# An algorithm for determining the composition of a musical chord using Fourier transformations

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## 1 Rationale

As someone who is interested in signals, music and mathematics it was obvious for me that I wanted my mathematics internal assessment to be related to music. Although I have known of and used the Fourier transformations before, I have never used Fourier transformations to find individual notes in sounds. I decided, therefore, to try and identify chords using Fourier transformations, and try to find which are the composing notes of a chord.

## 2 Introduction

Even now, in the 21st century, an algorithm<sup>1</sup> for identifying tones in music does not exist<sup>2</sup>. Research has been done which does allow simple chord progressions to be identified, but complex progressions remain a mystery.

The likes of Hausner, Kurnia and, of course, Wang, have developed algorithms which present themselves as being promising. Although currently, simple, with hhe exception of Wang's, these algorithms are powerful enough to follow chord progressions. However, with advancements in the field, it is likely that these algorithms would improve to a point where they could follow highly complex, highly noisy melodies, for example, in music like jazz and rock.

The research goal of this paper is to create a simple algorithm with decent performance which allows for recognizing chords and the individual tones<sup>3</sup> that make them up. This simple algorithm could prove useful as a piece of a bigger algorithm for identifying chord progressions in songs.

## 2.1 Musical background

A chord is a musical unit in which three or more tones are played simultaneously (Benward, and Saker). The most simple chords are triads, as they are composed of only three tones. Although they are modeled by three composing sine waves (for example, the C major chord is  $f_{C_4maj}(t) = sin(41.63t) + sin(52.46t) + sin(62.39t)$ ), it is very important to note that actual musical instruments do not produce sine waves, as there

<sup>&</sup>lt;sup>1</sup>A set of operations executed in steps

<sup>&</sup>lt;sup>2</sup>Trained models are able to identify tones, however, there doesn't exist a single untrained algorithm which can achieve this.

<sup>&</sup>lt;sup>3</sup>Air vibrations at a constant pitch

are physical events that prevent this - reverb<sup>4</sup>, distortion<sup>5</sup>, harmonic series<sup>6</sup>, etc.

Another major issue stands - the timbre of a musical instrument. *Timbre* does not have a clear-cut definition as it is quite abstract, but can be thought of as the color of the sound a musical instrument produces, ie. it is the property (or, rather *properties*) that allow us to discern musical instruments from one another.

All of these effects hinder the possibility of easily and effectively fitting a function using the least squares method. To find the individual frequencies composing a chord, it is necessary to convert the chord signal from the time domain to the frequency domain. For this, the *Fourier transformation* is used.

### 2.2 Definitions

For further development, a clear set of definitions is necessary.

Let codomain Y be the set of all values which a function f for a domain of X outputs values.

$$f:X\to Y$$

Let a periodic function be a function such that the values of its codomain repeat at certain, constant intervals in its domain. Let waveform, in this paper, refer to a periodic function. Such a function has a period T and a frequency  $\nu$ , with the property of  $\nu = \frac{1}{T}$ .

Let an *even* function be such a function in the real domain for which the values of the function of the negative subset of the domain are equal to the values of the positive subset of the domain:

$$f(-x) = f(x)$$

Let an *odd* function be a function in the real domain such that:

$$f(-x) = -f(x)$$

Let an integral of a function f(t) over a period T and therefore the interval  $\left[-\frac{T}{2}, \frac{T}{2}\right]$  be denominated as  $\int_T f(x)dx$ :

$$\int_{-\frac{T}{2}}^{\frac{T}{2}} f(t)dt = \int_{T} f(t)dx$$

$$A(a) = \begin{cases} -4 & \text{when } a < -4 \\ a & \text{when } -4 < a < 4 \\ 4 & \text{when } a > 4 \end{cases}$$

<sup>&</sup>lt;sup>4</sup>Reverberation is the effect of reflecting sound waves. When sound waves reflect off of surfaces many of them build up and then decay as they become absorbed (Valente et al.). It is what allows us to discern from big halls to small studios when hearing audio.

