Quantum Field Theory on a Highly Symmetric Lattice

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Why Lattice Quantum Chromodynamics?

In quantum field theory scattering amplitudes in the form

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What is a Lattice?

Definition: Lattice Λ

 $\Lambda = \{ \sum_{i=1}^{n} a_i e_i \mid a_i \in \mathbb{Z} \}, \text{ with } \{e_i\}$ any basis of \mathbb{R}^n

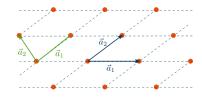


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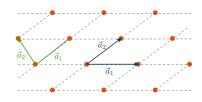


Figure: A bidimensional lattice.

Hypercubic lattice

 $\{e_i\}$ is the canonical basis of \mathbb{R}^n a is called *lattice spacing*.



Figure: A square lattice.

Basic idea

Fields can take values only in given parts of the lattice, $x \to n \in \Lambda$.

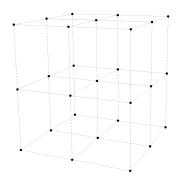


Figure: A (hyper)cubic lattice in \mathbb{R}^3 .

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• Scalar fields $\Phi(x) \to \Phi(n)$ on sites

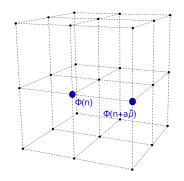


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Parallel Transporter

$$U_{\mu}(x) = \exp(igaA_{\mu}(x))$$

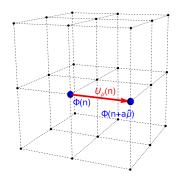


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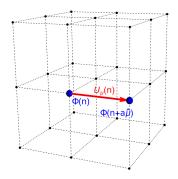


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Beware!

Spinorial fields are trickier to be discretized.

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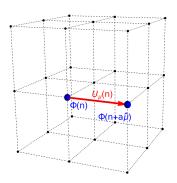


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Gauge-Invariant Observables and Wilson Action

The Yang-Mills continuum action is $S_E = \frac{1}{4} \int d^4x F^{a\mu\nu}(x) F^a_{\mu\nu}(x)$.

On the lattice, every closed path is gauge-invariant.

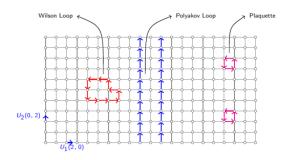


Figure: Gauge-invariant paths on a bidimensional lattice. [Sigdel, 2016]

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Definition: Plaquette $U_{\mu\nu}(n)$

$$U_{\mu}(n)U_{
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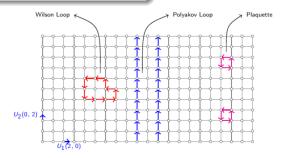


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Wilson's Idea

$$S = rac{eta}{2N} \sum_{n,\mu,
u} \mathfrak{Re} \operatorname{Tr} (\mathbb{1} - U_{\mu
u}(n))$$

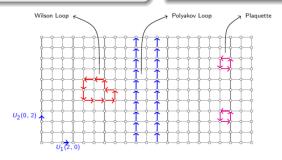


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Polyakov Loops and Potential

If the time coordinate is taken to be periodic, more closed paths arise.

Polyakov Loop

$$P(n) = \operatorname{Tr} \prod_{t=0}^{T-1} U_t(n)$$

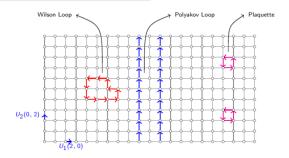


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The expectation value of two Polyakov loops is the potential.

Potential

$$V(R) = -\frac{1}{T}\log \langle P(0)P^{\dagger}(R) \rangle$$

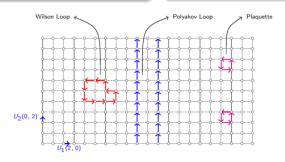


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$$x^{\mu} \rightarrow R^{\mu}_{\nu} x^{\nu} \quad R \in SO(4)$$

$$\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow \qquad \qquad \uparrow \qquad \qquad \Gamma \in T$$

T: group of rotations of multiples of 90° around any axis.

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Important:

Rotational invariance seems to be broken.

Rotational Invariance Restoration - Lang and Rebbi

Equipotential surfaces become spheres as the continuum limit is approached.

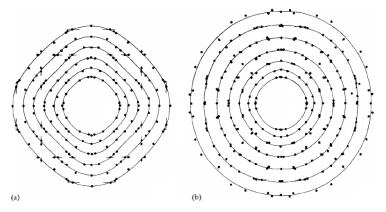
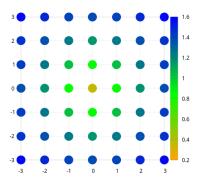


Figure: Restoration of rotational invariance from (a) $\beta=2$, $n_s=8$, $n_t=4$ to (b) $\beta=2.25$, $n_s=16$, $n_t=6$; the curves represent equipotential curves. [Lang and Rebbi, 1982]

Rotational Invariance Restoration

Values of β are slightly different from Lang and Rebbi's because $a(\beta) \approx \Lambda e^{-b_0 \beta}$, with Λ , $b_0 > 0^1$.



