

Improved analysis of impurity transport coefficient profiles

M.A. Chilenski, M. Greenwald, Y. Marzouk,^{*} N.T. Howard,
J.E. Rice, A.E. White

MIT PSFC/Alcator C-Mod
^{*}MIT Aero/Astro, Uncertainty Quantification Group

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Careful analysis is needed to properly measure impurity transport coefficient profiles

- Validation of simulations requires rigorous inference of the experimental quantities used for comparison.
- Measuring impurity transport coefficients is a very challenging nonlinear inverse problem.
- In particular, it is not sufficient to compute merely one reasonable value for the transport coefficients – there can be multiple, dramatically different solutions that describe the data equally well.
- This lack of uniqueness in the solution *must* be taken into account when computing the uncertainty in the solution and comparing to simulations.

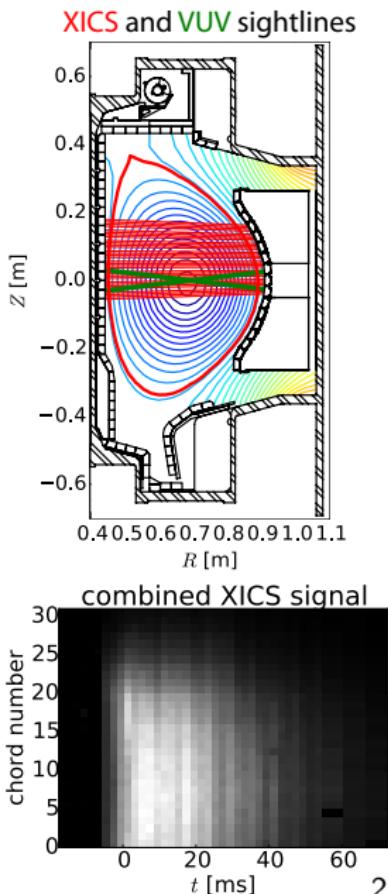
Alcator C-Mod is uniquely equipped to make detailed measurements of impurity transport

Multipulse laser blow-off impurity injector provides controlled impurity injections [1]

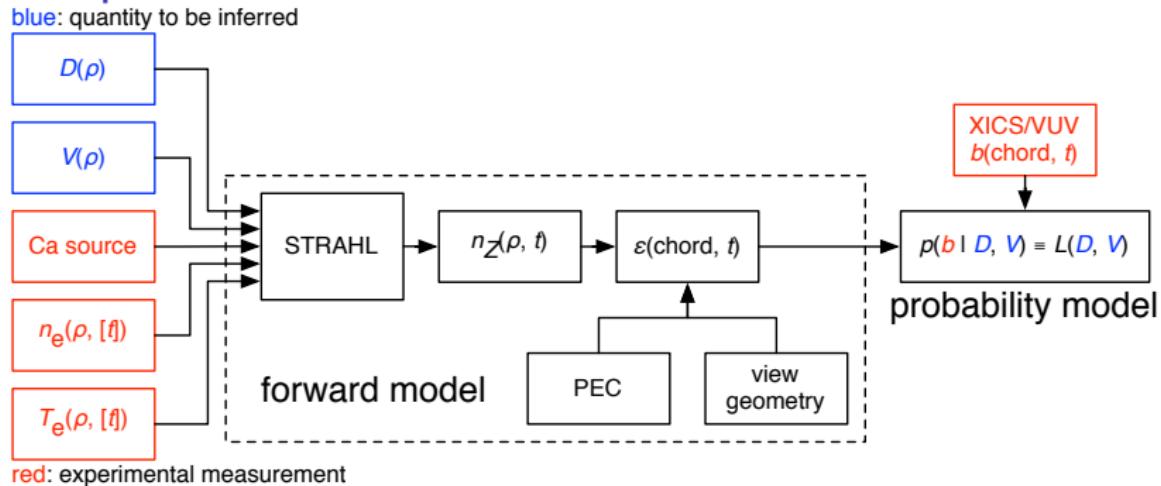
- Multiple injections per shot: up to 10 Hz
- Typically inject CaF_2 : calcium is *non-intrinsic* and *non-recycling*

X-ray imaging crystal spectrometer [2] and VUV spectrometers [3] track the impurities

- XICS observes spatial profile of a *single charge state* (Ca^{18+}): more direct interpretation than unresolved soft x-rays
- Two single-chord VUV spectrometers measure Ca^{16+} , Ca^{17+}



Inferring impurity transport coefficients is a nonlinear inverse problem



- This is an inverse problem: given the forward model $b = F(D, V)$, the objective is to find D, V profiles that best reproduce the observed brightness b on each of the diagnostics.
- Key issues are existence, uniqueness and stability of the solution.

