

Quantum Field Theory on a Highly Symmetric Lattice

Marco Aliberti

Università degli Studi di Torino

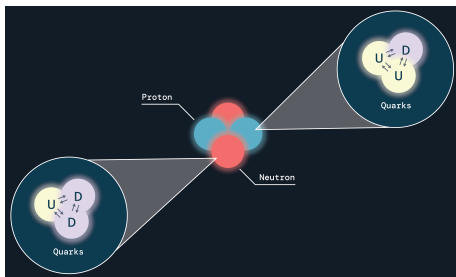
23rd October, 2023

The Strong Interaction

Matter is made of Atoms

Atoms are made of Nuclei and Electrons

Nuclei are made of Protons and Neutrons,
composed of Quarks and Gluons

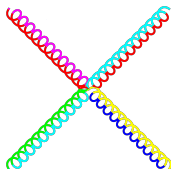


Described by Quantum Chromodynamics (QCD)

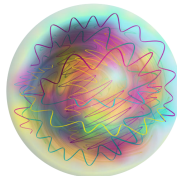
Image credits: NASA^[StrongForceImg]

Perturbative vs Non-Perturbative

Perturbative



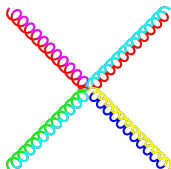
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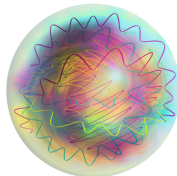
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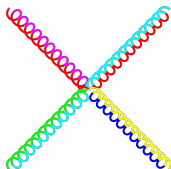
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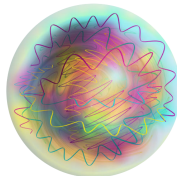
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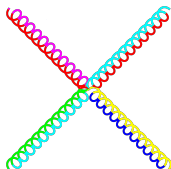
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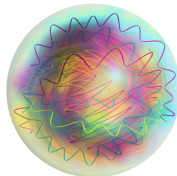
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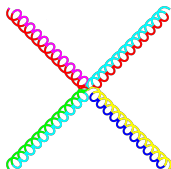
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- Straightforward series expansion in powers of small $g \Leftrightarrow$ Feynman diagrams with n loops
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- No straightforward approach

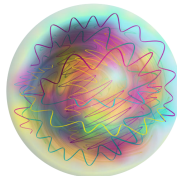
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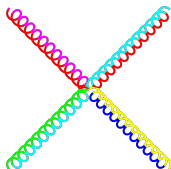
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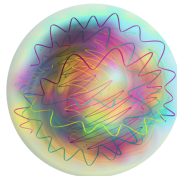
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What is a Lattice?

Definition: Lattice Λ

$\Lambda = \{ \sum_{i=1}^n a_i e_i \mid a_i \in \mathbb{Z} \}$, with $\{e_i\}$ any basis of \mathbb{R}^n

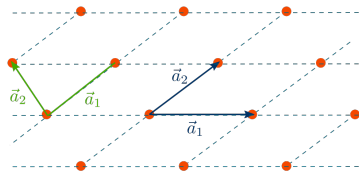


Figure: A bidimensional lattice.

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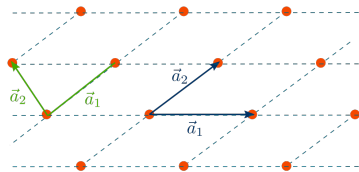


Figure: A bidimensional lattice.

Hypercubic lattice

$\{e_i\}$ is the canonical basis of \mathbb{R}^n
 a is called *lattice spacing*.

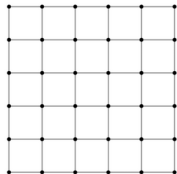


Figure: A square lattice.

Basic idea

Fields can take values only in given parts of the lattice, $x \rightarrow n \in \Lambda$.

