

Lecture 8: Stability, Padé approximations

- Stability of Runga-Kutta methods
 - **A-stability**
 - Aliasing
 - **L-stability**
 - Padé approximations
 - (Nonlinear stability: AN-, B- and algebraic stability)

Book: 14.6

Explicit Runge-Kutta (ERK) methods

- IVP: $\dot{y} = f(y, t), \quad y(0) = y_0$
- ERK:

$$k_1 = f(y_n, t_n)$$

$$k_2 = f(y_n + ha_{21}k_1, t_n + c_2h)$$

$$k_3 = f(y_n + h(a_{31}k_1 + a_{32}k_2), t_n + c_3h)$$

$$\vdots$$

$$k_\sigma = f(y_n + h(a_{\sigma,1}k_1 + a_{\sigma,2}k_2 + \dots + a_{\sigma,\sigma-1}k_{\sigma-1}), t_n + c_\sigma h)$$

$$y_{n+1} = y_n + h(b_1k_1 + b_2k_2 + \dots + b_\sigma k_\sigma)$$
- Butcher array:

c	A				
	b^T				
0					
c ₂	a ₂₁				
c ₃	a ₃₁	a ₃₂			
⋮	⋮	⋮	⋱		
c _σ	a _{σ,1}	a _{σ,2}	⋯	a _{σ,σ-1}	
	b ₁	b ₂	⋯	b _{σ-1}	b _σ

Recap: Test system, stability function

- One step method (typically: Runge-Kutta):

$$y_{n+1} = y_n + h\phi(y_n, t_n)$$

- Apply it to scalar test system:

$$\dot{y} = \lambda y$$

- We get:

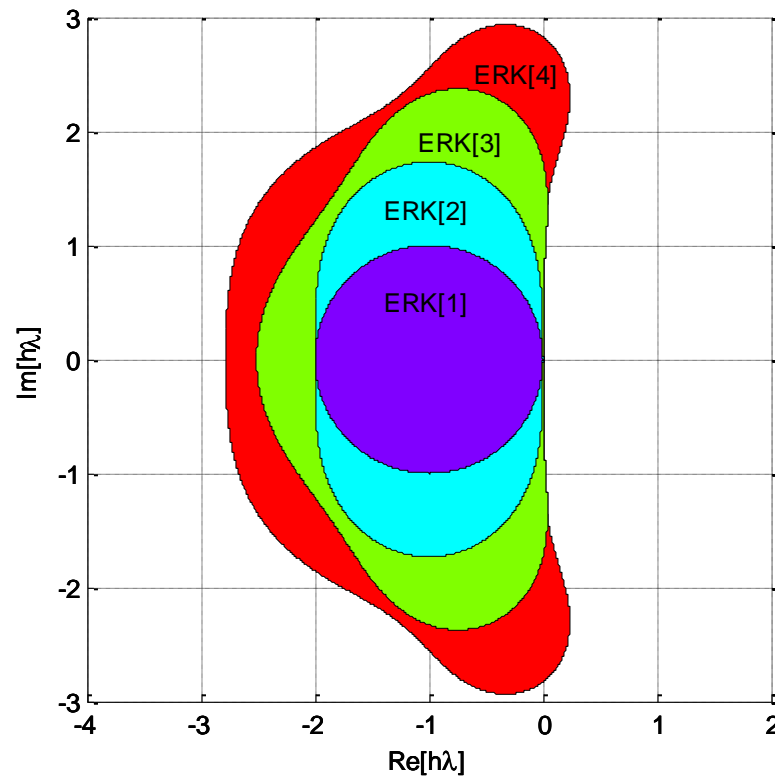
$$y_{n+1} = R(h\lambda)y_n$$

where $R(h\lambda)$ is stability function

- The method is stable (for test system!) if

$$|R(h\lambda)| \leq 1$$

Stability regions for ERK methods



Implicit Runge-Kutta (IRK) methods

- IVP: $\dot{y} = f(y, t), \quad y(0) = y_0$
- IRK:

$$k_1 = f(y_n + h(a_{1,1}k_1 + a_{1,2}k_2 + \dots + a_{1,\sigma}k_\sigma), t_n + c_1h)$$

$$k_2 = f(y_n + h(a_{2,1}k_1 + a_{2,2}k_2 + \dots + a_{2,\sigma}k_\sigma), t_n + c_2h)$$

$$k_3 = f(y_n + h(a_{3,1}k_1 + a_{3,2}k_2 + \dots + a_{3,\sigma}k_\sigma), t_n + c_3h)$$

$$\vdots$$

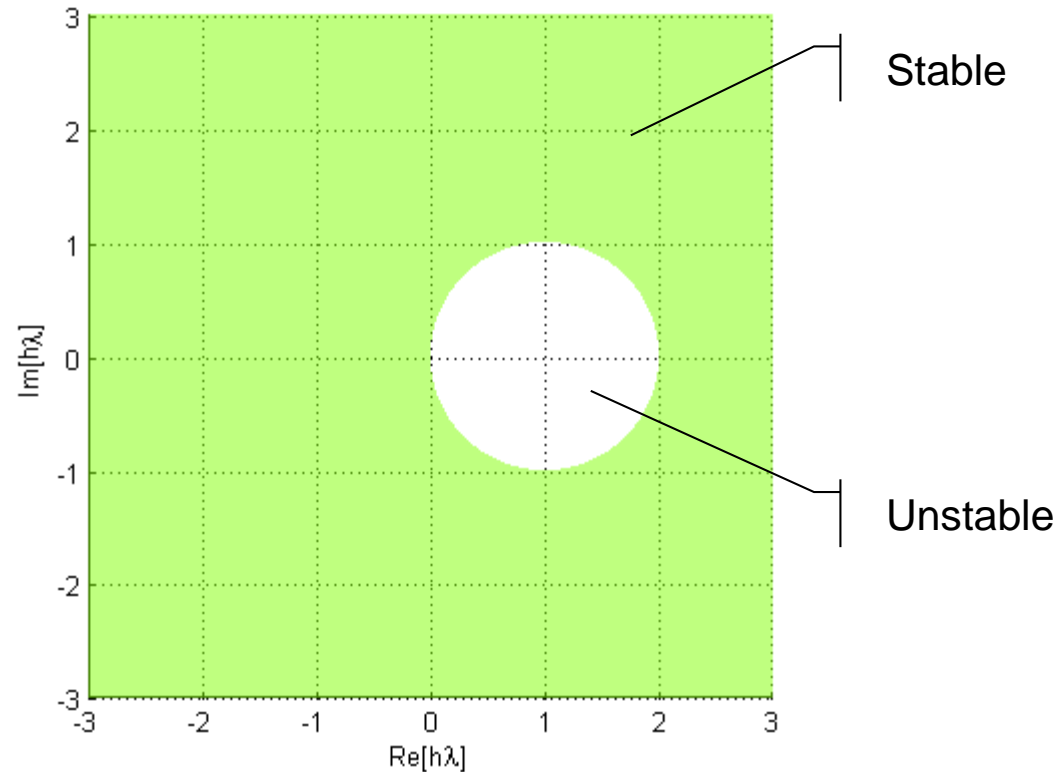
$$k_\sigma = f(y_n + h(a_{\sigma,1}k_1 + a_{\sigma,2}k_2 + \dots + a_{\sigma,\sigma}k_\sigma), t_n + c_\sigma h)$$

$$y_{n+1} = y_n + h(b_1k_1 + b_2k_2 + \dots + b_\sigma k_\sigma)$$
- Butcher array:

