa.s}$, the Prandtl number $\text{Pr} = 0.71$, and adiabatic bulk modulus κP_0 , where the atmospheric pressure is $P_0 = 1.013 \times 10^5 \text{ Pa}$. The total two-layer thickness is thus $b = b^{(1)} + b^{(2)} = 53 \text{ mm}$. $M^{(1)} = 5$ and $M^{(2)} = 11$ collocation points are employed respectively in the poroelastic and elastic layers.

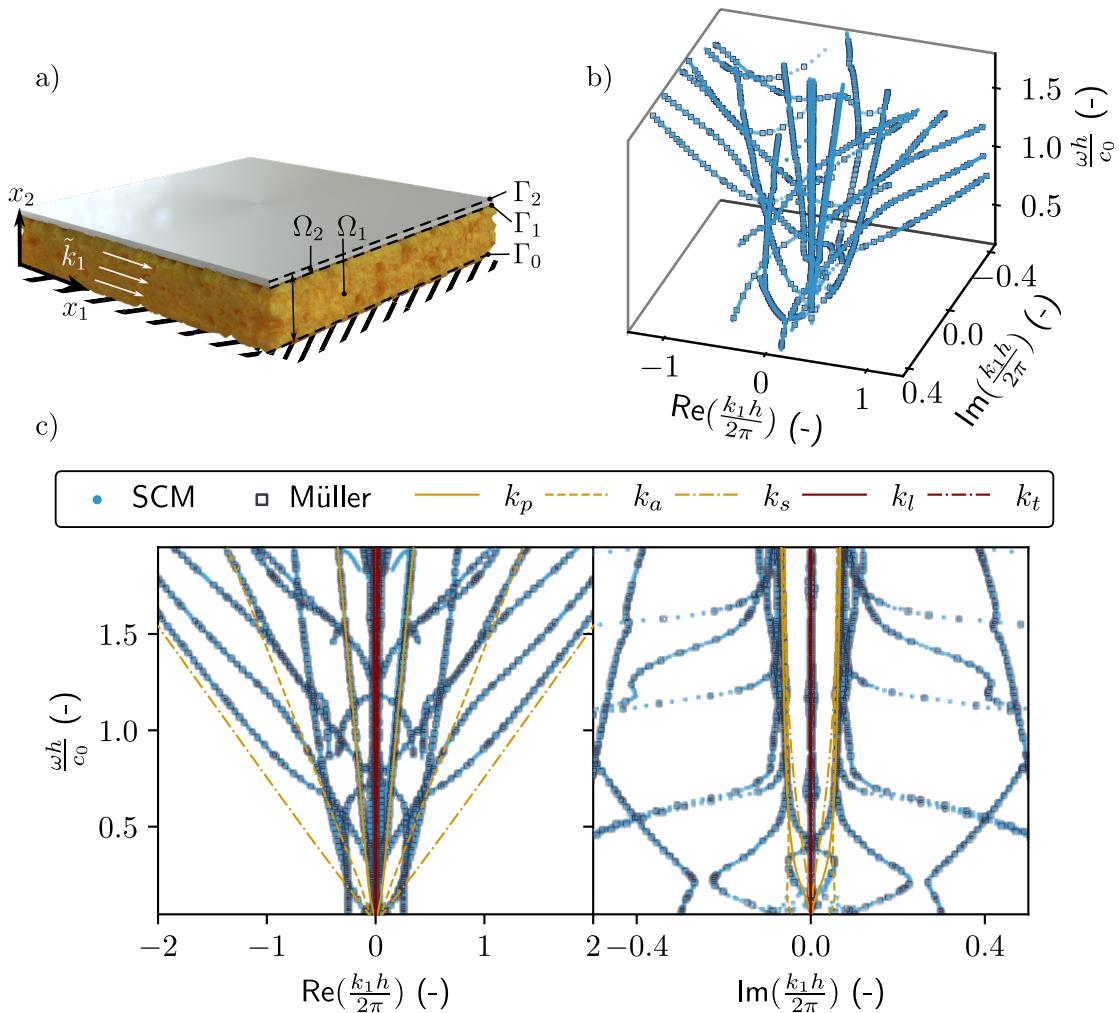


Figure 3.2: a) Sketch of the two-layer structure geometry. b) 3D view of the complex wavenumber-real frequency dispersion relation and c) real and imaginary parts of the dispersion relation. Blue dots represents the results as calculated with the SCM and black open markers represents those as calculated with the Müller method. Bulk wavenumbers in the poroelastic and elastic materials are highlighted in yellow and dark red respectively.

Table 3.1: Parameters of the poroelastic material (melamine foam)

Parameter	ϕ	ρ_1	R_f	Λ	Λ'	α_∞	ν	N_s
Unit	-	kg.m^{-3}	kPa.s.m^{-2}	μm	μm	-	-	kPa
Melamine	0.98	6.5	5.6	214	214	1	0.24	$11.96(1 + i0.07)$

The 3D view of the complex dispersion diagram is depicted Fig. 3.2b. To ease readability, real and imaginary parts are depicted separately in Fig. 3.2c. Please note that this last representation can be misleading because only the modes that lie in the complex wavenumber range depicted in Fig. 3.2b are represented in Fig. 3.2c. Cut-on frequencies are thus fictitious. The results calculated using the SCM method, plotted in blue markers, are compared with those obtained using the Müller method, plotted with black open markers. Both methods provide identical results, up to around four digits, therefore validating the present approach. Note that the SCM solutions are used as initial guesses of the Müller algorithm to shorten computational time. In this way, fewer iterations are needed for the results of the root-finding method to converge. The main advantages of SCM are that no initial guesses are needed, i.e. as long as the discretization is sufficient to model the full wave behavior in the structure, the dispersion relation is guaranteed to be calculated. In addition, the matrices used for the numerical model are very small compared to those required by other methods such as FEM, and are therefore much less demanding in terms of calculation.

Wavenumbers are perfectly symmetric with both $\text{Re}(k_1 b / 2\pi) = 0$, and $\text{Im}(k_1 b / 2\pi) = 0$ axis, which is a feature of reciprocal systems. Wavenumbers having a very large slope and a low imaginary part correspond to modes mostly propagating in the aluminium plate. They are very close to the bulk longitudinal k_l and transverse k_t wavenumbers in the aluminum. This is due to the very large contrast between the elastic properties of the aluminum and those of the poroelastic frame. The other branches asymptotically tend toward the bulk acoustic, k_a , compression, k_p , or shear, k_s , wavenumbers of the poroelastic medium at high frequency. At low frequencies, these modes are highly dispersive.

If we compare these results with the dispersion relation depicted in 3.6.Bc) for the guided waves in a single poroelastic plate twice as thick and in the absence of the elastic plate, we find some essential differences . The branches associated with the elastic plate and the branch clearly associated with the coupling between the elastic and poroelastic layers (referenced as (1) in Fig.3.6.Bc)) do not exist. The other branches are similar to those of this single poroelastic layer, although slightly modified by the presence of the elastic plate. This proves a weak coupling between the elastic and poroelastic layers, particularly at high frequencies. At lower frequencies, the coupling between the two layers is stronger. Effectively, the wavelength is much larger than the thickness of at least one of the two layers in these frequency ranges, which favors coupling. The A_0 branch observed in the case of a single poroelastic layer disappear in the two-layer system rigidly backed, because this rigid-backing does not allow the existence of such a mode.