The forward model is built around the STRAHL code

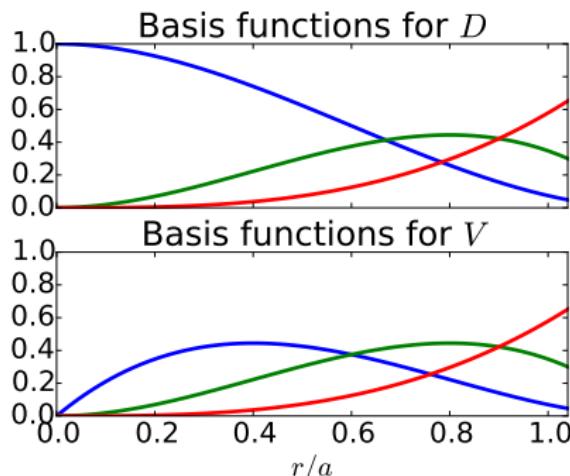
- Assume impurity flux of the form $\Gamma_Z = -D\nabla n_Z + Vn_z$.
- STRAHL [4] computes the temporal evolution of the impurity charge state densities $n_{Z,j}(\rho, t)$ for given profiles of D , V .
- Photon emission coefficients from ADAS [5] are used to convert $n_{Z,j}(\rho, t)$ to spectral line emissivity profiles.
- TRIPPy tomography code [6] is used to perform line integrations to obtain $b(\text{chord}, t)$.
- Uncertainty in data is taken to be Gaussian:

$$p(b_{\text{obs}} | D, V, n_e, T_e) =$$

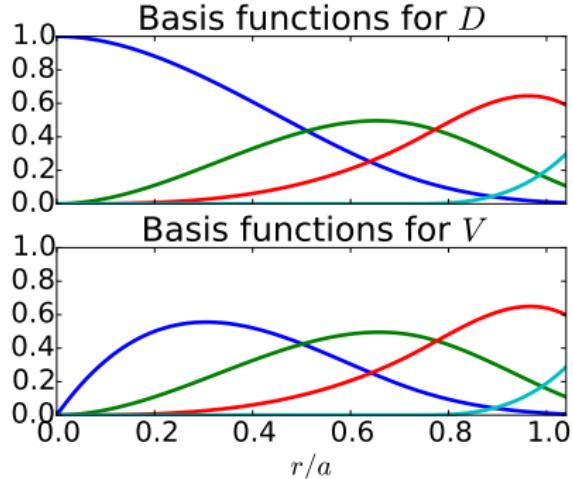
$$\prod_i \frac{1}{\sigma_i \sqrt{2\pi}} \exp \left\{ -\frac{[b_{\text{obs},i} - b_i(D, V, n_e, T_e)]^2}{2\sigma_i^2} \right\}$$

B-spline basis functions are used to obtain a smooth profile, impose constraints

Basis functions for 3 coefficient case,
no free knots



Basis functions for 4 coefficient case,
1 free knot



- $dD/dr = 0$ at $r/a = 0$
- $D \geq 0$ everywhere
- $V(0) = 0$

There are several challenges when solving this inverse problem

- The forward model is somewhat expensive to evaluate: $\sim 1\text{ s}$ per run on a typical workstation.
 - This reduces the practicality of sampling-based inference techniques, which may require $\sim 10^7$ samples to find the mode(s) of the posterior distribution and fully characterize the parameter space.
- The nonlinear relationship between the inputs D, V and the outputs $b(\text{chord}, t)$ introduces the possibility that there are multiple profiles of D, V which describe the data equally well.
 - In statistical terms, this means that the posterior distribution may be multimodal.
 - **Failure to account for multiple modes can lead to a dramatic underestimation of the uncertainty in D, V .**

Current approaches: maximum likelihood estimate (MLE)

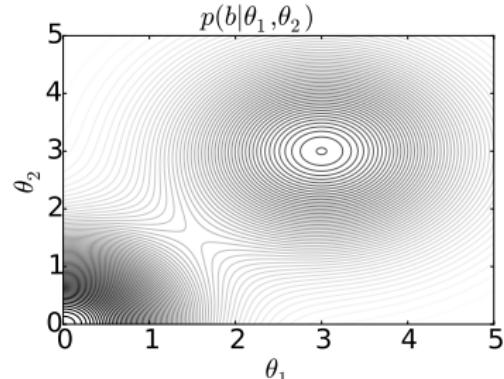
MLE is a standard approach to handle this problem...

$$\hat{D}, \hat{V} = \arg \max_{D, V} p(b|D, V)$$

- Pick D , V profiles which make the observations most likely.
- Use standard optimization techniques: assumption of Gaussian noise makes this a “least squares” problem.
- Need basis functions to represent the profiles with a finite number of variables: typically piecewise linear functions with fixed knots.

... but it has some potential shortcomings

- Point estimate:
 - Risk of underestimating uncertainty.
 - Not valid when there are multiple extrema.
- Propagation of uncertainty in n_e , T_e profiles requires an additional step.



