

Is Phase Really Needed for Weakly-Supervised Dereverberation ?

Detailed mathematical proofs

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1 Introduction

This document compiles some detailed mathematical proofs as supplementary material to the paper *Is Phase Really Needed for Weakly-Supervised Dereverberation ?* [3]. Section 2 recalls the statistical modeling and notations we use, as well as some additional details on the *Statistical Wave Field Theory* (hypotheses and physical interpretation). Then, complete descriptions of proofs are presented in Section 3.

2 Framework and notations

Introduced in [1], the Statistical Wave Field Theory (SWFT) describes the statistics of reverberation in the asymptotic regime of high frequencies and long times. Following the framework of SWFT, the statistics of a Room Impulse Response (RIR) are obtained by the randomization of the emitting source's position. Precisely, if h is the RIR, H its Fourier transform and V a 3D bounded domain representing the room, the following hypotheses are considered:

- (H1) *High frequencies/long times:* $H(f)$ is studied for $f \gg f_s$ where f_s is the Schroeder frequency. Similarly, $h(t)$ is studied for $t \gg t_m$ where t_m is the mixing time.
- (H2) *Uniform source distribution:* The source's position x_0 is a random variable, with uniform distribution over V . Moreover, the emitted impulse is isotropic.
- (H3) *Spatial stationarity:* The image sources' distribution (or more specifically, the B -function [1]) has stationary and isotropic 1st and 2nd order statistics over space.

Hypotheses (H1), (H2) and (H3) are linked to the properties of *mixing rooms* [2]. In practice, those assumptions are generally satisfied except for rooms presenting strong symmetries (shoebox, spherical or cylindrical rooms, etc).

Solving the waves equation under **(H1)**, **(H2)** and **(H3)**, SWFT leads to a time-frequency description of the RIR statistics, generalizing the so-called Polack model [2]. In formal words:

- $\forall t, h(t)$ is real, centered and Gaussian,
- The Wigner distribution of h is

$$W_h(t, f) := \int_{\mathbb{R}} \mathbb{E} \left[h \left(t + \frac{\tau}{2} \right) h \left(t - \frac{\tau}{2} \right) \right] e^{-2i\pi f\tau} d\tau \quad (1)$$

$$= \mathbb{1}_{\mathbb{R}^+}(t) [B(\cdot) e^{-\alpha(\cdot)t} \star \text{sinc}(4t\cdot)](f) \quad (2)$$

where B is a frequency-dependent amplitude, α a frequency-dependent energy decay rate, and $\text{sinc}(x) := \frac{\sin(\pi x)}{\pi x}$.

The sinc term is a border effect due to the discontinuity of h at $t = 0$. The following equation holds exactly for $t < 0$ and approximately for large values of $t > 0$:

$$W_h(t, f) \simeq \mathbb{1}_{\mathbb{R}^+}(t) B(f) e^{-\alpha(f)t}. \quad (3)$$

By definition of the Wigner distribution, the autocorrelation function is the inverse Fourier transform of W_h over f :

$$\gamma_h(t, \tau) = \mathbb{1}_{\mathbb{R}^+}(t) \int_{\mathbb{R}} B(f) e^{-\alpha(f)t + 2i\pi f\tau} df \quad (4)$$

Moreover, because h is Gaussian and centered, any scalar product $\langle h, \varphi \rangle := \int h\varphi$ is a centered gaussian random variable in \mathbb{R} , and the following covariance formula holds:

$$\text{Cov}(\langle h, \varphi \rangle, \langle h, \psi \rangle) := \iint_{\mathbb{R}^2} \varphi \left(t + \frac{\tau}{2} \right) \psi \left(t - \frac{\tau}{2} \right) \gamma_h(t, \tau) d\tau dt \quad (5)$$

$$= \int_0^{+\infty} \int_{\mathbb{R}} \int_{\mathbb{R}} \varphi \left(t + \frac{\tau}{2} \right) \psi \left(t - \frac{\tau}{2} \right) B(f) e^{-\alpha(f)t + 2i\pi f\tau} df d\tau dt \quad (6)$$

3 Detailed Proofs

3.1 Asymptotic distribution of the RIR in the Fourier domain

The aim of this section is to provide detailed computations that stand behind the main result presented in [3].

Proposition 1. *If we consider the RIR h to be randomly sampled according to the generalized Polack model, as in Section 2, with parameters $\alpha(f)$ and $B(f)$, and H its Fourier transform, then, asymptotically when $f \rightarrow +\infty$,*

$$H(f) \sim \mathcal{N}_{\mathbb{C}} \left(0, \frac{B(f)}{\alpha(f)} \right). \quad (7)$$