$\begin{array}{c c} \mathbf{c} & \mathbf{A} \\ \hline & \mathbf{b}^\top \end{array}$	c_1	a_{11}	a_{12}	\cdots	$a_{1,\sigma-1}$	$a_{1,\sigma}$
	c_2	a_{21}	a_{22}	\cdots	$a_{2,\sigma-1}$	$a_{2,\sigma}$
	\vdots	\vdots	\vdots	\ddots	\vdots	\vdots
	c_σ	$a_{\sigma,1}$	$a_{\sigma,2}$	\cdots	$a_{\sigma,\sigma-1}$	$a_{\sigma,\sigma}$
		b_1	b_2	\dots	$b_{\sigma-1}$	b_σ

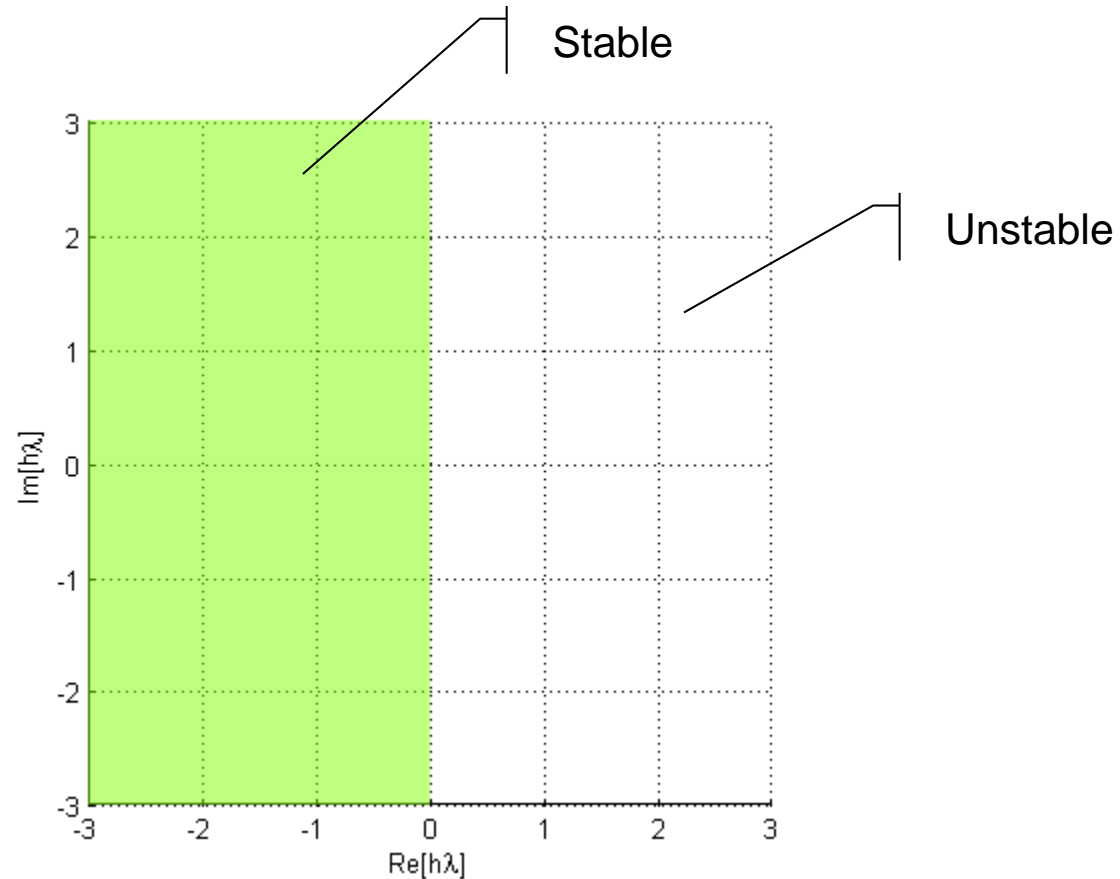
Stability regions for implicit methods

- Implicit Euler



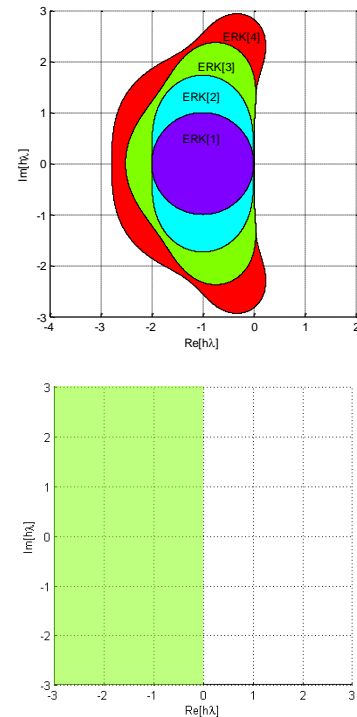
Stability regions for implicit methods

- Trapezoidal rule and implicit midpoint rule:



Why use IRK methods?

- IRK methods are much more complex (since we have to solve a set of nonlinear equations for each step) than ERK methods, so why use them?
 - Accuracy? Stability?
- Not because of accuracy
 - Even if an IRK method may have higher accuracy for a given number of stages, it is easy and cheap to achieve same accuracy for ERK by increasing the number of stages
- It's because of the much larger stability region!
- When is this important?
 - **Stiff systems!**



