

Quantum Field Theory on a Highly Symmetric Lattice

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Università degli Studi di Torino

23rd October, 2023

The Strong Interaction

Matter is made of Atoms

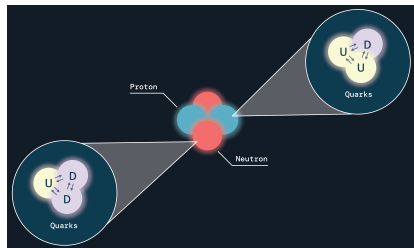


Image credits: NASA^[NASA, 2022]

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Atoms are made of Nuclei and Electrons

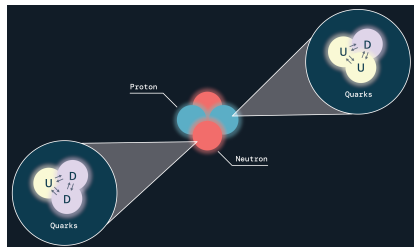


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Nuclei are made of Protons and Neutrons,
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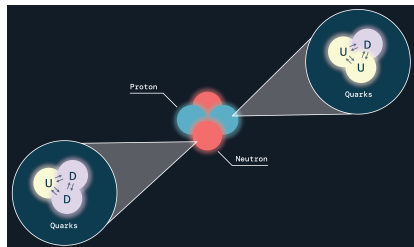


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Quarks and **Gluons**

Described by Quantum
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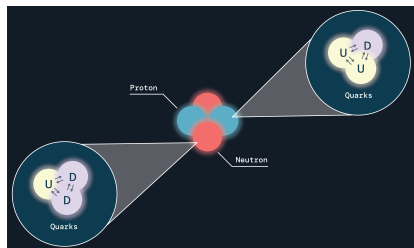


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Quantum Chromodynamics (QCD)

Described by an $SU(3)$ Yang-Mills theory

$$S = \frac{1}{4} \int d^4x F_{\mu\nu}^a(x) F^{a\mu\nu}(x)$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a - gf_{bc}^a A_\mu^b A_\nu^c$$

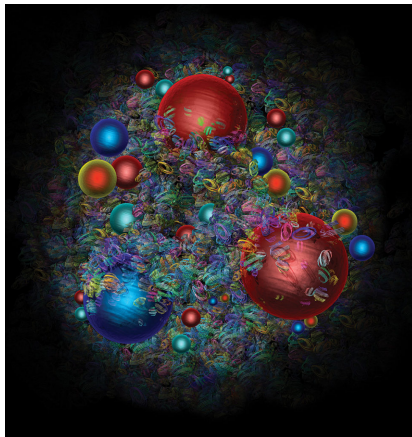


Figure: An artist's representation of a proton^[CERN, 2019].

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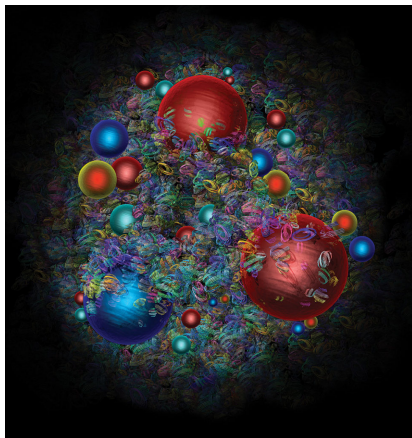


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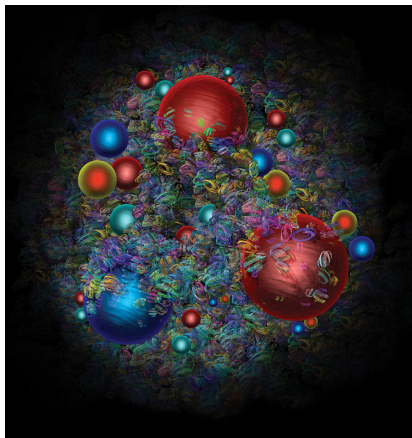