Proof. First, we decompose the Fourier transform of h in real and imaginary parts:

$$H(f) = \langle h, c_f \rangle - i \langle h, s_f \rangle \quad (8)$$

where $c_f(t) := \cos(2\pi ft)$ and $s_f(t) := \sin(2\pi ft)$. As a bi-dimensional random variable in the complex plane, $H(f)$ is Gaussian, centered, with a 2×2 covariance matrix $\Sigma(f) = \begin{pmatrix} \sigma_+^2(f) & C(f) \\ C(f) & \sigma_-^2(f) \end{pmatrix}$. Its coefficients can be computed thanks to the covariance formula (6). For instance:

$$\sigma_+^2(f) = \text{Var}(\langle h, c_f \rangle) = \int_0^{+\infty} \int_{\mathbb{R}} \int_{\mathbb{R}} c_f\left(t + \frac{\tau}{2}\right) c_f\left(t - \frac{\tau}{2}\right) B(\xi) e^{-\alpha(\xi)t + 2i\pi\xi\tau} d\xi d\tau dt. \quad (9)$$

Using the formula $\cos(a)\cos(b) = \frac{1}{2}[\cos(a+b) + \cos(a-b)]$, it yields

$$\int_{\mathbb{R}} c_f\left(t + \frac{\tau}{2}\right) c_f\left(t - \frac{\tau}{2}\right) e^{2i\pi\xi\tau} d\tau = \frac{1}{2} \left[\int_{\mathbb{R}} c_f(\tau) e^{-2i\pi\xi\tau} d\tau + \int_{\mathbb{R}} c_f(2t) e^{-2i\pi\xi\tau} d\tau \right] \quad (10)$$

$$= \frac{1}{4} [\delta(\xi + f) + \delta(\xi - f) + 2c_f(2t)\delta(\xi)]. \quad (11)$$

Injecting this expression in (9), and knowing that $B(-f) = B(f)$ and $\alpha(-f) = \alpha(f)$, $\sigma_+^2(f)$ boils down to:

$$\sigma_+^2(f) = \frac{B(f)}{2} \int_0^{+\infty} e^{-\alpha(f)t} dt + \frac{B(0)}{2} \int_0^{+\infty} e^{-\alpha(0)t} c_f(2t) dt. \quad (12)$$

The first integral can be computed directly: $\int_0^{+\infty} e^{-\alpha(f)t} dt = \frac{1}{\alpha(f)}$. Similarly, since $c_f(2t) = \frac{e^{4i\pi ft} + e^{-4i\pi ft}}{2}$, we obtain for the second integral:

$$\int_0^{+\infty} e^{-\alpha(0)t} c_f(2t) dt = \frac{1}{2} \left(\frac{1}{\alpha(0) - 4i\pi f} + \frac{1}{\alpha(0) + 4i\pi f} \right) \quad (13)$$

$$= \frac{\alpha(0)}{\alpha^2(0) + (4\pi f)^2} \quad (14)$$

$$= \frac{1}{\alpha(0) \left[1 + \left(\frac{4\pi f}{\alpha(0)} \right)^2 \right]}. \quad (15)$$

In the end :

$$\sigma_+^2(f) = \frac{B(f)}{2\alpha(f)} + \frac{B(0)}{2\alpha(0) \left[1 + \left(\frac{4\pi f}{\alpha(0)} \right)^2 \right]}. \quad (16)$$

With a similar reasoning, one can obtain:

$$\sigma_-^2(f) = \frac{B(f)}{2\alpha(f)} - \frac{B(0)}{2\alpha(0) \left[1 + \left(\frac{4\pi f}{\alpha(0)} \right)^2 \right]} \quad (17)$$

and

$$C(f) = \frac{2\pi f B(0)}{\alpha^2(0) \left[1 + \left(\frac{4\pi f}{\alpha(0)} \right)^2 \right]}. \quad (18)$$

All in all, when $f \rightarrow +\infty$, we have:

- $\sigma_{\pm}^2(f) \simeq \frac{B(f)}{2\alpha(f)}$,
- $C(f) \simeq 0$,

which means that $\Sigma(f) \simeq \frac{B(f)}{2\alpha(f)} I$, where I is the 2×2 identity matrix, and ends the proof. \square

3.2 Autocorrelation in the Fourier domain

In this section is demonstrated the second result used in [3]:

Proposition 2. *Let h be a general Polack RIR, with acoustics parameters α and B . Then its Fourier transform's autocorrelation function is equal to:*

$$\Gamma_H(f, \xi) := \mathbb{E} \left[H \left(f + \frac{\xi}{2} \right) H^* \left(f - \frac{\xi}{2} \right) \right] = \frac{B(f)}{\alpha(f)} \cdot \frac{1 - i \frac{2\pi\xi}{\alpha(f)}}{1 + \left(\frac{2\pi\xi}{\alpha(f)} \right)^2} \quad (19)$$

Proof. To prove this proposition, we use the properties of the Wigner distribution:

$$\Gamma_H(f, \xi) = \int_{\mathbb{R}} W_h(t, f) e^{-2i\pi\xi t} dt \quad (20)$$

$$= B(f) \int_0^{+\infty} e^{-(\alpha(f) + 2i\pi\xi)t} dt \quad (21)$$

$$= \frac{B(f)}{\alpha(f) + 2i\pi\xi} \quad (22)$$

$$= \frac{B(f)}{\alpha(f)} \cdot \frac{1 - i \frac{2\pi\xi}{\alpha(f)}}{1 + \left(\frac{2\pi\xi}{\alpha(f)} \right)^2} \quad (23)$$

\square

References

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