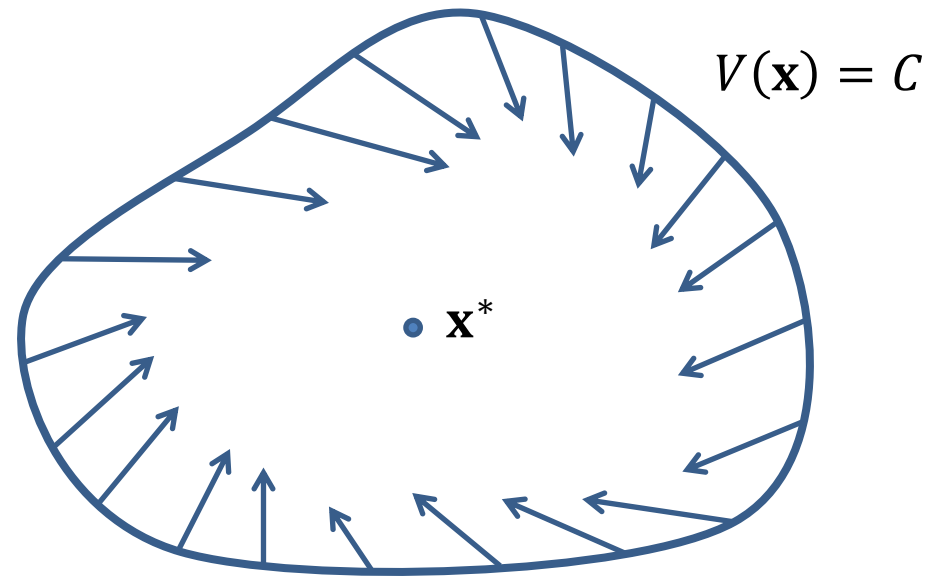


Lecture 4: Lyapunov Functions

- Suppose there exists a connected orientable region (meaning there is an inside and outside) defined by $\{\mathbf{x}: V(\mathbf{x}) \leq C\}$ surrounding an equilibrium point \mathbf{x}^* so that all flows crossing the boundary point remain inside the region.
- Once inside the region, the flow cannot escape



Nested boundaries

Consider a nested sequence of surfaces defined by a reducing set of constants. The increasing normal to each surface is given by $\underline{\nabla}V$

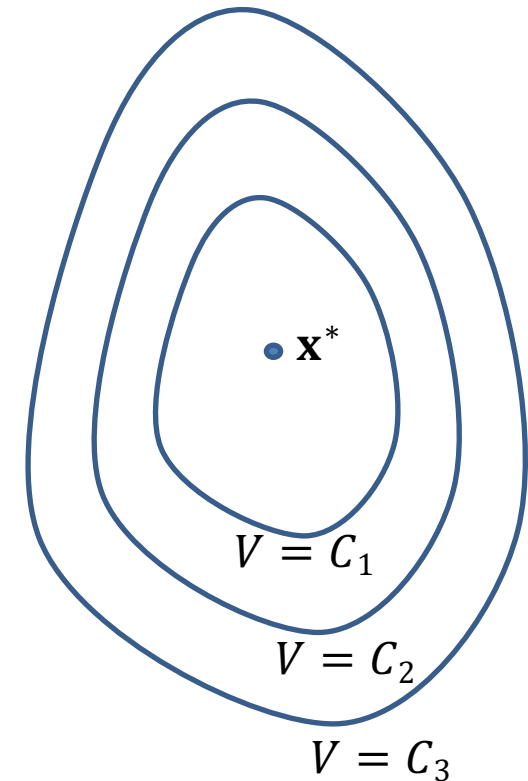
We require the flow to point inwards, i.e. $\underline{\nabla}V \cdot \dot{\mathbf{x}} \leq 0$ But $\dot{\mathbf{x}} = f(\mathbf{x})$. We thus require

$$\underline{\nabla}V \cdot f(\mathbf{x}) \leq 0$$

But

$$\frac{dV}{dt} = \sum \frac{\delta V}{\delta x_i} \frac{dx_i}{dt} = \underline{\nabla}V \cdot \dot{\mathbf{x}} \leq 0$$

so flows go downhill and end up at the bottom



Lyapunov's Theorem

- Let \mathbf{x}^* be an equilibrium point of $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x})$, i.e. $\mathbf{f}(\mathbf{x}^*) = 0$. Let D be an open set surrounding \mathbf{x}^* and let $V(\mathbf{x}): D \rightarrow \mathbb{R}$ be a continuously differentiable function on D such that

