

# ELECTRIC CURRENT

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# FROM STATIC CHARGES TO FLOWING ELECTRICITY

So far, we've focused on the fascinating world of **electrostatics**:

- Charges are at rest or moving at constant velocity.
- Electric fields are static.
- Potential energy is stored in configurations (like capacitors).

But how do we **use** electricity? How does a light bulb turn on, or a phone charge? The answer isn't in static charges, but in their **movement**.

To truly harness electricity and build the technology that surrounds us, we must understand how electric charge is **transported** through materials. This shift from static configurations to the **continuous flow of charge** is our next big step!

# ELECTRIC CURRENT: THE FLOW OF CHARGE

While electric current describes the flow of charge, we first consider the **current density vector  $\mathbf{J}$** , which provides a more fundamental, microscopic description of this flow.

- $\mathbf{J}$  describes how charges move. It allows **us** to find the amount of charge flowing per unit time through a unit cross-sectional area, perpendicular to the flow.
- It is a **vector quantity**, and its direction is defined as the direction of flow of positive charge.
- **Units:** Amperes per square meter ( $\text{A}/\text{m}^2$ ).

# CURRENT DENSITY AND CHARGE CARRIER MOTION

Think of a material where some charges (called charge carriers) are moving with a velocity known as drift velocity. Given the standard notation

- $n$ : Number of charge carriers per unit volume (charge carrier density,  $\text{m}^{-3}$ )
- $q$ : Charge of a single carrier (Coulombs, C)
- $\mathbf{v}_d$ : Average drift velocity of the charge carriers (meters per second, m/s). This is the average velocity component of the charged particles in the direction of the electric force acting on them.

The current density  $\mathbf{J}$  the charge carriers within the material is

$$\mathbf{J} = nq\mathbf{v}_d$$

To gain a clearer understanding, imagine a material where the charge carrier density ( $n$ ) and drift velocity ( $v_d$ ) are uniform. The flux of  $\mathbf{J}$  through a cross-sectional area  $A$  perpendicular to  $\mathbf{J}$  is:

$$\Phi_A(\mathbf{J}) = \int_A \mathbf{J} \cdot \hat{\mathbf{n}} dA = nq v_d A$$

Giving this formula some thought, we will realize it is just the charge per unit time that crosses the surface. This is what people usually call **electric current  $I$** .

### NOTE

For a current to exist, there must be charge carriers free to move (like electrons in a metal or ions in an electrolyte) and a force acting on them.

The **electric current ( $I$ )** is the total rate at which charge flows through a given cross-sectional area. It's the macroscopic manifestation of current density.

- For any surface  $A$  through which charge flows, the flux of  $\mathbf{J}$  is the current that passes through a surface

$$I = \int \mathbf{J} \cdot d\mathbf{A}.$$

If the current density  $\mathbf{J}$  is uniform and perpendicular to the cross-sectional area  $A$ , then:

$$I = JA$$

## Units:

The SI unit for electric current is the **Ampere (A)**, where 1 Ampere = 1 Coulomb per second (1 A = 1 C/s).

# DRIFT VELOCITY VS. SIGNAL SPEED

It's important to distinguish between the very slow drift velocity of individual charge carriers and the nearly instantaneous speed at which an electrical signal propagates.

- **Drift Velocity ( $v_d$ )**: This is typically very slow, on the order of millimeters per second in metals. Individual electrons literally "drift" through the conductor.
- **Speed of Electric Signal**: When you flip a light switch, the light turns on almost instantly. This is because the electric field, which pushes the electrons, propagates through the wires at a speed close to the speed of light ( $c$ ).

## ANALOGY:

Think of a water pipe already full of water. When you open the tap, water immediately comes out the other end, even though any single water molecule takes a long time to travel through the pipe. The pressure wave travels quickly.

# WHAT DRIVES CURRENT? THE ELECTRIC FIELD

In electrostatics, we learned that the electric field inside a conductor in equilibrium is zero. However, for a current to flow, there must be an electric field present within the conductor.

- An external electric field ( $\mathbf{E}$ ) exerts a force on the free charge carriers within the conductor.
- This force causes the charge carriers to accelerate, leading to their drift.
- Due to collisions with the atoms of the conductor, this acceleration is not continuous, and they reach an average drift velocity  $\mathbf{v}_d$ .
- Therefore, a sustained electric current implies a sustained electric field within the conductor.

## CONTRAST WITH ELECTROSTATICS

Remember: in electrostatic equilibrium, charges arrange themselves to cancel out any internal electric field. With current, charges are **constantly moving**, implying a continuous electric field driving them.