<sup>&</sup>lt;sup>5</sup>Distortion is most easily modeled as clipping. If a signal function is defined as  $f(t) = A(a)sin(\omega t)$  then clipping will, for example, occur when:

A(a) is called a transfer function, where a is an arbitrary value representing the signal's amplitude before clipping has occurred. The transfer function A(a) represents the signal's amplitude after clipping. This is an example of hard clipping, and is the simplest example, but other, more complicated examples exist.

<sup>&</sup>lt;sup>6</sup>The harmonic series is the result of multiple parts of a vibrating body vibrating themselves (Benward, and Saker). A tone, therefore, is not composed of only one sine wave, but multiple sine waves of varying amplitudes. Usually, the longest harmonic is the strongest.

Let the symbol  $\propto$  denominate a linear function:

$$f(x) = kx + c \mid x, k, c \in \mathbb{R} \iff f \propto x$$

#### 2.3 Fourier transformation

The Fourier transformation is based on the Fourier series. As such, an understanding of the Fourier series is necessary.

Before further explanation, a set of premises is established.

### 2.3.1 Premises

Orthogonality of functions There exist functions such that:

$$\int_{B} f(x)g(x)dx = 0$$

Let these functions be denominated as orthogonal (Weisstein).

**Orthogonality of sines** Two sinusoidal functions are orthogonal to each other in their period if their frequencies are not equal:

$$\omega_1 \neq \omega_2 \Rightarrow \int_T \sin(\omega_1 t) \sin(\omega_2 t) dt = 0$$

**Orthogonality of cosines** Cosinusoidal functions also have the property of orthogonality for differing frequencies.

$$\omega_1 \neq \omega_2 \Rightarrow \int_T \cos(\omega_1 t) \cos(\omega_2 t) dt = 0$$

**Negative frequency** Negative frequency is a signed counterpart to standard frequency (Smith III). It is, indeed, contrary to intuition, possible for negative frequency to exist:

$$\exists \cos(-f \cdot t) \text{ where } f \in \mathbb{R}^+$$

$$\exists \sin(-f \cdot t) \text{ where } f \in \mathbb{R}^+$$

The following properties exist:

$$cos(-ft) = cos(ft)$$
$$sin(-ft) = sin(ft + \pi)$$

**Principal premise** The principal premise of the Fourier series is that any waveform can be expressed as a series of sine and cosine waves.

$$f(t) = \sum_{n=0}^{\infty} (a_n \sin(n\nu t) + b_n \cos(nft)) : \nu \in \mathbb{R}$$

The terms of the series are referred to, in this paper, as *constituting* sine and cosine waves.

#### 2.3.2 Remarks

This paper will employ radial frequency in its notation, meaning that all of the frequency values will be multiplied by  $2\pi$ :

$$\sin(2\pi ft)$$

This is done because of the higher clarity in the field of application. It allows us to use the frequency in Hz. However, in no shape or form does the employment of this influence the final results, and where it does, it is noted so.

#### 2.3.3 Fourier series

Let us first consider an even function.

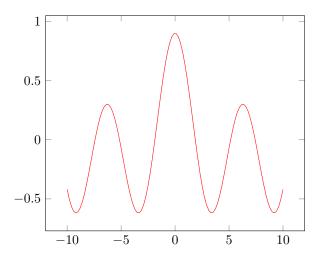


Figure 1: An arbitrary even function

From the principal premise, it is inferred that it is possible to express an even periodic signal such as the one given in figure 1 as a sum of cosine waves:

$$f(t) = \sum_{n=0}^{\infty} a_n \cos(2\pi\nu nt)$$
  
=  $a_0 + a_1 \cos(2\pi\nu t) + a_2 \cos(4\pi\nu t) + a_3(6\pi\nu t) + \dots$ 

Now the question of what the individual values of  $a_n$  are is posed.