1. $V(\mathbf{x}^*) = 0$ and $V(\mathbf{x}) > 0$ for all $\mathbf{x} \neq \mathbf{x}^*$
2. $\dot{V}(\mathbf{x}) = \underline{\nabla} V \cdot \mathbf{f}(\mathbf{x}) \leq 0$

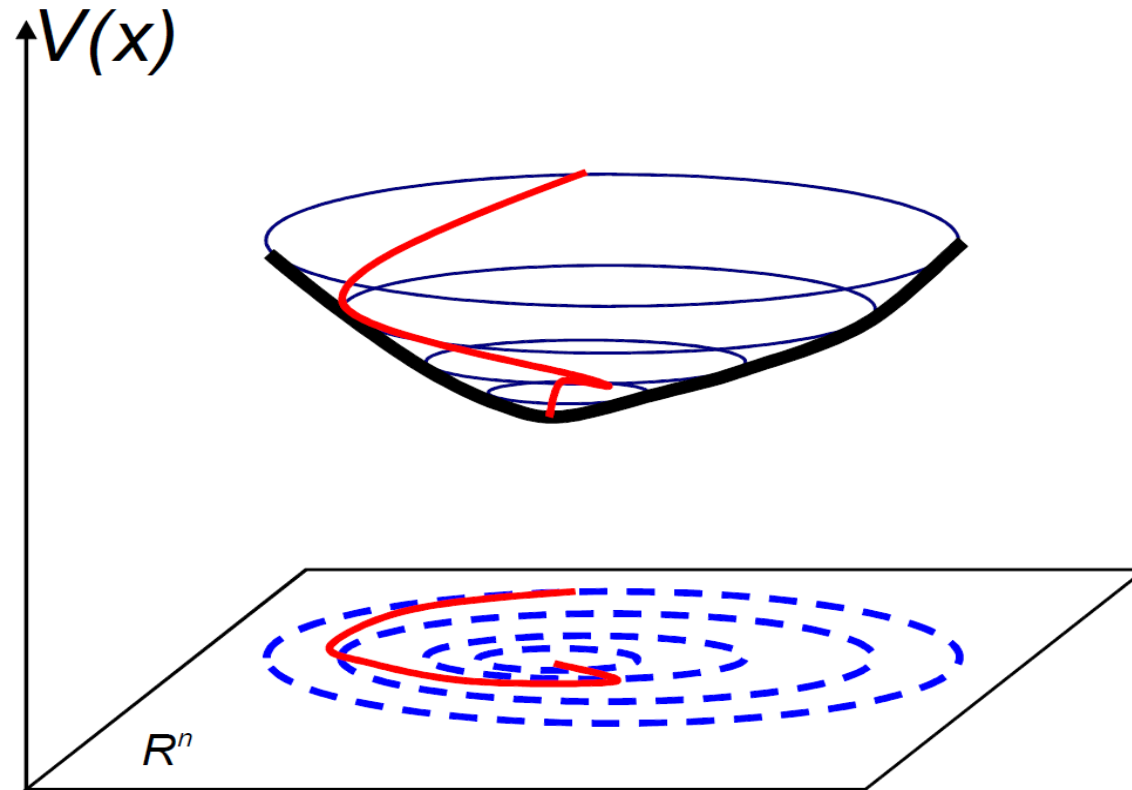
then \mathbf{x}^* is **stable**. If, in addition