Bayesian statistics provide a framework to overcome the shortcomings of MLE

Use Bayes' rule to obtain the posterior distribution $p(D, V|b)$, including constraints/prior knowledge $p(D, V)$:

$$\underbrace{f(D, V|b)}_{\text{posterior}} = \frac{\overbrace{f(b|D, V)}^{\text{likelihood}} \overbrace{f(D, V)}^{\text{prior}}}{\underbrace{f(b)}_{\text{evidence}}}$$

- **Likelihood**: Probability of observing the data b given D, V , assumed to be Gaussian.
- **Prior**: Distribution encoding any prior assumptions about D, V (positivity, typical values, etc.)
- **Evidence**: Probability of the data under the model. Just a normalization constant for parameter estimation.
- **Posterior**: Probability distribution for D, V given the data b : **contains all information which is known about D, V .**

Markov chain Monte Carlo (MCMC) sampling enables a complete accounting of uncertainty

- MCMC draws samples from unnormalized probability distribution such as $D^{(i)}, V^{(i)} \sim p(D, V|b) \propto p(b|D, V)p(D, V)$.
- Histogram to view $p(D, V|b)$ directly: nonuniqueness can be identified immediately.
- Allows for better point estimates, such as posterior mean and variance:

$$\mathbb{E}[D|b] = \int D p(D|b) dD \approx \frac{1}{N} \sum_{i=1}^N D^{(i)}$$

$$\text{var}[D|b] = \int (D - \mathbb{E}[D|b])^2 p(D|b) dD \approx \frac{1}{N-1} \sum_{i=1}^N (D^{(i)} - \mathbb{E}[D|b])^2$$

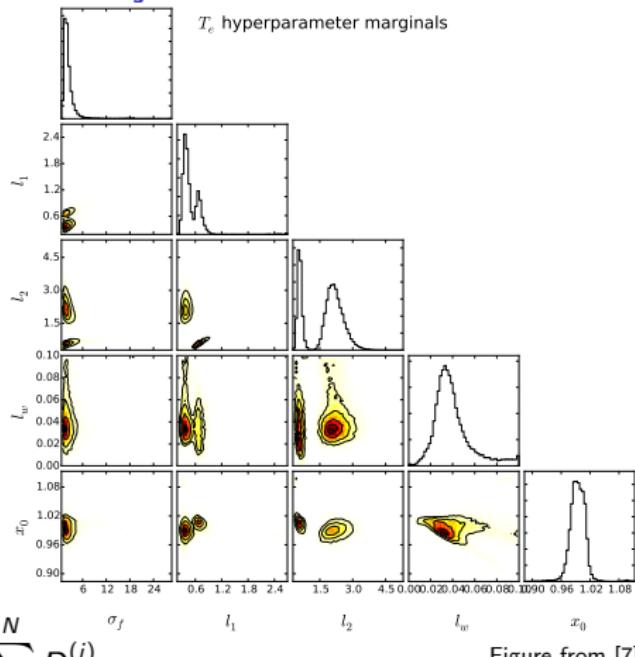
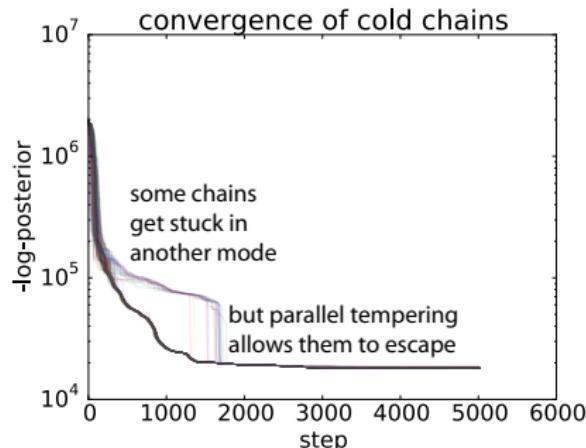


Figure from [7]

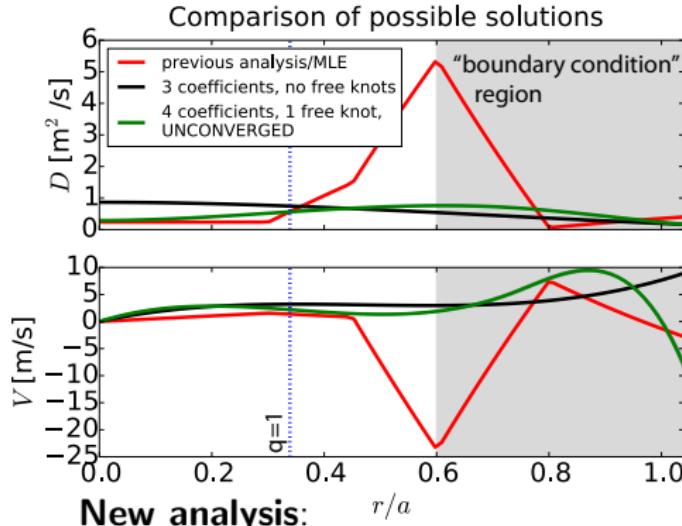
Multimodal posterior necessitates advanced MCMC

- Affine-invariant ensemble sampler (ES) [8, 9]
 - Eliminates need to tune proposal distribution.
 - But, cannot efficiently sample distributions with well-separated modes.
- Parallel tempering (PT) [10]
 - Sample from $p(b|D, V)^{1/T} p(D, V)$ for multiple values of $1 \leq T \leq \infty$.
 - Exchange of information between adjacent T lets chains move between modes.
- Adaptive parallel tempering (APT) [11]
 - Automatically tune T ladder.



- APT with ES in each temperature.
- 200 walkers per temperature, 25 temperatures.
- Plot shows $-\ln p(D, V|b)$ on a log scale: lower value = better fit.