A formula for calculating the values of  $a_n$  is presented, as well as a derivation.

First, let us multiply the left-hand and the right-hand side by  $\cos(2\pi\nu mt)$ , where  $m\in\mathbb{R}^+$ .

$$f(t)\cos(2\pi\nu mt) = \sum_{n=0}^{\infty} a_n \cos(2\pi\nu nt)\cos(2\pi\nu mt)$$

Then, integration is done on both sides of the equation.

$$\int_{T} f(t) \cos(2\pi\nu mt) dt = \int_{T} \sum_{n=0}^{\infty} a_n \cos(2\pi\nu nt) \cos(2\pi\nu mt) dt$$

$$= \int_{T} a_0 \cos(2\pi\nu mt) + a_1 \cos(2\pi\nu t) \cos(2\pi\nu mt) + a_2 \cos(4\pi\nu t) \cos(2\pi\nu mt) + \dots dt$$

$$= \int_{T} a_0 \cos(2\pi\nu mt) dt + \int_{T} a_1 \cos(2\pi\nu t) \cos(2\pi\nu mt) dt + \dots$$

Since the values of  $a_n$  are constant real numbers, it follows that the integral of the sum can be expressed as follows:

$$\int\limits_T f(t)\cos(2\pi\nu mt)dt = \sum_{n=0}^\infty \left(a_n\int\limits_T \cos(2\pi\nu nt)\cos(2\pi\nu mt)dt\right)$$

Let us now employ the trigonometric product identity:

$$\int_{T} \cos(2\pi\nu nt)\cos(2\pi\nu mt)dt = \frac{1}{2}\sum_{n=0}^{\infty} \left( a_n \int_{T} \cos((m+n)2\pi\nu t) + \cos((m-n)2\pi\nu t) dt \right)$$
(1)

Since  $m \ge 0$ , and since m is an integer (making m + n an integer), it follows that  $\cos((m + n)2\pi\nu t)$  has m + n oscillations in a period. Since this is true, when this function is integrated, due to the orthogonality of cosines:

$$\int_{T} \cos((m+n) 2\pi \nu t) dt = 0$$

Equation 1 is then simplified to:

$$\int_{T} f(t)\cos(2\pi\nu mt)dt = \frac{1}{2}\sum_{n=0}^{\infty} a_n \int_{T} \cos((m-n)2\pi\nu t)dt$$
(2)

Since the cosine function is even, in one period (ie. the interval of the integration - from  $-\frac{T}{2}$  to  $\frac{T}{2}$ ) the following is true:

$$\cos((m-n)2\pi\nu t) = \cos((n-m)2\pi\nu t)$$

When m = n:

$$\cos((m-n)2\pi\nu t) = \cos(0) = 1$$

To consider the case of m-n, we must acknowledge the fact that the cosine function will complete |m-n| oscillations in a single period. Since the interval of integration is one period, it follows that:

$$\int_{T} \cos((m-n)2\pi\nu t)dt = \int_{T} 1 \cdot dt$$

Summing up:

$$\int_{T} \cos((m-n)2\pi\nu t) = \begin{cases} \int_{T} \cos((m-n)2\pi\nu t)dt = 0 & \text{when } m \neq n \\ \int_{T} 1 \cdot dt = T & \text{when } m = n \end{cases}$$

Let us observe the summation in equation 2. In the summation, it can be observed that every term, except for the one where m = n is equal to zero, and does not contribute to the series. Then, the summation can be reduced exclusively to the case when m = n.

$$\int_{T} f(t)\cos(m2\pi\nu t)dt = \frac{1}{2}a_{m}T$$

Then  $a_m$  is:

$$a_m = \frac{2}{T} \int_T f(t) \cos(m2\pi\nu t) dt$$

We can replace m with n, since m did not contribute to the original equation:

$$a_n = \frac{2}{T} \int_T f(t) \cos(m2\pi\nu t) dt \tag{3}$$

Substituting with m=0 for  $a_0$  in equation 1, we get:

$$\int_{T} f(t)\cos(2\pi\nu mt)dt = \sum_{n=0}^{\infty} a_n \int_{T} \cos(2\pi\nu nt)dt$$

$$\int_{T} f(t)dt = Ta_0$$

This gives the final result for the y axis offset:

$$a_0 = \frac{1}{T} \int_T f(t)dt \tag{4}$$

For even functions, we have derived a formula for the amplitudes of the terms in the series.