3. $\dot{V}(\mathbf{x}) < 0$ for all $\mathbf{x} \neq \mathbf{x}^*$

then \mathbf{x}^* is **asymptotically stable**

- $V(\mathbf{x})$ is called a Lyapunov function
- If $\lim_{\|\mathbf{x}\| \rightarrow \infty} V(\mathbf{x}) = \infty$ and $D = \mathbb{R}^n$ then \mathbf{x}^* is **globally asymptotically stable**

Illustration



$V(\mathbf{x})$ decreases along solution trajectories

Example 1

Consider the dynamical system

$$\dot{x} = y$$

$$\dot{y} = -x + \epsilon x^2 y$$

Equilibrium: $(0,0)$ has Jacobian: $\begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$ with eigenvalues: $\lambda = \pm j$

so Hartman-Grobman doesn't apply

Let $V(x, y) = \frac{1}{2}(x^2 + y^2)$,

$$\frac{dV}{dt} = \underline{\nabla} V \cdot \dot{\mathbf{x}} = x\dot{x} + y\dot{y} = xy - xy + \epsilon x^2 y^2 = \epsilon x^2 y^2$$

If $\epsilon < 0$ then $(0,0)$ is stable

Example 2

$$\dot{x}_1 = -2x_2 + x_2x_3$$

$$\dot{x}_2 = x_1 - x_1x_3$$

$$\dot{x}_3 = x_1x_2$$

- Equilibrium point: $(0, 0, 0)$ is a linear centre. Let

$$V(\mathbf{x}) = \frac{1}{2}(c_1x_1^2 + c_2x_2^2 + c_3x_3^2)$$

Then

$$\begin{aligned}\dot{V} &= \underline{\nabla} V \cdot \dot{\mathbf{x}} = c_1x_1(-2x_2 + x_2x_3) + c_2x_2(x_1 - x_1x_3) + c_3x_3x_1x_2 \\ &= (c_1 - c_2 + c_3)x_1x_2x_3 + (-2c_1 + c_2)x_1x_2\end{aligned}$$

