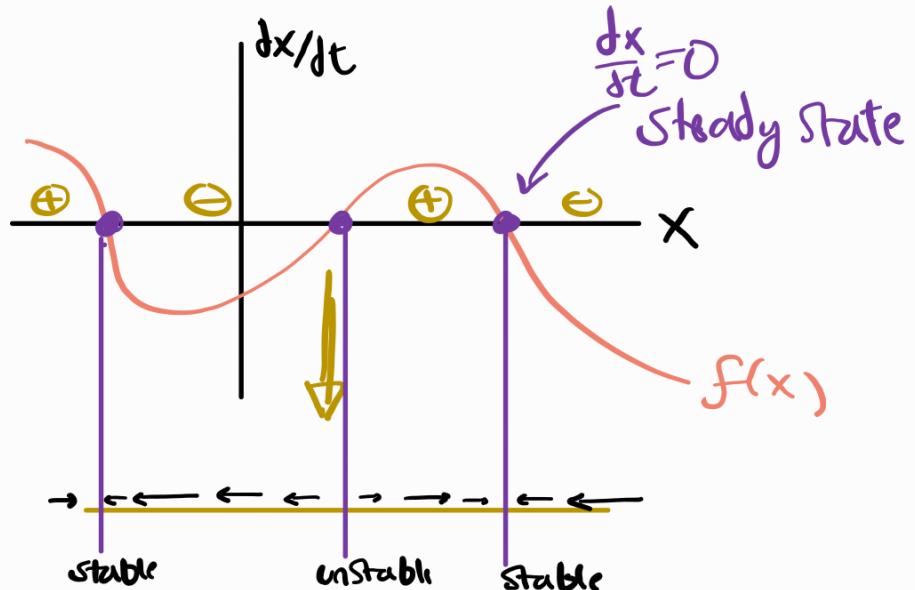


Last Week: Autonomous 1D ODEs

$$\frac{dx}{dt} = f(x) \quad \rightarrow$$



Linear Stability Analysis

We can determine fixed pt stability analytically:

⇒ let x^* be a fixed pt ($f(x^*) = 0$)
⇒ small perturbation

$$\varepsilon(t) = x(t) - x^* \rightarrow \text{how does it evolve?}$$

$$\begin{aligned} \frac{d}{dt}\varepsilon(t) &= \frac{dx}{dt} = f(x(t)) = f(\varepsilon(t) + x^*) \\ &= \underline{\underline{f(x^*)}} + \varepsilon f'(x^*) + O(\varepsilon^2) \end{aligned} \quad \text{↓ Taylor expansion}$$

$$\frac{d\varepsilon}{dt} \approx \varepsilon f'(x^*) \xrightarrow{\text{soln}} \varepsilon(t) = \varepsilon_0 \exp(f'(x^*) t)$$

$f'(x^*) > 0 \Rightarrow$ perturbation grows \Rightarrow unstable

$f'(x^*) < 0 \Rightarrow$ perturbation shrinks \Rightarrow stable

[also, we see perturbations evolve w/ timescale $|1/f'(x^*)|$]

Example

$$\frac{dN}{dt} = rN\left(1 - \frac{N}{K}\right) = f(N) \quad \begin{matrix} r > 0 \\ K > 0 \end{matrix}$$

steady States: $N^* = 0, N^* = K$

Stability: $f'(N) = r\left(1 - \frac{N}{K}\right) - \frac{rN}{K}$

$$\Rightarrow f'(0) = r > 0 \quad f'(K) = -r < 0$$

Unstable Stable

Autonomous ODEs in 2D

General form in higher dimensions

$$\left\{ \begin{array}{l} \frac{dx_1}{dt} = F_1(x_1, x_2, \dots, x_N) \\ \frac{dx_2}{dt} = F_2(x_1, x_2, \dots, x_N) \\ \vdots \\ \frac{dx_N}{dt} = F_N(x_1, x_2, \dots, x_N) \end{array} \right.$$

2D Case

$$\left\{ \begin{array}{l} \frac{dx}{dt} = f(x, y) \\ \frac{dy}{dt} = g(x, y) \end{array} \right.$$