Odd functions are represented with sine waves:

$$f_o(t) = \sum_{n=1}^{\infty} b_n \sin(2\pi\nu nt)$$

It can similarly be shown that the amplitude of each constituting sine wave  $b_n$  is given by:

$$b_n = \frac{2}{T} \int_T f_o(t) \sin(2\pi\nu nt) dt$$

The derivation is similar to that of even functions and is therefore omitted.

Finally, the individual even and odd constituents' formulae can be combined to satisfy the requirement of the original premise that any waveform can be represented as a sum of sine and cosine waves:

$$f(t) = a_0 + \sum_{n=1}^{\infty} (a_n \cos(2\pi\nu nt) + b_n \sin(2\pi\nu nt))$$

where

$$a_0 = \frac{1}{T} \int_T f(t)dt$$

$$a_n = \frac{2}{T} \int_T f(t) \cos(2\pi\nu nt)dt, n \neq 0$$

$$b_n = \frac{2}{T} \int_T f(t) \sin(2\pi\nu nt)dt$$

Euler's formulae allow for expressing the Fourier series in complex form:

$$\cos\theta = \frac{e^{i\theta} + e^{-i\theta}}{2}$$

$$sin\theta = \frac{e^{i\theta} - e^{-i\theta}}{2i}$$

$$f(t) = a_0 + \sum_{n=1}^{\infty} (a_n \cos(2\pi\nu nt) + b_n \sin(2\pi n\nu t))$$

$$= a_0 + \sum_{n=1}^{\infty} \left( a_n \frac{e^{i2\pi\nu nt} + e^{-i2\pi\nu nt}}{2} + b_n \frac{e^{i2\pi\nu nt} - e^{-i2\pi\nu nt}}{2i} \right)$$

$$= a_0 + \sum_{n=1}^{\infty} \frac{a_n - ib_n}{2} e^{2\pi\nu nt} + \sum_{n=1}^{\infty} \frac{a_n + ib_n}{2} e^{-2\pi\nu nt}$$

To solve the issue of having different complex exponentials let us set n to be n = -n in the first summation:

$$f(t) = a_0 + \sum_{n=-1}^{-\infty} \frac{a_{-n} - ib_{-n}}{2} e^{-2\pi\nu nt} + \sum_{n=1}^{\infty} \frac{a_n + ib_n}{2} e^{-2\pi\nu nt}$$

$$(5)$$

Let us also define the following:

$$c_0 = a_0$$

$$c_{n+} = \frac{a_n + ib_n}{2}$$

$$c_{n-} = \frac{a_n - ib_n}{2}$$

Then, with these defined, we can rewrite equation 5 as<sup>7</sup>:

$$f(t) = \sum_{n = -\infty}^{\infty} c_n e^{-i2\pi\nu nt} \tag{6}$$

 $c_n$  is expressed in the following<sup>8</sup>:

$$c_{n} = \frac{1}{2} \cdot \left(\frac{2}{T} \int_{T} f(t) \cos(2\pi\nu n t) dt - i\frac{2}{T} \int_{T} f(t) \sin(2\pi\nu n t) dt\right)$$

$$= \frac{1}{T} \int_{T} f(t) \cos(2\pi n \nu t) - if(t) \sin(2\pi\nu n t) dt$$

$$= \frac{1}{T} \int_{T} f(t) \left(\frac{e^{i2\pi\nu n t} + e^{-i2\pi\nu n t}}{2} - i\frac{e^{i2\pi\nu n t} - e^{-i2\pi\nu n t}}{2i}\right)$$