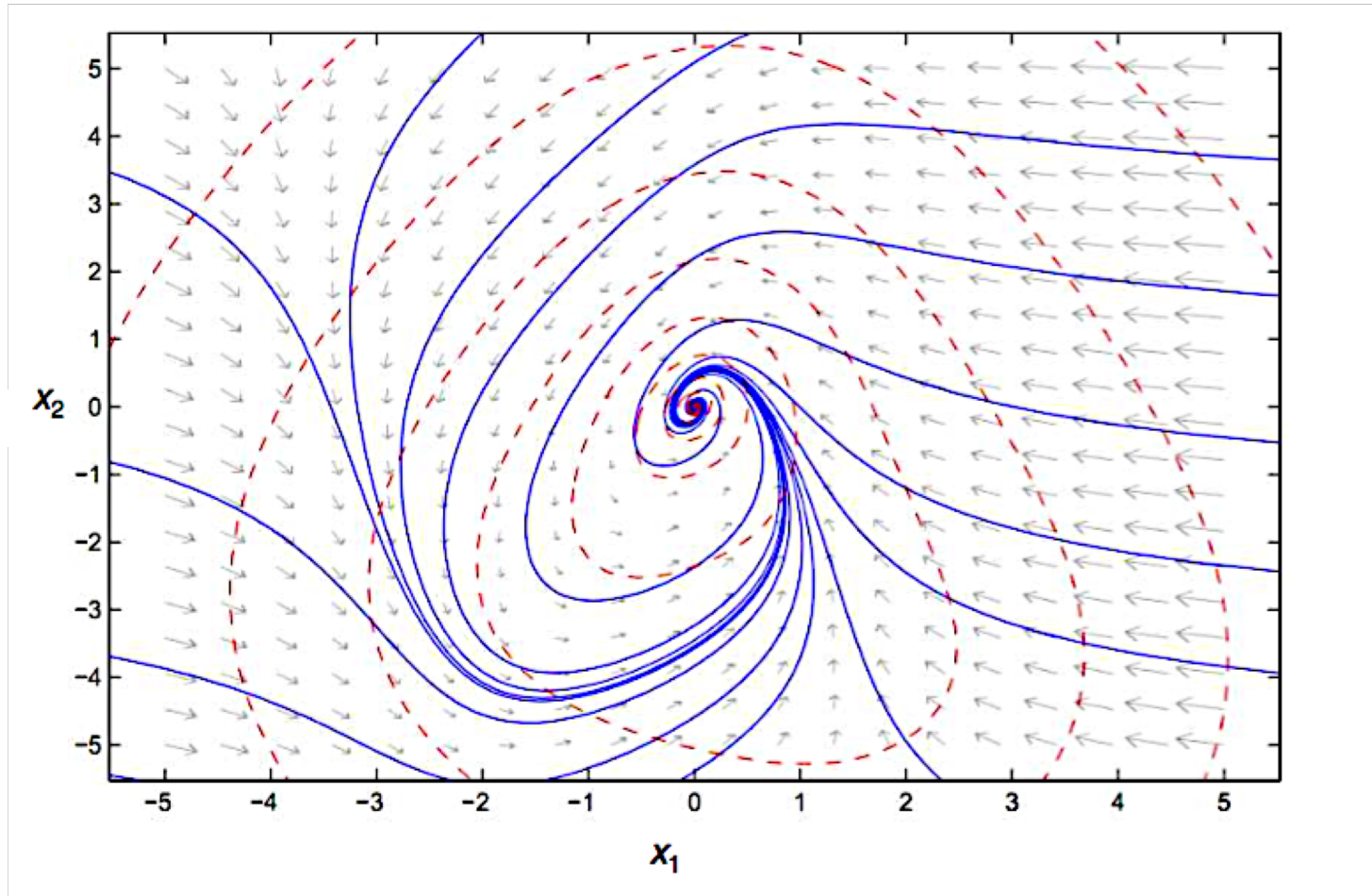
- Choose $c_2 = 2c_1 > 0$ and $c_3 = c_1$, then $\dot{V} = 0$, so equilibrium is stable
- $\dot{V} = 0$ on $V(\mathbf{x}) = \frac{1}{2}(x_1^2 + 2x_2^2 + x_3^2)$ so $\mathbf{x}(t)$ lies on $V(\mathbf{x}) = \text{const}$

Jet Engine Example

$$\begin{aligned}\dot{x}_1 &= -x_2 + 1.5x_1^2 - 0.5x_1^3 \\ \dot{x}_2 &= 3x_1 - x_2\end{aligned}$$

- Equilibrium point: $(0,0)$ has Jacobian $\begin{bmatrix} 0 & -1 \\ 3 & -1 \end{bmatrix}$ with $\lambda = \frac{-1 \pm j\sqrt{11}}{2}$ so is a linear stable focus.
- Hartman-Grobman theorem states that the non-linear system is stable (but only close to the origin)
- Lyapunov functions can extend this result globally using specially constructed functions – see the lecture notes

Jet Engine Example



Level curves of the Lyapunov function showing global stability of the Jet engine model

Vector fields possessing an integral

- Consider the flow associated with the solution of $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x})$ as a vector field

- This is said to have an integral $I(\mathbf{x})$ (a scalar function) if

$$\frac{dI(\mathbf{x})}{dt} = \frac{\partial I(\mathbf{x})}{\partial \mathbf{x}} \dot{\mathbf{x}} = \frac{\partial I(\mathbf{x})}{\partial \mathbf{x}} \mathbf{f}(\mathbf{x}) = 0$$

- $\frac{\partial I(\mathbf{x})}{\partial \mathbf{x}}$ is the gradient vector of $I(\mathbf{x})$
- $I(\mathbf{x})$ defines level sets which contain the flow