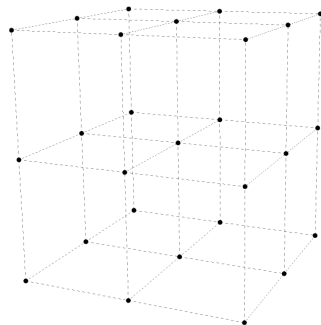


Figure: A (hyper)cubic lattice in \mathbb{R}^3 .

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Examples:

- **Scalar fields** $\Phi(x) \rightarrow \Phi(n)$ on sites

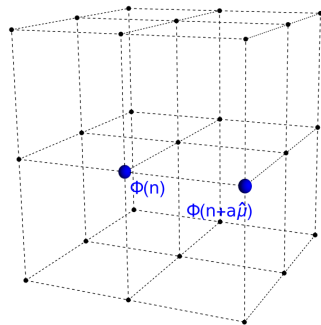


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Parallel Transporter

$$U_\mu(x) = \exp(igaA_\mu(x))$$

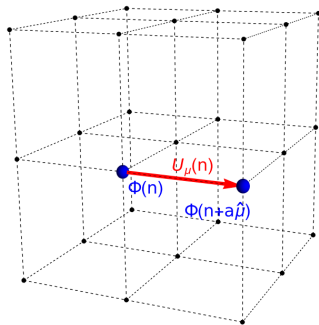


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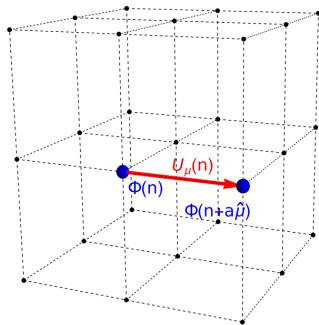


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Beware!

Spinorial fields are trickier to be discretized.

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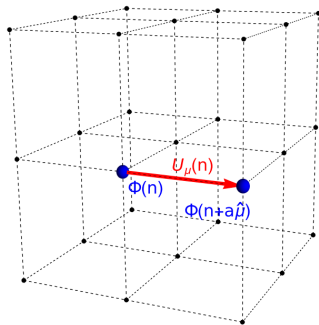


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Gauge-Invariant Observables and Wilson Action

The Yang-Mills continuum action is
$$S_E = \frac{1}{4} \int d^4x F^{a\mu\nu}(x) F_{\mu\nu}^a(x).$$

On the lattice, every closed path is gauge-invariant.

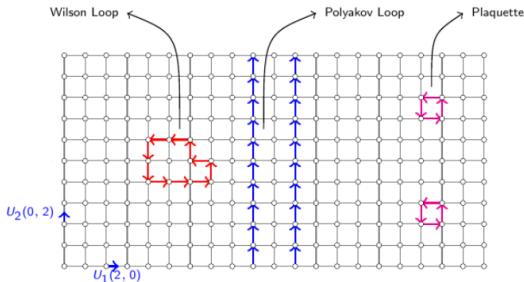


Figure: Gauge-invariant paths on a bidimensional lattice. [Sigdel, 2016]

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$$U_\mu(n) U_\nu(n + \mu) U_\mu^\dagger(n + \nu) U_\nu^\dagger(n)$$

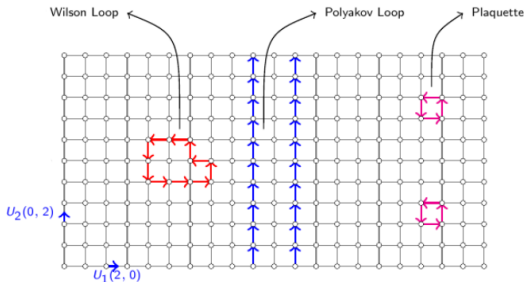


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Wilson's Idea

$$S = \frac{\beta}{2N} \sum_{n,\mu,\nu} \Re \text{Tr} (1 - U_{\mu\nu}(n))$$

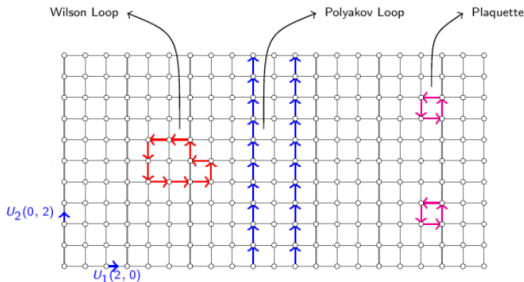


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Polyakov Loops and Potential

