An ALE Finite Element Method for 2D Navier-Stokes Equation with Species Transport Equation

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Outline



- 1. Introduction
- 2. Mathematical Model
- 3. Computational Code
- 4. Validation
- 5. Results
- 6. Conclusion

Introduction

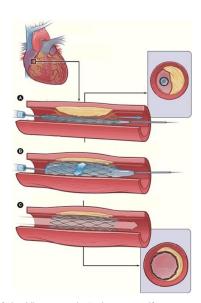


Motivation:

► Ischaemic heart disease and stroke have remained the leading death causes globally in the last 15 years [1]

Goals:

- ► To develop a ALE-FE code for 2D Navier-Stokes Equation with Species Transport Equation
- ➤ To create new drug-eluting stent design patent



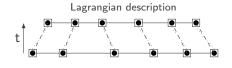


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Arbitrary Lagrangian-Eulerian (ALE)



The Arbitrary Lagrangian-Eulerian combines the classical motion descriptions, while it provides [2]:



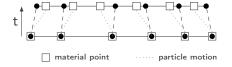
Advantages:

 Simulations in fluid-structure and moving boundary problems

Eulerian description

Disadvantages:

► The computational mesh requires an extensive topological treatment



node

ALE description

[2] Donea, J., Huerta, A., Ponthot, J.-P. and Rodríguez-Ferran, A. (2004). Arbitrary Lagrangian–Eulerian Methods. In Encyclopedia of Computational Mechanics doi:10.1002/0470091355.ecm009

mesh motion

Governing Equations



Assumptions [3]:

- 1. Continuum hypothesis
- 2. Homogeneous and Isotropic
- 3. Incompressible
- 4. Newtonian
- 5. Constant Mass Difusivity
- 6. Single-phase Flow
- 7. Two-dimensional flow

$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} - \hat{\mathbf{v}}) \cdot \nabla \mathbf{v} = -\nabla p + \frac{1}{Re} \nabla^2 \mathbf{v}$$

$$\nabla \cdot \mathbf{v} = 0$$

$$\frac{\partial e}{\partial t} + (\mathbf{v} - \hat{\mathbf{v}}) \cdot \nabla e = \frac{1}{ReSc} \nabla^2 e$$

▶ If the mesh velocity field $\hat{\mathbf{v}} = \mathbf{v} (Lagrangian)$ or $\hat{\mathbf{v}} = 0 (Eulerian)$

[3] Panton, R. (2013). Incompressible Flow John Wiley & Sons, Ltd

Semi-Lagrangian Method



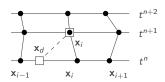
The convective term was replaced by material derivative in the direction of characteristic trajectory, that is: $\frac{D(\cdot)}{Dt} = \frac{\partial(\cdot)}{\partial t} + (\mathbf{v} - \mathbf{\hat{v}}) \cdot \nabla(\cdot)$

$$\frac{D\mathbf{v}}{Dt} = -\nabla p + \frac{1}{Re}\nabla^2 \mathbf{v}$$
$$\nabla \cdot \mathbf{v} = 0$$

 $\frac{De}{Dt} = \frac{1}{ReSc} \nabla^2 e$

The departure node is calculated by $x_d^n = x_i^{n+1} - (\mathbf{v} - \mathbf{v}) \Delta t$. Then, a searching procedure is required to find x_d^n using barycentric coordina-

tes



Semi-Lagrangian Method



The implicit semi-Lagrangian time discretization provides [4]:

Advantages:

- ► Symmetric linear systems
- Unconditionnal stability

Disadvantages:

- Numerical Diffusion
- Searching procedure may lead to excessive computational cost if it is not well designed

$$\frac{\mathbf{v}_i^{n+1} - \mathbf{v}_d^n}{\Delta t} = -\nabla p^{n+1} + \frac{1}{Re} \nabla^2 \mathbf{v}^{n+1}$$

$$\nabla \cdot \mathbf{v} = 0$$

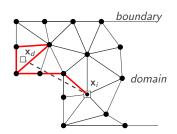
$$\frac{e_i^{n+1} - e_d^n}{\Delta t} = \frac{1}{ReSc} \nabla^2 e^{n+1}$$