Pendulum example

$$\begin{aligned}\dot{q} &= p \\ \dot{p} &= -\frac{g}{l} \sin q\end{aligned}$$

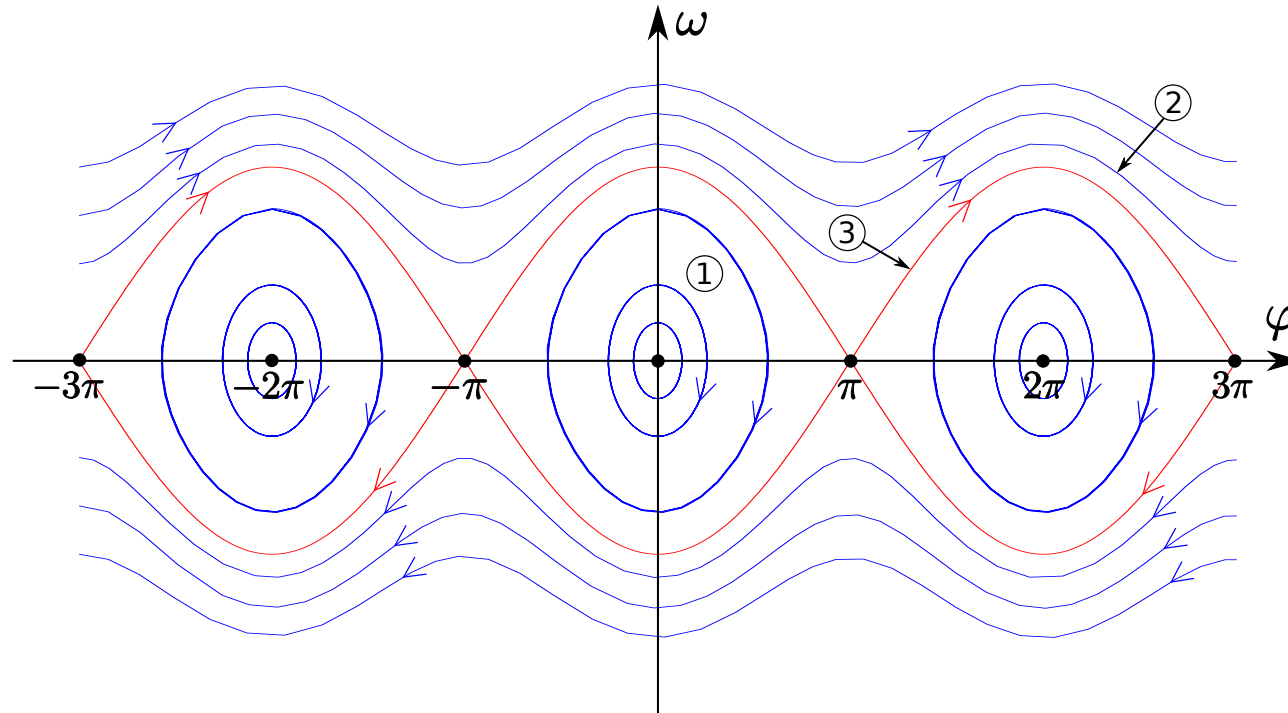
The total stored energy E is conserved

$$E = \frac{1}{2}p^2 - \frac{g}{l} \cos q$$

i.e.