$$= \frac{1}{T} \int_{T} f(t) \frac{2e^{i2\pi\nu n t}}{2} dt$$

$$c_{n} = \frac{1}{T} \int_{T} f(t) e^{i2\pi\nu n t} dt$$

$$(7)$$

## 2.4 The Fourier transformation

The Fourier transformation, unlike the Fourier series which operates on periodic functions, operates on non-periodic functions. This is achieved based on the fact that a non-periodic function is expressible as a periodic function whose interval is  $\infty$ .

Its derivation is given in this subsection. Let us start from the Fourier series.

Next, let  $\omega$  be:

$$\omega = 2\pi\nu n$$

<sup>&</sup>lt;sup>7</sup>This is often referred to as the *synthesis* equation.

<sup>&</sup>lt;sup>8</sup>This equation is referred to as the *analysis* equation.

This expression allows us to express equation 7 as a function of  $\omega$ :

$$\frac{1}{T} \int_{T} f(t)e^{i2\pi\nu nt}dt = \frac{1}{T} \int_{T} f(t)e^{i\omega t}dt$$

$$g(\omega) = \frac{1}{T} \int_{T} f(t)e^{i\omega t}dt$$
(8)

Rewriting equation 6 we get:

$$f(t) = \sum_{n = -\infty}^{\infty} g(\omega)e^{-i\omega t}$$

Note that this equation connects the time and frequency domains.

Substituting:

$$f(t) = \sum_{n = -\infty}^{\infty} \left( \frac{1}{T} \left( \int_{-\frac{T}{2}}^{\frac{T}{2}} f(t) e^{i\omega t} dt \right) e^{-i\omega t} \right)$$
 (9)

Let us now find the difference between the two terms  $\Delta\omega$ :

$$\Delta\omega = \omega_{n+1} - \omega n - 1$$

$$= ((n+1)2\pi\nu) - (n2\pi\nu)$$

$$= \pi\nu$$

$$= \frac{2\pi}{T}$$

Let us solve equation 9 for  $T = \frac{2\pi}{\Delta\omega}$ .

$$f(t) = \sum_{n = -\infty}^{\infty} \left( \frac{\Delta \omega}{\pi} \left( \int_{-\frac{T}{2}}^{\frac{T}{2}} f(t) e^{i\omega t} dt \right) e^{-i\omega t} \right)$$

Now, let T approach  $\infty$  as per the definition of the Fourier transformation:

$$T \to \infty$$

Conversely:

$$\lim_{T \to \infty} \Delta \omega = 0$$

Since  $\Delta\omega \to 0$ , the function becomes continuous. Because of this, we can express the sum as a Riemann sum. Note that  $\Delta\omega$  becomes  $d\omega$ :

$$f(t) = \sum_{n = -\infty}^{\infty} \left( \frac{\Delta \omega}{\pi} \left( \int_{-\infty}^{\infty} f(t)e^{i\omega t} dt \right) e^{-i\omega t} \right)$$

$$= \int_{-\infty}^{\infty} \left( \int_{-\infty}^{\infty} f(t)e^{i\omega t} dt \right) e^{-i\omega t} \frac{d\omega}{\pi}$$

$$= \frac{1}{\pi} \int_{-\infty}^{\infty} \left( \int_{-\infty}^{\infty} f(t)e^{i\omega t} dt \right) e^{-i\omega t} d\omega$$

$$(10)$$

Trivially, it can be shown that equation 10 is equal to:

$$f(t) = \frac{1}{\pi} \int_{-\infty}^{\infty} g(\omega)e^{-i\omega t}dt$$

The operation  $\hat{f}(t)$  is defined to be  $g(\omega)$ :

$$\hat{f}(t) = g(\omega) = \int_{-\infty}^{\infty} f(t)e^{i\omega t}dt$$

This is the Fourier transformation.