X-nullcline

$$\{(x, y) \mid f(x, y) = 0\}$$

Y-nullcline

$$\{(x, y) \mid g(x, y) = 0\}$$

Intersections are fixed points

2D phase line \Rightarrow 2D phase plane

Example

$$\begin{cases} \frac{dx}{dt} = (x-1)(y+1) \\ \frac{dy}{dt} = (y+x)(x+2) \end{cases} \Rightarrow \text{plot nullclines in the phase plane}$$

X-nullcline

$x=1, y: \text{any}$

OR

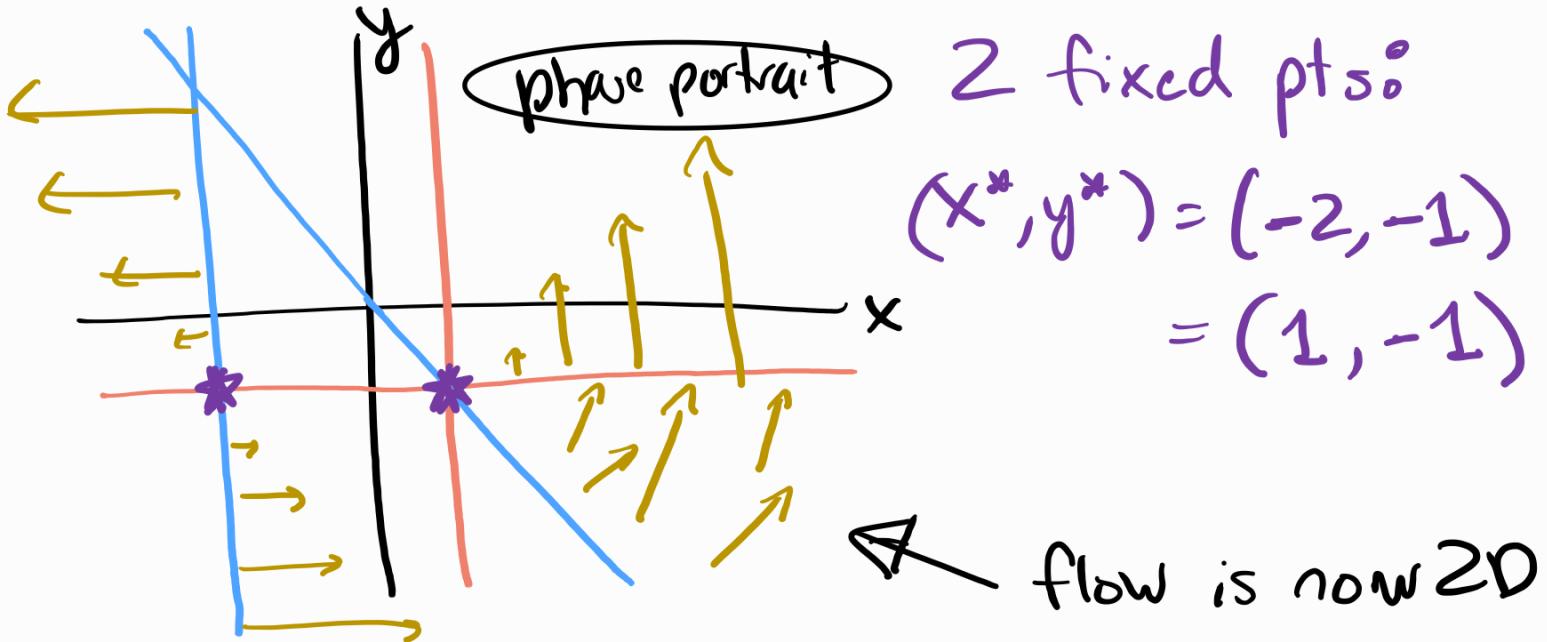
$y=-1, x: \text{any}$

Y-nullcline

$y = -x$

OR

$x=-2, y: \text{any}$



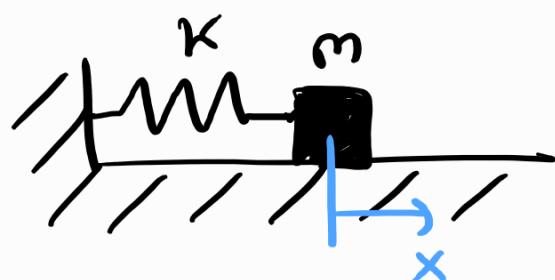
Direction field can help indicate stability, but for nonlinear problems the flow can change dramatically close vs. far

