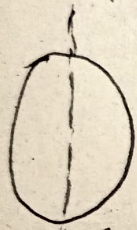
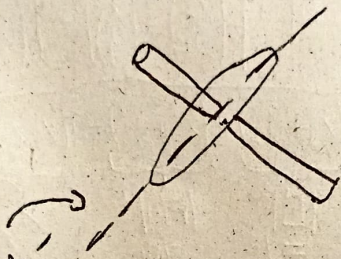


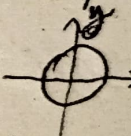
a)  $I_0$  será el momento de inercia al rededor de la dirección azimutal.



eje de rotación.



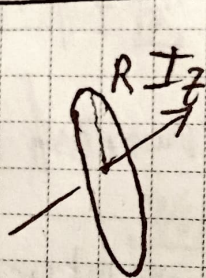
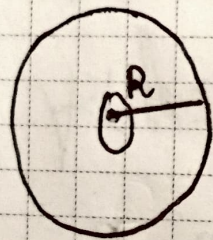
En este caso planteamos para el momento de Inercia sobre cada pedazo infinitesimal del disco:  $dI = r^2 dm$  (Esto es el momento de inercia generado por un pedazo de masa casi puntual.)

∴ Nótese que por la simetría del disco,  $I_x = I_y$  para  así que vale la pena hacerla para solo un eje. Elegimos y. Planteamos la integral que resulta en  $\frac{\pi}{4} R^4$ . Así pues  $I_y = I_x = \text{or } \frac{\pi}{4} R^4 = \frac{\pi}{4} R^4 \left( \frac{M}{\pi R^2} \right) = \frac{1}{4} M R^2$ .

No obstante, sobre este eje no hay ninguna rotación. Por este motivo usamos el teorema de ejes paralelos para que el eje se desplace hasta cortar con el eje de rotación real (sobre la ligadura).

Según ejes paralelos:  $I_0 = I_x + M d^2$ . Para el palo de largo  $D$ :  $I_0 = \frac{1}{4} M R^2 + M D^2$





$$\int dI = \int r^2 dm$$

$$\frac{M}{A} = \sigma = \frac{dm}{dA}$$

$$I_{cm} = \int_0^R r^2 \sigma dA = \int_0^R r^2 \sigma 2\pi r dr$$

$$dA = 2\pi r dr$$

$$I_{cm} = \int_0^R 2r^3 \pi \frac{M}{A} dr$$

$$\Rightarrow I_{cm} = 2\pi \frac{M}{A} \int_0^R r^3 dr = \frac{2\pi M}{\pi R^2} \cdot \frac{r^4}{4} \Big|_0^R = \frac{2M}{R^2} \cdot \frac{R^4}{4} = \frac{MR^2}{2}$$

$$\therefore I_z = \frac{1}{2} MR^2$$

$$\frac{1}{2} I_2 (\dot{\phi}^2 \cos^2 \theta + 2 \dot{\phi} \cos \theta \dot{\psi} + \dot{\psi}^2)$$

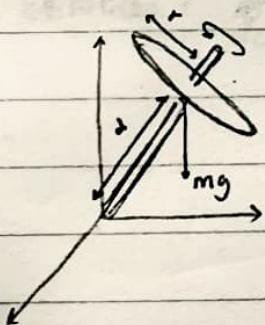
c)

$$L = \frac{1}{2} I_0 (\dot{\theta}^2 + \dot{\phi}^2 \sin^2 \theta) + \frac{1}{2} I_2 (\dot{\phi} \cos \theta + \dot{\psi})^2 - mgd \cos \theta$$

$$\begin{aligned} \star \frac{\partial L}{\partial \dot{\phi}} &= I_0 \dot{\phi} \sin^2 \theta + \frac{I_2}{2} (2 \dot{\phi} \cos^2 \theta + 2 \cos \theta \dot{\psi}) \\ &= I_0 \dot{\phi} \sin^2 \theta + I_2 (\dot{\phi} \cos^2 \theta + \cos \theta \dot{\psi}) \\ &= \dot{\phi} (I_0 \sin^2 \theta + I_2 \cos^2 \theta) + I_2 \cos \theta \dot{\psi} = p_{\phi} // \end{aligned}$$

$$\star \frac{\partial L}{\partial \dot{\psi}} = \frac{1}{2} I_2 (2 \dot{\phi} \cos \theta + 2 \dot{\psi}) = I_2 (\dot{\phi} \cos \theta + \dot{\psi}) = p_{\psi} //$$

$$\begin{aligned} \star I_0 \ddot{\theta} &= \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) \rightarrow \frac{\partial L}{\partial \theta} = I_0 \dot{\phi}^2 \sin \theta \cdot \cos \theta + I_2 \dot{\phi}^2 \cos \theta \sin \theta \\ &\quad - I_2 \dot{\phi} \dot{\psi} \sin \theta + mgd \sin \theta \\ &= I_0 \dot{\phi}^2 \sin \theta \cos \theta - I_2 \dot{\phi} \sin \theta (\dot{\phi} \cos \theta + \dot{\psi}) + mgd \sin \theta // \end{aligned}$$



Por Euler-Lagrange:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{\theta}} \right) = \frac{\partial L}{\partial \theta}$$

$$I_0 \ddot{\theta} = I_0 \dot{\phi}^2 \sin \theta \cos \theta - I_2 \dot{\phi} \sin \theta (\dot{\phi} \cos \theta + \dot{\psi})$$

$$+ mgd \sin \theta //$$

$$\ddot{\theta} = \dot{\phi}^2 \sin \theta \cos \theta - \frac{I_2}{I_0} \dot{\phi} \sin \theta (\dot{\phi} \cos \theta + \dot{\psi}) + \frac{mgd}{I_0} \sin \theta$$