

### Università degli Studi di Torino

Corso di Laurea Magistrale in Astrofisica e Fisica Teorica

### Titolo

TESI DI LAUREA MAGISTRALE

Relatore:

Prof.

Panero Marco

Candidato:

Aliberti Marco Matricola 855766

Anno Accademico 2022/2023

## Abstract

## Contents

T	QCD on the Lattice					
	1.1 The QCD Continuum Action	2				
	1.1.1 Spinor Fields					
	1.1.2 Quantum Electrodynamics	2				
	1.1.3 Nonabelian Gauge Theories	9				
2	Computer Simulation of Gauge Theories	4				
3	Gauge Theories Simulation on non-hypercubic lattice F4	5				
4	Simulation Results	6				
5	Conclusions	7				

## QCD on the Lattice

#### 1.1 The QCD Continuum Action

In order to write the action of QCD on the lattice, I must first recall how the theory is formulated in the continuum.

#### 1.1.1 Spinor Fields

Let us take into consideration a (free) quantum field theory describing a fermion, such as a quark or a lepton, in a 4-dimensional spacetime with metric  $\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$ . Its action (in natural units, where  $c = \hbar = 1$ ) can be written as:

$$S_F[\psi(x), \overline{\psi}(x)] = \int d^4x \left( i\overline{\psi} \partial \!\!\!/ \psi - m\overline{\psi} \psi \right)$$
 (1.1.1)

from which, upon the application of the variational principle, the Dirac equation follows:

$$(i\partial - m) \psi(x) = 0 \tag{1.1.2}$$

It can now be easily checked by direct computation that this action is invariant under a rigid (global) phase transformation, also called global U(1) transformation:

$$\psi \to \psi' = e^{-i\alpha}\psi \tag{1.1.3}$$

$$\overline{\psi} \to \overline{\psi}' = \overline{\psi}e^{i\alpha}$$

where  $\alpha$  is a constant that does not depend on the spacetime coordinate x, while if  $\alpha = \alpha(x)$  the action (1.1.1) would not be invariant because of the kinetic term.

#### 1.1.2 Quantum Electrodynamics

As the free field theory itself is non interacting, it does not provide any real-world prediction, so it is useful to write an interacting action where the spinor field is coupled, for instance, to a vector field  $A_{\mu}$ , i.e. the photon. One way to implement this interaction is to ask for local, instead of global, invariance of the action (1.1.1) under the phase transformation (1.1.3), where now  $\alpha = \alpha(x)$ . In order to do so, one has to define the covariant derivative as follows:

$$D_{\mu} \equiv (\partial_{\mu} + igA_{\mu}) \tag{1.1.4}$$

where g is the coupling constant.<sup>1</sup>

The vector field's kinetic term is written in terms of its field-strength, namely:

$$F_{\mu\nu} \equiv -\frac{\mathrm{i}}{g} \left[ D_{\mu}, D_{\nu} \right] =$$

$$= -\frac{\mathrm{i}}{g} \left( D_{\mu} \left( \partial_{\nu} + \mathrm{i} g A_{\nu} \right) - D_{\nu} \left( \partial_{\mu} + \mathrm{i} g A_{\mu} \right) \right) =$$

$$= -\frac{\mathrm{i}}{g} \left( \partial_{\mu} \partial_{\nu} + \mathrm{i} g \partial_{\mu} A_{\nu} - g^{2} A_{\mu} A_{\nu} - \partial_{\nu} \partial_{\mu} - \mathrm{i} g \partial_{\nu} A_{\mu} + g^{2} A_{\nu} A_{\mu} \right) =$$

$$= \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu} + \mathrm{i} g \left[ A_{\mu}, A_{\nu} \right] =$$

$$= \partial_{\mu} A_{\nu} - \partial_{\nu} A_{\mu}$$

$$(1.1.5)$$

<sup>&</sup>lt;sup>1</sup>Usually g is called e, the electron charge, though I will be using g in analogy to nonabelian gauge theories.

where the term with the commutator is = 0 in the abelian theory.

Two different fields  $A_{\mu}$  and  $A'_{\mu}$  describe the same physics if one can be obtained from another throug a gauge transformation:

$$A'_{\mu}(x) = A_{\mu}(x) - \frac{1}{g}\partial_{\mu}\alpha(x)$$

$$F'_{\mu\nu} = F_{\mu\nu} - \frac{1}{g}(\partial_{\mu}\partial_{\nu} - \partial_{\nu}\partial_{\mu})\alpha(x) = F_{\mu\nu}$$

$$(1.1.6)$$

Thus, the free action for the vector field is:

$$S_{EM} = -\frac{1}{4} \int d^4x F_{\mu\nu} F^{\mu\nu} \tag{1.1.7}$$

That is also gauge invariant, i.e. invariant under (1.1.6), as  $F_{\mu\nu}$  is gauge invariant. The term that broke the local phase invariance of the action (1.1.1) can now be "absorbed" by  $A_{\mu}$  through a gauge transformation (1.1.6), thus making the full action gauge invariant:

$$S = \int d^4x \left( i\overline{\psi} \partial \psi - m\overline{\psi}\psi - g\overline{\psi} A\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu} \right)$$
 (1.1.8)

#### 1.1.3 Nonabelian Gauge Theories

# Computer Simulation of Gauge Theories

# Gauge Theories Simulation on non-hypercubic lattice F4