Preliminary results do not match previous analysis



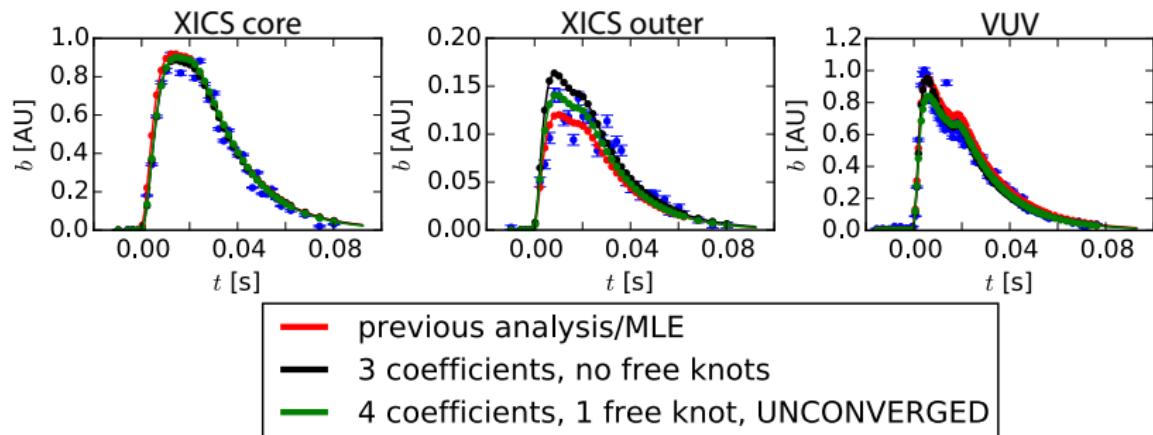
New analysis:

- Cubic B-spline basis functions.
- APT to handle multiple maxima, width of posterior distribution.
- Uncertainty estimate in $r/a > 0.6$ still too small to be consistent with assumed lack of knowledge there.
- *Cases shown are likely overconstrained.*
- Models with more free parameters failed to burn in, even after many thousands of CPU-hours.

Previous analysis [12]:

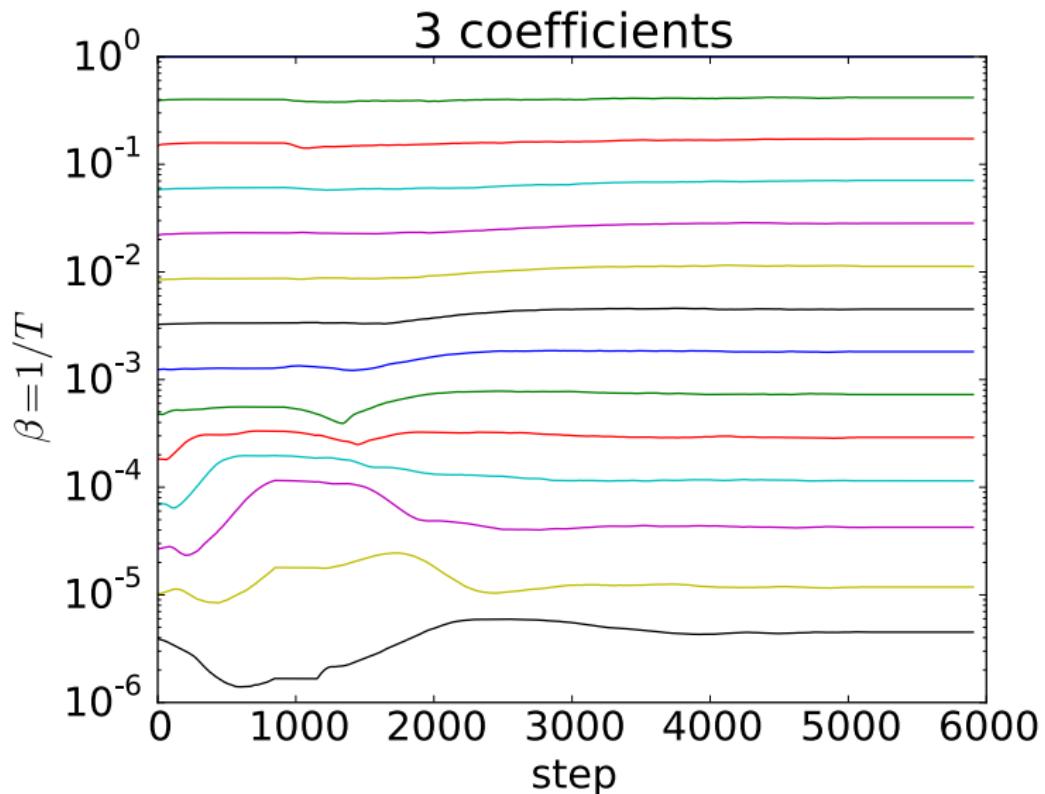
- Piecewise linear basis functions.
- MLE without estimate of width of posterior distribution.
- Behavior in $r/a > 0.6$ thought to be only weakly constrained.
- But, uncertainty there too small to be consistent with this.