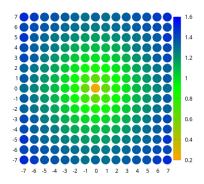


Figure: Potential from $\beta = 2.20$, $n_s = 8$, $n_t = 4$.

Figure: Potential from $\beta = 2.35$, $n_s = 16$, $n_t = 6$.

¹The simulation code is based on the code presented in refs [Panero, 2009; Mykkänen, Panero, and Rummukainen, 2012]

Higher Symmetry Lattices

Other, more rotational-symmetric, lattices have been used:

Body Centered Tesseract [Celmaster, 1982]

- 24 nearest neighbours
- 1152-element symmetry group

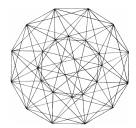


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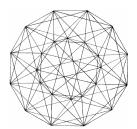


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F₄ coroots lattice [Neuberger, 1987]

- 48 nearest neighbours
- 2304-element symmetry group

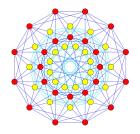


Figure: Two-dimensional projection of a F_4 coroots lattice. [Wikipedia, 2023]

Obtained from the roots lattice of the exceptional Lie algebra F₄ and its dual;

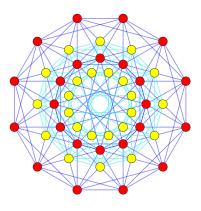


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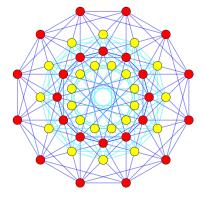


Figure: The 24 roots (red) of the F_4 lattice, projected on a bidimensional plane.

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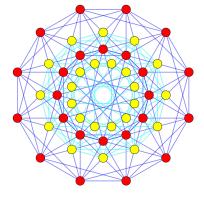


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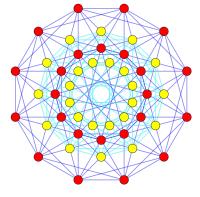


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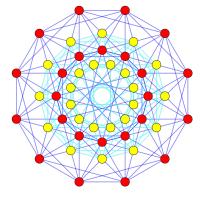


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- Has been used only to simulate scalar fields, in [Neuberger, 1987].

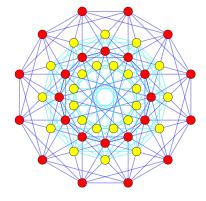
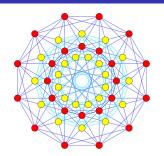


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Work in Progress

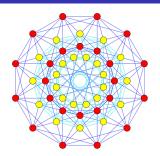
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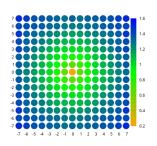


Work in Progress

 Implement the F₄ lattice in the simulation program and make efficiency studies;

 Make a rotational invariance study on the new lattice, hoping to get better results than the Simple Hypercubic lattice.





Thank you for your attention

Bibliography I

- Sigdel, Dibakar (2016). "Two Dimensional Lattice Gauge Theory with and without Fermion Content". In: FIU Electronic Theses and Dissertations 3224. DOI: 10.25148/etd.FIDC001748. URL: https://digitalcommons.fiu.edu/etd/3224?utm_source=digitalcommons.fiu.edu%2Fetd%2F3224&utm_medium=PDF&utm_campaign=PDFCoverPages.
- Lang, C. B. and C. Rebbi (1982). "Potential and Restoration of Rotational Symmetry in SU(2) Lattice Gauge Theory". In: *Phys. Lett.* B115. [, 322 (1982)], p. 137. DOI: 10.1016/0370-2693(82)90813-9.
- Panero, Marco (2009). "Thermodynamics of the QCD plasma and the large-N limit". In: *Phys. Rev. Lett.* 103, p. 232001. DOI: 10.1103/PhysRevLett.103.232001. arXiv: 0907.3719 [hep-lat].

Bibliography II

- Mykkänen, Anne, Marco Panero, and Kari Rummukainen (2012). "Casimir scaling and renormalization of Polyakov loops in large-N gauge theories". In: *JHEP* 1205, p. 069. DOI: 10.1007/JHEP05(2012)069. arXiv: 1202.2762 [hep-lat].
- Celmaster, William (1982). "Gauge Theories on the Body Centered Hypercubic Lattice". In: *Phys. Rev.* D26, p. 2955. DOI: 10.1103/PhysRevD.26.2955.
 - Neuberger, Herbert (1987). "SPINLESS FIELDS ON F(4) LATTICES". In: *Phys. Lett. B* 199, pp. 536–540. DOI: 10.1016/0370-2693(87)91623-6.
- Wikipedia (2023). URL: https://en.wikipedia.org/wiki/F4_%28mathematics%29 (visited on 08/29/2023).