Lecture 5: Asymptotic behaviour

Global trajectories:

- Suppose $\mathbf{x}(t) = \phi(t, \mathbf{x}_0)$ is a solution of $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x})$ such that $\mathbf{x}(0) = \mathbf{x}_0$ (i.e. for some set D, $\frac{\partial \phi(t, \mathbf{x})}{\partial t} = \mathbf{f}(\mathbf{x}) \ \forall \ \mathbf{x} \in D$, with $\phi(0, \mathbf{x}_0) = \mathbf{x}_0 \in D$)
- This solution defines (for t < 0 as well as $t \ge 0$) a path or trajectory: $\Gamma_{\mathbf{x}_0} = \{\mathbf{x} \in D: \ \mathbf{x} = \phi(t, \mathbf{x}_0), t \in \mathbb{R}\}$
- We want to determine the asymptotic behaviour of this solution Hence define the α and ω limit points of the trajectory

Limit points

Definition: A point $\mathbf{p} \in D$ is called an ω **limit point** of the trajectory $\phi(t, \mathbf{x})$ if there exists a sequence of times $\{t_i\}$, $t_i \to \infty$, such that $\lim_{i \to \infty} \phi(t_i, \mathbf{x}) = \mathbf{p}$

this point is denoted $\omega(\mathbf{x})$.

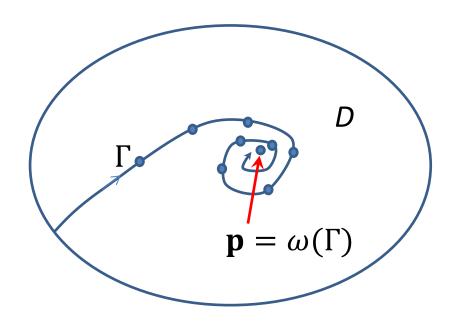
- Note that we may need to choose the times $\{t_i\}$ carefully in order to get a limit point (e.g. consider $\phi(t, \mathbf{x}) = \mathrm{e}^{-2t} + \cos(t)$)
- An α limit point is defined in the same way, but with $t_i \to -\infty$. This point is denoted $\alpha(\mathbf{x})$

Limit points

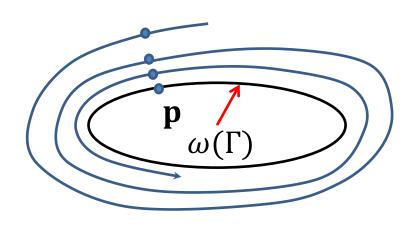
Definition: $\alpha(\Gamma)$ and $\omega(\Gamma)$ are called the α -limit set and ω -limit set respectively. These are the sets of all α limit points and ω limit points for the trajectory Γ .

- Hence $\alpha(\Gamma)$ is the set of points from which the trajectory Γ originates (at $t=-\infty$) and $\omega(\Gamma)$ is the set of points to which it tends (at $t=\infty$)
- The set of all limit points is called the **limit set of** Γ

Example limit sequences



A sequence of points leading to an isolated point ω -limit set



A carefully chosen sequence of points to an ω -limit point when the ω -limit set is not an isolated point

Equilibrium points

- An equilibrium point \mathbf{x}^* is its own α and ω limit point. Conversely, if a trajectory has a unique ω limit point \mathbf{x}^* , then \mathbf{x}^* must be an equilibrium point.
- Not all ω limit points are equilibrium points (e.g. see the previous slide). If a point \mathbf{p} is a limit point and $\dot{\mathbf{p}} \neq 0$ then this trajectory is a closed orbit. Note that we had to choose the sequence of points on the trajectory carefully in order to find a limit point and that there are infinitely many points on the ω limit set.

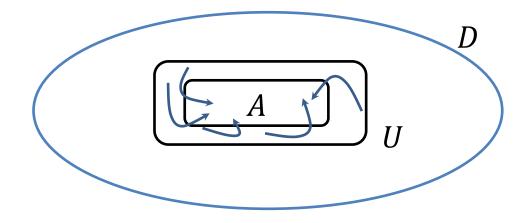
Invariance

Definition: If the differential equation is defined on an open set D with the flow $\phi(t, \mathbf{x}) \stackrel{\text{def}}{=} \phi_t(D)$, then the set $S \subset D$ is called **positively invariant** with respect to the flow if $\phi_t(S) \subset S$ for all $t \geq 0$.