This equation gives a continuous relationship between frequency and amplitude, which allows for identifying individual constituting waveforms inside complex waveforms. Its applications are almost limitless - from analogue signal, to digital signal processing, image processing, seismic analysis and even medicine.

# 3 The chord identification algorithm

## 3.1 Definitions

A few more definitions are necessary for a clearer understanding of the algorithm.

A *sample*, in this context, is a discrete measurement of a physical waveform and its storage in a digital format. Usually, computers sample a signal 44000 times every second.

The sampling frequency  $F_s$  of a digital audio recording is the number of samples that is created by a digital recording device in a second.

Let us define a convenience operation #(v) such that it returns the number of dimensions of the vector space containing the vector v. To exemplify:

$$\# \left( \begin{bmatrix} a \\ b \\ c \end{bmatrix} \right) = 3$$

$$\# \left( \begin{bmatrix} a \\ b \\ c \\ d \end{bmatrix} \right) = 4$$

where a, b, c and d are real numbers.

Let us define our final convenience operation  $v = \operatorname{append}(v, b)$  which appends a value b to a vector v. For example:

$$v = \begin{bmatrix} 2 \\ 8 \\ 3 \end{bmatrix}$$

append 
$$(v, 4) = \begin{bmatrix} 2\\8\\3\\4 \end{bmatrix}$$

Let a *symmetric vector* be such a vector:

$$a = \begin{bmatrix} a_0 \\ a_1 \\ \vdots \\ a_{\left(\frac{n}{2}-1\right)} \\ a_{\frac{n}{2}} \\ a_{\left(\frac{n}{2}+1\right)} \\ \vdots \\ a_{n-1} \\ a_n \end{bmatrix}$$

with the following set of properties:

$$\left\{a_{\left(\frac{n}{2}-1\right)} = a_{\left(\frac{n}{2}+1\right)} \dots a_1 = a_{n-1}, a_0 = a_n, n = 2l+1 : l \in \mathbb{Z}^+\right\}$$

Let us define another convenience function  $\max(v)$ , which for a vector v returns the numerically largest dimension. For example:

$$\max\left(\begin{bmatrix} 5\\9\\2 \end{bmatrix}\right) = 9$$

A Discrete Fourier transformation can be thought of as a discrete counterpart to the continuous Fourier transformation previously discussed. Unlike the Fourier series, it can operate on non-periodic signals, and for a sufficiently small value of  $\Delta\omega$  the error produced is negligible. Let it be denoted as:

a = dft(b, c) where b is the domain of the waveform and c is the codomain

Such a function returns values of the amplitudes of the constituting waveforms and stores them in a vector here denoted as a.

## 3.2 The algorithm

Given an input signal such as the one given in figure 2, the algorithm operates in 4 distinct stages.

The computer is initially instructed to count the number of samples n in the digital recording.

The computer is instructed to load the samples of the signal in a column vector A.

$$A = \begin{bmatrix} a_1 \\ a_2 \\ a_3 \\ \vdots \\ a_n \end{bmatrix}$$
 where  $a$  is the sample value and  $n$  is the number of samples

Note that these are the values represented on the y axis in figure 2.

To get the values of the x axis, ie. the time domain, it is necessary to create a column vector as such.

$$t = \frac{1}{F_s} \cdot \begin{bmatrix} 0 \\ 1 \\ 2 \\ \vdots \\ n-1 \end{bmatrix}$$

These two vectors possess the property #(A) = #(t).

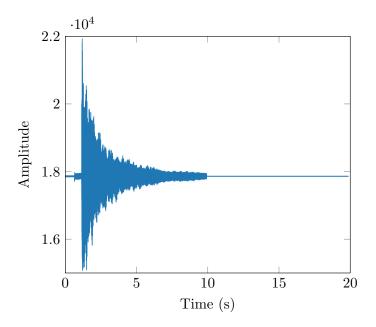


Figure 2: An example of an input signal - in this case a  $C_{maj}$  chord

With these values, it is possible to proceed in the algorithm.