Predicted brightnesses are similar between all three cases



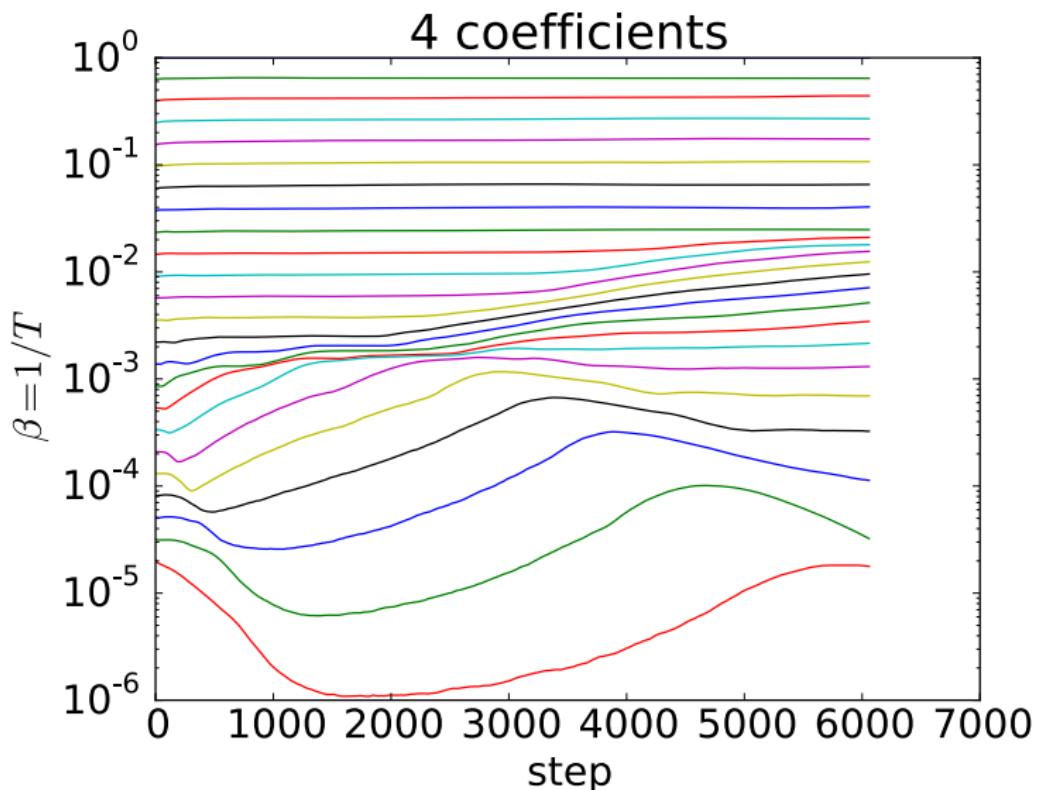
- Agreement on core XICS chords is good in all cases.
- Agreement on outer XICS chords shows widest variation – **4 coefficient** case seems to do best job.
- Agreement on VUV spectrometer is reasonable in all cases.
- *This shows the importance of accounting for the possibility of multiple solutions.*

Temperature ladder adaptation for 3 coefficient case



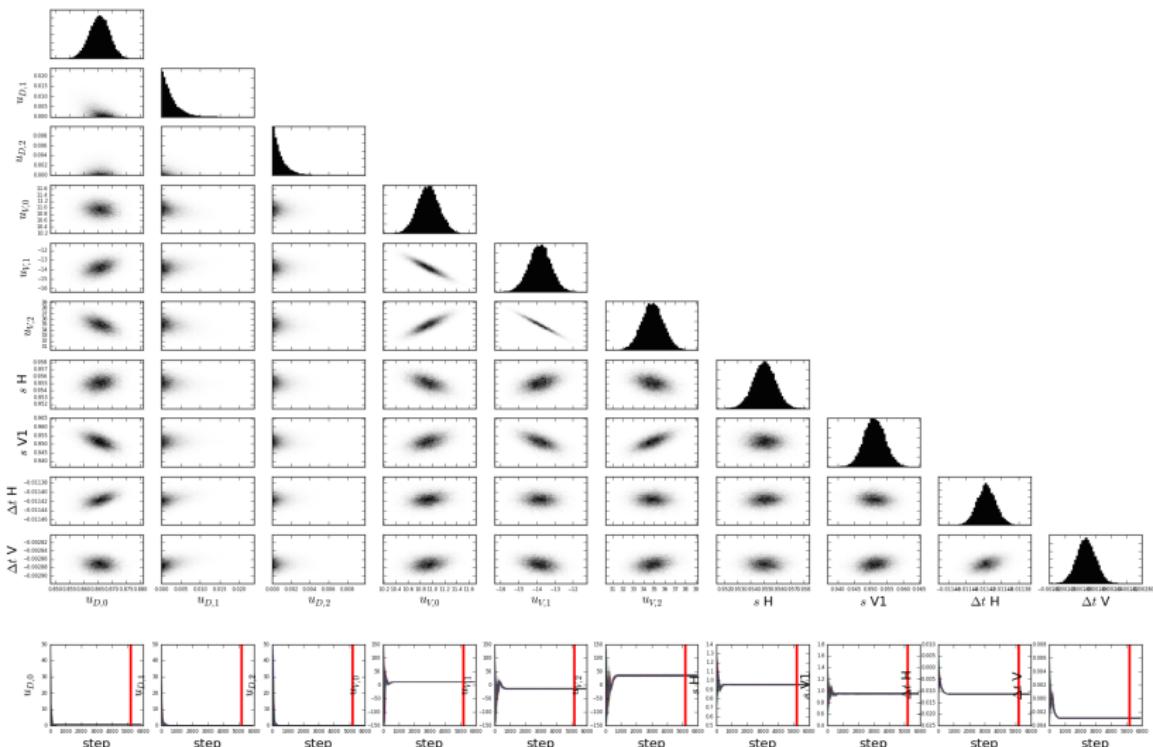
Appears to have settled down after about 5000 steps.