- All points in *S* stay in *S* under the action of the flow the solution cannot 'escape' from *S*. We saw an example of this in the case of the stable and unstable invariant manifolds in earlier lectures.
- If a region M is positively invariant, closed and bounded, then the ω limit set is not empty (i.e. you have to go somewhere!). The limit set itself is positively invariant (since once there, you stay in the set).

Attraction

Definition: An invariant set $A \subset D$ is **attracting** if there is some neighbourhood U of A which is positively invariant and all trajectories starting in U tend to A as $t \to \infty$. U is called the **trapping region** of A.



Neighbourhood: A set surrounding a point \mathbf{x} so that the distance to all points in the neighbourhood from the point \mathbf{x} is less than some positive number $\boldsymbol{\epsilon}$.

Attractors

Definition: An **attractor** is an invariant attracting set that contains a dense orbit (such as a limit cycle or equilibrium point), i.e. there is no subset of the attractor that is itself an invariant attracting set.

Example: A stable node or focus is the ω limit set of all trajectories passing through points in a neighbourhood of the equilibrium point – the equilibrium point is an attractor. A saddle point on its own cannot be an attractor as trajectories leave the saddle point's neighbourhood.

Example

$$\dot{x} = -y + x(1 - x^2 - y^2)$$

$$\dot{y} = x + y(1 - x^2 - y^2)$$

In polar co-ords:

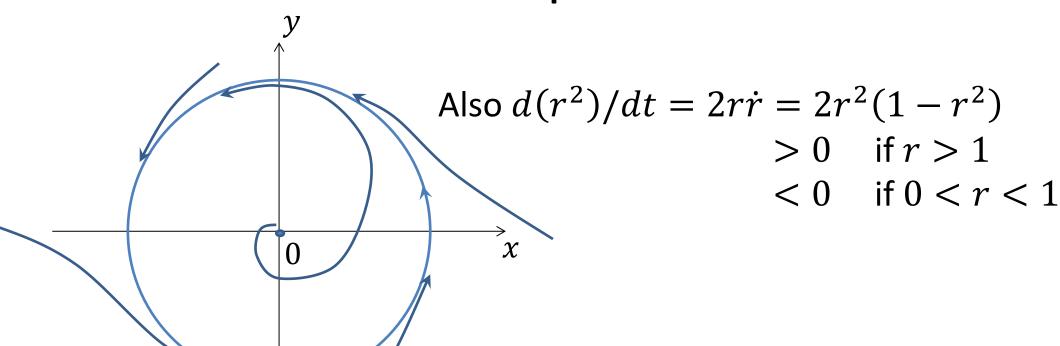
$$\dot{r} = r(1 - r^2)$$

$$\dot{\theta} = 1$$

The origin is an unstable hyperbolic equilibrium point.

Also r=1 is a centre ($\dot{r}=0$), called a limit cycle. If we perturb r about r=1 by δr we get $\dot{\delta r}\approx -2\delta r$, which is stable. Thus r=1 is a stable limit cycle.

Example

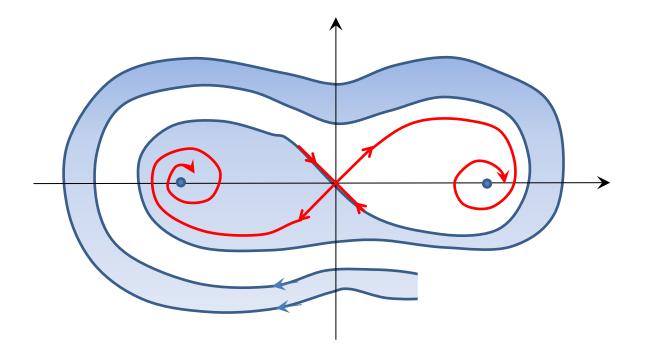


Hence r=1 is the ω limit set for all points in the plane except for the origin (which is its own ω limit set). The trapping region U is the whole of the plane minus the origin.