## 3.3 The Fourier transformation

A discrete Fourier transformation is done on the vectors

$$F = dft(t, A)$$

The Fourier transformation, for real inputs, gives a symmetric vector F. Due to the unnecessary values in F it is cut off at the middle part onwards.

$$F' = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \\ \vdots \\ F_{\frac{\#(F)-1}{2}} \end{bmatrix}$$

Since the numbers given by the Fourier transformation are complex, to get the amplitudes of the frequencies it is necessary to get absolute values for all of the values in the vector given by the Fourier transformation.

$$F'' = \begin{bmatrix} |F_1'| \\ |F_2'| \\ |F_3'| \\ \vdots \\ |F_{\#(F')}'| \end{bmatrix} = \begin{bmatrix} \sqrt{\text{Re}(F_1')^2 + \text{Im}(F_1')^2} \\ \sqrt{\text{Re}(F_2')^2 + \text{Im}(F_2')^2} \\ \sqrt{\text{Re}(F_3')^2 + \text{Im}(F_3')^2} \\ \vdots \\ \sqrt{\text{Re}\left(F_{\#(F')}'\right)^2 + \text{Im}\left(F_{\#(F')}'\right)^2} \end{bmatrix}$$

A vector f is created to store the domain of the Fourier transformation, such that it is a linear space from 0 to 1 with the norm being equal to the one returned Fourier transformation, multiplied by the sampling

frequency of the original file. This is in turn, creates the frequency domain.

$$f = F_s \cdot \begin{bmatrix} 0 \\ 1 \cdot \frac{1}{\#(F'')} \\ 2 \cdot \frac{1}{\#(F'')} \\ \vdots \\ (\#(F'') - 1) \cdot \frac{1}{\#(F'')} \\ \#(F'') \cdot \frac{1}{\#(F'')} \end{bmatrix}$$

## 3.4 Fourier transform peak identification

In this section of the algorithm the computer tries to identify which peaks play a significant role in composing a chord.

Due to the possibility of different waveforms having different amplitudes, and therefore the possibility of their respective Fourier transformations having different amplitudes, it is impossible to compare them before normalizing them. To normalize them (ie. make their maximal values the same) the existing vector is multiplied by the reciprocal of the largest dimension in it:

$$F''' = \frac{1}{\max(F'')} \cdot F''$$

It can be assumed that there are insignificant frequencies that do not play a role in a chord, but may be present (noise, instrument insignificant noises, like the moving of the piston on a trumpet, or the actual picking of the string). So, to improve the algorithm, it was decided to take values which are numerically larger than a certain constant value, named cutoff value and denoted as  $c_o$ . This was achieved using the following logical condition.

$$p = \emptyset$$

$$\forall v \in F''' : \left(v > c_o \Rightarrow p = \text{append}(p, v)\right)$$

As an example, the normalized Fourier transformation, alongside the cutoff value of  $c_o = 0.15$ , is plotted in figure 3.

Due to the fact that the Fourier transform is continuous, the apparent peaks of the Fourier transform are not discrete values, but are rather reminiscent of the peaks shown in figure 4.

To resolve this issue, the average of neighboring values of a peak is calculated. Neighboring values w are defined to be neighboring to a value v if the absolute difference between them is less than a leeway value l.

$$\forall v \in p : \left( \left( \left( w \notin p' : |v - w| < l \right) \to p' = \text{append}\left( p', w \right) \right) \land \left( \left( w \in p' : |v - w| < l \right) \to w = \frac{\left( w + v \right)}{2} \right) \right)$$

At this point, the vector p' contains the frequencies of the played chord that are more intense than a given cutoff value.