If the time coordinate is taken to be periodic, more closed paths arise.

Polyakov Loop

$$P(n) = \text{Tr} \prod_{t=0}^{T-1} U_t(n)$$

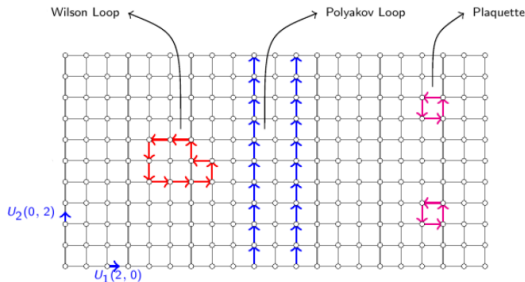


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The expectation value of two Polyakov loops is the potential.

Polyakov Loop

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Potential

$$V(R) = -\frac{1}{T} \log \langle P(0) P^\dagger(R) \rangle$$

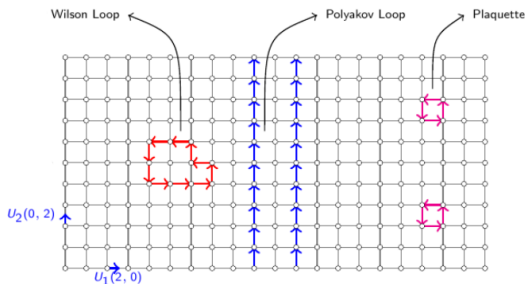


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$$x^\mu \rightarrow R^\mu_\nu x^\nu \quad R \in SO(4)$$

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$$n \rightarrow \Gamma n \quad \Gamma \in T$$

T : group of rotations of multiples of 90° around any axis.

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Important:

Rotational invariance seems to be broken.

Rotational Invariance Restoration - Lang and Rebbi

Equipotential surfaces become spheres as the continuum limit is approached.

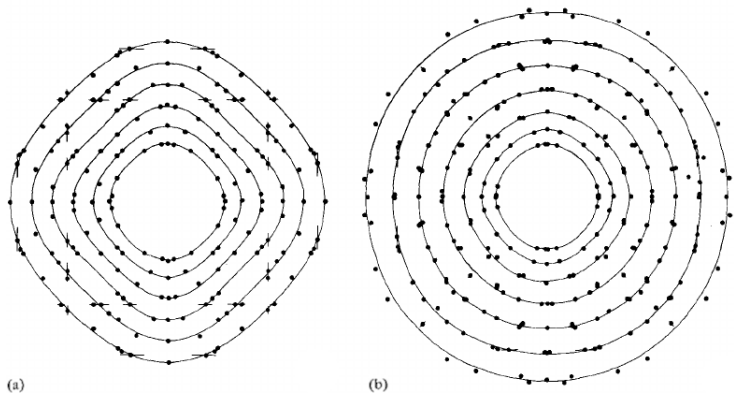


Figure: Restoration of rotational invariance from (a) $\beta = 2$, $n_s = 8$, $n_t = 4$ to (b) $\beta = 2.25$, $n_s = 16$, $n_t = 6$; the curves represent equipotential curves.
[Lang and Rebbi, 1982]

Rotational Invariance Restoration

Values of β are slightly different from Lang and Rebbi's because $a(\beta) \approx \Lambda e^{-b_0 \beta}$, with $\Lambda, b_0 > 0$ ¹.