3.4.3. Mode shapes and energy fluxes

Beyond the SCM numerical efficiency, a significant advantage of the method is that the eigenvectors associated to each eigenvalue directly provides the associated mode shape. Since these mode shapes, i.e. eigenvectors, are calculated from an eigenvalue problem, their amplitude are defined up to a constant. They are normalized to the value of the pressure field at the bottom of the poroelastic layer, i.e. at the location of the rigid backing, because this condition imposes $\mathbf{n}^{(1)}(0) = 0$ (and in particular $n_2^{(1)}(0) = 0$), which is equivalent to a maximum pressure field. From these normalized eigenvectors, any physical fields can be calculated as detailed in 3.6.C.

Let us consider two solutions $k_1^\pm = \pm|\text{Re}(k_1)|$ symmetric in the dispersion diagram with respect to the axis $\text{Re}(k_1) = 0$ as shown in Figs. 3.3a-b. The displacement fields and total energy fluxes in each direction are depicted in Fig. 3.3c-f. The results calculated using the SCM method, plotted in continuous lines, match those obtained using the Müller method, plotted with marker lines. These two modes are identical but propagate in opposite directions, i.e., the mode k_1^+ propagates towards positive x_1 while the mode k_1^- propagates towards negative x_1 . Displacement fields and energy fluxes are identical along the x_2 axis but are of equal modulus but opposite sign along the x_1 axis (see Fig. 3.3c-d and e-f). Both displacement field and energy flux amplitudes are two to three orders of magnitude smaller in the elastic plate (for $b_1 < x_2 < b$) than in the poroelastic

plate. The aluminum coating is almost acting as a rigid boundary condition, as also testified by the fact that the x_1 components of the displacement field and energy fluxes are almost symmetric with respect to the axis $x_2 = b_1/2$, while their x_2 components are almost antisymmetric along the x_2 axis. Note that the fluid phase Poynting vectors depicted in Fig. 3.3f are of equal modulus but opposite sign, symmetric with respect to the axis $x_2 = b_1/2$, and more than a hundred times smaller than total energy flux in amplitude. The energy of these modes are thus mainly localized in the solid phase of the poroelastic plate.

To go a step further, we evaluated the energy transport velocity $\bar{V}_e = \bar{P}_t/\bar{U}$ as the ratio between the average energy flux in each layer \bar{P}_t over the mean total energy \bar{U} . The energy transport velocity of the mode highlighted in orange in Fig. 3.3a-b is compared to the group velocity $v_g = \partial\omega/\partial\text{Re}(k_1)$ [45] in Fig. 3.3g. It is numerically computed by doing finite differences on a branch of the dispersion relation. Both velocities match for frequency ranges where the imaginary part of k_1 is low, i.e., for weak attenuation modes [46–49], but are different for strong attenuation. In particular, the group velocity diverges around the cut-on frequency. This is the reason why the energy transport velocity is usually preferred to study absorbing structures. The energy transport velocity \bar{V}_e of the mode highlighted in green in Fig. 3.3a-b is depicted in Fig. 3.3h. This energy transport velocity is always positive, while k_1 crosses the $\text{Re}(k_1) = 0$ axis. When $\text{Re}(k_1) < 0$ the phase velocity is negative, while the $\bar{V}_e > 0$. This branch (and the symmetric one) has also a negative group velocity as also testified by the fact $\text{Re}(k_1) < 0$ while $\text{Im}(k_1) > 0$ in the shaded area in Fig. 3.3h. This branch has also an exponentially growing amplitude.

3.4.4. Experimental validation

The experimental set-up is depicted in Fig. 3.4a. The sample is the two-layer structure considered in the previous subsection. The aluminum plate is glued on top of the poroelastic layer. The latter layer is glued on a 3 cm-thick aluminium plate, which mimic the rigid backing. The three elements are 45 cm wide and 85 cm long. The surrounding medium and the saturating fluid is air, at regular pressure and temperature conditions. The sample is positioned perpendicularly to the ground on its longest side as depicted in Fig. 3.4a to ease the mode excitation in the poroelastic layer. The excitation is provided by a shaker (Brüel and Kjaer type 4810), which is rigidly attached to the sample with a threaded steel rod fixed to the shaker on one side and glued on a 1mm thick aluminum plate of width 10 cm and height 1.5 cm. This plate is cut at the edge opposite to the threaded steel rod and glued to the porous sample, creating a line source at the edge of the sample. The resonance of this part was measured to be 4500 Hz, which is the upper limit of the measured frequency range. The excitation signal is a swept-sine function with 400 points ranging from 1 kHz to 3 kHz. The general layout of the experiments is similar to that used originally in [50]. The key difference in our case is that the excitation is operated directly to the poroelastic layer and not to the top of the aluminum plate, because the modes that corresponds to propagating solutions in the poroelastic layer are encapsulated in between two almost rigid plates (see Section 3.4.3). The normal displacement is measured with a laser Doppler vibrometer (Polytec VibroFlex Neo) on a line of length $L = 40$ cm, along the side of the thickness of the porous layer. These measurements are averaged 100 times. The excitation signal and the measured field are interfaced through a computer via a Zürich Instruments acquisition card.

The SLaTCoW method [50] is used to analyze the space-frequency displacement measurements and recover the complex wavenumber-real frequency dispersion relation. This method relies on the Laplace spatial transform of these measurements, which gives access to the real and imaginary components of the complex wavenumber. The difference between this Laplace transform and a correctly chosen ansatz function is finally minimized for each frequency, to recover the dispersion relation. + détails ? The complex wavenumber-real frequency dispersion relation in Fig. 3.4b-c depicts the wavenumbers recovered experimentally k_s and simulated k_n . Four modes were sought during the minimization with the SLaTCoW method: three of them correspond to the three bulk waves due to excitation and the fourth is depicted. The experimentally recovered mode jumps from one branch to the other because of the large modal density. The cut-on frequencies of the modes are clearly visible, especially on the imaginary part of the wavenumber. The lower real wavenumber parts of the numerical dispersion relation are not retrieved with the experimental results. This is because they correspond to a high attenuation of the mode, a part that is inherently difficult to measure in this experimental conditions. Generally, the recovered branches are those around the bulk solid compression wavenumber, corresponding to the wave that we most probably excite in the structure. The challenges of exciting homogeneously all the waves