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- Heavily non-perturbative nature (except for specific regimes)

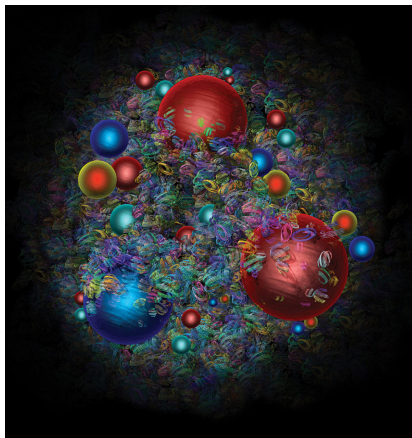


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Lattice Field Theory

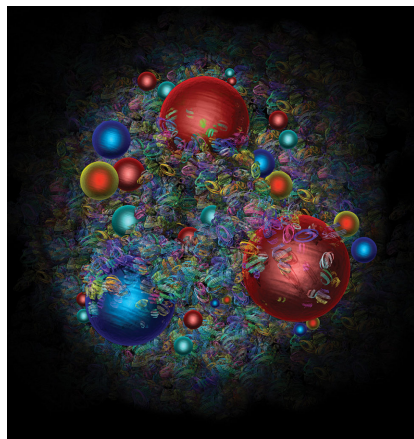


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What is a Lattice?

Definition: Lattice Λ

$\Lambda = \{ \sum_{i=1}^n a_i \mathbf{e}_i \mid a_i \in \mathbb{Z} \}$, with $\{\mathbf{e}_i\}$ any basis of \mathbb{R}^n

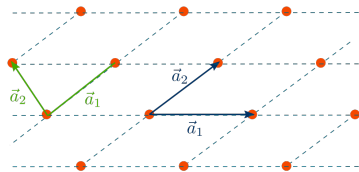


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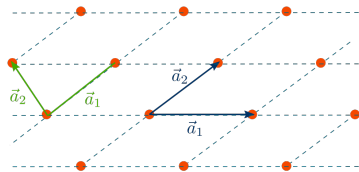


Figure: A bidimensional lattice.

Hypercubic lattice

$\{e_i\}$ is the canonical basis of \mathbb{R}^n
 a is called *lattice spacing*.

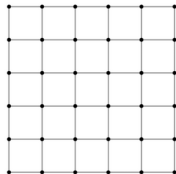


Figure: A square lattice.

Basic idea

Fields can take values only in given parts of the lattice, $x \rightarrow n \in \Lambda$.

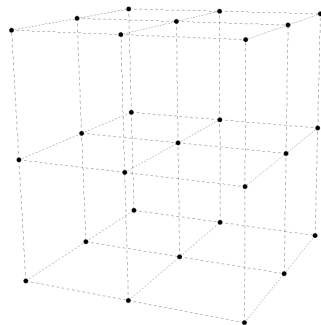


Figure: A (hyper)cubic lattice in \mathbb{R}^3 .

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Examples:

- **Scalar fields** $\Phi(x) \rightarrow \Phi(n)$ on sites

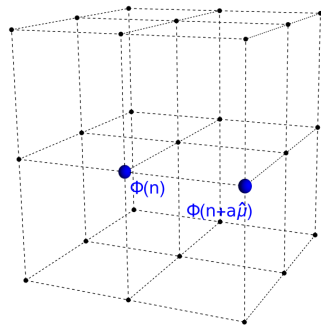


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Parallel Transporter

$$U_\mu(x) = \exp(igaA_\mu(x))$$

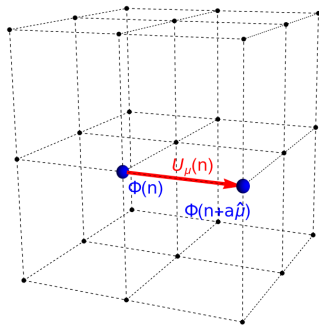


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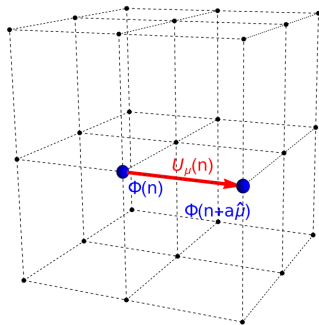


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Beware!

