It before, we oriented the discretized eg as an gorox for the continuum equations — now we will do the reverse: start with the discrete problem, I derive continuum problem as a limit or approximation

(1)

* Balls and springs Maises (m) sliding without / mm @ with on - frictun, displacement: u, with springs (k) (Un = 0 at equilibrium) Fn+1 chang in length net force on un: k(un+1-un) - k(un-un-1) (Hooke's) "Fn+1" $= \left\{ \left(u_{n+1} - 2u_n + u_{n-1} \right) \right\}$ (looks like 2 d2 without the ex? !

more systematically:

(i) get Fn+1/2 's from k × (differences in un's)

(ii) set net force from differences in Faty 5

1

(i) in matrix form:
$$\begin{cases}
F_{1} \\
F_{3/2}
\end{cases} = K$$

$$\begin{cases}
N+1 \\
\text{componenty}
\end{cases}$$

$$\Rightarrow \overrightarrow{F} = K D \overrightarrow{u} \cdot \Delta X$$

$$Same D as fir FD approx.$$

(ii) in motion form:

m
$$\vec{u}$$
 = net force = $\begin{pmatrix} -1 & 1 \\ -1 & 1 \end{pmatrix}$

N components

= - DX DT

Aron before.

$$\Rightarrow \overrightarrow{u} \overrightarrow{u} = -\frac{K}{m} \Delta x^2 \overrightarrow{D} \overrightarrow{D} \overrightarrow{u} = \overrightarrow{A} \overrightarrow{u}$$

$$A \qquad real-symmetric$$

negative definite

= $\frac{K}{m} \Delta x^2 \begin{pmatrix} -2 & 1 \\ 1 & -2 & 1 \\ \hline & & & \\ & & \\ & & & \\ &$

 $\vec{u}(0) = \underbrace{\leq}_{n} \alpha_{n} \otimes \vec{u}_{n} \Rightarrow \underbrace{\vec{u}_{n}}_{n} \vec{u}(0)$ $\vec{u}(0) = \underbrace{\leq}_{n} \omega_{n} \beta_{n} \vec{u}_{n} \Rightarrow \underbrace{\vec{u}_{n}}_{n} \vec{u}(0)$ $\underbrace{\vec{u}_{n}}_{n} = \underbrace{\vec{u}_{n}}_{n} \vec{u}(0)$

oscillating with frequencies wh

3

examples .

A
$$N=1$$
: $M=0$ $M=1$: $M=0$ $M=1$: $M=0$ $M=0$: M

(4

$$N = 2 : \frac{1}{\sqrt{N}} = \frac{1}{\sqrt{N$$

$$A = \frac{k}{m} \begin{pmatrix} -2 & 1 & 0 \\ 1 & -2 & 1 \\ 0 & \hline{1 & -2} \end{pmatrix}$$

$$\vec{u}_1 = \frac{1}{2} \begin{pmatrix} \vec{v}_2 \\ \vec{v}_1 \end{pmatrix} \longrightarrow \omega_1 \approx 0.765 \sqrt{\frac{E}{m}}$$

4

* The Continuum Limit N-> 00

· make ox smaller + smaller

density (per length)

· shortening springs increases k!

$$\left(\begin{array}{cccc} -W & W & = & -W & -W & = & -W & \\ k_1 & k_2 & \left(\frac{1}{k_1} + \frac{1}{k_2}\right)^{-1} & 2k & 2k & k \end{array}\right)$$

$$\implies \vec{i} = -\frac{k}{m} \Delta x^{2} \vec{D}^{T} \vec{D} \vec{x} = -\frac{c}{m} \vec{D}^{T} \vec{D} \vec{x}$$

$$\frac{\partial^2 u(x,+)}{\partial x^2} = + \frac{C}{\rho} \frac{\partial^2 u(x,+)}{\partial x^2} \qquad \frac{\partial^2 u(x,+)}{\partial x^2}$$

scalar wave equation!

$$\ddot{U} = A U$$

fixed ends:
$$u(0,+)$$

$$= u(L,+) = 0$$

$$\hat{A} = + \frac{c}{8} \frac{\delta^2}{8x^2}$$
 negative definite $\frac{d}{dx} = \frac{d}{8x^2}$ to self-adjoint

* Inhomogeneous materials:

- suppose each m, k is different:

→ (i) F = K Dù x

(N+1) × (N+1)

diagonal matrix

of K's

<a, \$> = \$\tilde{x}^* M\$

a negative-det

for m, 1<>0

i = - M Ax2 D K D i = Ai

(/m, /mz = NXN digjoral matrix of 1/m 's

A = - AX2 MIDTKD = A*

 $\hat{A} = \frac{1}{\rho(x)} \frac{1}{\delta x} c(x) \frac{1}{\delta x}$

= A* under <u, v>0 = Spuv

4 negative-Lefinite for p, c > 0

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