[4] Pironneau, O. On the transport-diffusion algorithm and its applications to the Navier-Stokes equations. Numer. Math. 38, 309–332 (1982). https://doi.org/10.1007/BF01396435

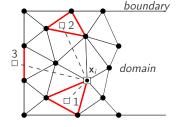
Searching Procedure



The start searching procedure was used. This procedure has presented an acceptable computational cost, in addition to the concave domains can be used

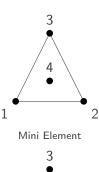


The searching procedure must be prepared for three situations that may occur. After the departure node is found, the interpolation is performed



Galerkin FE Method





$$\begin{bmatrix} \frac{\mathsf{M}}{\Delta t} + \frac{\mathsf{K}}{Re} & -\mathsf{G} \\ \mathsf{D} & 0 \end{bmatrix} \begin{bmatrix} \mathsf{v}_i^{n+1} \\ p \end{bmatrix} = \begin{bmatrix} \frac{\mathsf{M}}{\Delta t} \mathsf{v}_d^n \\ 0 \end{bmatrix} + bc$$

$$\left[\frac{\mathsf{M}}{\Delta t} + \frac{\mathsf{K}}{ReSc}\right] e_i^{n+1} = \frac{\mathsf{M}}{\Delta t} e_d^n + bc_e$$

Mesh Velocity



The computational mesh velocity $\hat{\mathbf{v}}$ is defined as a linear combination of other velocities, such as:

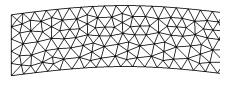
$$\mathbf{\hat{v}} = \beta_1 \mathbf{v}_1 + \beta_2 \mathbf{v}_2$$

where, \mathbf{v}_1 is the Lagrangian velocity, \mathbf{v}_2 is the Laplacian Smoothing velocity, β_1 and β_2 are the parameters control the Lagrangian and Laplacian motion respectively.

Laplacian Smoothing



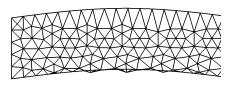
To avoid the fast degradation of the computational elements due to ALE description, it was used the Laplacian Smoothing Method [5] The new node position \hat{x}_i can be approximated by:



with Laplacian Smoothing

$$\hat{\mathbf{x}}_{\mathbf{i}} = w_{ij} \sum_{j}^{N_1} (\mathbf{x}_j - \mathbf{x}_i)$$

where, N_1 is the 1-ring neighbors of the node, w_{ij} is the weight and was calculated by the inverse distance from neighbors vertices



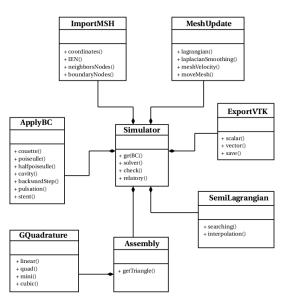
no Laplacian Smoothing



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Code Framework





Simplified Class Diagram

Computational Cost

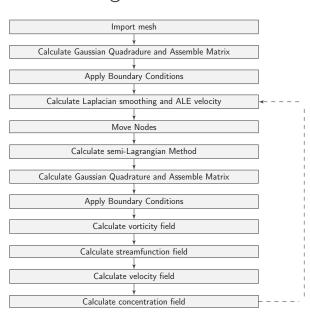


Process	AVG Computational Cost (%)
Mesh import	1.24
Assembly	73.87
BC Apply	6.70
Mesh update	10.05
Semi-Lagrangian	2.27
Vorticity Solver	5.51
VTK export	0.36

Average computational cost for the linear triangular element.