\Rightarrow first, linear examples

\Rightarrow then, linear Stability analysis

Example

Consider a mass on a spring on a smooth table
Let x be the deviation in position from rest.

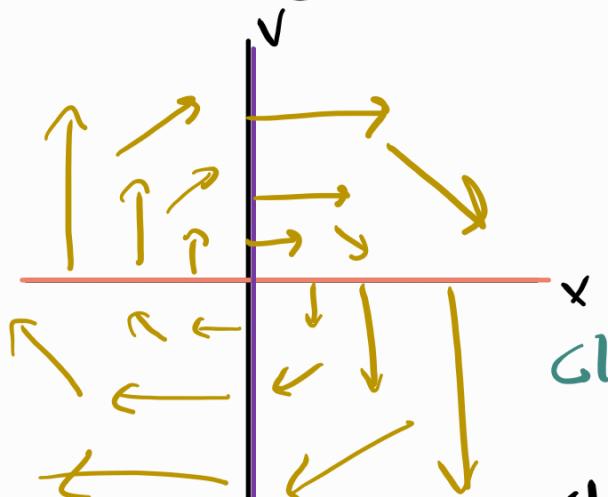


Newton's law tells us

$$m\ddot{x} = -kx \quad \text{← Spring force}$$

We can rewrite this as

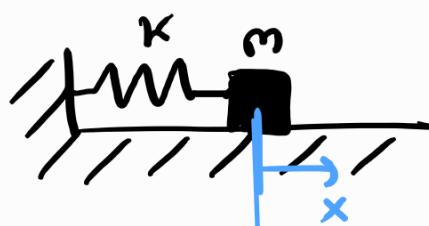
$$\begin{cases} \frac{dx}{dt} = v & \text{velocity} \\ \frac{dv}{dt} = -\omega^2 x & \text{acceleration} \end{cases} \quad \omega^2 = k/m$$



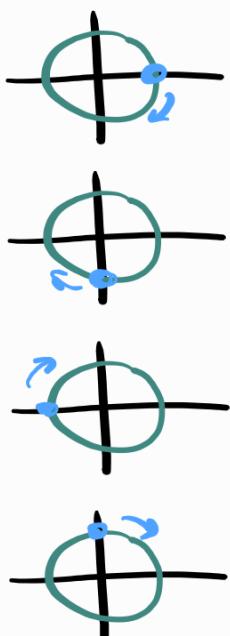
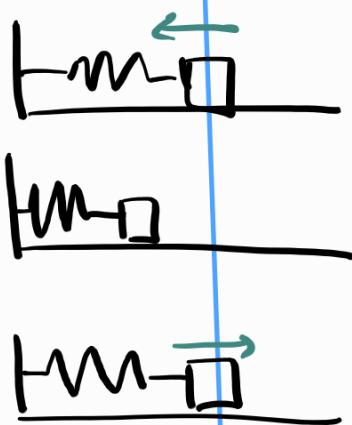
x nullcline: $v=0$

v nullcline: $x=0$

Closed orbits around fixed points?