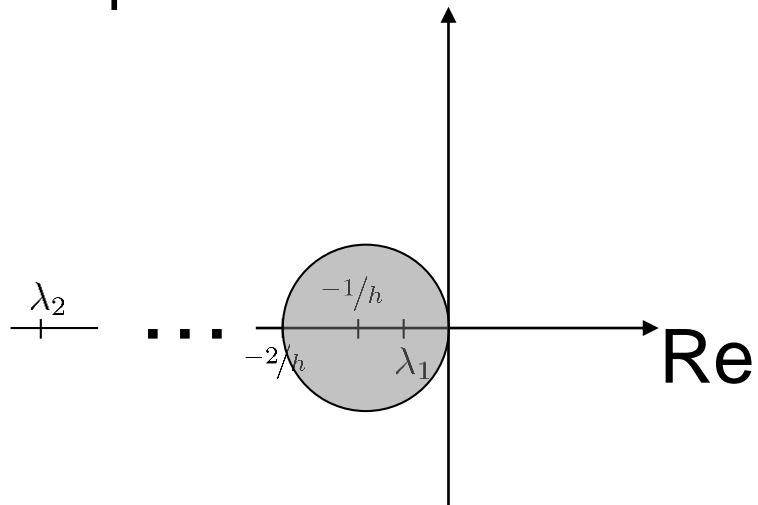
Stiff system

$$\dot{y}_1 = \lambda_1 y_1$$

$$\dot{y}_2 = \lambda_2 y_2$$

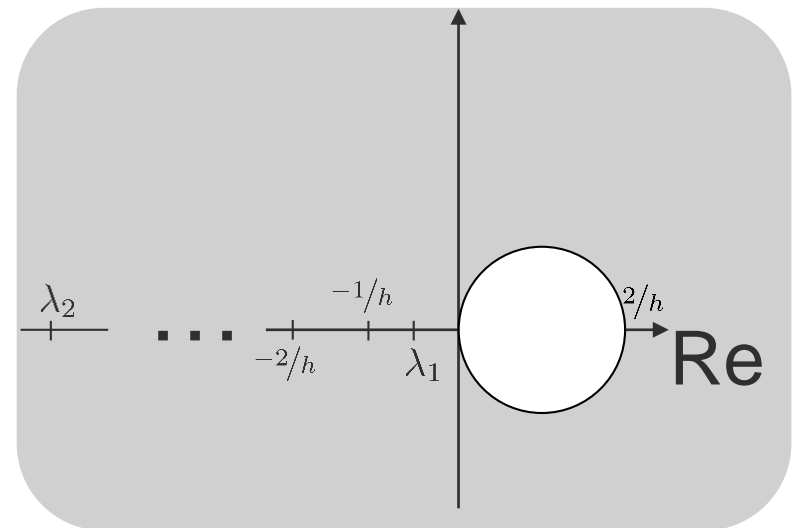
$$\lambda_2 \ll \lambda_1 < 0$$

Explicit Euler Im



must choose a small h for stability

Implicit Euler Im

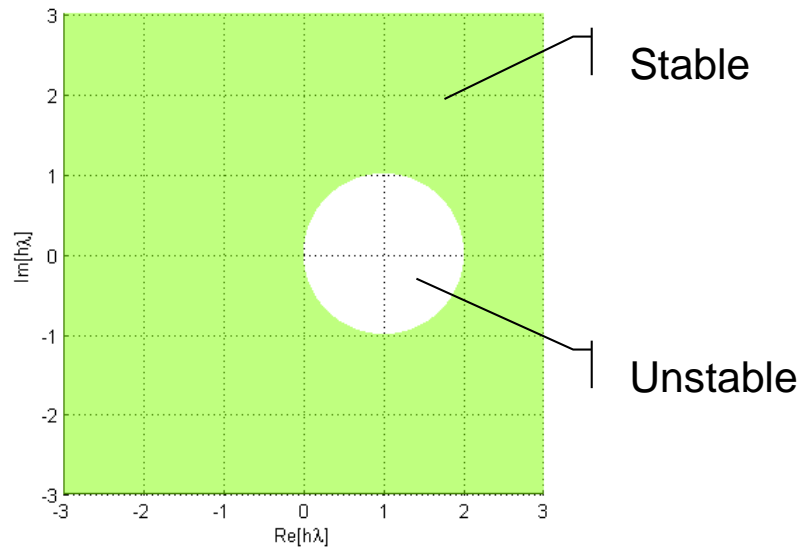


can choose all h for stability

Example 221

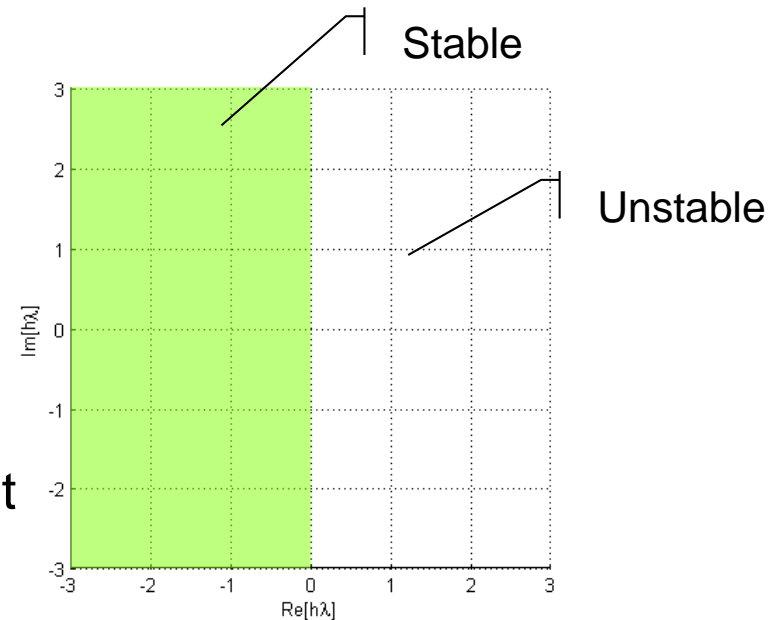
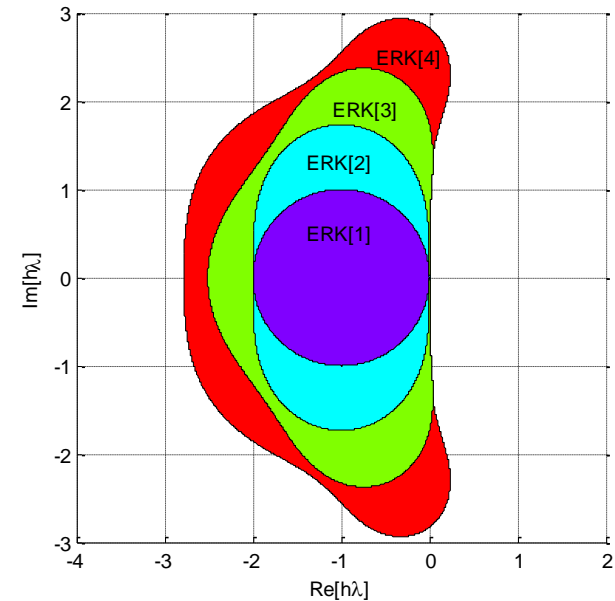
ERK vs IRK stability regions

Implicit Euler



Trapezoidal/
Implicit Midpoint

ERK-methods



Aliasing I

Aliasing II

- Assume:

Test system I

$$\dot{y} = \lambda y ; \quad \lambda = \sigma + j\omega$$

$$y_L(t_n; t_{n+1}) = e^{h\lambda} y_n$$

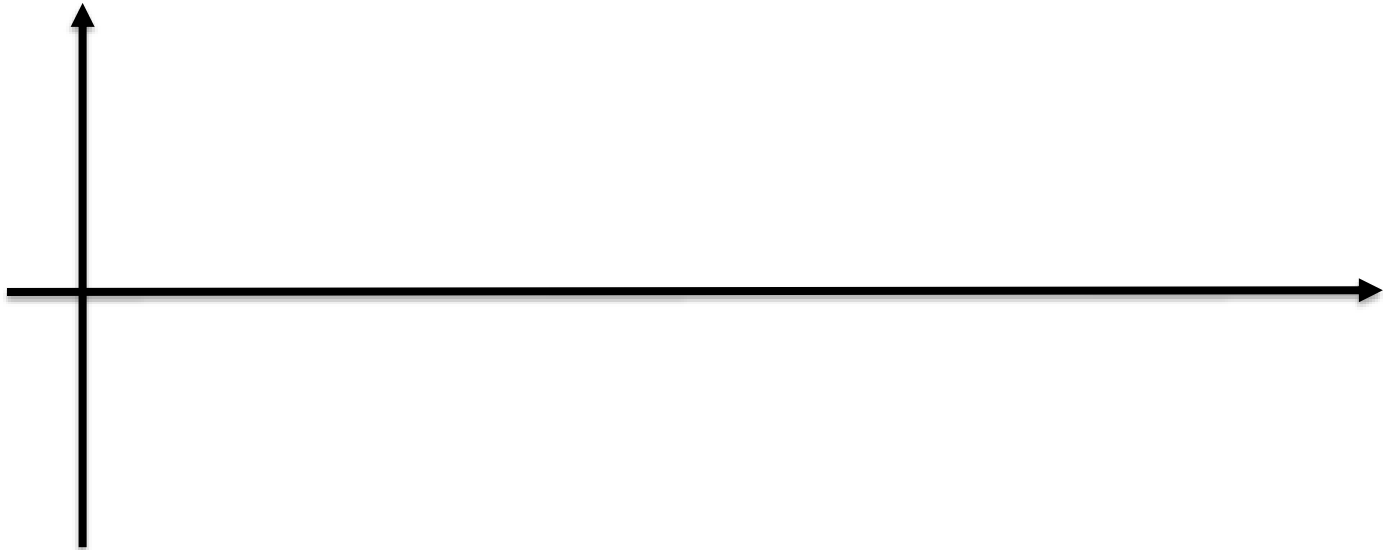
Test system II

$$\dot{y} = \mu y$$

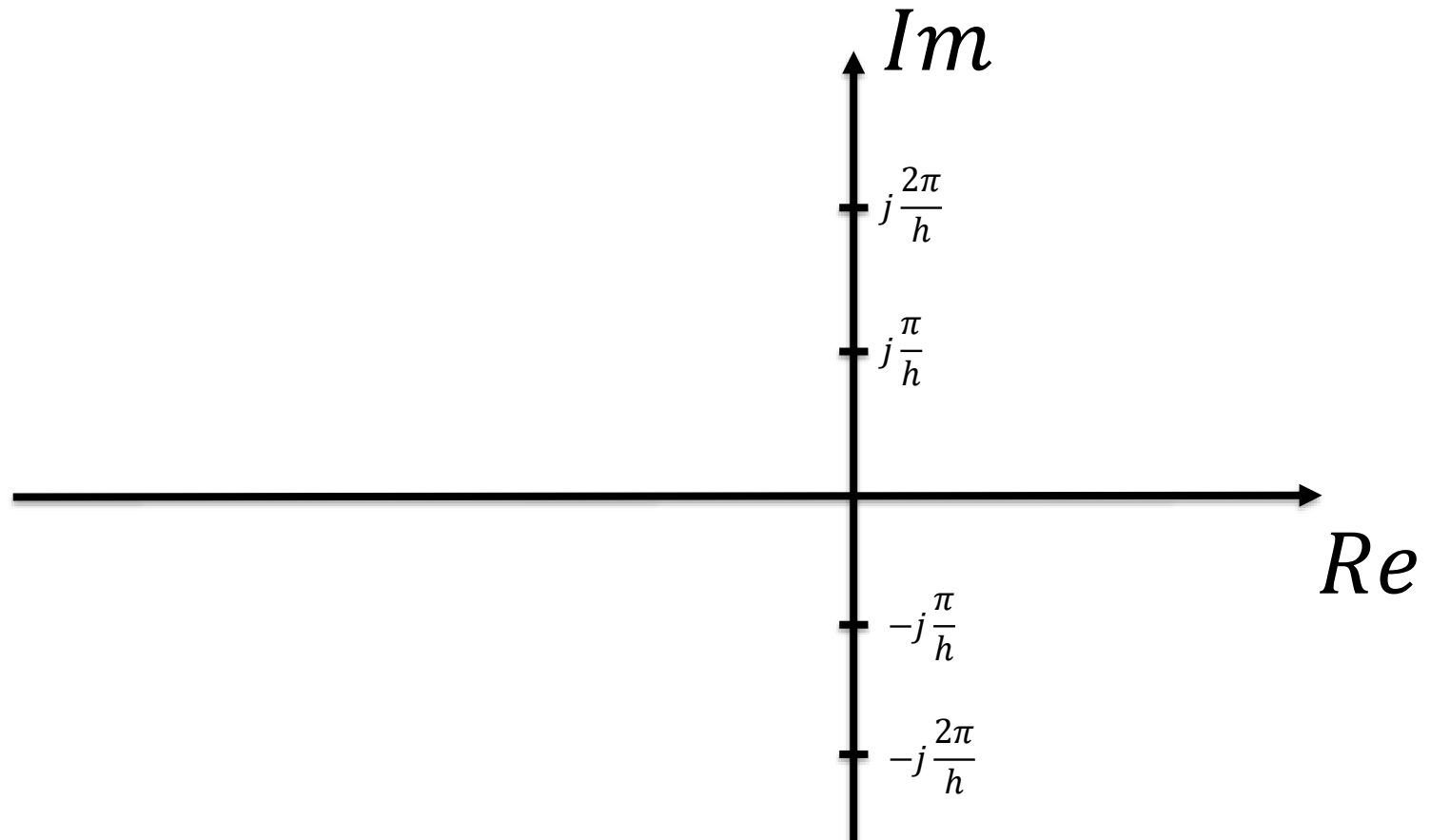
$$y_L(t_n; t_{n+1}) = e^{h\mu} y_n$$

When are these two systems the same?

Visualisation I

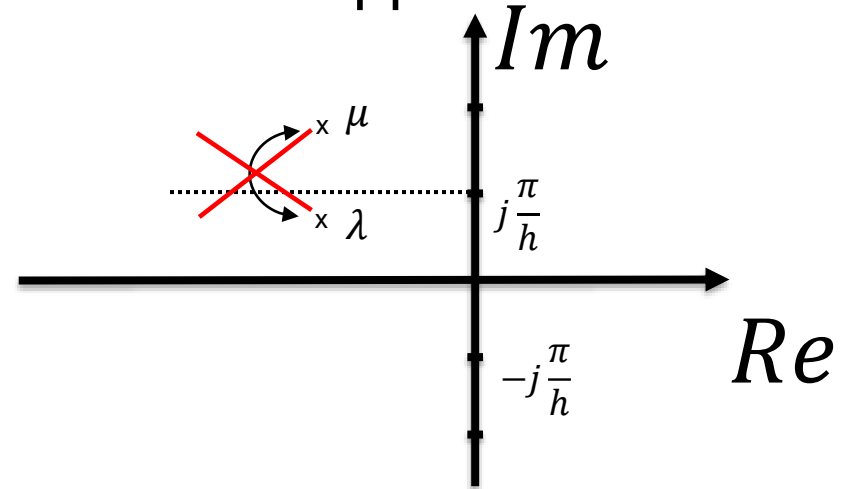


Visualisation II



L-stability I

- Assume μ eigenvalues and we want to suppress:



L-stability II

Example – L-stability

Padé-approximation I

Test system: $\dot{y} = \lambda y$

exact

$$y_L(t_n; t_{n+1}) = e^{h\lambda} y_n$$

RK

$$y_{n+1} = R(h\lambda) y_n$$

«good» methods fulfill $e^s \approx R(s)$ $s = h\lambda$

Padé-approximation II

Def: Padé approximation $P_m^k(s)$ of e^s is the rational function:

$$P_m^k(s) = \frac{1 + \beta_1 s + \cdots + \beta_k s^k}{1 + \gamma_1 s + \cdots + \gamma_m s^m},$$

which approximates e^s the best given k and m

Padé approximations to e^s

$\begin{smallmatrix} k \\ m \end{smallmatrix}$	0	1	2	3
0	$\frac{1}{1}$	$\frac{1+s}{1}$	$\frac{1+s+\frac{1}{2}s^2}{1}$	$\frac{1+s+\frac{1}{2}s^2+\frac{1}{6}s^3}{1}$
1	$\frac{1}{1-s}$	$\frac{1+\frac{1}{2}s}{1-\frac{1}{2}s}$	$\frac{1+\frac{2}{3}s+\frac{1}{6}s^2}{1-\frac{1}{3}s}$	$\frac{1+\frac{3}{4}s+\frac{1}{4}s^2+\frac{1}{24}s^3}{1-\frac{1}{4}s}$
2	$\frac{1}{1-s+\frac{1}{2}s^2}$	$\frac{1+\frac{1}{3}s}{1-\frac{2}{3}s+\frac{1}{6}s^2}$	$\frac{1+\frac{1}{2}s+\frac{1}{12}s^2}{1-\frac{1}{2}s+\frac{1}{12}s^2}$	$\frac{1+\frac{3}{5}s+\frac{3}{20}s^2+\frac{1}{60}s^3}{1-\frac{2}{5}s+\frac{1}{20}s^2}$
3	$\frac{1}{1-s+\frac{1}{2}s^2-\frac{1}{6}s^3}$	$\frac{1+\frac{1}{4}s}{1-\frac{3}{4}s+\frac{1}{4}s^2-\frac{1}{24}s^3}$	$\frac{1+\frac{2}{5}s+\frac{1}{20}s^2}{1-\frac{3}{5}s+\frac{3}{20}s^2-\frac{1}{60}s^3}$	$\frac{1+\frac{1}{2}s+\frac{1}{10}s^2+\frac{1}{120}s^3}{1-\frac{1}{2}s+\frac{1}{10}s^2-\frac{1}{120}s^3}$

- $m = 0$: Explicit Runge-Kutta methods with $p = \sigma$
- $m = k$: Gauss, Lobatto IIIA/IIIB (incl. implicit mid-point, trapezoidal)
- $m = k+1$: Radau-methods (incl. implicit Euler)
- $m = k+2$: Lobatto IIIC

Padé approximations to e^s

$\begin{smallmatrix} k \\ m \end{smallmatrix}$	0	1	2	3
0	$\frac{1}{1}$	$\frac{1+s}{1}$	$\frac{1+s+\frac{1}{2}s^2}{1}$	$\frac{1+s+\frac{1}{2}s^2+\frac{1}{6}s^3}{1}$
1	$\frac{1}{1-s}$	$\frac{1+\frac{1}{2}s}{1-\frac{1}{2}s}$	$\frac{1+\frac{2}{3}s+\frac{1}{6}s^2}{1-\frac{1}{3}s}$	$\frac{1+\frac{3}{4}s+\frac{1}{4}s^2+\frac{1}{24}s^3}{1-\frac{1}{4}s}$
2	$\frac{1}{1-s+\frac{1}{2}s^2}$	$\frac{1+\frac{1}{3}s}{1-\frac{2}{3}s+\frac{1}{6}s^2}$	$\frac{1+\frac{1}{2}s+\frac{1}{12}s^2}{1-\frac{1}{2}s+\frac{1}{12}s^2}$	$\frac{1+\frac{3}{5}s+\frac{3}{20}s^2+\frac{1}{60}s^3}{1-\frac{2}{5}s+\frac{1}{20}s^2}$
3	$\frac{1}{1-s+\frac{1}{2}s^2-\frac{1}{6}s^3}$	$\frac{1+\frac{1}{4}s}{1-\frac{3}{4}s+\frac{1}{4}s^2-\frac{1}{24}s^3}$	$\frac{1+\frac{2}{5}s+\frac{1}{20}s^2}{1-\frac{3}{5}s+\frac{3}{20}s^2-\frac{1}{60}s^3}$	$\frac{1+\frac{1}{2}s+\frac{1}{10}s^2+\frac{1}{120}s^3}{1-\frac{1}{2}s+\frac{1}{10}s^2-\frac{1}{120}s^3}$



L-stable



L-stable



A-stable

- $m = 0$: Explicit Runge-Kutta methods with $p = \sigma$
- $m = k$: Gauss, Lobatto IIIA/IIIB (incl. implicit mid-point, trapezoidal)
- $m = k+1$: Radau-methods (incl. implicit Euler)
- $m = k+2$: Lobatto IIIC

Padé approximations as stability function

1. Assume $k \leq m \leq k + 2$

In that case $|P_m^k(s)| \leq 1$ if $\operatorname{Re}[s] \leq 0$

2. $|P_m^m(j\omega)| = 1$ if $\omega \rightarrow \infty$

3. Assume $m > k$

In that case $|P_m^k(s)| \rightarrow 0$ if $\omega \rightarrow \infty$

One step methods with stability function $R(s) = P_m^k(s)$ are:

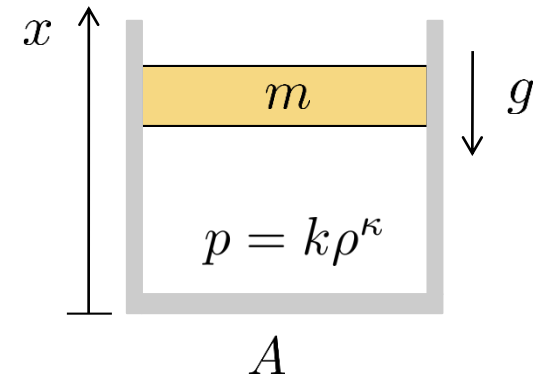
- A-stable if $k = m$
- L-stable if $m = k + 1$ or $m = k + 2$

Pneumatic spring example, again

- Model from Newton's 2nd law:

$$\ddot{x} + g(1 - x^{-\kappa}) = 0$$

"mass-spring-damper with nonlinear spring"



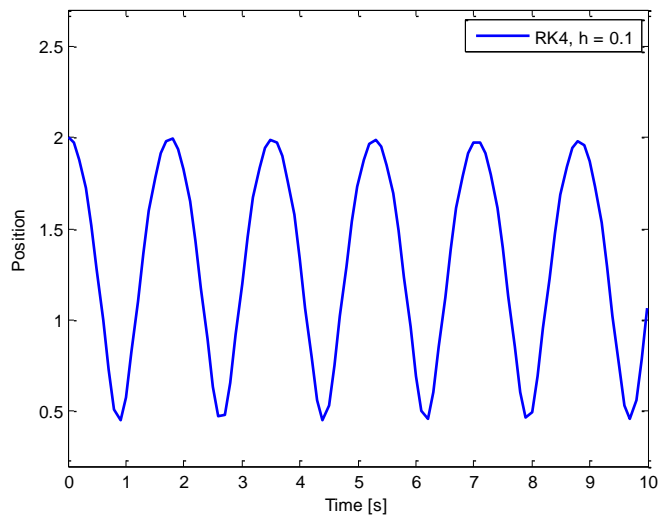
- On state-space form $\dot{y} = f(y, t)$

$$\begin{pmatrix} \dot{y}_1 \\ \dot{y}_2 \end{pmatrix} = \begin{pmatrix} y_2 \\ -g(1 - y_1^{-\kappa}) \end{pmatrix}$$