Solution Algorithm





Repeat the procedure for the next time step until the steady state



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Validation - Poiseuille Flow



Inflow condition: $v_x = u_{analytical}$, $v_y = 0$

Top plate: $v_x=0$, $v_y=0$, $\psi=1$

Bottom plate: $v_x=0$, $v_y=0$, $\psi=0$

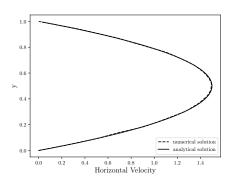
Vorticity condition: $\omega = \nabla \times \mathbf{v}$

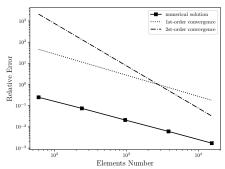
$$\beta_1 = 0 \qquad \beta_2 = 1$$



Nodes: 3835 Flements: 7299

Relative Error: 0.4%





Validation - Symmetric Poiseuille Flow



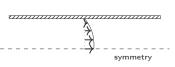
Inflow condition: $v_x = u_{analytical}, v_y = 0$

Top plate: $v_{\scriptscriptstyle X}=$ 0, $v_{\scriptscriptstyle V}=$ 0, $\psi=1$

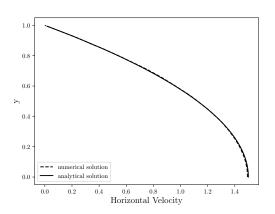
Symmetric axis: $v_y=0$, $\psi=0$

Vorticity condition: $\omega = \nabla \times \mathbf{v}$

$$\beta_1 = 0 \qquad \beta_2 = 1$$



Nodes: 3835 Elements: 7299



Validation - Lid Driven Cavity Flow

COPPE ,,UFRJ

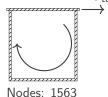
Bottom and side plates: $v_x=0$, $v_y=0$ e $\psi=0$

Top plate: $\emph{v}_{\it x}=$ 1, $\emph{v}_{\it y}=$ 0 e $\psi=$ 0

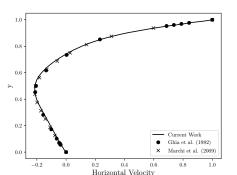
Vorticity condition: $\omega = \nabla \times \mathbf{v}$

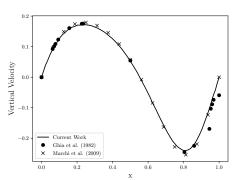
$$\beta_1 = 0 \qquad \beta_2 = 1$$

Re = 100



Nodes: 1563 Elements: 2988



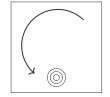


Validation - Pure Advection Flow



Assumptions

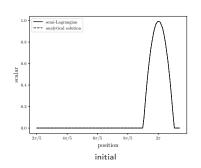
- 1. Continuum hypothesis
- 2. Homogeneous and Isotropic
- 3. Incompressible
- 4. Newtonian
- 5. Two-dimensional Flow
- 6. Single-phase Flow
- 7. High Reynolds number ($Re = \infty$)

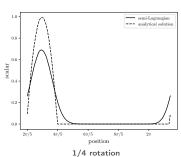


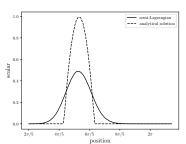
$$\frac{\partial \alpha}{\partial t} + \mathbf{v} \cdot \nabla \alpha = 0$$

Linear Triangular Element

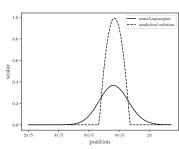






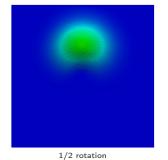


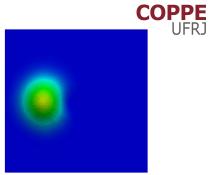
1/2 rotation



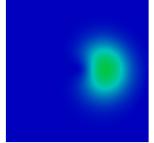
Linear Triangular Element

initial





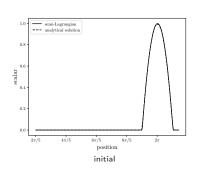
1/4 rotation

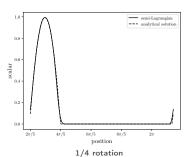


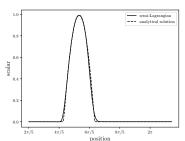
3/4 rotation

Quadratic Triangular Element

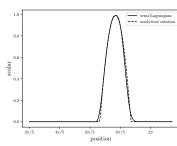




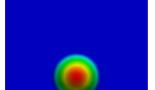


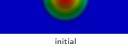


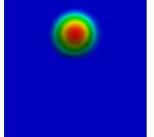
1/2 rotation



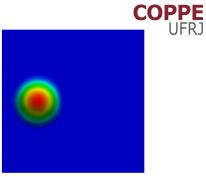
Quadratic Triangular Element



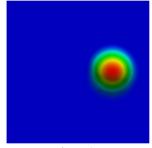




1/2 rotation 3/4 rotation



1/4 rotation



25/34



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Results





Non-dimensional symmetric geometry for blood flow in coronary artery with drug-eluting stent placed by Wang et al. (2017): (a) Curved Channel with Stent (b) Real Channel with Stent.