Spinorial fields are trickier to be discretized.

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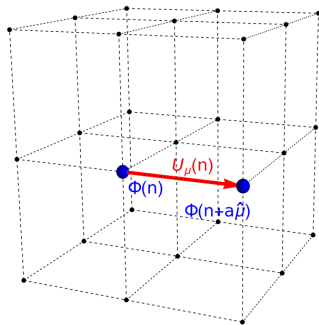


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Gauge-Invariant Observables and Wilson Action

The Yang-Mills continuum action is
$$S_E = \frac{1}{4} \int d^4x F^{a\mu\nu}(x) F_{\mu\nu}^a(x).$$

On the lattice, every closed path is gauge-invariant.

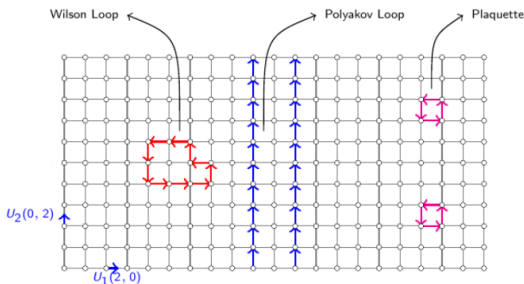


Figure: Gauge-invariant paths on a bidimensional lattice [Sigdel, 2016].

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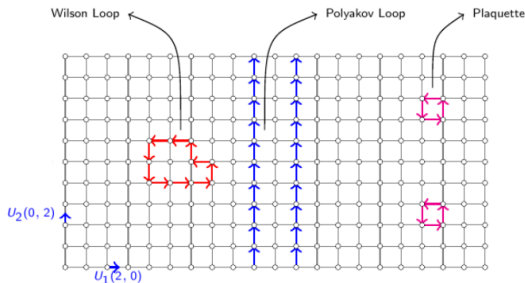


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Wilson's Idea

$$S = \frac{\beta}{2N} \sum_{n,\mu,\nu} \Re \text{Tr} (1 - U_{\mu\nu}(n))$$

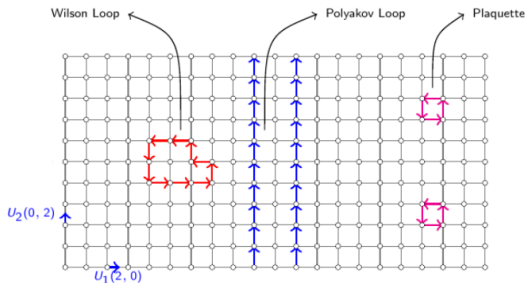


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Polyakov Loops and Potential

If the time coordinate is taken to be periodic, more closed paths arise.

Polyakov Loop

$$P(n) = \text{Tr} \prod_{t=0}^{T-1} U_t(n)$$

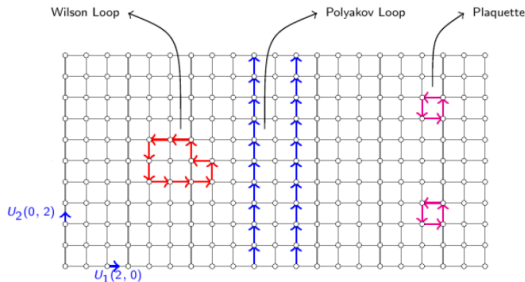


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Polyakov Loops and Potential

If the time coordinate is taken to be periodic, more closed paths arise.