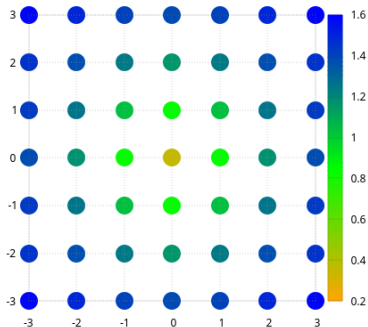


Figure: Potential from $\beta = 2.20$,
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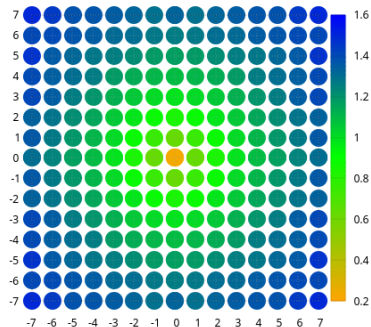


Figure: Potential from $\beta = 2.35$,
 $n_s = 16$, $n_t = 6$.

¹The simulation code is based on the code presented in
refs. [Panero, 2009; Mykkänen, Panero, and Rummukainen, 2012].

Higher Symmetry Lattices

Other, more rotational-symmetric, lattices have been used:

Body Centered Tesseract

[Celmaster, 1982]

- 24 nearest neighbours
- 1152-element symmetry group

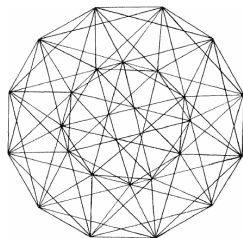


Figure: Two-dimensional projection of a BCT. [Celmaster, 1982]

The SH lattice has 8 nearest neighbours and a 288-element symmetry group.

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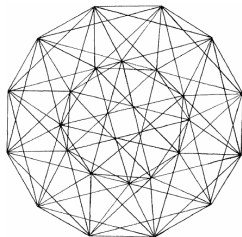


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F_4 coroots lattice [Neuberger, 1987]

- 48 nearest neighbours
- 2304-element symmetry group

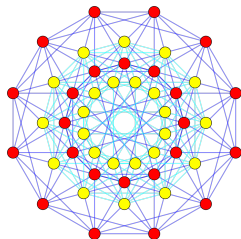


Figure: Two-dimensional projection of a F_4 coroots lattice. [Wikipedia, 2023]

The E_8 lattice has 8 nearest neighbours and a 288-element symmetry group.

F_4 Coroots Lattice

- Obtained from the roots lattice of the exceptional Lie algebra F_4 and its dual;

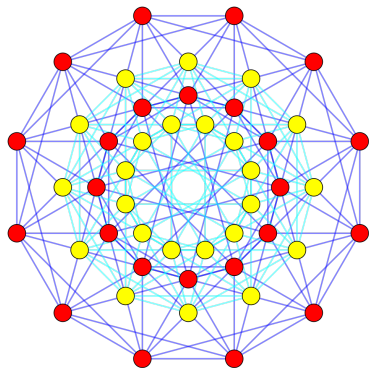


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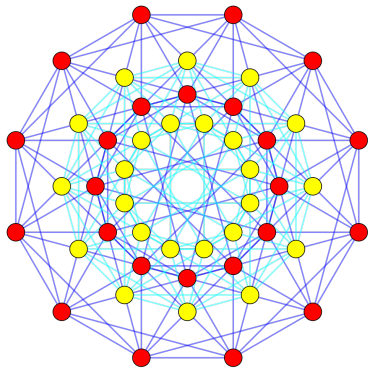


Figure: The 24 roots (red) of the F_4 lattice, projected on a bidimensional plane.