## 3.5 Note identification

Notes are identified in this step by finding the closest value in a look-up table of frequencies<sup>9</sup>. In this subsection the methodology for finding the note is presented. In this section, let us denote any value in p' as v. Let us also denote the lookup table as a column vector r. We start from the premise that v lies between two closest values in r:

$$r_n < v < r_{n+1}$$

<sup>&</sup>lt;sup>9</sup>This table is available at the following link: https://pages.mtu.edu/~suits/notefreqs.html.

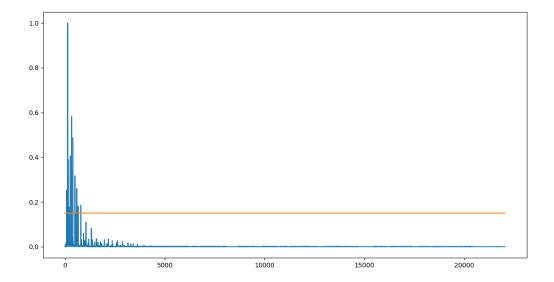


Figure 3: The normalized Fourier transformation of the signal given in figure 2, ie. a  $C_{maj}$  chord.



Figure 4: An example of how the peaks are (Is There A One-Line Function That Generates A Triangle Wave?)

The closest value is evaluated to be the expected note:

$$(|v - r_n| > |v - r_{n+1}| \Rightarrow e = r_{n+1}) \land (|v - r_n| < |v - r_{n+1}| \Rightarrow e = r_n)$$

where e is the final note.

The full algorithm is available at the following link: https://github.com/markovejnovic/Chordy.

## 3.6 Chord identification

Since this algorithm is only designed for chord triads, it was not difficult to design an algorithm which identifies chords from the previously identified notes. A simple look-up table was created and a simple searching algorithm was implemented.

# 4 Evaluation

This paper defines an arbitrary success rate as the ration between the number correctly analyzed chords (as per the chord played on an instrument) and the number of chords analyzed in total.

To identify how well the algorithm performs tests were created. 10 guitar chord samples were downloaded from the internet (*vide* the works cited). These chord samples were then analyzed using the algorithm. Since it was known which chords were played, it was possible to determine the success rate of the algorithm.

For a value of  $c_o = 0.15$ , the algorithm's success rate was 60%. However, when the value of  $c_o$  was chosen to be a value greater or lesser than 0.15 different success rates were achieved. A plot of the relationship between  $c_o$  and the success rate of the algorithm is given in figure 5.

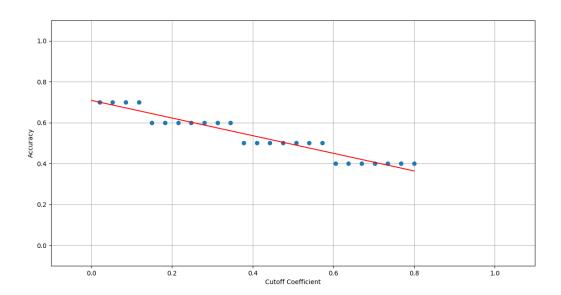


Figure 5: The relationship between  $c_o$  and the success rate of the algorithm.

It is obvious that the maximal success rate is achieved for the minimal value of  $c_o = 0.02$ . This success rate is equal to 70%. Note, however, that the chords played in the samples might have not been the ones that the authors stated they were. It is possible that the guitars the chords were played on were out of tune, or the chords were played improperly, or the tune of the guitars was not one of  $A_4 = 440$ , or it is even possible that the guitars used were not precisely manufactured.

Figure 5 also contains a negative linear fit which predicts that at  $c_o = 0$  the success rate is 0.7086, ie. 70.9%.

## 5 Conclusion

The algorithm seems effective in achieving the required result. The evaluation of the algorithm was incomplete, due to a lack of chord samples. For more accurate evaluation it is necessary to use better chord samples.

With a success rate higher than 70%, this algorithm can be used for roughly estimating chords, but cannot be used for rigorous musical analysis.

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