$t=0$



Example (Linear, Saddle pt)

$$\begin{cases} \frac{dx}{dt} = x + y \\ \frac{dy}{dt} = 4x - 2y \end{cases} \Rightarrow \frac{d\bar{x}}{dt} = \underline{\underline{A}} \bar{x} \text{ with } \underline{\underline{A}} = \begin{bmatrix} 1 & 1 \\ 4 & -2 \end{bmatrix}$$

\bar{x} -nullcline: $y = -x$ y -nullcline: $y = 2x$

Q) What do solns look like?

direction
↓

⇒ Consider ansatz $\bar{x}(t) = \exp(\lambda t) \bar{v}$

$$\frac{d\bar{x}}{dt} = \cancel{\lambda \exp(\lambda t) \bar{v}} = \underline{\underline{A}} \bar{x} = \cancel{\exp(\lambda t) \underline{\underline{A}}} \bar{v} \quad \begin{matrix} \uparrow \\ \text{growth rate} \end{matrix}$$

$$\Rightarrow \underline{\underline{A}} \bar{v} = \lambda \bar{v} \quad \begin{matrix} \text{eigenvalue} \\ \text{eigenvector} \end{matrix} \quad \text{pair!} \quad \begin{matrix} v_1, \lambda_1 \\ v_2, \lambda_2 \end{matrix}$$

If you have initial condition

$$\bar{x}_0 = c_1 \bar{v}_1 + c_2 \bar{v}_2$$

\Downarrow

Note: \bar{v}_1, \bar{v}_2
are linearly
independent

$$\bar{x}(t) = c_1 \exp(\lambda_1 t) \bar{v}_1 + c_2 \exp(\lambda_2 t) \bar{v}_2$$