Temperature ladder adaptation for 4 coefficient case



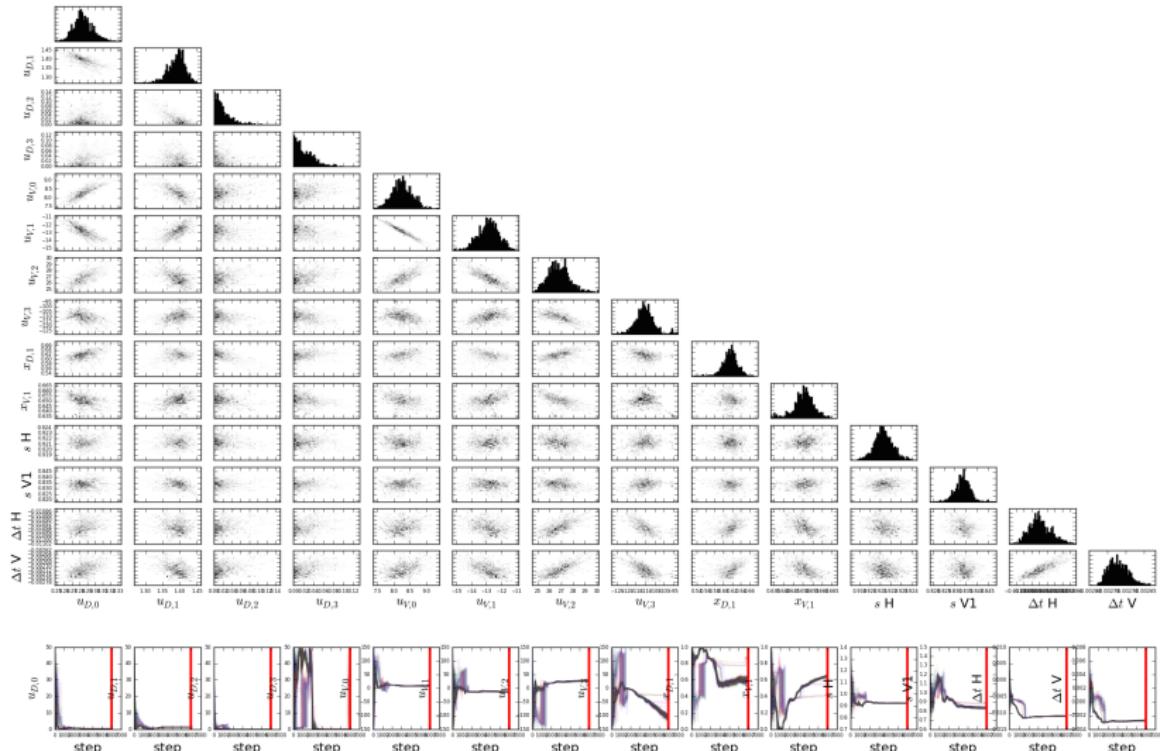
Has not settled down even after 6000 steps.

Posterior distribution for 3 coefficient case



Appears to be unimodal and is fully burned-in.

Posterior distribution for 4 coefficient case



Appears to be unimodal *but has not burned in even after 6000 steps.*

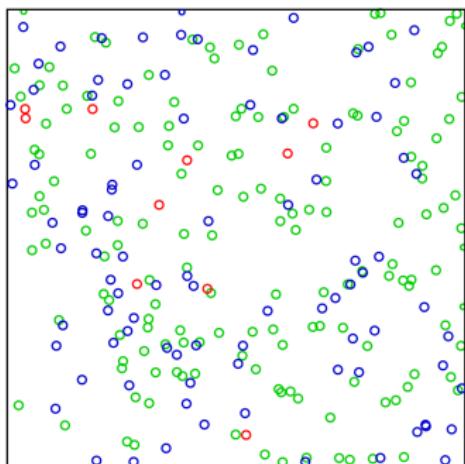
Convergence of MCMC/APT depends strongly on quality of starting conditions

- 4 coefficient case above used 3×10^7 calls to STRAHL – about 9000 CPU-hours.
- More flexible basis functions take even longer to burn in.
- Ideally, the sampler will be initialized with most walkers already near the posterior modes.
- This necessitates the use of global optimization techniques to find “all” of the posterior modes:
 - PaGMO/PyGMO [13] enables parallelization of genetic algorithms-based global optimizers through use of the Generalized Island Model [14].
 - Also includes tools for efficiently searching a parameter space for local extrema.