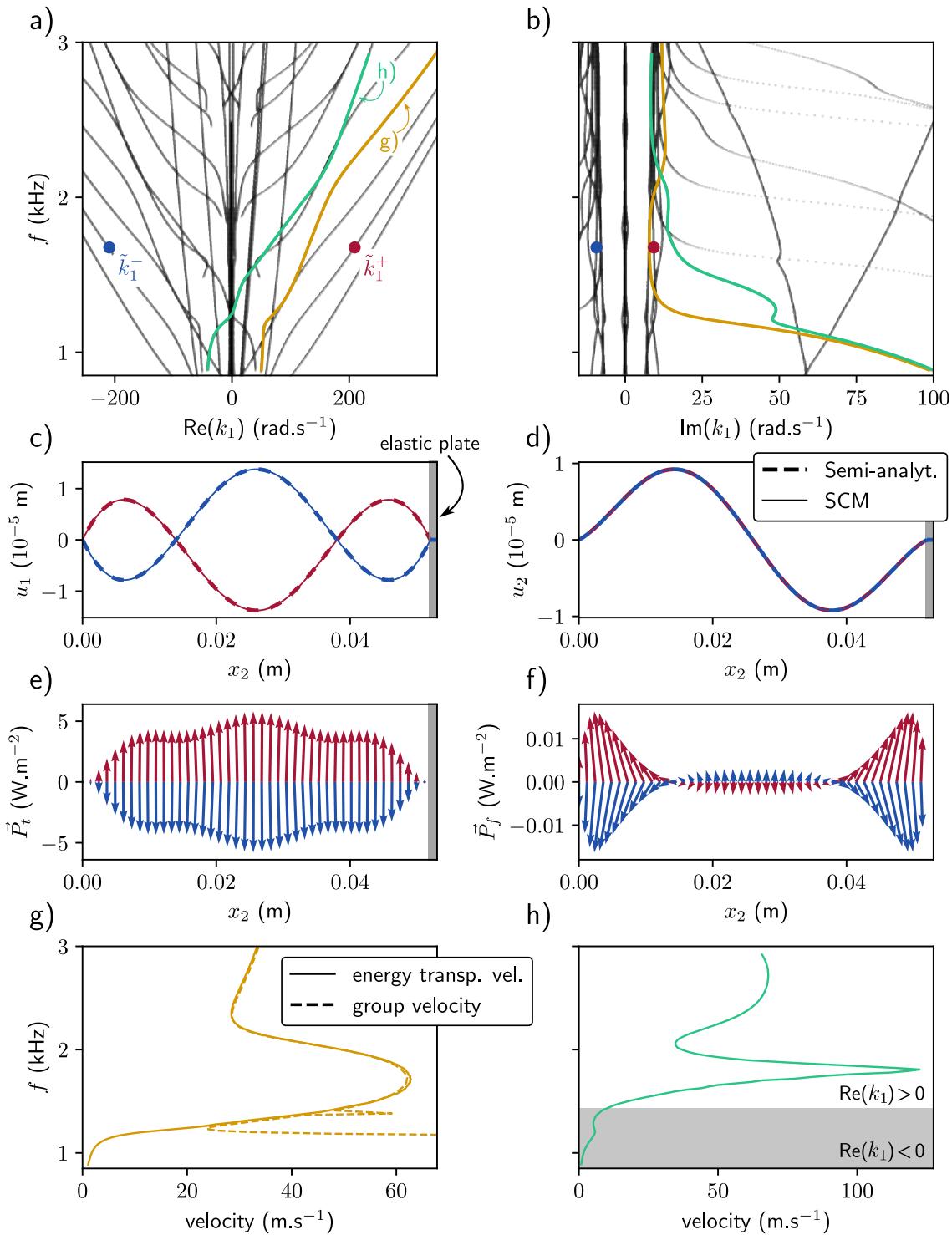


Figure 3.3: a) and b) represent the real and imaginary parts of the dispersion relation of the two-layer structure. Two branches are highlighted in yellow and green, as well as two specific solutions (blue and red) symmetric with respect to the $\text{Re}(k_1) = 0$ axis. Displacement fields in the x_1 (c) and x_2 direction d) associated with the forward (red) and backward (blue) solutions displayed above. e) Total Poynting vector and f) fluid Poynting vector in the poroelastic material distributions along the thickness of the structure for these two specific solutions. g) Group and energy transport velocities for the yellow branch. h) Energy transport velocity for the green branch. The gray area represents the frequency range over which this velocity is negative.

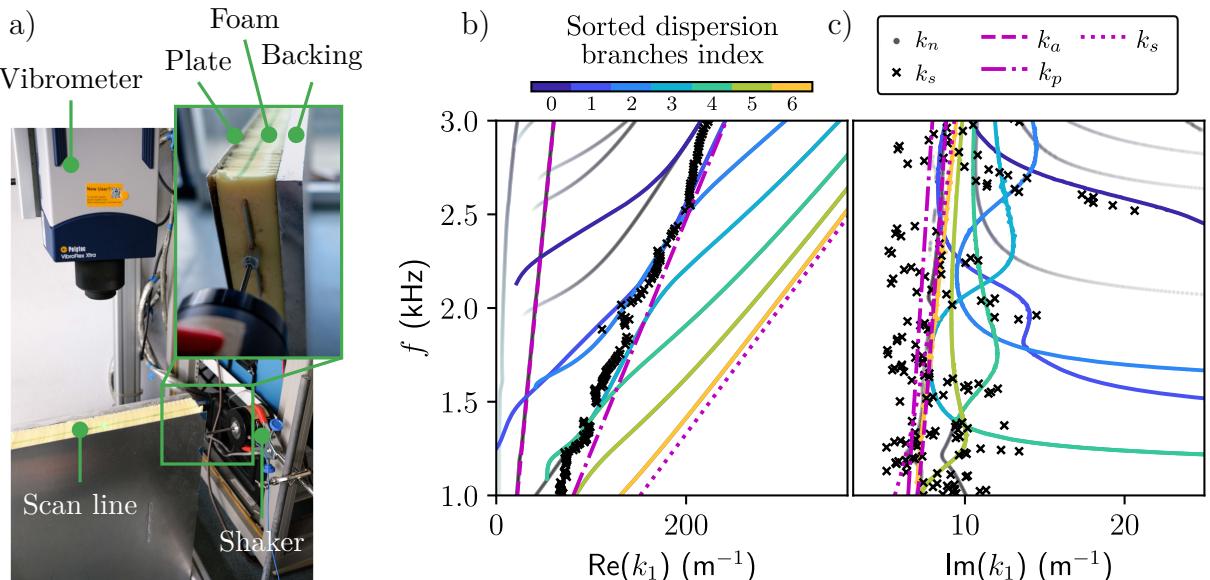


Figure 3.4: a) Experimental set-up used for the measurement of the normal displacement field of the two-layer structure, with a close-up of the sample excitation. b) Real and c) imaginary parts of the dispersion relations recovered experimentally from the SLaTCoW method (black crosses) as calculated with the SCM (colored lines) and the bulk wavenumbers in the poroelastic material (magenta lines).

that is supported by this configuration explain the other missing regions of the dispersion relation. + détails ? Nevertheless, the agreement between experimental results and simulation is quite good, thus validating the method.

3.5 | Conclusion

A spectral collocation method (SCM) is proposed to calculate dispersion relations for guided acoustic wave propagation in almost all multilayer dissipative structures Changer cette phrase ?. The formalism is general enough to be able to consider arbitrary arrangements of (visco-) elastic, poroelastic, and (viscothermal-) fluid layers. Outgoing fluid radiation, leading to Leaky modes, can be accounted for. The fields in each layer is approximated by Chebyshev polynomials and discretized on the corresponding nodes. Interface conditions are implemented to couple layers together. The solution provides the full complex wavenumber-real frequency dispersion relation of the considered structure. In addition, the physical fields in each layer are directly evaluated since they are the eigenvectors associated with the eigenvalue problem giving the dispersion relation. The dispersion diagram calculated with the present SCM is validated against experimental results and that calculated with a root-finding method (Müller method) in a simple two-layer structure consisting of a rigidly backed poroelastic layer covered by a thin aluminum plate radiating in a fluid half-space. The semi-analytical approach underlying the root-finding method is taken a step further to solve for the wave amplitudes as well, providing a reference for the mode shapes calculated with the SCM. The dispersion properties in such structures, as well as the direction of the energy density in each phase of the poroelastic layers using a decoupled expression for the Poynting vectors is analyzed. Results on energy transport velocity are also reported.