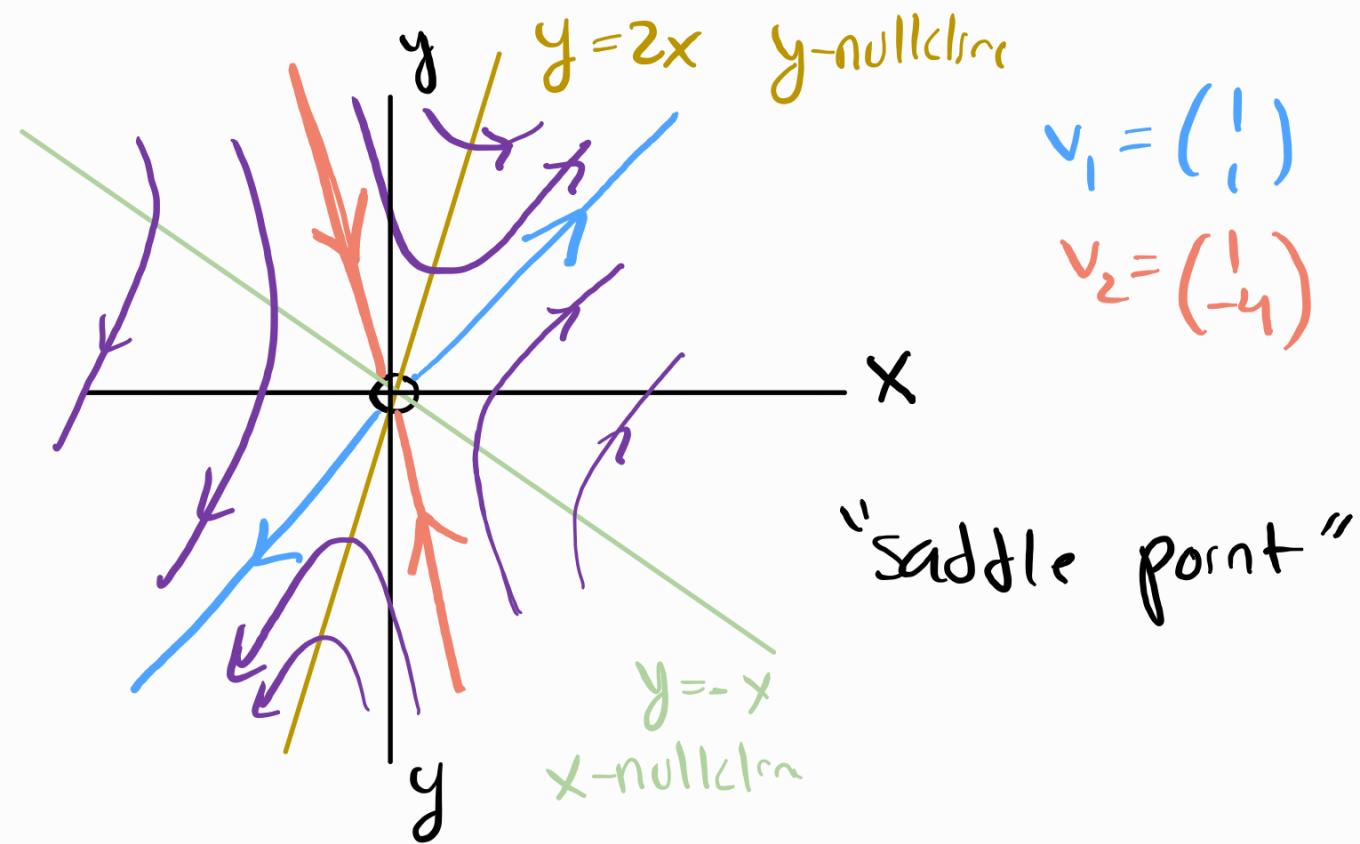
⇒ By uniqueness, the only soln

$$\lambda_1 = 2, \bar{v}_1 = \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

growing

$$\lambda_2 = -3, \bar{v}_2 = \begin{bmatrix} 1 \\ -4 \end{bmatrix}$$

decaying



Suppose matrix $\underline{A}^{2 \times 2}$ has eigenvalues λ_1, λ_2

\rightsquigarrow we know from lin alg that:

$$\begin{cases} \text{tr}(\underline{A}) = T = \lambda_1 + \lambda_2 \\ \det(\underline{A}) = \Delta = \lambda_1 \lambda_2 \end{cases} \quad \begin{matrix} \text{note: both} \\ \underline{T}, \underline{\Delta} \in \mathbb{R} \end{matrix}$$

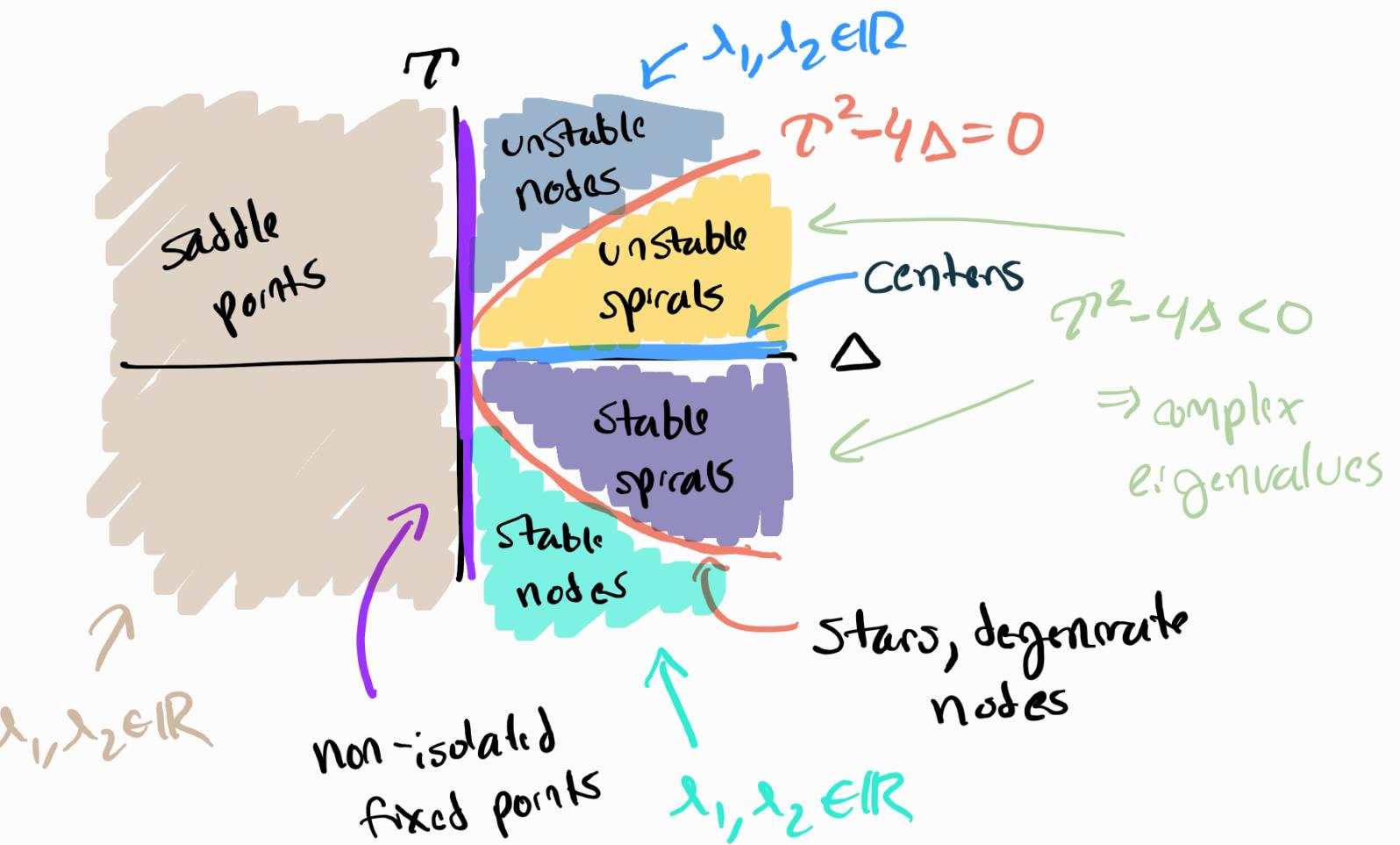
\rightarrow Both λ_1, λ_2 and T, Δ classify S.S.

