FYS4150 Project 4: The Ising model in 2 dimensions

Marie Foss (# 56), Maria Hammerstrøm (# 59)

Abstract

Lorem ipsum dolor sit amet, consectetuer adipiscing elit. Ut purus elit, vestibulum ut, placerat ac, adipiscing vitae, felis. Curabitur dictum gravida mauris. Nam arcu libero, nonummy eget, consectetuer id, vulputate a, magna. Donec vehicula augue eu neque. Pellentesque habitant morbi tristique senectus et netus et malesuada fames ac turpis egestas. Mauris ut leo. Cras viverra metus rhoncus sem. Nulla et lectus vestibulum urna fringilla ultrices. Phasellus eu tellus sit amet tortor gravida placerat. Integer sapien est, iaculis in, pretium quis, viverra ac, nunc. Praesent eget sem vel leo ultrices bibendum. Aenean faucibus. Morbi dolor nulla, malesuada eu, pulvinar at, mollis ac, nulla. Curabitur auctor semper nulla. Donec varius orci eget risus. Duis nibh mi, congue eu, accumsan eleifend, sagittis quis, diam. Duis eget orci sit amet orci dignissim rutrum.

Github: https://github.com/mariahammerstrom/Project4

1 Introduction

The project deals with the Ising model in two dimensions, without an external magnetic field. The Ising model is a model to study phase transitions at finite temperature for magnetic systems. In its simplest form the **energy** for a specific microstate i is expressed as

$$E_i = -J \sum_{\langle kl \rangle}^N s_k s_l,\tag{1}$$

with $s_k = \pm 1$, N is the total number of spins and J is a coupling constant expressing the strength of the interaction between neighboring spins. The symbol < kl > indicates that we sum over nearest neighbors only. We will assume that we have a ferromagnetic ordering, meaning J > 0. We will use periodic boundary conditions and the Metropolis algorithm only.

We will be using the **Boltzmann probability distribution** (PDF) defined as

$$w_i = \frac{1}{Z} e^{-\beta E_i},\tag{2}$$

where $\beta = 1/kT$ with k the Boltzmann constant and T the temperature, and with Z as the **partition function** for the canonical ensemble defined as

$$Z = \sum_{i=1}^{\infty} e^{-\beta E_i}.$$
 (3)

summing over all micro states i. w_i expresses the probability of finding the system in a given microstate i.

The magnetic moment of a given microstate is

$$M_i = \sum_j s_j. (4)$$

Some quantities we are interested in calculating, are the **expection values** for the energy $\langle E \rangle$ and magnetic moment $\langle M \rangle$:

$$\langle E \rangle = \frac{1}{Z} \sum_{i} E_{i} e^{-\beta E_{i}}$$

$$\langle M \rangle = \frac{1}{Z} \sum_{i} M_{i} e^{-\beta E_{i}},$$
(5)

as well as the variances for the energy σ_E^2 and for the magnetic moment σ_M^2 :

$$\sigma_E^2 = \langle E^2 \rangle - \langle E \rangle^2 \sigma_M^2 = \langle M^2 \rangle - \langle M \rangle^2,$$
 (6)

where

$$\sigma_E^2 = \frac{1}{Z} \sum_{i} E_i^2 e^{-\beta E_i} - \left[\frac{1}{Z} \sum_{i} E_i e^{-\beta E_i} \right]^2$$
 (7)

First we will assume that we only have two spins in each dimension, that is L=2, where L is the lattice length. In this case, we can find a closed form expression for the partition function in Eq. (3) and the corresponding expectation values for E, $|\mathcal{M}|$, the specific heat C_V and the susceptibility χ as functions of T using periodic boundary conditions. **Periodic boundary conditions** means that the neighbor to the right of a given spin s_N in Eq. (1) takes the value of s_1 . Similarly, the neighbor to the left of s_1 takes the value of s_N .

In the case L = 2 we can write Eq. (1) as:

$$E_{i} = -J \sum_{\langle kl \rangle}^{N} s_{k} s_{l} = -J \sum_{k=1}^{2} \sum_{l=1}^{2} s_{k} s_{l} = -J(s_{1} s_{1} + s_{1} s_{2} + s_{2} s_{1} + s_{2} s_{2})$$
(8)

... EXPRESSIONS!

2 Methods

We write a code for the Ising model which computes the mean energy E, mean magnetization $|\mathcal{M}|$, the specific heat C_V and the susceptibility χ as functions of T using periodic boundary conditions for in the x and y directions. Using the Ising model in two dimensions, the number of configurations is given by 2^N with $N = L \times L$ number of spins for a lattice of length L.

We will use the **Metropolis algorithm**. The algorithm goes as follows:

- 1. Generate a random configuration in the lattice to create an initial state with energy E_b .
- 2. Change the initial configuration by flipping for example one spin only. Compute the energy of this state E_t .
- 3. Calculate $\Delta E = E_t E_b$.
- 4. The Metropolis test: If $\Delta E \leq 0$ the new configuration is accepted, meaning the energy is lowered and we are moving towards the energy minimum at a given temperature. If $\Delta E > 0$, calculate $w = e^{-\beta \Delta E}$. If w < r where r is a random number, accept the new configuration. If else, keep the old configuration.
- 5. Compute the new energy $E' = E_t 0 \Delta E$. Calculate various expectation values using this energy.
- 6. Repeat steps 2-5 for the chosen number of Monte Carlo cycles, meaning the number of times we should sweep through the lattice and summed over all spins.
- 7. Compute the various expectation values by dividing by the total number of cycles (and possibly the number of spins).

3 Results

Our numerical results from computing the mean energy E, mean magnetization $|\mathcal{M}|$, the specific heat C_V and the susceptibility χ as functions of T using periodic boundary conditions, can be compared with expressions given in Eq. (8) - (?) for a temperature T = 1.0 (in units of kT/J). This gives the following results:

Quantity	Closed form	Numerical
Mean energy E	-	=
Mean magnetization $ \mathcal{M} $	-	-
Specific heat C_V	-	-
Susceptibility χ	-	-

The number of Monte Carlo cycles needed to achieve good agreement is XXX.

Next we looked at the case of L=20 spins in the x and y directions. In this case we wanted to perform a study of the time (or number of Monte Carlo cycles) we need before reaching an equilibrium situation and can start computing the various expectation values.

A rough and study is done by plotting the various expectation values as a function of number of Monte Carlo cycles, for both T = 1.0 kT/J and T = 2.4 kT/J, as shown in Fig. XXX. PLOT + COMMENTS!

We also study the total number of accepted configurations as a function of the total number of Monte Carlo cycles, shown in Fig. XXX. PLOT + COMMENTS!

...

By simply counting the number of times a given energy appears in our computation we can compute the probability w_i . MORE.

4 Conclusions

•••

5 List of codes

The codes developed for this project are:

main.cpp - main program (C++).

plotting.m – plotting program which makes plots to study the total number of accepted configurations as a function of total number of Monte Carlo cycles and temperature T (Python).