This method paves the way of complex wavenumber-real frequency dispersion relation calculation and analysis of more complex structures, like layers with embedded periodic inclusions (elastic or resonant), i.e. metaporolastic surfaces.

3.5 | Conflict of interest

The authors have no conflicts to disclose

3.5 | Author Contributions

Mathieu Maréchal Formal analysis (lead); Investigation (equal); Methodology (equal); Software (lead); Visualization (lead); Writing – original draft (lead). **Alan Geslain** Writing – review & editing (supporting); Investigation (equal). **Jean-Philippe Groby** Formal analysis (equal); Methodology (equal); Supervision (equal); Writing – review & editing (lead). **Vicent Romero-García** Formal analysis (supporting); Methodology (supporting); Supervision (equal); Writing – review & editing (equal). **Olivier Dazel** Formal analysis (equal); Investigation (equal); Methodology (equal); Supervision (equal); Writing – review & editing (equal).

3.5 | Data Availability Statement

The data that support the findings of this study are available from the corresponding author upon reasonable request.

3.6 | Appendices

3.6.A. Expression of the Biot elastic coefficients and effective parameters in poroelastic materials

The complex and frequency dependent density and bulk modulus of the fluid phase, which accounts respectively for the viscous and thermal losses are [22, 23] as,

$$\begin{aligned}\rho_{eq}(\omega) &= \frac{\rho_f \alpha_\infty}{\phi} \left(1 + \frac{iR_f \phi}{\rho_f \alpha_\infty \omega} \sqrt{1 - i\omega \rho_f \mu_f \left(\frac{2\alpha_\infty}{R_f \phi \Lambda} \right)^2} \right), \\ K_{eq} &= \frac{\kappa P_0}{\phi} \left(\kappa - (\kappa - 1) / \left(1 + \frac{iR_f \phi}{\rho_f \alpha_\infty \text{Pr} \omega} \sqrt{1 - i\text{Pr} \omega \rho_f \mu_f \left(\frac{2\alpha_\infty}{R_f \phi \Lambda'} \right)^2} \right) \right)^{-1},\end{aligned}\quad (3.27)$$

where ϕ is the porosity, α_∞ is the tortuosity, Λ and Λ' are respectively the viscous and thermal characteristic lengths, R_f is the flow resistivity, μ_f is the dynamic viscosity, κ is the heat capacity ratio, ρ_f the density of the saturating fluid, and P_0 is the ambient pressure.

The Biot elastic coefficients, P , Q , and R , are more commonly expressed in terms of the frame, K_b , effective fluid, K_{eq} , and shear, N_s , moduli when saturated by a light fluid as

$$P = K_b + \frac{4}{3} N_s + (1 - \phi)^2 \frac{K_{eq}}{\phi}; \quad Q = K_f (1 - \phi), \quad R = \phi K_{eq}, \quad (3.28)$$

while the apparent densities are

$$\rho_{22} = \phi^2 \rho_{eq}, \quad \rho_{12} = \phi \rho_f - \rho_{22}, \quad \rho_{11} = (1 - \phi) \rho_s - \rho_{12}, \quad (3.29)$$

with ρ_s the density of the solid phase. The coefficients in Eq. (3.13) are thus

$$\hat{A} = K_b - \frac{2}{3} N_s + \frac{K_{eq}}{\phi} ((1 + (1 - \phi)), \quad \hat{\rho} = \rho_1 + \phi \rho_f - 2\phi^2 \rho_{eq} - \frac{\rho_f^2}{\rho_{eq}}, \quad \beta = \frac{\rho_f}{\rho_{eq}} - 2\phi - 1, \quad (3.30)$$

where $\rho_1 = \phi \rho_f + (1 - \phi) \rho_s$ is the apparent density of the frame. Finally, the frame bulk modulus and shear modulus are linked by the Poisson ratio ν with $K_b = 2N_s(1 + \nu)/3(1 - 2\nu)$.