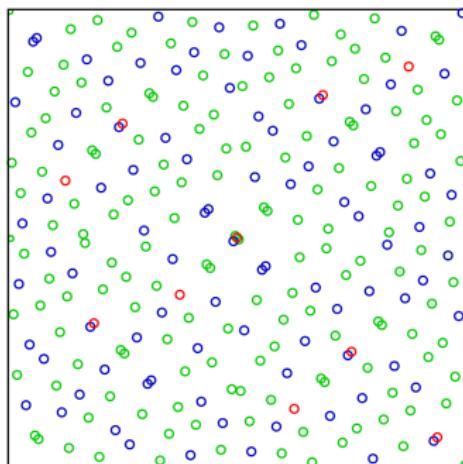
The Sobol sequence provides a systematic, efficient way of exploring the parameter space

Sobol sequence efficiently fills the parameter space, leaving fewer holes than pseudorandom sampling, better statistical properties than a uniform grid [15, 16].

Pseudorandom sequence:



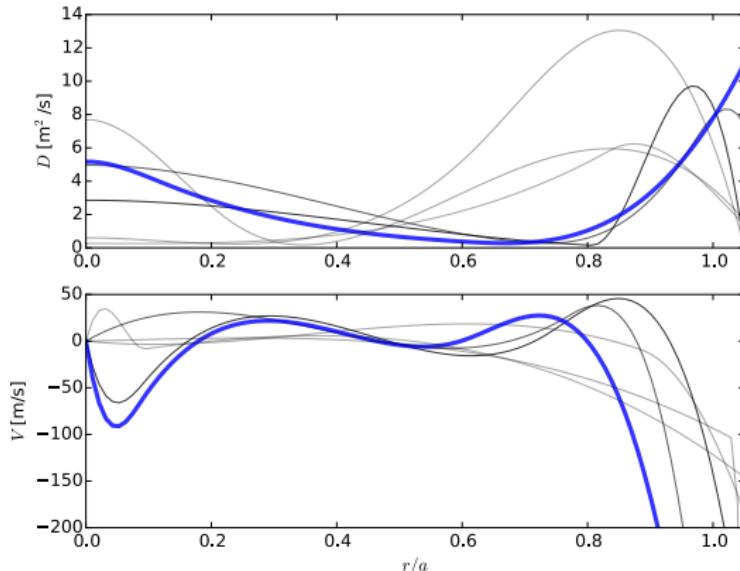
Sobol quasirandom sequence:



Images from [17, 18]

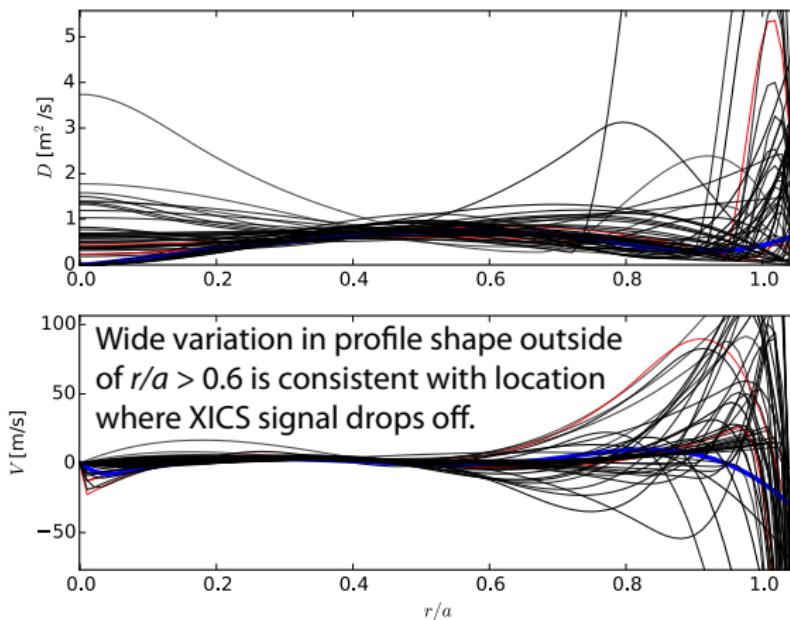
Systematic exploration of the parameter space using local optimizers is underway

- Sampled parameter space using 5×10^5 point Sobol sequence then started local optimizers at the 100 best points found.
- Discarded solutions which ended up too close to bounds.
- This left 6 possible solutions, best solution is shown in blue, remainder are shaded according to χ^2 .

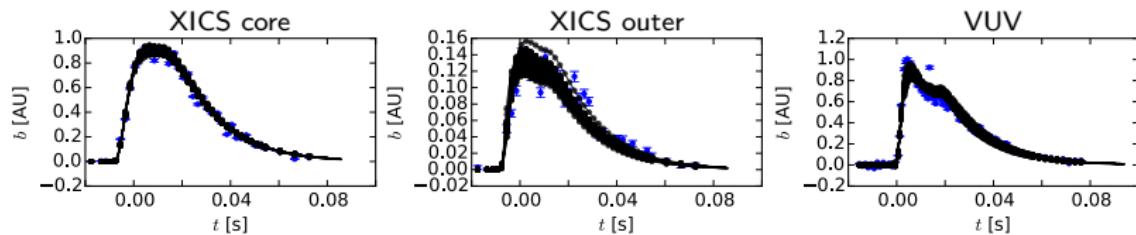


A more brute-force approach was also attempted

- Launched local optimizers at ~ 1000 points, again using a Sobol sequence to efficiently sample the space.
- Repeatedly restarted optimizers from the previous solutions, periodically pruning bad/stuck solutions.
- Ended up with 46 solutions.