Note also:

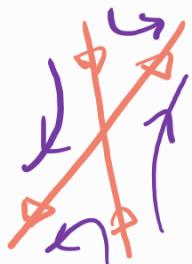
$$\lambda_{1,2} = \frac{1}{2} (T \pm \sqrt{T^2 - 4\Delta})$$

\Rightarrow if $T^2 - 4\Delta < 0 \Rightarrow$ complex eigenvalues

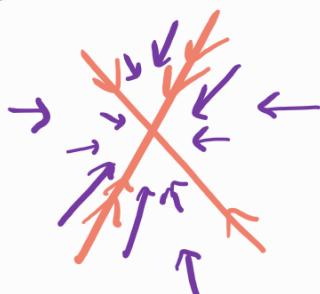
\Rightarrow if $T=0, \Delta > 0 \Rightarrow$ purely imaginary eigenvalues



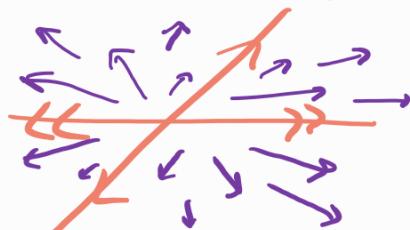
Saddle



stable node



unstable node



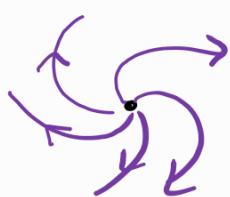
center



stable spiral

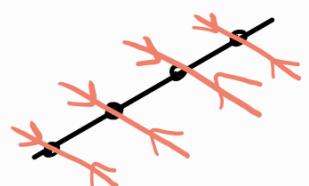


unstable spiral



non-isolated f.p.

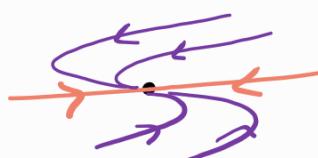
$\Delta = 0 \rightarrow$ at least one $\lambda \equiv 0$



stars



degenerate node



$\lambda_1 = \lambda_2 = \frac{\tau}{2} \neq 0$ \longrightarrow \oplus one distinct eigenvector
 \oplus two distinct eigenvectors

all trajectories
are straight lines

as $t \rightarrow \infty$,
trajectories are \parallel
to eigenvector

Linearization (analogous to 1D)

Let

$$u = x - x^* \quad v = y - y^*$$

u, v small

$$\Rightarrow \frac{\partial u}{\partial t} = \frac{\partial x}{\partial t} = f(x, y) = f(u+x^*, v+y^*)$$