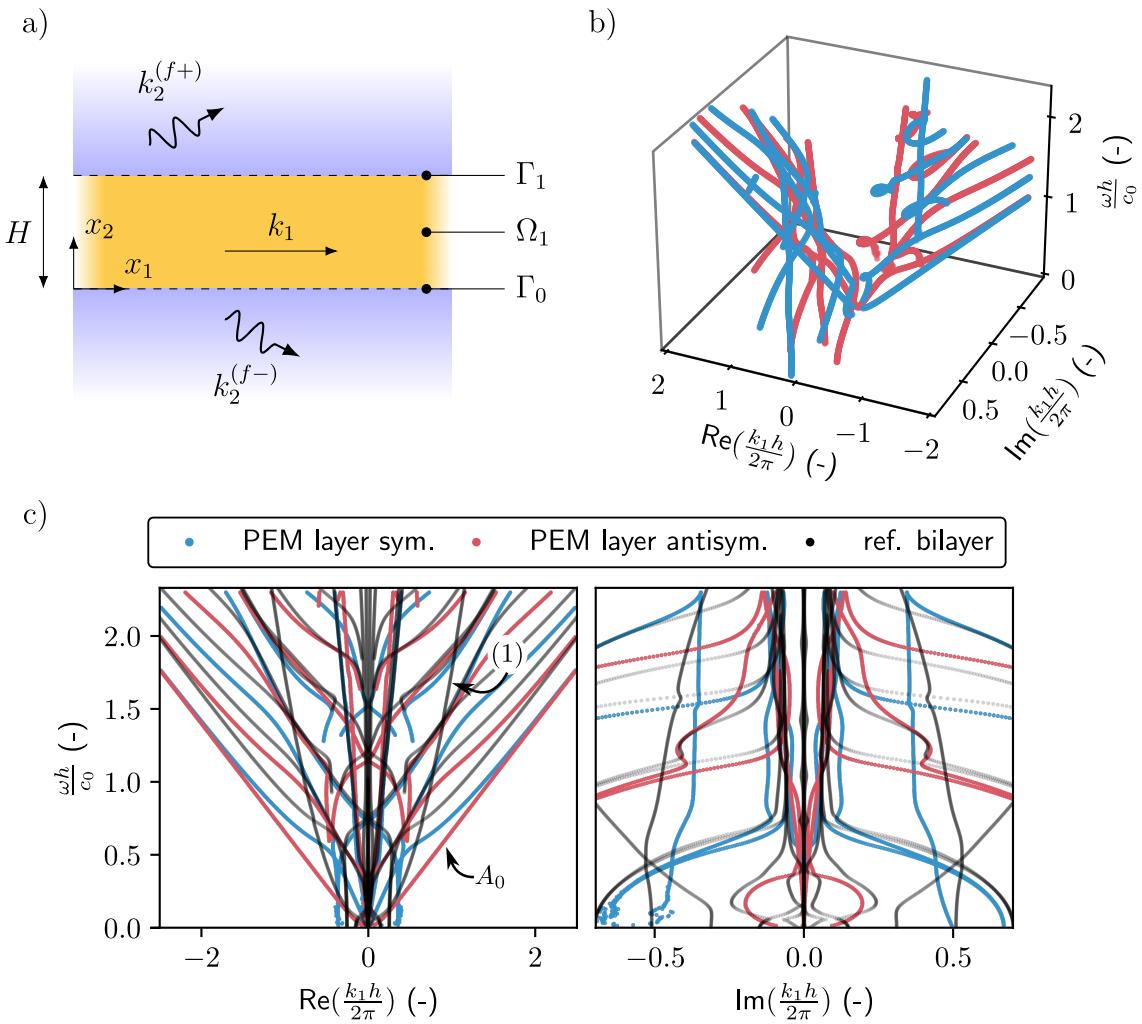


Figure 3.5: a) Sketch of the single layer poroelastic structure geometry. b) 3D view of the complex wavenumber-real frequency dispersion relation and c) real and imaginary parts of the dispersion relation. The branches of the symmetric and antisymmetric modes are represented by blue and red curves respectively. The dispersion relation of the two-layer elatic-poroelastic system studied in Section 3.4 are reminded with black lines.

3.6.B. Application to a single poroelastic layer surrounded by two fluid half-spaces

The dispersion relation of guided waves in a single layer of poroelastic material surrounded by two identical fluid half-spaces is considered (See Fig. 3.5a). The poroelastic material properties are those considered in Section 3.4.2 and the layer thickness is $2h^{(1)} = 104$ mm, i.e. twice the thickness of the poroelastic layer considered in the two-layer system.

The discretized equations of motion is that provided in Eq.(3.16). $M^{(1)} = 11$ collocation points are considered. Only the boundary conditions at the lower, i.e. Γ_0 at $x_2 = 0$, and upper, i.e. Γ_1 at $x_2 = h$, interfaces are modified. They read as [44]

$$\begin{aligned} \left(1 + \phi + \phi \frac{\rho_{12}}{\rho_{22}}\right) u_2^{(1)} + \frac{\phi^2}{\rho_{22}\omega^2} \partial_2 p^{(1)} &= u_2^{(0\pm)}, \quad p^{(1)} = p^{(0\pm)}, \\ \hat{\sigma}_{22}^s &= \left(1 - \phi \left(1 + \frac{Q}{R}\right)\right) p^{(0\pm)}, \quad \hat{\sigma}_{12}^s = 0. \end{aligned} \tag{3.31}$$

The discretized form of these boundary conditions is

$$\underline{\underline{C}}_1 = \begin{pmatrix} \frac{\phi^2}{\rho_{22}} \underline{D}_2^+ & 0 & \omega^2 \left(1 + \phi \left(1 + \frac{\rho_{12}}{\rho_{22}} \right) \right) \underline{I}^+ & \frac{ik_2^{(f)}}{\rho_f} \\ \underline{I}^+ & 0 & 0 & -1 \\ 0 & ik_1(\hat{P} - 2N_s)\underline{I}^+ & \hat{P}\underline{D}_2^+ & 1 - \phi \left(1 + \frac{Q}{R} \right) \\ 0 & N_s\underline{D}_2^+ & ik_1N_s\underline{I}^+ & 0 \end{pmatrix}. \quad (3.32)$$

These matrices are then arranged in the same manner as Eq.(3.10) and the wavenumbers are calculated. Fig. 3.5b-c depicts the results.

In a similar way as Lamb modes, the different modes can be sorted in symmetric and antisymmetric modes. This is done by comparing the sign of the radiated amplitudes at each edge of the layer. Results are shown in Fig. 3.5b. All modes necessarily have an imaginary part, both because of the radiation condition, i.e. leakage, and because of viscothermal dissipation as can be noticed in Fig. 3.5c.

3.6.C. Evaluation of the physical fields from the SCM solutions

The discrete sets of eigenvectors resulting from the SCM are the spatial Laplace transforms of some physical fields at collocation points, i.e. the locations where the discrete equations of motion (residue) are exactly satisfied. These discrete fields are interpolated into a continuous form using

$$\underline{s}^{(n)}(x_2) = \left[\underline{\psi}^{-1}(\xi) \underline{s}_m^{(n)}(\xi) \right] \psi(\xi^{(n)}) \quad (3.33)$$

where the terms in brackets correspond to the coefficients $\alpha_m^{(n)}$ given in Eq. (3.8). The Chebyshev polynomials ψ is the basis that maps from the discrete grid ξ back to the physical space.

The stress tensor components depend on the direct expressions of the fields and their first order derivative in x_2 , denoted with a prime symbol in the following. The various fields, differentiated up to any order, can be calculated by replacing ψ by its derivative in Eq. (3.33). Explicit expressions are given in the following,

- in the elastic layer, the spatial Laplace transforms of stress tensor components are computed from $\underline{\underline{u}}$, with

$$\begin{aligned} \sigma_{11} &= (\lambda + 2\mu)ik_1\underline{\underline{u}}_1 + \lambda\underline{\underline{u}}'_2, & \sigma_{12} &= \mu(\underline{\underline{u}}'_1 + ik_1\underline{\underline{u}}_2), \\ \sigma_{22} &= \lambda ik_1\underline{\underline{u}}_1 + (\lambda + 2\mu)\underline{\underline{u}}'_2. \end{aligned} \quad (3.34)$$

- in the poroelastic layer, the full solid stress tensor is expressed as $\sigma_s = \underline{\underline{\sigma}}_s - \phi(Q/R)\underline{\underline{p}}$, with $\underline{\underline{\sigma}}_s$ the spatial Laplace transform of the *in vacuo* solid stress tensor introduced in Section 3.4. The components of this full solid stress tensor are,

$$\begin{aligned} \sigma_{s,11} &= P ik_1 \underline{\underline{u}}_{s,1} + \hat{A} \underline{\underline{u}}'_{s,2} - \phi \frac{Q}{R} \underline{\underline{p}}, & \sigma_{s,12} &= N_s \left(\underline{\underline{u}}'_{s,1} + ik_1 \underline{\underline{u}}_{s,2} \right), \\ \sigma_{s,22} &= \hat{A} ik_1 \underline{\underline{u}}_{s,2} + P \underline{\underline{u}}'_{s,1} - \phi \frac{Q}{R} \underline{\underline{p}}. \end{aligned} \quad (3.35)$$

The Laplace transform of the fluid stress tensor is calculated as

$$\sigma_f = -\phi \underline{\underline{p}} \delta_{ij}, \quad (3.36)$$

and the fluid displacement components are

$$\underline{\underline{u}}_{f,1} = \frac{\phi}{\rho_{22}\omega^2} ik_1 \underline{\underline{p}} - \frac{\rho_{12}}{\rho_{22}} \underline{\underline{u}}_{s,1}, \quad \underline{\underline{u}}_{f,2} = \frac{\phi}{\rho_{22}\omega^2} \underline{\underline{p}}' - \frac{\rho_{12}}{\rho_{22}} \underline{\underline{u}}_{s,2}. \quad (3.37)$$

All physical fields $\theta(\underline{x})$ can finally be calculated by the inverse spatial Laplace transform with $\theta(\underline{x}) = \underline{\theta}(k_1, x_2) e^{ik_1 x_1}$.

