

Formelblad

Medelhastighet: $\langle \mathbf{v} \rangle = \frac{1}{A} \int_A \mathbf{v} \cdot \mathbf{n} dA$

Volumetrisk flödeshastighet: $Q = \langle \mathbf{v} \rangle A$

Acceleration: $\boldsymbol{\alpha} = \frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v}$

Newton viskositetslag: $\tau_{ij} = \mu \left(\frac{\partial v_i}{\partial x_j} + \frac{\partial v_j}{\partial x_i} \right); \quad \tau_{xy} = \mu \frac{dv_x}{dy}$

Kontinuitetsekvationen: $\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0; \quad \frac{\partial \rho}{\partial t} = - \left(\frac{\partial(\rho v_x)}{\partial x} + \frac{\partial(\rho v_y)}{\partial y} + \frac{\partial(\rho v_z)}{\partial z} \right)$

Navier-Stokes ekvation för en inkompressibel fluid i rektangulära koordinater:

$$\begin{aligned} \rho \left(\frac{\partial v_x}{\partial t} + v_x \frac{\partial v_x}{\partial x} + v_y \frac{\partial v_x}{\partial y} + v_z \frac{\partial v_x}{\partial z} \right) &= - \frac{\partial p}{\partial x} + \mu \left[\frac{\partial^2 v_x}{\partial x^2} + \frac{\partial^2 v_x}{\partial y^2} + \frac{\partial^2 v_x}{\partial z^2} \right] + \rho g_x \quad \text{i } x\text{-riktningen} \\ \rho \left(\frac{\partial v_y}{\partial t} + v_x \frac{\partial v_y}{\partial x} + v_y \frac{\partial v_y}{\partial y} + v_z \frac{\partial v_y}{\partial z} \right) &= - \frac{\partial p}{\partial y} + \mu \left[\frac{\partial^2 v_y}{\partial x^2} + \frac{\partial^2 v_y}{\partial y^2} + \frac{\partial^2 v_y}{\partial z^2} \right] + \rho g_y \quad \text{i } y\text{-riktningen} \\ \rho \left(\frac{\partial v_z}{\partial t} + v_x \frac{\partial v_z}{\partial x} + v_y \frac{\partial v_z}{\partial y} + v_z \frac{\partial v_z}{\partial z} \right) &= - \frac{\partial p}{\partial z} + \mu \left[\frac{\partial^2 v_z}{\partial x^2} + \frac{\partial^2 v_z}{\partial y^2} + \frac{\partial^2 v_z}{\partial z^2} \right] + \rho g_z \quad \text{i } z\text{-riktningen} \end{aligned}$$

Bernoullis lag: $\frac{1}{2} \rho v^2 + p + \rho g z = \text{konstant}$

Inloppslängd: $L_e = 0.058 D \text{Re}$ (för cylindriskt rör)

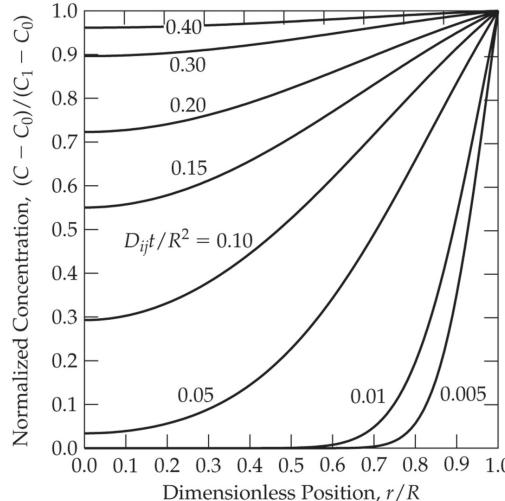
Poiseuilles lag: $Q = \frac{\Delta p \pi R^4}{8 \mu L}$ (för cylindriskt rör)

Ficks första lag: $J_{ix} = -D_{ij} \frac{dc_i}{dx}; \quad$ Ficks andra lag: $\frac{dc_i}{dt} = D_{ij} \frac{d^2 c_i}{dx^2}$

Stoke-Einstiens ekvation: $D_{ij} = \frac{k_B T}{f}; \quad f = 6\pi\mu R$ för en sfär

Instationär diffusion

(i en sfär med radien R)



Dimensionslösa tal:

$Re = \frac{\rho v L}{\mu}$	$De = \sqrt{\delta} Re ; \delta = \frac{a}{R}$	$Pe = \frac{vL}{D_{ij}}$	
$\alpha = R\sqrt{\omega/v} ; v = \frac{\mu}{\rho}$	$Nu = \frac{h_f L}{k}$	$Sh = \frac{k_f L}{D_{ij}}$	$St = \frac{L_x}{T < v_x >}$
$Pe_T = \frac{< v > L}{\alpha}$	$Pr = \frac{v}{\alpha} = \frac{c_p \mu}{k}$	$Sc = \frac{v}{D_{ij}}$	

Darcys lag med gravitation: $\mathbf{v} = -K(\nabla p - \rho g)$

Permeabilitet: $k = K\mu$

Brinkmans ekvation: $\mu \nabla^2 \mathbf{v} - \frac{1}{K} \mathbf{v} - \nabla p = 0$

Konvektion-diffusionsekvationen: $\frac{\partial C}{\partial t} + \nabla \cdot (f \mathbf{v}_f C) = D_{eff} \nabla^2 C + \Phi_B - \Phi_L + Q$

Retardationskoefficienten: $f = \frac{v_s}{v_f}$

Starlings filtreringslag: $J_v = L_p S (\Delta p - \sigma_s \Delta \pi)$

Energiekvationen:

$$\rho C_p \left(\frac{\partial T}{\partial t} + v_x \frac{\partial T}{\partial x} + v_y \frac{\partial T}{\partial y} + v_z \frac{\partial T}{\partial z} \right) = k \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \right) + \dot{Q}_p + \dot{W}_t + \dot{\Phi}_v$$

$$\frac{\partial T}{\partial t} = \alpha \nabla^2 T \quad (\text{endast konduktiv värmetransport})$$

Fouriers lag: $\mathbf{q} = -k \nabla T$

Differentiering, cartesiska koordinater:

$$\begin{aligned} \nabla \psi &= \hat{\mathbf{x}} \frac{\partial \psi}{\partial x} + \hat{\mathbf{y}} \frac{\partial \psi}{\partial y} + \hat{\mathbf{z}} \frac{\partial \psi}{\partial z} \\ \nabla^2 \psi &= \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \end{aligned}$$

Differentiering, cylindriska koordinater:

$$\begin{aligned} \nabla \psi &= \hat{\mathbf{r}}_c \frac{\partial \psi}{\partial r_c} + \hat{\phi} \frac{1}{r_c} \frac{\partial \psi}{\partial \phi} + \hat{\mathbf{z}} \frac{\partial \psi}{\partial z} \\ \nabla^2 \psi &= \frac{1}{r_c} \frac{\partial}{\partial r_c} \left(r_c \frac{\partial \psi}{\partial r_c} \right) + \frac{1}{r_c^2} \frac{\partial^2 \psi}{\partial \phi^2} + \frac{\partial^2 \psi}{\partial z^2} \end{aligned}$$

Differentiering, sfäriska koordinater:

$$\begin{aligned} \nabla \psi &= \hat{\mathbf{r}} \frac{\partial \psi}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial \psi}{\partial \theta} + \hat{\phi} \frac{1}{r \sin \theta} \frac{\partial \psi}{\partial \phi} \\ \nabla^2 \psi &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \\ &= \frac{1}{r} \frac{\partial^2}{\partial r^2} (r \psi) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \end{aligned}$$