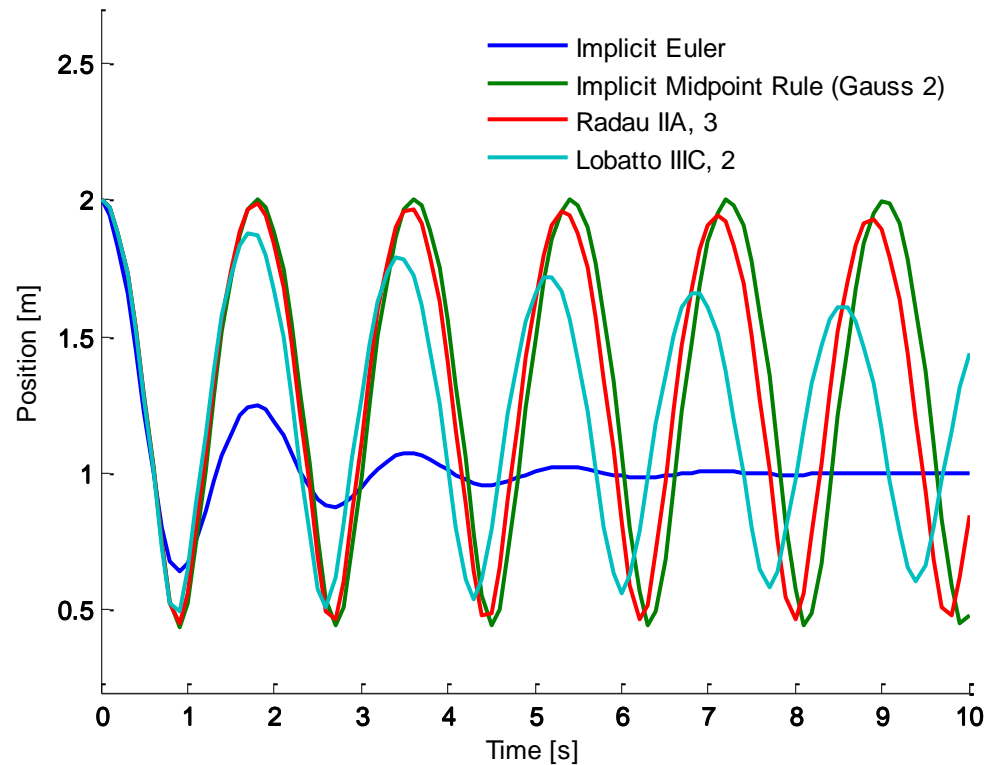
- Linearization about equilibrium:

$$\frac{\partial f}{\partial y} = \begin{pmatrix} 0 & 1 \\ -g\kappa & 0 \end{pmatrix}, \quad \lambda_{1,2} = \pm j\omega_0, \quad \omega_0 = \sqrt{g\kappa} \approx 3.7$$

Simulation



$h = 0.5$ (stability limit for RK4)



Padé approximations to e^s

$\begin{smallmatrix} k \\ m \end{smallmatrix}$	0	1	2	3
0	$\frac{1}{1}$	$\frac{1+s}{1}$	$\frac{1+s+\frac{1}{2}s^2}{1}$	$\frac{1+s+\frac{1}{2}s^2+\frac{1}{6}s^3}{1}$
1	$\frac{1}{1-s}$	$\frac{1+\frac{1}{2}s}{1-\frac{1}{2}s}$	$\frac{1+\frac{2}{3}s+\frac{1}{6}s^2}{1-\frac{1}{3}s}$	$\frac{1+\frac{3}{4}s+\frac{1}{4}s^2+\frac{1}{24}s^3}{1-\frac{1}{4}s}$
2	$\frac{1}{1-s+\frac{1}{2}s^2}$	$\frac{1+\frac{1}{3}s}{1-\frac{2}{3}s+\frac{1}{6}s^2}$	$\frac{1+\frac{1}{2}s+\frac{1}{12}s^2}{1-\frac{1}{2}s+\frac{1}{12}s^2}$	$\frac{1+\frac{3}{5}s+\frac{3}{20}s^2+\frac{1}{60}s^3}{1-\frac{2}{5}s+\frac{1}{20}s^2}$
3	$\frac{1}{1-s+\frac{1}{2}s^2-\frac{1}{6}s^3}$	$\frac{1+\frac{1}{4}s}{1-\frac{3}{4}s+\frac{1}{4}s^2-\frac{1}{24}s^3}$	$\frac{1+\frac{2}{5}s+\frac{1}{20}s^2}{1-\frac{3}{5}s+\frac{3}{20}s^2-\frac{1}{60}s^3}$	$\frac{1+\frac{1}{2}s+\frac{1}{10}s^2+\frac{1}{120}s^3}{1-\frac{1}{2}s+\frac{1}{10}s^2-\frac{1}{120}s^3}$



L-stable



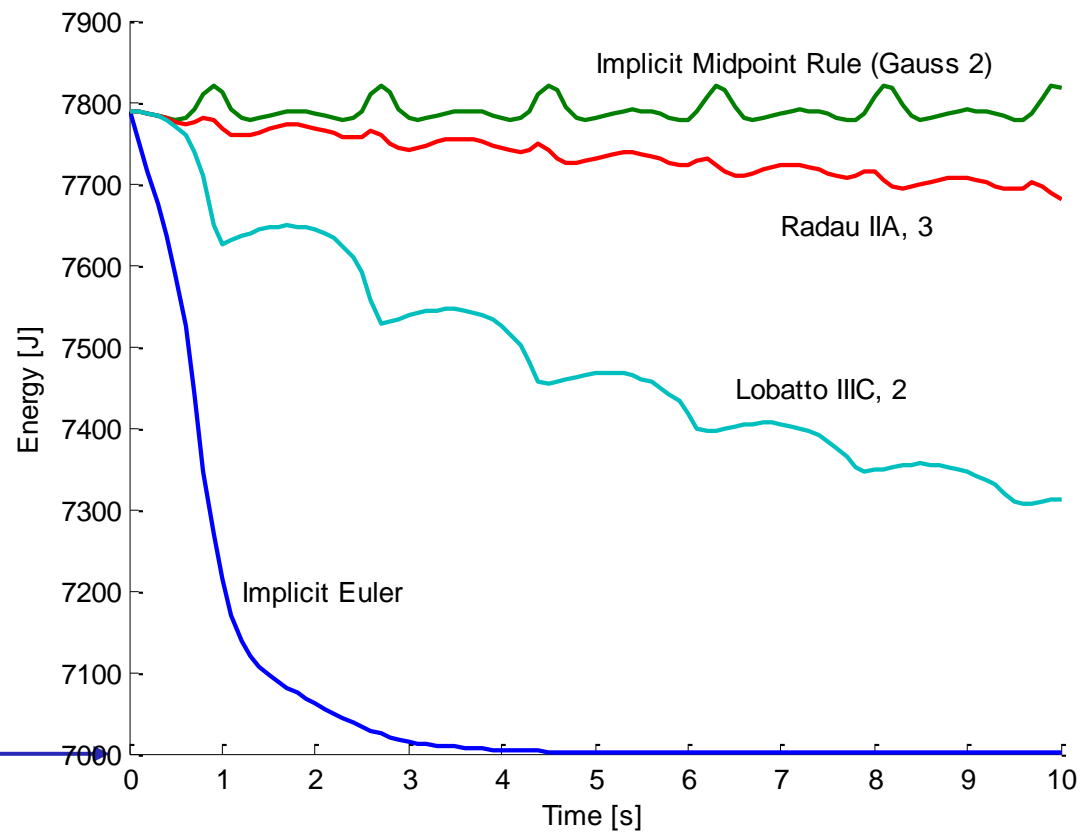
L-stable



A-stable

- $m = 0$: Explicit Runge-Kutta methods with $p = \sigma$
- $m = k$: Gauss, Lobatto IIIA/IIIB (incl. implicit mid-point, trapezoidal)
- $m = k+1$: Radau-methods (incl. implicit Euler)
- $m = k+2$: Lobatto IIIC