3.6.D. Energy flux in the multilayer structure - Poynting theorem in a poroelastic medium

Energy conservation is commonly conducted in the time domain. Let us consider a volum Ω , bounded by $\partial\Omega$, with \mathbf{n} the outgoing normal vector. The energy balance or Poynting theorem in a poroelastic materials reads as [18, 51]

$$\begin{aligned} - \int_{\partial\Omega} (-\sigma_s \cdot \mathbf{v}_s^*) \cdot \mathbf{n} + (-\sigma_f \cdot \mathbf{v}_f^*) \cdot \mathbf{n} \, dS &= \frac{1}{2} \int_{\Omega} \rho_1 \partial_t |\mathbf{v}_s|^2 - \rho'_{12} \partial_t (\mathbf{v}_f - \mathbf{v}_s)^2 + \rho_2 \partial_t |\mathbf{v}_f|^2 \, d\Omega \\ &\quad + \int_{\Omega} \sigma_s : \partial_t \epsilon_s^* + \sigma_f : \partial_t \epsilon_f^* - b (\mathbf{v}_s - \mathbf{v}_f)^2 \, d\Omega, \end{aligned} \quad (3.38)$$

where $\mathbf{v}_s = \partial_t \mathbf{u}_s$ and $\mathbf{v}_f = \partial_t \mathbf{u}_f$ are the velocity fields, σ_s and σ_f are the stress tensors, $\epsilon_s = 1/2 (\nabla \mathbf{u}_s + \nabla^T \mathbf{u}_s)$ and $\epsilon_f = 1/2 (\nabla \mathbf{u}_f + \nabla^T \mathbf{u}_f)$ are the strain tensors in the solid frame and in the fluid phase respectively, and \star denotes the complex conjugate. Note that a similar expression is provided in [49], where the alternative 1962 Biot formulation is used.

The left-hand side contains the energy density flux $\mathbf{P}_s = -\sigma_s \cdot \mathbf{v}_s^*$ and $\mathbf{P}_f = -\sigma_f \cdot \mathbf{v}_f^*$ in the elastic frame and fluid phase. The right-hand side contains the kinetic energy K [42] and the internal forces $P_{\text{int}} = \partial_t W + P_{\text{dis}}$. The latter is the sum of the time derivative of the strain energy W and the dissipated power P_{dis} by the elastic damping and the viscous and thermal losses [51]. In the frequency domain, the kinetic and strain energies read as

$$\begin{aligned} K &= \frac{1}{2} \rho_1 \partial_t |\mathbf{v}_s|^2 - \frac{1}{2} \rho'_{12} \partial_t (\mathbf{v}_f - \mathbf{v}_s)^2 + \frac{1}{2} \rho_2 \partial_t |\mathbf{v}_f|^2, \\ W &= \frac{1}{2} \sigma_s : \epsilon_s^* + \frac{1}{2} R \nabla \cdot (Q \mathbf{u}_s + R \mathbf{u}_f), \end{aligned} \quad (3.39)$$

Finally, the time average total Poynting vector is $\langle \mathbf{P}_t \rangle = \langle \mathbf{P}_s \rangle + \langle \mathbf{P}_f \rangle$, with $\langle \mathbf{P}_s \rangle = -\text{Re}(\sigma_s \cdot \mathbf{v}_s^*)/2$ and $\langle \mathbf{P}_f \rangle = -\text{Re}(\sigma_f \cdot \mathbf{v}_f^*)/2$ and the total energy density is $\langle U \rangle = \langle K \rangle + \langle W \rangle = \text{Re}(K)/2 + \text{Re}(W)/2$. Note that a similar procedure can be followed for isotropic elastic or fluid medium [52]. In particular, the elastic Umov-Poynting vector is $\langle \mathbf{P} \rangle = \text{Re}(\sigma \cdot \mathbf{v}^*)/2$, where σ is the stress tensor and \mathbf{v} is the velocity in the elastic medium.

The energy transport velocity V_e is evaluated from the latter expressions

$$V_e = \frac{\langle \mathbf{P}_t \rangle \cdot \mathbf{n}}{\langle U \rangle}. \quad (3.40)$$

When a multi-layer configuration is considered, the latter velocity is averaged over the structure thickness $\overline{V}_e = \overline{\overline{P}_t}/\overline{\overline{U}}$, with the average Poynting vector and the average total energy defined as

$$\overline{P}_t = \sum_n \frac{1}{b_n} \int_{b_{n-1}}^{b_n} \langle \mathbf{P}_t^{(n)} \rangle \cdot \mathbf{e}_{x_2} \, dx_2, \quad \overline{U} = \sum_n \frac{1}{b_n} \int_{b_{n-1}}^{b_n} \langle U^{(n)} \rangle \, dx_2. \quad (3.41)$$

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Appendices

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A

Implémentation de SAFE pour des modes de Lamb

A.1. Formalism

1. Variational form for elasticity equations

$$-\rho\omega^2 \int_{\Omega} u \delta u \, d\Omega + \int_{\Omega} \sigma(u) : \epsilon(\delta u) \, d\Omega = \int_{\partial\Omega} \sigma(\delta u) \cdot \mathbf{n} \, d\Gamma, \quad \forall \delta u \quad (42)$$

2. Displacement fields $\mathbf{u}(x) = \mathbf{u}(x_2)e^{ik_1x_1}e^{-i\omega t}$

3. Double dot product $\sigma(u) : \epsilon(\delta u^\star) = \sigma_{ij}(u)\epsilon_{ij}(\delta u^\star)$

$$\begin{aligned} \sigma_{11}\epsilon_{11} &= k_1^2(\lambda + 2\mu)u_1\delta u_1 + -ik_1\lambda\nabla u_2\delta u_1 \\ \sigma_{12}\epsilon_{12} &= \frac{\mu}{2} \left(-ik_1\nabla u_1\delta u_2 + \nabla u_1\delta\nabla u_1 + k_1^2u_2\delta u_2 + ik_1u_2\nabla\delta u_1 \right) \\ \sigma_{22}\epsilon_{22} &= ik_1\lambda u_1\nabla\delta u_2 + (\lambda + 2\mu)\nabla u_2\nabla\delta u_2 \end{aligned} \quad (43)$$

4. Explicit form for the weak form of elasticity equations,

$$\begin{aligned} \int_{\Omega} -\rho\omega^2(u_1\delta u_1 + u_2\delta u_2) + k_1^2(\lambda + 2\mu)u_1\delta u_1 - ik_1\lambda\nabla u_2\delta u_1 - ik_1\mu\nabla u_1\delta u_2 + \mu\nabla u_1\delta\nabla u_1 \\ + \mu k_1^2 u_2\delta u_2 + ik_1\mu u_2\nabla\delta u_1) + ik_1\lambda u_1\nabla\delta u_2 + (\lambda + 2\mu)\nabla u_2\nabla\delta u_2 \, d\Omega = 0 \quad \forall \delta u \end{aligned} \quad (44)$$

where ∇u_i denotes differentiation of the i -th component of the displacement along x_2 . Note that the surface term initially expressed vanishes from the no-traction boundary condition.

5. Discretized fields

$$\mathbf{u}(\mathbf{x}) = \mathbf{N}(\xi) \mathbf{U} e^{ik_1 x_1} e^{-i\omega t}, \quad (45)$$

with \mathbf{N} containing the second-order Lagrange shape functions, and \mathbf{U} is the vector of displacement amplitudes.