Basin of attraction

The domain or basin of attraction of an attracting set A is the union of all trajectories forming a trapping region of A.

The domain of attraction of the left hand equilibrium point of the Duffing oscillator



La Salle's Invariance Principle

- Let $\dot{\mathbf{x}} = \mathbf{f}(\mathbf{x}), \mathbf{x} \in \mathbb{R}^n$ and let $D \subset \mathbb{R}^n$ be a positively invariant set (all points starting in D remain in D). If the boundary of D is differentiable and D has a non-empty interior, then D is a trapping region.
- Suppose there exists a $V(\mathbf{x})$ that satisfies $V \leq 0$ on D and consider the following two sets:

$$E = \{ \mathbf{x} \in D : \dot{V}(\mathbf{x}) = 0 \}$$

 $M = \{ \text{union of all positively invariant sets in } E \}$

The Principle: Every trajectory starting at $\mathbf{x} \in D$ tends to M as $t \to \infty$.

Example: Duffing Oscillator

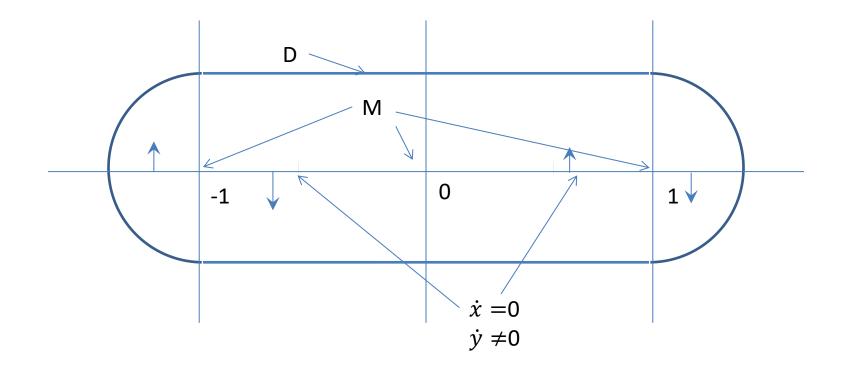
$$\dot{x} = y$$

$$\dot{y} = x - x^3 - \gamma y, \qquad \gamma > 0$$

Let
$$V(x,y) = \frac{y^2}{2} - \frac{x^2}{2} + \frac{x^4}{4}$$
, then $\dot{V} = -\gamma y^2$

- Here $E = \{(x, y) : y = 0\}$ and $M = \{(-1,0), (0,0), (1,0)\}$.
- Let $D = \{(x, y) : V(x, y) = c\}$ for large c, then D contains M and is positively invariant since $\dot{V} \leq 0$.
- LaSalle's principle implies that all trajectories starting in D converge to M, and hence to one of the three equilibria.

The Duffing Oscillator and LaSalle



Types of orbits – Preparation for The Poincaré Bendixson Theorem

The Poincaré Bendixson concerns the range of attractors that can exist in the **Phase Plane**. We will consider possible attractors, but first define some terms:

- A **homoclinic orbit** is a trajectory that joins a saddle point equilibrium point to itself (move out on an unstable manifold and come back on a stable manifold).
- A **heteroclinic orbit** (or **heteroclinic connection**) joins two different equilibrium points.
- A **separatrix cycle** partitions phase space into two regions with different characteristics and there are many ways to construct such cycles.