The expectation value of two Polyakov loops is the potential.

Polyakov Loop

$$P(n) = \text{Tr} \prod_{t=0}^{T-1} U_t(n)$$

Potential

$$V(R) = -\frac{1}{T} \log \langle P(0) P^\dagger(R) \rangle$$

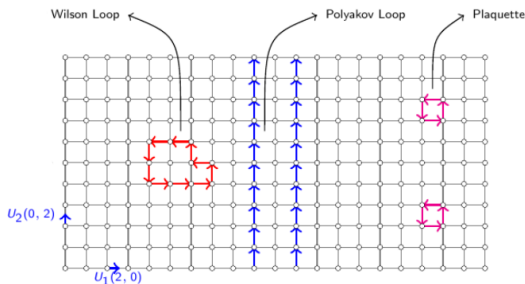


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Monte Carlo Simulations

Computers are used to simulate Lattice Field Theories



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Monte Carlo Simulations

Computers are used to simulate Lattice Field Theories

- Random configurations of link variables are generated.
- Proper Monte Carlo algorithms evolve the configurations towards minimums of the action.
- A great number of observables is evaluated and then their mean value is computed.

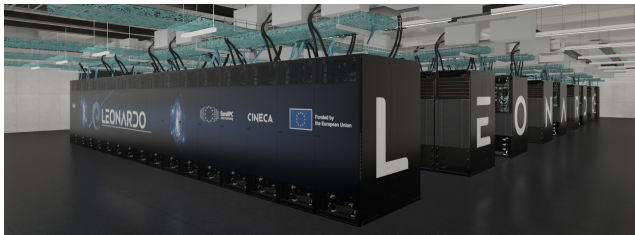


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Poincaré Group can be divided in:

Translations

Rotations

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Rotations

$$x^\mu \rightarrow R^\mu_\nu x^\nu \quad R \in SO(4)$$

\Downarrow

$$n \rightarrow \Gamma n \quad \Gamma \in G_{\Lambda_{SH}}$$

$G_{\Lambda_{SH}}$: group of rotations of multiples of 90° around any axis.

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$$\Gamma \not\rightarrow R \text{ for } a \rightarrow 0$$

Important:

Rotational invariance seems to be broken.

Rotational Invariance Restoration - Lang and Rebbi

Equipotential surfaces become spheres as the continuum limit is approached^[Lang and Rebbi, 1982].

The gauge group used was the discrete icosahedral subgroup $\tilde{Y} \subset SU(2)$.

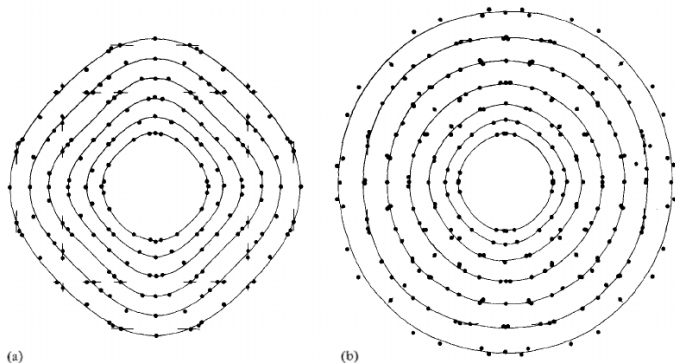


Figure: Restoration of rotational invariance from (a) $\beta = 2$, $n_s = 8$, $n_t = 4$ to (b) $\beta = 2.25$, $n_s = 16$, $n_t = 6$; the curves represent equipotential curves.

Rotational Invariance Restoration

Results of simulations for gauge group $SU(2)$ with 20000 measurements each¹. Approach slightly different than Lang and Rebbi's.