# OHM'S LAW: MICROSCOPIC FORM

The relationship between the electric field that drives the current and the resulting current density is fundamental.

For many materials (especially metals) at a constant temperature, the current density  $\mathbf{J}$  is directly proportional to the applied electric field  $\mathbf{E}$  [Ohm's Law]:

$$\mathbf{J} = \sigma \mathbf{E}$$

Where  $\sigma$  (sigma) is the electrical conductivity of the material.

The reciprocal of conductivity is resistivity ( $\rho$ ):  $\rho = 1/\sigma$ .

So, Ohm's Law can also be written as:

$$\mathbf{E} = \rho \mathbf{J}$$

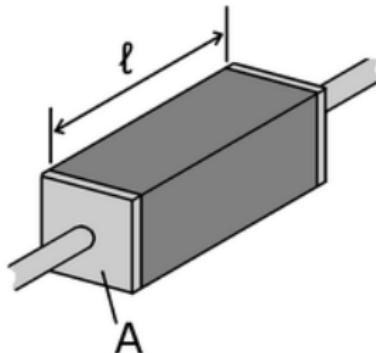
## UNITS:

- Conductivity ( $\sigma$ ) Siemens per meter (S/m) or ( $Om\cdot m$ ) $^{-1}$
- Resistivity ( $\rho$ )  $Om\cdot m$

# RESISTANCE ( $R$ ) - THE MACROSCOPIC PROPERTY

While resistivity ( $\rho$ ) is an intrinsic property of a material, **resistance ( $R$ )** is a macroscopic property of a specific object (like a wire) that depends on the material it is made of, its shape, and size.

Think of a conductor of uniform cross-sectional area  $A$  and length  $\ell$ , carrying a uniform current density  $\mathbf{J}$  driven by a uniform electric field  $\mathbf{E}$ :



Under the described circumstances:

- The potential difference across the conductor is  $V = E\ell$ .
- The current through the conductor is  $I = JA$ .

Using Ohm's Law in microscopic form ( $\mathbf{E} = \rho\mathbf{J}$ ), so

$$E\ell = \rho JL$$

therefore

$$V = \rho \frac{I}{A} \ell$$

which yields

$$V = \left( \frac{\rho \ell}{A} \right) I$$

## OHM'S LAW: MACROSCOPIC FORM

The previous derivation, naturally suggests defining the **resistance** ( $R$ ) of the conductor as:

$$R = \frac{\rho L}{A}$$

Clearly, the SI unit for resistance is the **Ohm ( $\Omega$ )**.

The  $V - I$  relation across the conductor can now be cast as

$$V = IR ,$$

which is the well known form of **Ohm's Law**:

# POWER DISSIPATION IN RESISTORS

When current flows through a resistor, electrical potential energy is continuously converted into other forms of energy, typically thermal energy. This is known as **Joule heating**. The rate at which energy is dissipated (or converted) is the **power ( $P$ )**

- Consider a tiny charge  $dQ$  moving through a potential difference  $V$  in a short time  $dt$ .
- The work done on the charge by the electric field is  $dU = VdQ$
- The power is  $P = dU/dt = V(dQ/dt)$ .
- Since  $I = dQ/dt$ :

$$P = IV$$

- Using Ohm's Law ( $V = IR$  or  $I = V/R$ ), we can express power in alternative forms:

$$P = I^2R \quad \text{or} \quad P = \frac{V^2}{R}$$

- A little reminder, the SI unit for power is the Watt (W) = 1 Joule per second.



# SUMMARY OF ELECTRIC CURRENT

We've explored the fundamental aspects of electric current:

- Current density vector:  $\mathbf{J} = nq\mathbf{v}_d$ , naturally defined from microscopic charge flow.
- Electric Current, flux of the current vector through some surface ( $I = \int \mathbf{J} \cdot \hat{\mathbf{n}} dA$ ).
- The crucial role of the Electric Field (**E**) in driving current within conductors.
- Introduced conductivity and resistivity and Ohm's Law as a relation between fields ( $\mathbf{J} = \sigma \mathbf{E}$  or  $\mathbf{E} = \rho \mathbf{J}$ )

- Concept of Resistance ( $R = \rho L/A$ ).
- The familiar expression of Ohm's Law ( $V = IR$ ) as a current voltage relation.
- Finally, Power Dissipation ( $P = IV = I^2R = V^2/R$ ) in resistances (circuits).

This lays the groundwork for analyzing complete circuits involving current, voltage, and resistance!