$$\frac{dE}{dt} = p\dot{p} + \dot{q} \frac{g}{l} \sin q = 0$$

Pendulum example



Phase plane of pendulum and level sets of constant energy

Duffing Oscillator for $\delta=0$

$$\begin{aligned}\dot{x} &= y \\ \dot{y} &= x - x^3\end{aligned}$$

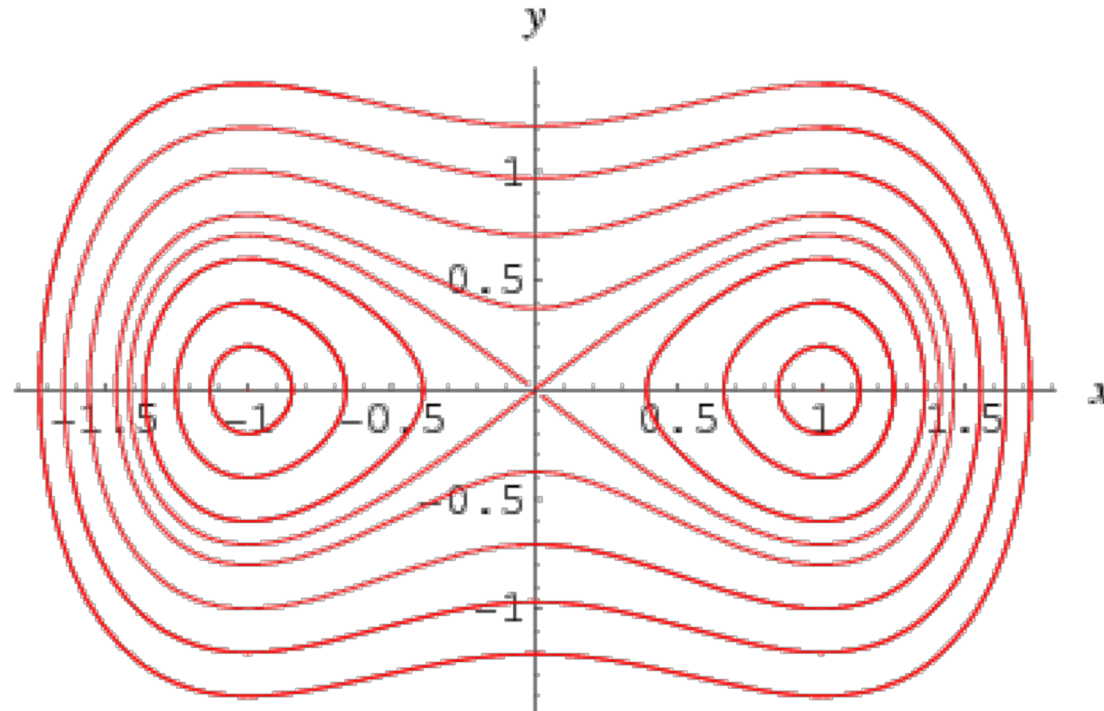
We require

$$\begin{aligned}\frac{dI}{dt} &= \frac{\partial I}{\partial x} \dot{x} + \frac{\partial I}{\partial y} \dot{y} = 0 \\ \frac{\partial I}{\partial x} y + \frac{\partial I}{\partial y} (x - x^3) &= 0\end{aligned}$$

So, for example

$$I = \frac{y^2}{2} - \frac{x^2}{2} + \frac{x^4}{4}$$

Duffing Oscillator for $\delta=0$



Level sets of $I(\mathbf{x})$ in the phase plane of the duffing oscillator

Hamiltonian systems

Hamiltonian systems have vector fields that possess an integral

Definition: Systems of the form

$$\begin{aligned}\dot{\mathbf{p}} &= \mathbf{f}(\mathbf{p}, \mathbf{q}) \\ \dot{\mathbf{q}} &= \mathbf{g}(\mathbf{p}, \mathbf{q})\end{aligned}$$

such that

$$\mathbf{f}(\mathbf{p}, \mathbf{q}) = \partial H(\mathbf{p}, \mathbf{q}) / \partial \mathbf{q}, \quad \mathbf{g}(\mathbf{p}, \mathbf{q}) = -\partial H(\mathbf{p}, \mathbf{q}) / \partial \mathbf{p}$$

are called **Hamiltonian Systems**.

- \mathbf{p} and \mathbf{q} are real vectors with n elements
- H is a twice differentiable function called the Hamiltonian
- \mathbf{q} is the vector of generalised positions, \mathbf{p} the vector of generalised momenta
- All Hamiltonian systems are conservative by construction

More on Hamiltonian systems

- If $(\mathbf{p}^*, \mathbf{q}^*)$ is an equilibrium and $H(\mathbf{p}, \mathbf{q}) > 0$ in a region surrounding the equilibrium, then the equilibrium is stable
- A Newtonian system of the form $\ddot{x} = f(x)$ can be written as a Hamiltonian system by summing the potential energy and kinetic energy

$$\begin{aligned}\dot{x} &= v \\ \dot{v} &= f(x) \\ H(x, v) &= \frac{v^2}{2} - \int_{x_0}^x f(s) ds\end{aligned}$$

Gradient Systems

Definition: Let $V(\mathbf{x})$ be a twice differentiable function in a region $D \subseteq \mathbb{R}^n$. The system

$$\dot{x}_i = -\frac{\partial V}{\partial x_i}$$

is called a **gradient** system.