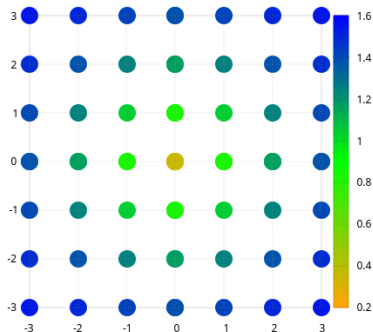


Figure: Potential from $\beta = 2.20$,
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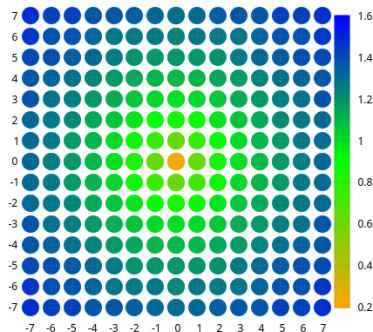


Figure: Potential from $\beta = 2.35$,
 $n_s = 16$, $n_t = 6$.

¹The simulation code is based on the code presented in refs. [Panero, 2009; Mykkänen, Panero, and Rummukainen, 2012].

Higher Symmetry Lattices

Other, more rotational-symmetric, lattices have been used:

Body Centered Tesseract

- 24 nearest neighbours
- 1152-element symmetry group

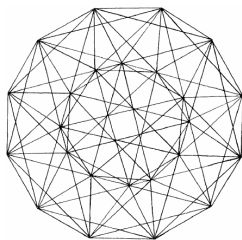


Figure: Two-dimensional projection of a BCT[Celmaster, 1982].

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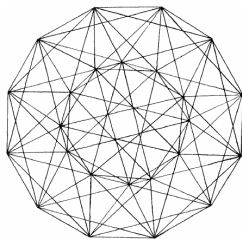


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F_4 coroots lattice^[Neuberger, 1987]

- 48 nearest neighbours
- 2304-element symmetry group

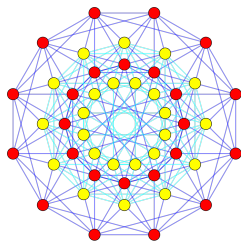


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Body Centered Tesseract (BCT)

- Obtained from Simple Hypercubic lattice considering also the centers;

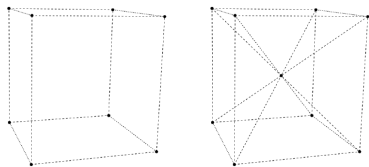


Figure: Cubic Cell (left) and BC Cubic Cell (right).

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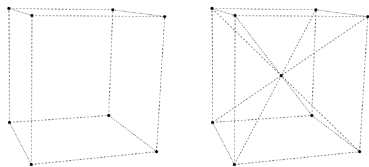


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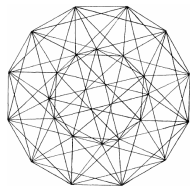


Figure: Bidimensional projection of the 24-cell.

Body Centered Tesseract (BCT)

- Obtained from Simple Hypercubic lattice considering also the centers;
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 - The 8 possible permutations of $(\pm 1, 0, 0, 0)$
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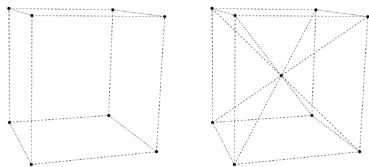


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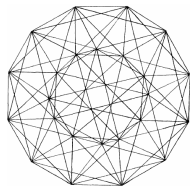


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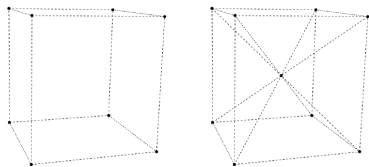


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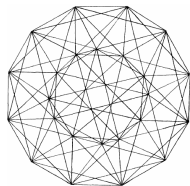


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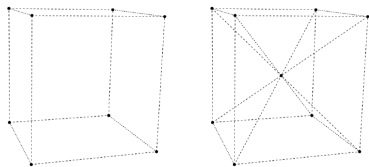


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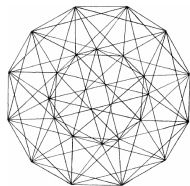


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- Plaquettes are triangular;
- Contains the Simple Hypercubic lattice;
- Has been used to simulate $SU(2)$ Yang-Mills theories, in [Celmaster, 1982].