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 - The 8 possible permutations of $(\pm 1, 0, 0, 0)$
 - The 16 vectors of the form $(\pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2})$

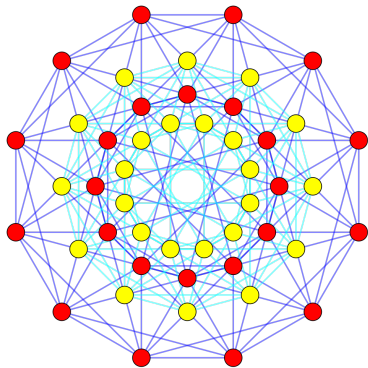


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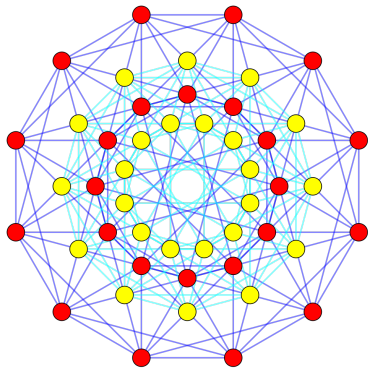


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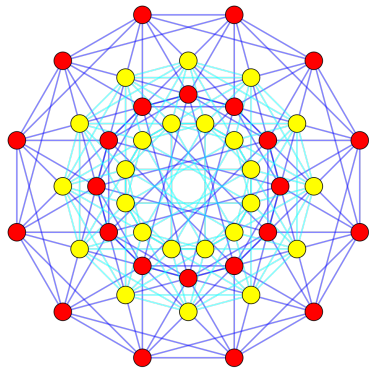


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- Exists only in 4 dimensions;
- Contains the Simple Hypercubic lattice and the BCT;
- Has been used only to simulate scalar fields, in [Neuberger, 1987].

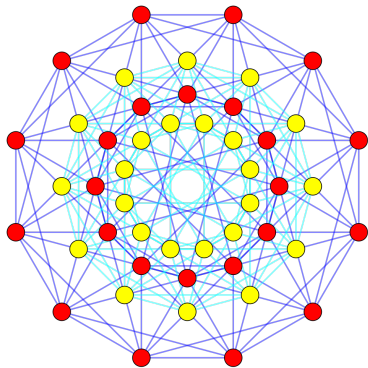
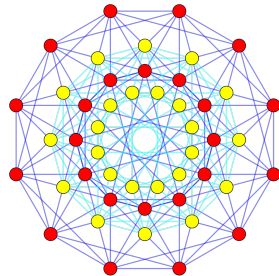
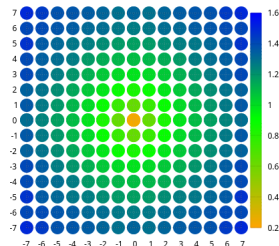
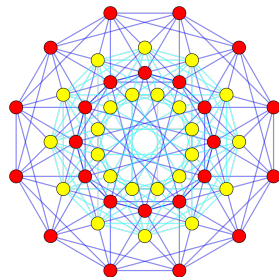


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- Make a rotational invariance study on the new lattice, hoping to get better results than the Simple Hypercubic lattice.



Thank you for your attention



Sigdel, Dibakar (2016). “Two Dimensional Lattice Gauge Theory with and without Fermion Content”. In: *FIU Electronic Theses and Dissertations* 3224. DOI: 10.25148/etd.FIDC001748. URL: https://digitalcommons.fiu.edu/etd/3224?utm_source=digitalcommons.fiu.edu%2Fetd%2F3224&utm_medium=PDF&utm_campaign=PDFCoverPages.



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Neuberger, Herbert (1987). “SPINLESS FIELDS ON $F(4)$ LATTICES”. In: *Phys. Lett. B* 199, pp. 536–540. DOI: [10.1016/0370-2693\(87\)91623-6](https://doi.org/10.1016/0370-2693(87)91623-6).



Wikipedia (2023). URL: https://en.wikipedia.org/wiki/F4_%28mathematics%29 (visited on 08/29/2023).