Inflow condition:
$$v_x = u_{analytical}$$
, $v_y = 0$;

Top plate:
$$v_x = 0$$
, $v_v = 0$, $\psi = 1$;

Symmetry condition:
$$v_y = 0$$
, $\psi = 0$;

Drug-eluting stent:
$$\emph{v}_{\emph{x}}=\emph{0}$$
, $\emph{v}_{\emph{y}}=\emph{0}$, $\psi=\emph{1}$, $\emph{e}=\emph{1}$

Vorticity condition:
$$\omega = \nabla \times \mathbf{v}$$

$$\beta_1 = 0 \qquad \beta_2 = 1$$

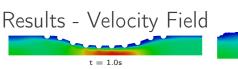
$$R = 0.0015 m$$

$$\mu = 0.0035 Pa.s$$

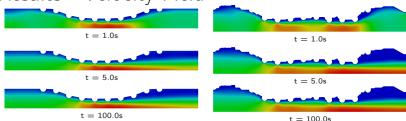
$$\rho=1060 kg/m^3$$

$$u = 12cm/s$$

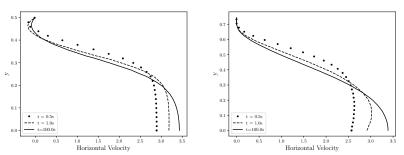
$$Re = 54.5$$







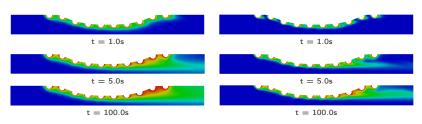
Evolution in time and space of velocity field: Curved Channel (left column) and Real Channel (right column)



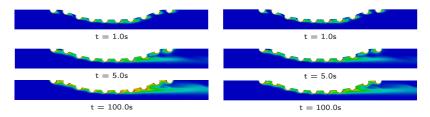
Evolution of velocity profile in centerline (x = 0.5L): (a) Curved Channel and (b) Real Channel

Results - Concentration Field





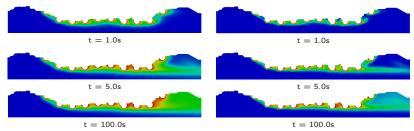
Evolution in time and space of concentration field in Curved Channel: Sc = 1 (left column) and Sc = 10 (right column)



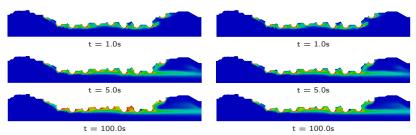
Evolution in time and space of concentration field in Curved Channel: Sc=100 (left column) and Sc=1000 (right column)

Results - Concentration Field





Evolution in time and space of concentration field in Real Channel: Sc = 1 (left column) and Sc = 10 (right column)



Evolution in time and space of concentration field in Real Channel: Sc=100 (left column) and Sc=1000 (right column)



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Conclusion



- The ALE description allows to perform simulations in the fluid-structure and moving boundary problems, in addition it is possible to assign several mesh velocity values in specific regions of the problem in order to improve the accuracy of the numerical results
- The searching procedure of the semi-Lagrangian Method used in this
 work was well designed since it corresponds to only 2% of average
 computational cost of the numerical simulation. Therefore, the methos
 showed to be very useful due to the unconditionnal stability and symmetric linear system
- 3. As expected, the chemical species transport in blood flow is directly influenced by Schmidt number, where the diffusion is increased as the Schmidt number decreases. It is possible that the density and viscosity of the blood are affected by chemical species diffusion. However, this influence is not considered in this work

Further developments:



- 1. Increase assembly performance
- 2. The use of primitive variables in the 3D Navier-Stokes equation
- 3. Blood flow model as a multiphase problem
- 4. Blood flow model as a non-Newtonian fluid



Thank you!

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