Quantum Field Theory on a Highly Symmetric Lattice

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Why Lattice Quantum Chromodynamics?

In quantum field theory scattering amplitudes in the form

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See a second sec



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 Feynman diagrams with n loops



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What is a Lattice?

Definition: Lattice Λ

$$\Lambda = \{ \sum_{i=1}^{n} a_i e_i \mid a_i \in \mathbb{Z} \}, \text{ with } \{e_i\}$$
 any basis of \mathbb{R}^n

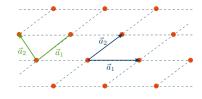


Figure: A bidimensional lattice.

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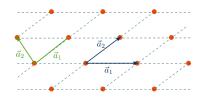


Figure: A bidimensional lattice.

Hypercubic lattice

 $\{e_i\}$ is the canonical basis of \mathbb{R}^n a is called *lattice spacing*.



Figure: A square lattice.

Basic idea

Fields can take values only in given parts of the lattice, $x \to n \in \Lambda$.

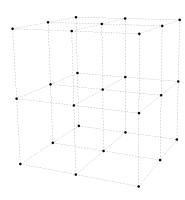


Figure: A (hyper)cubic lattice in \mathbb{R}^3 .

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Examples:

• Scalar fields $\Phi(x) \to \Phi(n)$ on sites

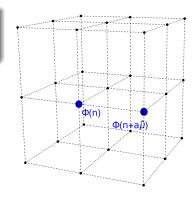


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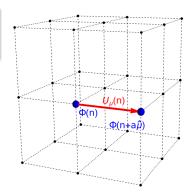


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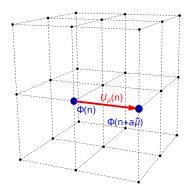


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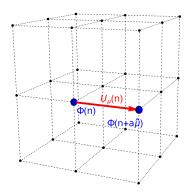


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Attention!

Spinorial fields are trickier to be discretized.

Gauge-Invariant Observables and Wilson Action

The Yang-Mills continuum action is
$$S = \frac{1}{4} \int \mathrm{d}^4 x F^{a\mu\nu}(x) F^a_{\mu\nu}(x).$$

On the lattice, every closed path is gauge-invariant.

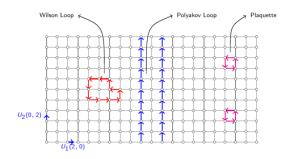


Figure: Gauge-invariant paths on a bidimensional lattice.[1]

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Definition: Plaquette $\overline{U_{\mu\nu}(n)}$

$$U_{\mu}(n)U_{
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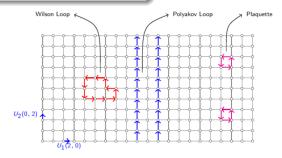


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Wilson's Idea

$$S = \frac{\beta}{2N} \sum_{n,\mu,\nu} \mathfrak{Re} \operatorname{Tr} \left(\mathbb{1} - U_{\mu\nu}(n) \right)$$

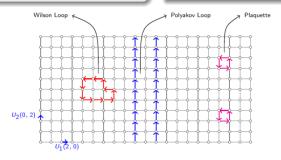


Figure: Gauge-invariant paths on a bidimensional lattice.[1]

Polyakov Loops and Potential

If the time coordinate is taken to be periodic, more closed paths arise.

Polyakov Loop

$$P(n) = \operatorname{Tr} \prod_{t=0}^{T-1} U_t(n)$$

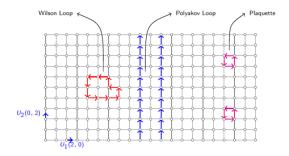


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The expectation value of two Polyakov loops is the potential.

Potential

$$V(R) = -\frac{1}{T}\log \langle P(0)P^{\dagger}(R) \rangle$$

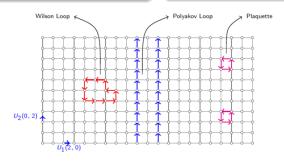


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Lattice symmetries

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