Taylor Series

$$\approx f(x^*, y^*) + u \frac{\partial f}{\partial x}(x^*, y^*) + v \frac{\partial f}{\partial y}(x^*, y^*)$$

$$+ \frac{u^2}{2} \frac{\partial^2 f}{\partial x^2}(x^*, y^*) + \cancel{u v} \frac{\partial^2 f}{\partial x \partial y}(x^*, y^*)$$

$$+ \frac{v^2}{2} \frac{\partial^2 f}{\partial y^2}(x^*, y^*) + O(u^3, uv^2, u^2v, v^3)$$

Small!

$$\Rightarrow \begin{cases} \frac{\partial u}{\partial t} \approx u \underbrace{\frac{\partial f}{\partial x}(x^*, y^*)}_{\text{constants!}} + v \underbrace{\frac{\partial f}{\partial y}(x^*, y^*)}_{\text{constants!}} \\ \frac{\partial v}{\partial t} \approx u \underbrace{\frac{\partial g}{\partial x}(x^*, y^*)}_{\text{constants!}} + v \underbrace{\frac{\partial g}{\partial y}(x^*, y^*)}_{\text{constants!}} \end{cases}$$

linear system

constants!

$$\frac{d}{dt} \begin{bmatrix} u \\ v \end{bmatrix} = J \Big|_{x^*, y^*} \begin{bmatrix} u \\ v \end{bmatrix} \quad \text{Jacobian}$$

$$J = \begin{bmatrix} \frac{\partial f}{\partial x} & \frac{\partial f}{\partial y} \\ \frac{\partial g}{\partial x} & \frac{\partial g}{\partial y} \end{bmatrix}$$

\Rightarrow find eigenvalues, classify them

Example

S.p. $(-2, -1)$ & $(1, -1)$

$$\begin{cases} \frac{dx}{dt} = (x-1)(y+1) \\ \frac{dy}{dt} = (y+x)(x+2) \end{cases} \Rightarrow J = \begin{bmatrix} y+1 & x-1 \\ y+2x+2 & x+2 \end{bmatrix}$$

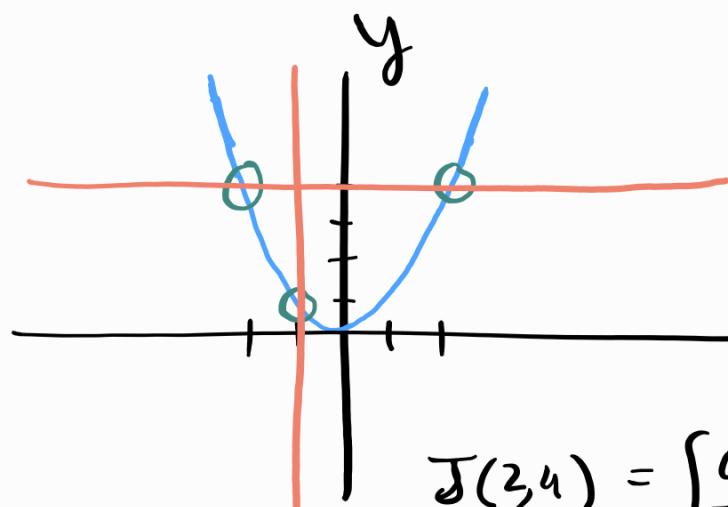
$$J(-2, -1) = \begin{bmatrix} 0 & -3 \\ -3 & 0 \end{bmatrix} \Rightarrow \begin{matrix} r=0 \\ \Delta = -9 \end{matrix} \rightarrow \text{saddle pt}$$

$$J(1, -1) = \begin{bmatrix} 0 & 0 \\ 3 & 3 \end{bmatrix} \Rightarrow \begin{matrix} r=3 \\ \Delta = 0 \end{matrix} \rightarrow \text{'non isolated S.p.'}$$