6. Elementary matrices

$$\begin{aligned} \mathbf{M}^{(e)} &= \frac{b}{2} \int_{-1}^1 \mathbf{N}^\top \mathbf{N} d\xi, & \mathbf{K}^{(e)} &= \frac{2}{b} \int_{-1}^1 \nabla \mathbf{N}^\top \nabla \mathbf{N} d\xi, \\ \mathbf{C}_1^{(e)} &= \int_{-1}^1 \mathbf{N}^\top \nabla \mathbf{N} d\xi, & \mathbf{C}_2^{(e)} &= \int_{-1}^1 \nabla^\top \mathbf{N}^\top \mathbf{N} \mathbf{U} d\xi, \end{aligned} \quad (46)$$

7. Discretization $\forall \delta \mathbf{U}$, weak form for one element and rearranging the terms

$$\begin{aligned} & \int_{\Omega_e} k_1^2 \left(\mu \delta \mathbf{U}_2^\top \mathbf{M}^{(e)} \mathbf{U}_2 + (\lambda + 2\mu) \delta \mathbf{U}_1^\top \mathbf{M}^{(e)} \mathbf{U}_1 \right) + ik_1 \left(\mu \delta \mathbf{U}_2^\top \mathbf{C}_1^{(e)} \mathbf{U} - \mu \delta \mathbf{U}_1^\top \mathbf{C}_2^{(e)} \mathbf{U}_2 \right. \\ & \quad \left. + \lambda \delta \mathbf{U}_1^\top \mathbf{C}_1^{(e)} \mathbf{U}_2 - \lambda \delta \mathbf{U}_2^\top \mathbf{C}_2^{(e)} \mathbf{U}_1 \right) + \mu \delta \mathbf{U}_1^\top \mathbf{K}^{(e)} \mathbf{U}_1 + (\lambda + 2\mu) \delta \mathbf{U}_2^\top \mathbf{K}^{(e)} \mathbf{U}_2 \\ & \quad - \rho \omega^2 \delta \mathbf{U}_1^\top \mathbf{M}^{(e)} \mathbf{U}_1 - \rho \omega^2 \delta \mathbf{U}_2^\top \mathbf{M}^{(e)} \mathbf{U}_2 d\Omega_e = 0 \end{aligned} \quad (47)$$

8. General form for the problem

$$\left(k_1^2 \mathbf{A}_2 + ik_1 \mathbf{A}_1 + \mathbf{A}_0 - \omega^2 \mathbf{M} \right) \mathbf{U} = 0. \quad (48)$$

Rewrites as a generalized eigenvalue problem,

$$\left[\begin{pmatrix} -\mathbf{A}_1 & \mathbf{A}_0 - \omega^2 \mathbf{M} \\ \mathbf{I} & \mathbf{0} \end{pmatrix} - k_1 \begin{pmatrix} \mathbf{A}_2 & \mathbf{0} \\ \mathbf{0} & \mathbf{I} \end{pmatrix} \right] \begin{pmatrix} k_1 \mathbf{U} \\ \mathbf{U} \end{pmatrix} = 0 \quad (49)$$

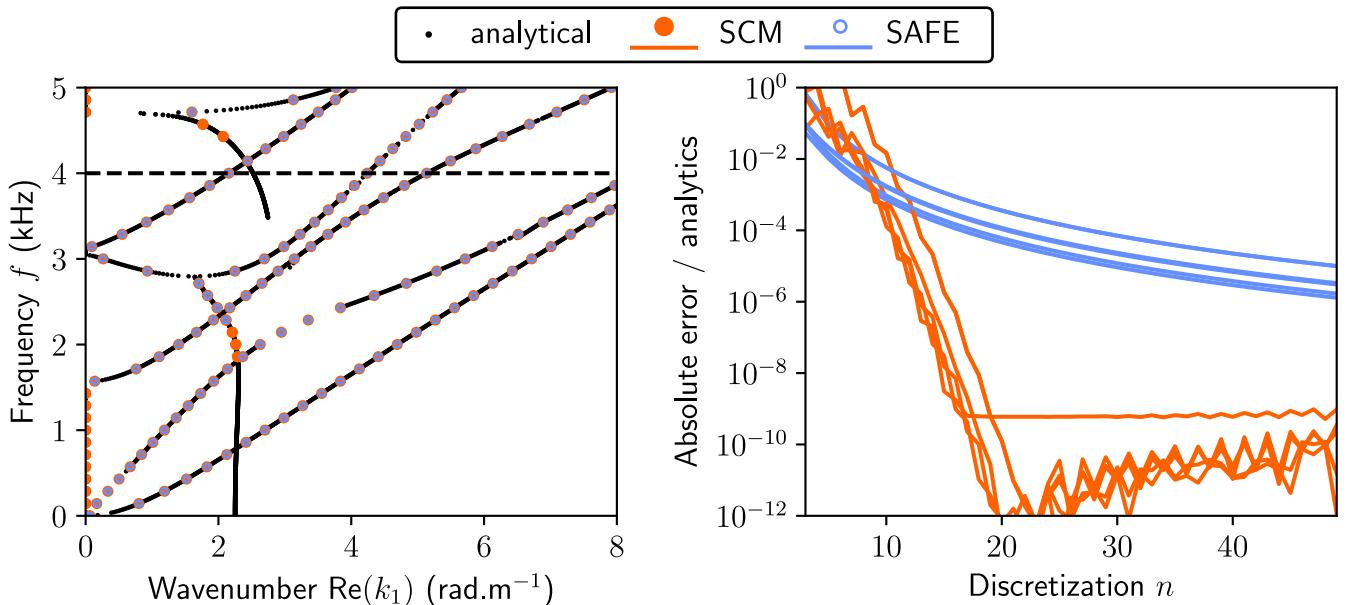


Figure 6: a) Dispersion relation for the Lamb modes with the root-finding results, alongside the spectral collocation and semi-analytical finite elements solutions. b) Convergence curve of the 2 methods at a given frequency $f = 4$ kHz. The multiple lines depicts the convergence of each wavenumber.

A.2. Results

Analytical results are obtained by applying the root-finding approach using the Müller method presented in the paper. Solutions are sought from the classical Lamb modes equations, derived in many books such as Royer. Figure 6a) shows the results of the three methods for a given frequency range. They all match well. In order to compare the two numerical methods in play, a brief convergence study has been done using the analytical results from the root-finding method as a reference.

The convergence curve of the results given in Fig. 6b) is unsurprisingly in agreement with the literature. It is calculated at a frequency $f = 4$ kHz: the solutions along the dashed line in Fig 6a) are selected. The parameter n used for the discretization is shared between the two methods and corresponds to the number of elements taken for the discretization in each method. SAFE matrices have a size $2(2n + 1)$, while SCM matrices have a size $2n$. The discretization is swept from $n = 3$ to $n = 50$.

The SCM convergence shows a steep slope and reaches a plateau around $n = 20$, while the SAFE curve follows the expected slope of a finite elements method that use second-order shape functions. When comparing at a given discretization order, the error from the SAFE is much lower than that of SCM by several order of magnitudes. In order to see comparable convergence, it would require to use higher-order shape functions or use spectral elements instead.