Despite dramatic differences in profile shape, the solutions from the brute-force local extrema search all provide reasonable fits to the brightness data



- Fit to core XICS, VUV chords is good in all cases.
- Fit to outer XICS chords again shows widest variation.
- **This indicates that the solution is not unique. Failing to account for the multiple possible solutions leads to an underestimation of the uncertainty in D , V .**

Choice of local optimizer has a critical effect on the quality of the solution obtained

algorithm	result
Nelder-Mead	17967
Subplex	19390
compass search	46328
COBYLA	60674
BOBYQA	84250

- Result shown is the negative log-posterior ($\propto \chi^2$): lower is better.
- Each algorithm started with the same initial guess.
- Limited each algorithm to 1000 iterations.
- Repeated 10 times, since this (surprisingly) seemed to deliver better performance than just running for 10000 iterations.

Next step: include uncertainty in n_e , T_e profiles

Form joint posterior distribution, now also conditional on the profile measurements d :

$$p(D, V, n_e, T_e | b, d) = p(D, V | n_e, T_e, b, d) p(n_e, T_e | b, d)$$

Use Gaussian processes for n_e , T_e [7]:

$$p(n_e | d) = \mathcal{N}(m(\rho), k(\rho, \rho))$$

Reduce dimension of parameter space by approximating this with truncated eigendecomposition:

$$n_e = Q\Lambda^{1/2}u + m(\rho), \quad u \sim \mathcal{N}(0, I), \quad k(\rho, \rho) = Q\Lambda Q^{-1}$$

Find marginal posterior distribution for D , V using MCMC:

$$p(D, V | b, d) = \int p(D, V, n_e, T_e | b, d) dn_e dT_e$$

Conclusions

- Rigorous quantification of the uncertainties in impurity transport coefficients is essential for validation of multichannel transport simulations.
- The computational expense, nonlinearity of the forward model make the problem difficult to solve and susceptible to multiple extrema.
- Work is underway to combine advanced optimization and inference tools to overcome these issues.

Future Work

- Deployment of more advanced optimizers to more efficiently identify local extrema.
- MCMC sampling using results from local extrema search.
- Incorporation of SXR data.
- Accounting for uncertainties in n_e , T_e data.

References

- [1] N. T. Howard, M. Greenwald, and J. E. Rice, *Review of Scientific Instruments* **82**, 033512 (2011).
- [2] A. Ince-Cushman et al., *Review of Scientific Instruments* **79**, 10E302 (2008).
- [3] M. L. Reinke et al., *Review of Scientific Instruments* **81**, 10D736 (2010).
- [4] R. Dux, STRAHL user manual, Technical Report IPP 10/30, Max-Planck-Institut für Plasmaphysik, 2006.
- [5] H. P. Summers, The ADAS user manual, version 2.6, <http://www.adas.ac.uk>, 2004.
- [6] I. C. Faust, Toroidal Radiation Inversion Protocol (Python), <https://github.com/icfaust/TRIPPy>.
- [7] M. A. Chilenski et al., *Nuclear Fusion* **55**, 023012 (2015).
- [8] J. Goodman and J. Weare, *Communications in Applied Mathematics and Computational Science* **5**, 65 (2010).
- [9] D. Foreman-Mackey, D. W. Hogg, D. Lang, and J. Goodman, *Publications of the Astronomical Society of the Pacific* **125**, 306 (2013).
- [10] D. J. Earl and M. W. Deem, *Physical Chemistry Chemical Physics* **7**, 3910 (2005).
- [11] W. Vousden, W. M. Farr, and I. Mandel, Dynamic temperature selection for parallel-tempering in Markov chain Monte Carlo simulations, <http://arxiv.org/abs/1501.05823>, 2015.
- [12] N. T. Howard et al., *Nuclear Fusion* **52**, 063002 (2012).
- [13] D. Izzo, PyGMO and PyKEP: Open source tools for massively parallel optimization in astrodynamics (the case of interplanetary trajectory optimization), in *5th ICATT: International Conference on Astrodynamics Tools and Techniques, ESTEC/ESA, The Netherlands, 29 May – 1 June 2012*, 2012.
- [14] D. Izzo, M. Ruciński, and F. Biscani, *Parallel Architectures and Bioinspired Algorithms*, volume 415 of *Studies in Computational Intelligence*, chapter The Generalized Island Model, pages 151–169, Springer-Verlag, Berlin Heidelberg, 2012.
- [15] I. M. Sobol, *USSR Computational Mathematics and Mathematical Physics* **7**, 86 (1967).
- [16] C. Lemieux, *Monte Carlo and Quasi-Monte Carlo Sampling*, Springer Series in Statistics, Springer Science+Business Media, New York, NY, 2009.
- [17] J. Heald, Pseudorandom sequence 2d, https://en.wikipedia.org/wiki/File:Pseudorandom_sequence_2D.svg, 2011.
- [18] J. Heald, Sobol sequence 2d, https://en.wikipedia.org/wiki/File:Sobol_sequence_2D.svg, 2011.

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