- Equilibrium points are the critical points of V . Away from critical points the trajectories are orthogonal to the level sets of V .
- If \mathbf{x}^* is a strict local minimum of V then $V(\mathbf{x}) - V(\mathbf{x}^*)$ is a Lyapunov function for \mathbf{x}^* , showing that \mathbf{x}^* is asymptotically stable. If \mathbf{x}^* is a strict local maximum, then the equilibrium is unstable.

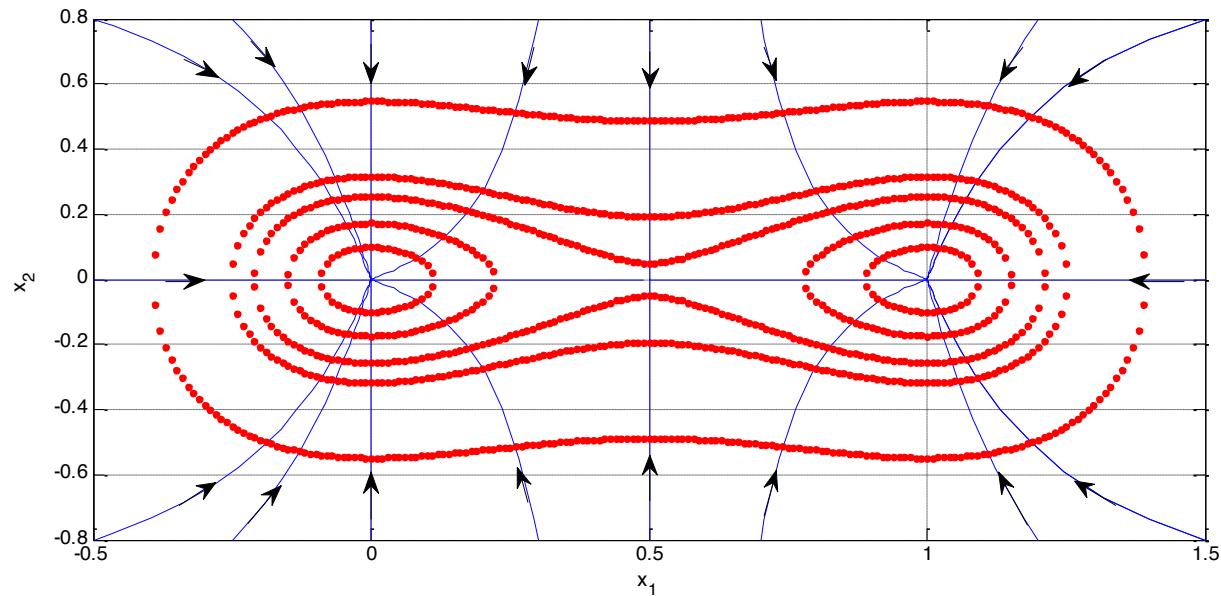
Example Gradient System

$$\dot{x} = -4x(x - 1)(x - 0.5)$$

$$\dot{y} = -2y$$

Has

$$V(x, y) = x^2(x - 1)^2 + y^2$$



Relationship between Gradient and Hamiltonian Systems

- The system

$$\begin{aligned}\dot{x} &= f(x, y) = \frac{\partial H}{\partial y} \\ \dot{y} &= g(x, y) = -\frac{\partial H}{\partial x}\end{aligned}$$

is orthogonal to

$$\begin{aligned}\dot{x} &= g(x, y) \\ \dot{y} &= -f(x, y)\end{aligned}$$

- They have the same equilibria, centres map to nodes, saddles to saddles and foci to foci.