* note: this classification is only local, because we linearized for small u, v

Example

$$\begin{cases} \frac{dx}{dt} = (x+1)(y-4) \\ \frac{dy}{dt} = y - x^2 \end{cases} \Rightarrow J = \begin{bmatrix} y-4 & x+1 \\ -2x & 1 \end{bmatrix}$$



$$(x^*, y^*) = (2, 4)$$

$$(x^*, y^*) = (-2, 4)$$

$$(x^*, y^*) = (1, 1)$$

$$J(2, 4) = \begin{bmatrix} 0 & 3 \\ -4 & 1 \end{bmatrix} \Rightarrow \begin{matrix} r=1 \\ \Delta=12 \end{matrix} \rightarrow \text{unstable spiral}$$

$$J(-1, 1) = \begin{bmatrix} -5 & 0 \\ 2 & 1 \end{bmatrix} \Rightarrow \begin{matrix} r=-4 \\ \Delta=-5 \end{matrix} \Rightarrow \text{saddle}$$

$$J(-2, 4) = \begin{bmatrix} 0 & -1 \\ 4 & 1 \end{bmatrix} \Rightarrow \begin{matrix} r=1 \\ \Delta=4 \end{matrix} \rightarrow \text{unstable spiral}$$

Why spiral?

Suppose $\lambda = \alpha + i\beta$:

$$\exp(\lambda t) = \exp(\alpha t) \exp(i\beta t)$$

$$= \underbrace{\exp(\alpha t)}_{\text{exponential magnitude}} \left[\underbrace{\cos(\beta t) + i \sin(\beta t)}_{\text{oscillatory}} \right]$$

~ if $\operatorname{Re}(\lambda) < 0 \Rightarrow$ Stable spiral

$\operatorname{Re}(\lambda) > 0 \Rightarrow$ unstable spiral

$\operatorname{Re}(\lambda) = 0 \Rightarrow$ center

note: neither \rightarrow stable nor unstable!

Quiz

Consider the following "Romeo-Juliet" models:

$$\left\{ \begin{array}{l} \frac{dR}{dt} = \alpha R \\ \frac{dJ}{dt} = 2R - J \end{array} \right.$$

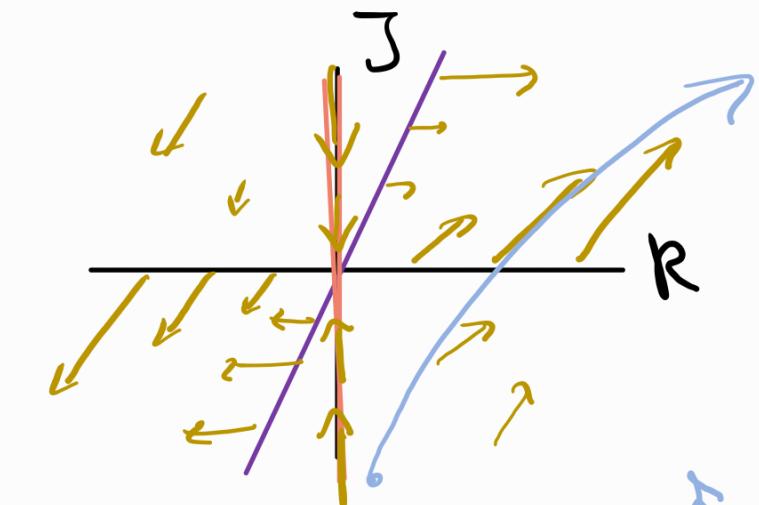
$$\left\{ \begin{array}{l} \frac{dR}{dt} = \alpha R \\ \frac{dJ}{dt} = 2R - J \end{array} \right.$$

- a) Juliet feels cautious of her own feelings but is responsive to Romeo's. Romeo does not care how Juliet feels.
 If $\alpha > 0$, he is bolstered by his own feelings. If $\alpha < 0$, he is cautious.
 If $\alpha = 0$, nothing could ever change how he feels.

- b) Find and sketch the nullclines
 when $\alpha = 2$

$$\frac{dR}{dt} = 0 \Rightarrow R = 0$$

$$\frac{dJ}{dt} = 0 \Rightarrow J = 2R$$



- c) Sketch the phase portrait ($\alpha = 2$)
 d) saddle node
 e) eventually, they will be madly in love