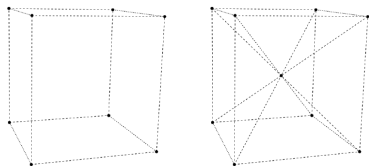


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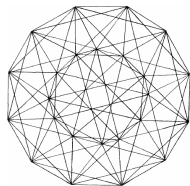


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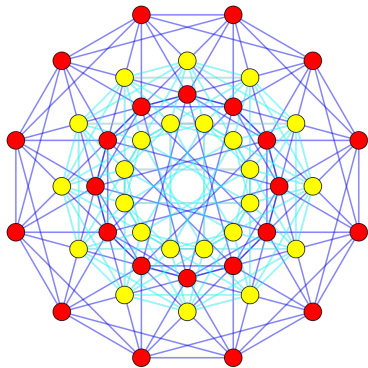


Figure: Bidimensional projection of the F_4 lattice.

F_4 Coroots Lattice

- Obtained from the roots lattice of the exceptional Lie algebra F_4 and its dual;
- Has 48 nearest neighbours:
 - The 24 roots are all possible permutations of coordinate positions of $(\pm 1, \pm 1, 0, 0)$

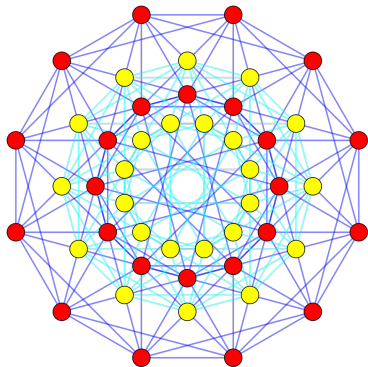


Figure: The 24 roots (red) of the F_4 lattice, projected on a bidimensional plane.

F_4 Coroots Lattice

- Obtained from the roots lattice of the exceptional Lie algebra F_4 and its dual;
- Has 48 nearest neighbours:
 - The 24 roots are all possible permutations of coordinate positions of $(\pm 1, \pm 1, 0, 0)$
 - The 24 dual roots (coroots) are:
 - The 8 possible permutations of $(\pm 1, 0, 0, 0)$
 - The 16 vectors of the form $(\pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2}, \pm \frac{1}{2})$

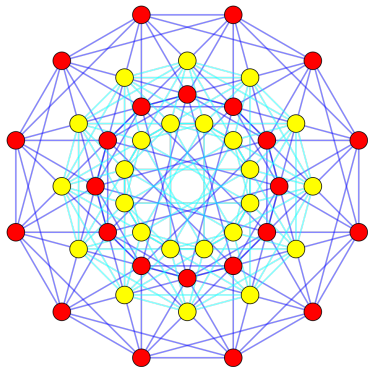


Figure: The 24 roots (red) and the 24 coroots (yellow) of the F_4 lattice, projected on a bidimensional plane.

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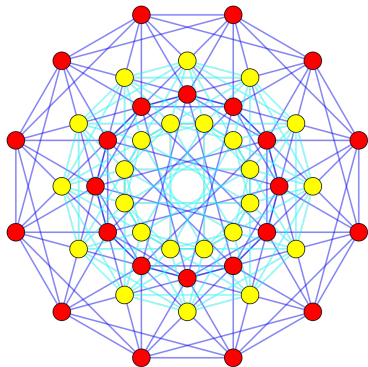


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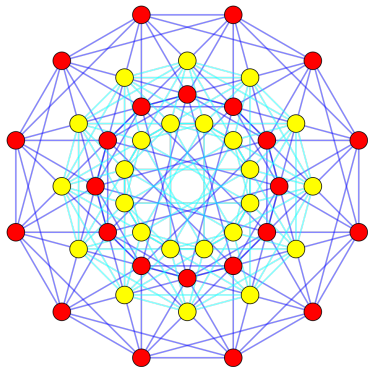


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- Exists only in 4 dimensions;
- Is a more symmetric version of the BCT;
- Has been used only to simulate scalar fields, in [Neuberger, 1987].