Homoclinic orbit example

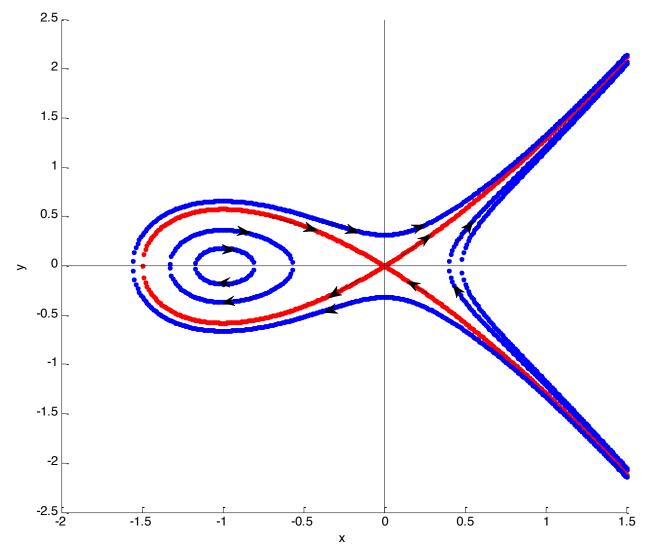
Consider the Hamiltonian system:

$$\dot{x} = y
\dot{y} = x + x^2
H(x, y) = \frac{y^2}{2} - \frac{x^2}{2} - \frac{x^3}{3}$$

- The solution curves are the level sets of the Hamiltonian (energy is constant) and are defined by $y^2-x^2-\frac{2}{3}x^3=c$
- If c=0, then $y^2=x^2+\frac{2}{3}x^3$, which goes through a saddle point at (x,y)=(0,0) (the Jacobian is $\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$ with eigenvalues 1, -1).

Homoclinic orbit example

One of trajectories (on the left) leaves the origin on the local unstable manifold and returns on the local stable manifold.



Heteroclinic example

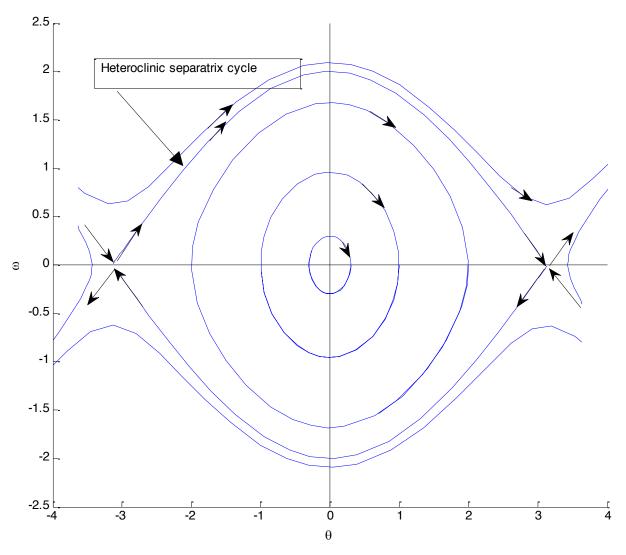
The undamped simple pendulum

$$\dot{\theta} = \omega$$

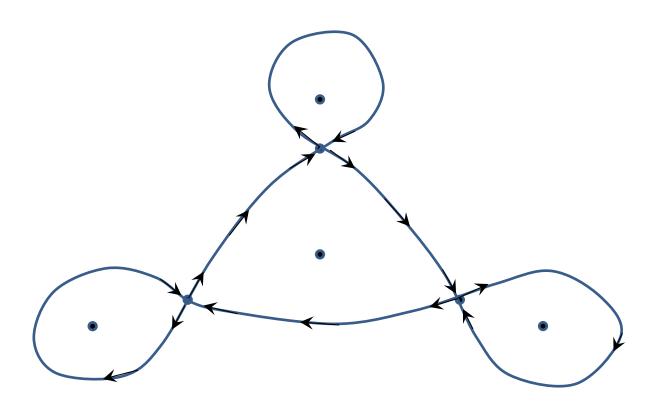
$$\dot{\omega} = -\sin\theta$$

- There are saddle points when the pendulum is pointing upwards these saddle points are connected by a heteroclinic orbit.
- The two heteroclinic orbits in the upper and lower half of the plane define a heteroclinic separatrix cycle.
- A number of 'compound' separatrix cycles are shown in the notes. Note: At any point where the trajectory appears to cross itself there must be an equilibrium point.