A.3. Point sur la biblio

- "Wave propagation along transversely periodic structures" cite [1]
- "Finite element model for waves guided along solid systems of arbitrary section coupled to infinite solid media" [2]
- "SAFE-PML approach for modal study of waveguides with arbitrary cross sections immersed in inviscid fluid" [3]
- "A coupled SAFE-2.5D BEM approach for the dispersion analysis of damped leaky guided waves in embedded waveguides of arbitrary cross-section" [4] + "Dispersion analysis of leaky guided waves in fluid-loaded waveguides of generic shape" and [5] "Dispersion analysis of leaky guided waves in fluid-loaded waveguides of generic shape"
- Livre J. Rose [6]
- Semi-Analytical Finite Element (SAFE) method for plotting Lamb waves dispersion curves of an aluminum plate and comparison with Disperse software [7]

B | Details on the Müller method

B.1. The root-finding algorithm

equations, iterative scheme etc.

[8]

$$\begin{aligned} f(x_0) &= a(x_0 - x_2)^2 + b(x_0 - x_2) \\ f(x_1) &= a(x_1 - x_2)^2 + b(x_1 - x_2) \\ f(x_2) &= a(x_2 - x_2)^2 + b(x_2 - x_2) \end{aligned} \quad (50)$$

Note that we have dropped the subscript “2” from the function for conciseness. Because we have three equations, we can solve for the three unknown coefficients, a , b , and c . Because two of the terms in Eq. (7.20) are zero, it can be immediately solved for $c = f(x_2)$. Thus, the coefficient c is merely equal to the function value evaluated at the third guess, x_2 . This result can then be substituted into Eqs. (7.18) and (7.19) to yield two equations with two unknowns:

$$\begin{aligned} f(x_0) - f(x_2) &= a(x_0 - x_2)^2 + b(x_0 - x_2) \\ f(x_1) - f(x_2) &= a(x_1 - x_2)^2 + b(x_1 - x_2) \end{aligned}$$

Algebraic manipulation can then be used to solve for the remaining coefficients, a and b . One way to do this involves defining a number of differences,

$$\begin{aligned} h_0 &= x_1 - x_0 & b_1 &= x_2 - x_1 \\ \delta_0 &= \frac{f(x_1) - f(x_0)}{x_1 - x_0} & \delta_1 &= \frac{f(x_2) - f(x_1)}{x_2 - x_1} \end{aligned}$$

These can be substituted into Eqs.(7.21) and (7.22) to give

$$\begin{aligned} (h_0 + h_1)b - (h_0 + h_1)^2a &= h_0\delta_0 + h_1\delta_1 \\ h_1 \quad b - \quad b_1^2 \quad a &= \quad h_1\delta_1 \end{aligned}$$

which can be solved for a and b . The results can be summarized as (7.24)

$$\begin{aligned} a &= \frac{\delta_1 - \delta_0}{h_1 + h_0} \\ b &= ah_1 + \delta_1 \\ c &= f(x_2) \end{aligned}$$

(7.25) (7.26) To find the root, we apply the quadratic formula to Eq. (7.17). However, because of potential round-off error, rather than using the conventional form, we use the alternative formulation [Eq.(3.13)] to yield

$$x_3 - x_2 = \frac{-2c}{b \pm \sqrt{b^2 - 4ac}}$$

or isolating the unknown x_3 on the left side of the equal sign,

$$x_3 = x_2 + \frac{-2c}{b \pm \sqrt{b^2 - 4ac}}$$

Note that the use of the quadratic formula means that both real and complex roots can be located. This is a major benefit of the method. In addition, Eq. (7.27a) provides a neat means to determine the approximate error. Because the left side represents the difference between the present (x_3) and the previous (x_2) root estimate, the error can be calculated as

$$\varepsilon_a = \left| \frac{x_3 - x_2}{x_3} \right| 100\%$$

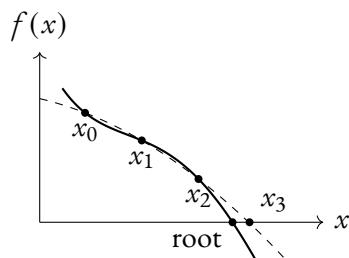


Figure 7: iteration muller

B.2. A complete routine for the automated solving of dispersion relations

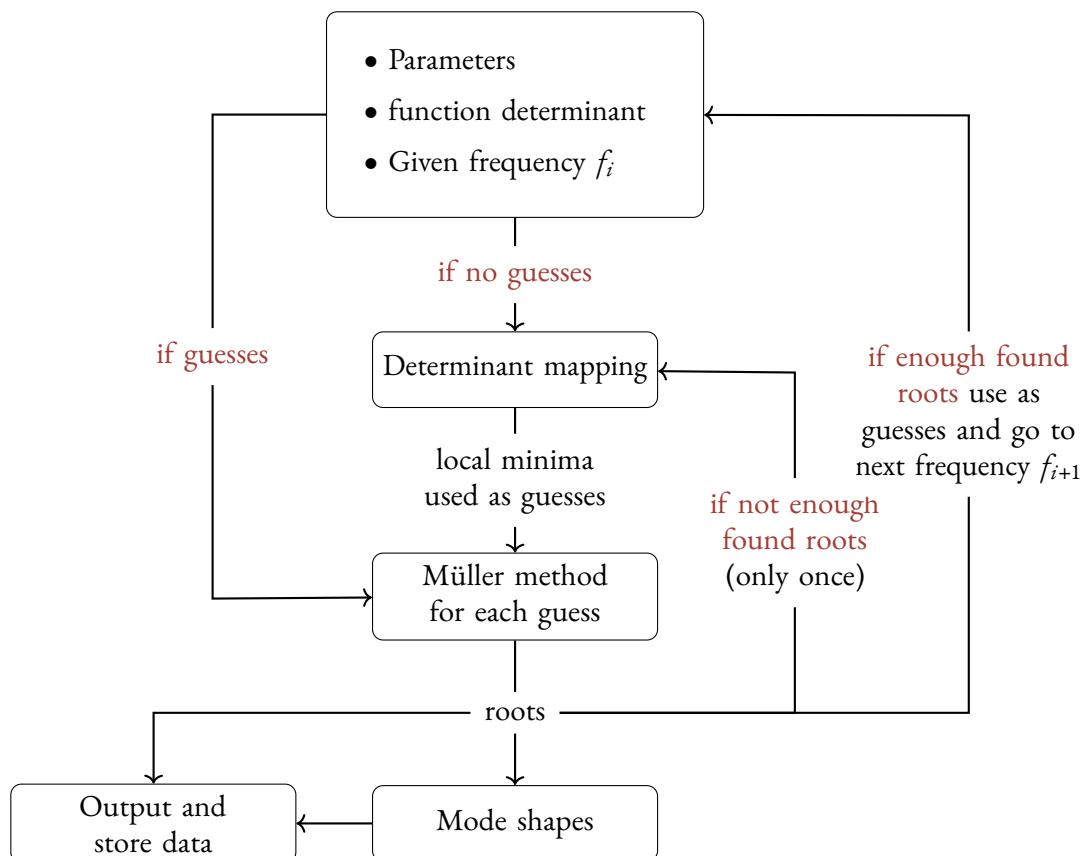


Figure 8: Description of the routine

C | References

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