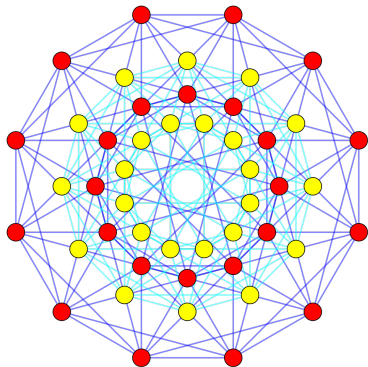


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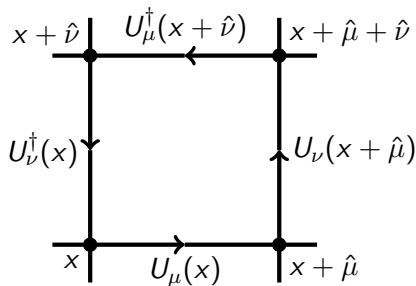
Simulations on SH Lattice

Wilson Action:

$$S_W = \frac{\beta}{2N} \sum_{x \in \Lambda} \sum_{\mu < \nu} \Re \text{Tr}[\mathbb{1} - U_{\mu\nu}(x)]$$

Plaquette:

$$U_{\mu\nu} = U_\mu(x) U_\nu(x + \hat{\mu}) U_\mu^\dagger(x + \hat{\nu}) U_\nu^\dagger(x)$$



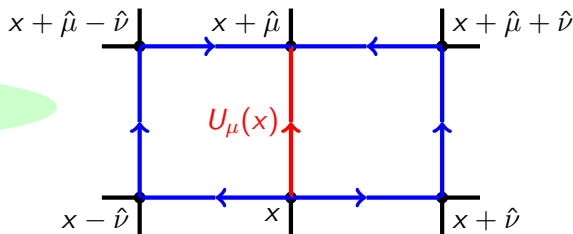
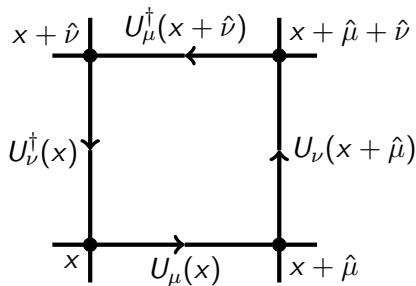
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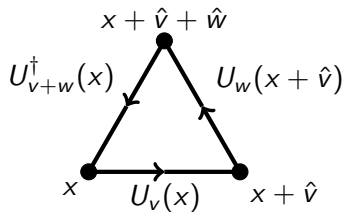
6 staples for each link

BCT Action:

$$S_{BCT} = \frac{\beta}{8} \sum_{\Delta} \Re \text{Tr } U_{\Delta}$$

Plaquette:

$$U_{\Delta} = U_v(x) U_w(x + \hat{v}) U_{v+w}^{\dagger}(x)$$



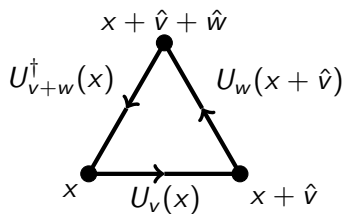
Simulations on BCT Lattice

BCT Action:

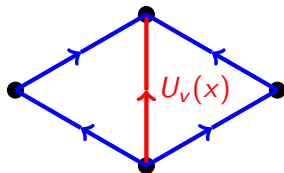
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Plaquette:

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8 staples for each link



Thank you for your attention

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