Simple pendulum phase portrait



Example compound separatrix cycle



Poincaré Bendixson Theorem in the plane

- Let M be a positively invariant region of a vector field in \mathbb{R}^2 containing only a finite number of equilibria. Let $\mathbf{x} \in M$ and consider $\omega(\mathbf{x})$. Then one of the following possibilities holds:
 - i. $\omega(\mathbf{x})$ is an equilibrium.
 - ii. $\omega(\mathbf{x})$ is a closed orbit.
 - iii. $\omega(\mathbf{x})$ consists of a finite number of equilibria $\mathbf{x}_1^*, \dots, \mathbf{x}_n^*$ and orbits γ with $\alpha(\gamma) = \mathbf{x}_i^*$ and $\omega(\gamma) = \mathbf{x}_j^*$. (Note: this defines a set of heteroclinic connections consider the pendulum phase portrait.)
- If there are only stable equilibria inside M, then there can only be one. If there are no equilibria, then there is a closed orbit inside it.

Example application of Poincaré Bendixson Thm.

Consider the Duffing oscillator:

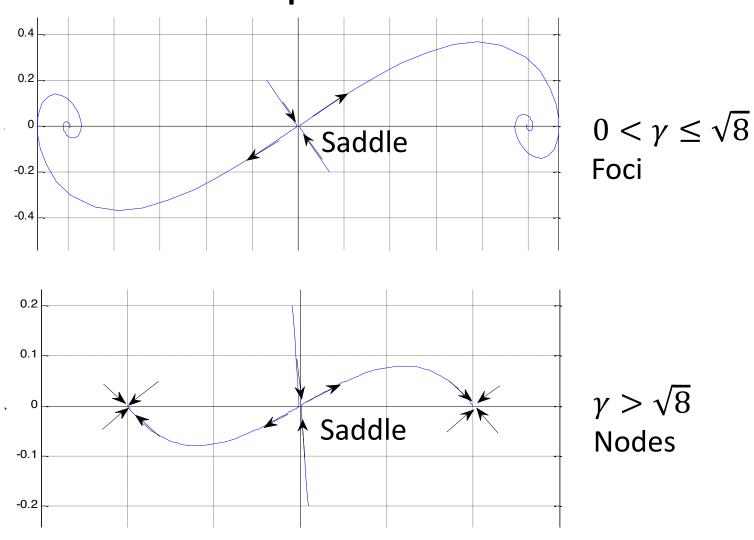
$$\dot{x}=y$$

$$\dot{y}=x-x^3-\gamma y, \qquad \gamma>0$$

The level sets of $V(x,y) = \frac{y^2}{2} - \frac{x^2}{2} + \frac{x^4}{4} = c$ are positively invariant since $\dot{V}(x,y) = -\gamma y^2 \le 0$.

- For c > 0, three equilibria lie inside $\{(x,y) : V(x,y) \le c\}$: an unstable equilibrium at (0,0), and two stable equilibria at (-1,0), (1,0).
- For c = 0, the level sets split into two with a common point at (0,0).
- Hence trajectories leaving the unstable equilibrium point at (0,0) must end up at the stable equilibria.

Duffing Oscillator example illustrating two types of stable equilibria



Further example

$$\dot{x}_1 = x_1 + x_2 - x_1(x_1^2 + x_2^2)$$

$$\dot{x}_2 = -x_1 + x_2 - x_2(x_1^2 + x_2^2)$$

The origin is an equilibrium point, linearisation gives Jacobian $\begin{bmatrix} 1 & 1 \\ -1 & 1 \end{bmatrix}$ with eigenvalues at $1 \pm i$ (an unstable spiral).

If we consider the region defined by

$$V = \frac{{x_1}^2}{2} + \frac{{x_1}^2}{2} = c$$

If c > 1, then $\dot{V} = c - c^2 < 0$ so $\{(x, y) : V \le c\}$ is an invariant region. As there is only an unstable equilibrium point inside the region, by Poincaré Bendixson the region must contain a stable limit cycle.

Unstable spiral within a stable limit cycle