Energy



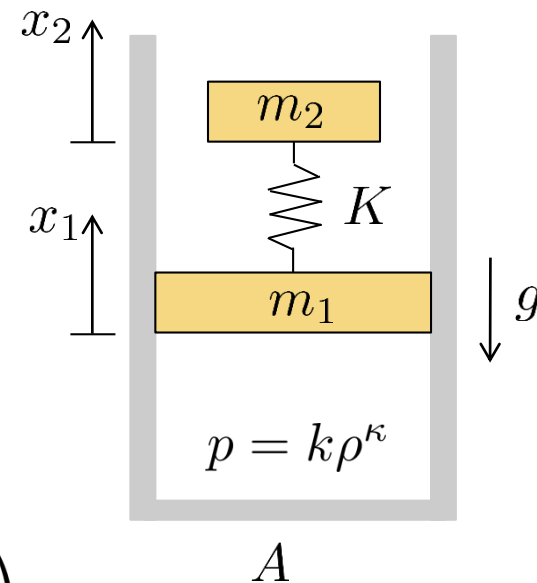
Equilibrium energy

Pneumatic spring with resonant load

- Equations of motion (Newton's law):

$$\ddot{x}_1 + g \left(1 - \frac{m_1 + m_2}{m_1} x_1^{-\kappa} \right) + \frac{\omega_2^2}{2} (x_1 - x_2) = 0$$

$$\ddot{x}_2 + g + \frac{\omega_2^2}{2} (x_2 - x_1) = 0$$



- Linearization around equilibrium:

$$J = \begin{pmatrix} 0 & 1 & 0 & 0 \\ -g \frac{m_1 + m_2}{m_1} \kappa (x_1^*)^{-(\kappa-1)} - \frac{\omega_2^2}{2} & 0 & \frac{\omega_2^2}{2} & 0 \\ 0 & 0 & 0 & 1 \\ \frac{\omega_2^2}{2} & 0 & -\frac{\omega_2^2}{2} & 0 \end{pmatrix}$$

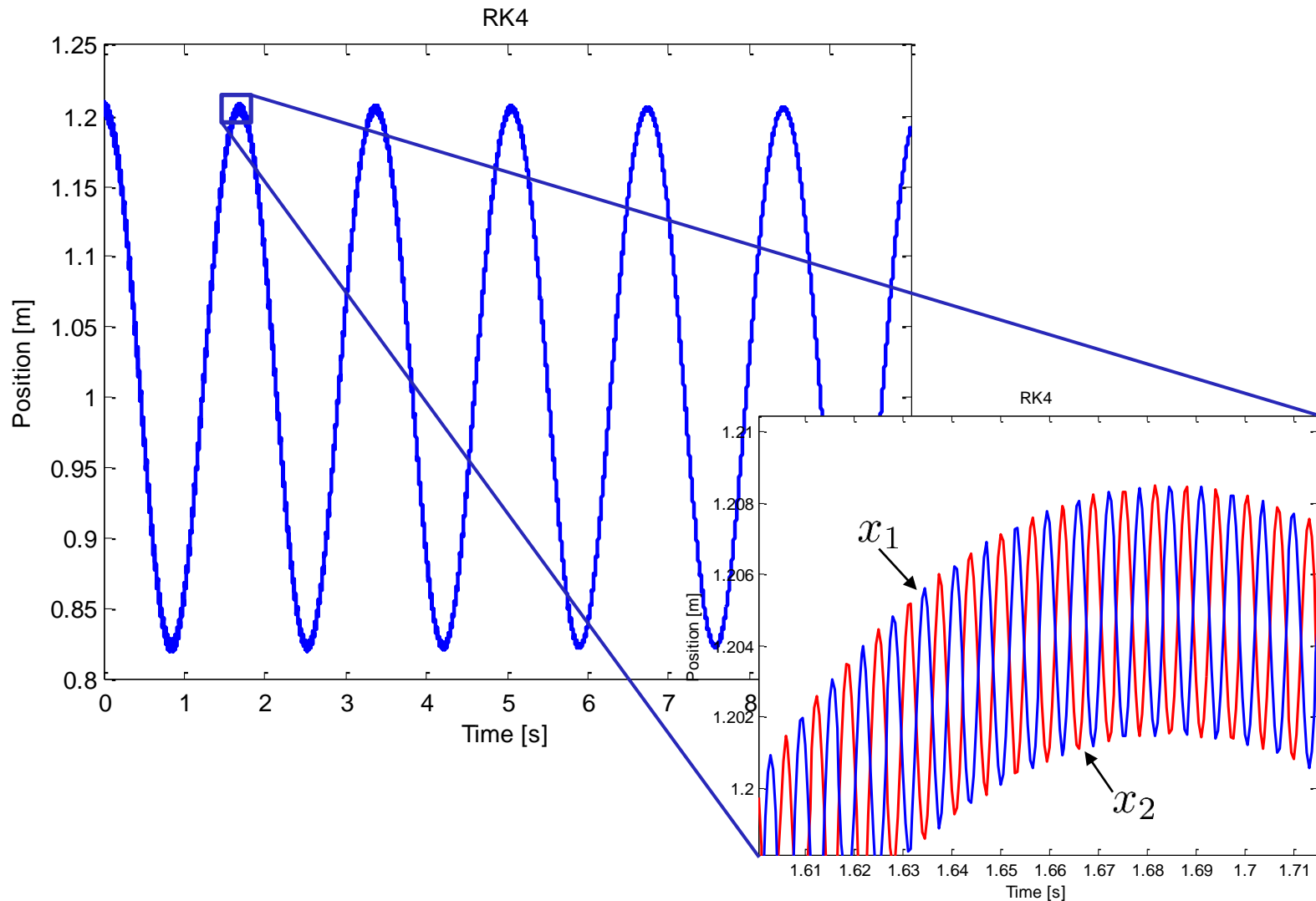
- Eigenvalues:

$$\lambda_{1,2} = \pm j\omega_1, \quad \omega_1 = 3.7 \text{ rad/s}$$

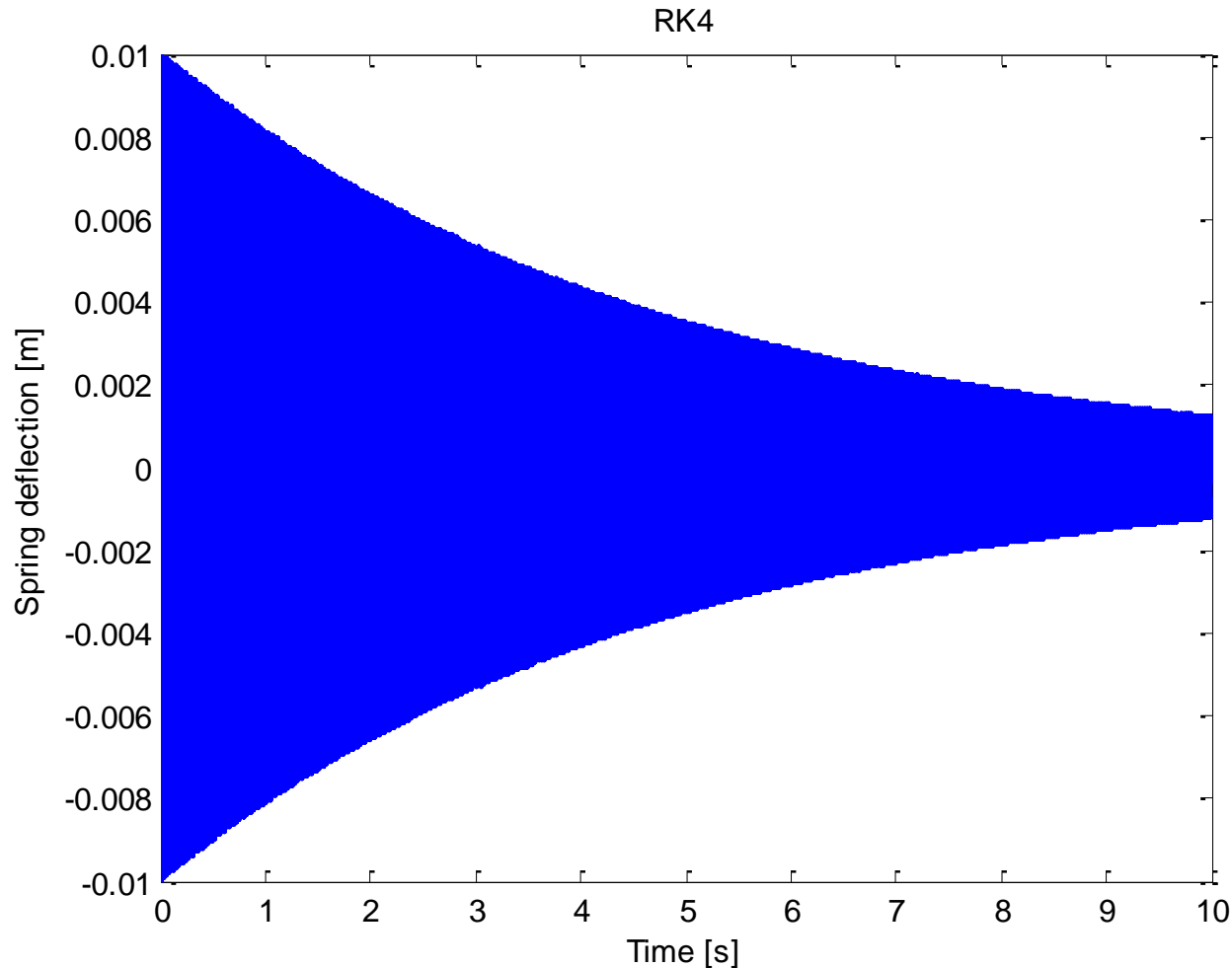
$$\lambda_{3,4} = \pm j\omega_2, \quad \omega_2 = 1000 \text{ rad/s}$$

Position of the two masses

RK4 with time step $h = 0.0005$

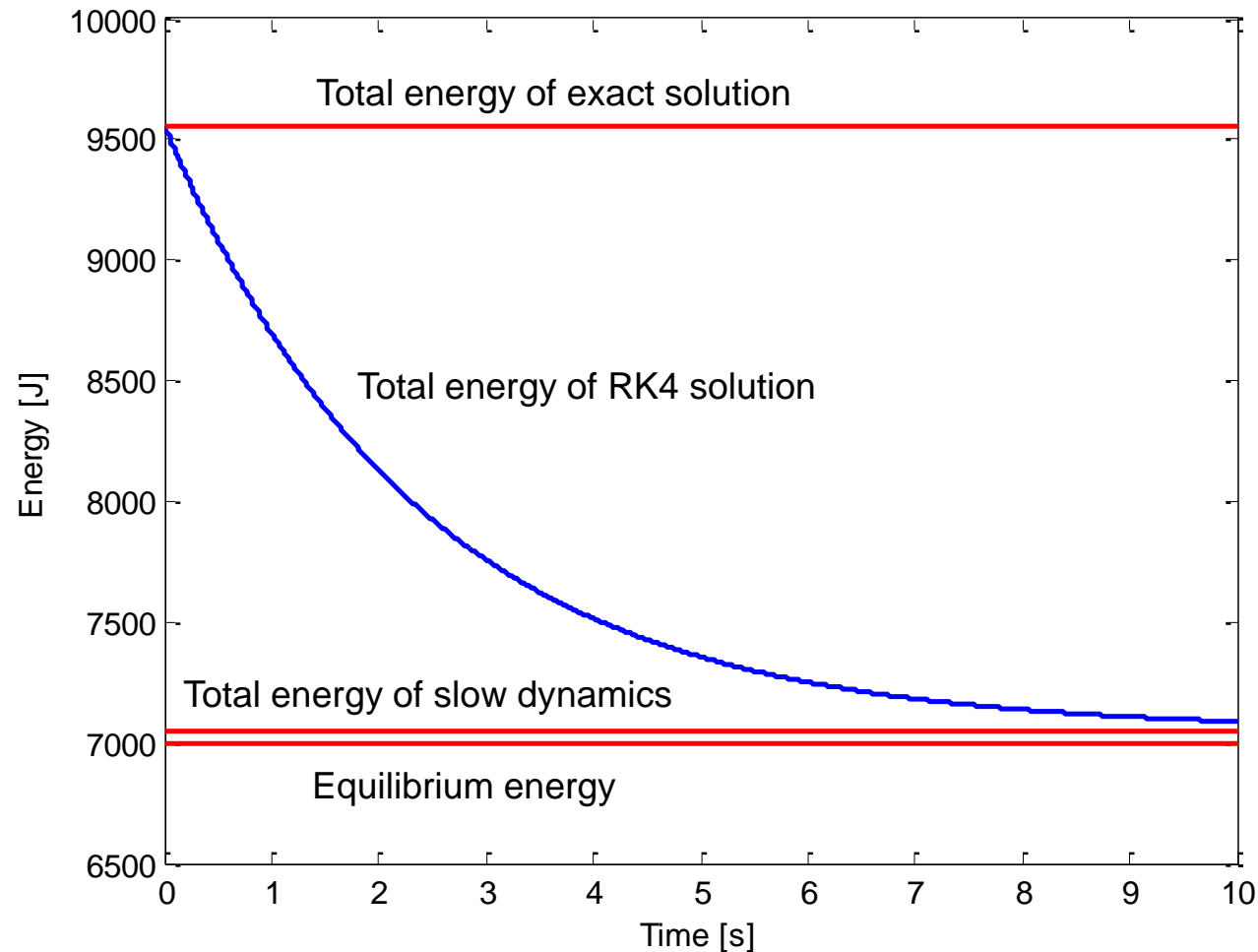


Spring deflection, RK4 with $h = 0.0005$



- Oscillation is lightly damped by integration method

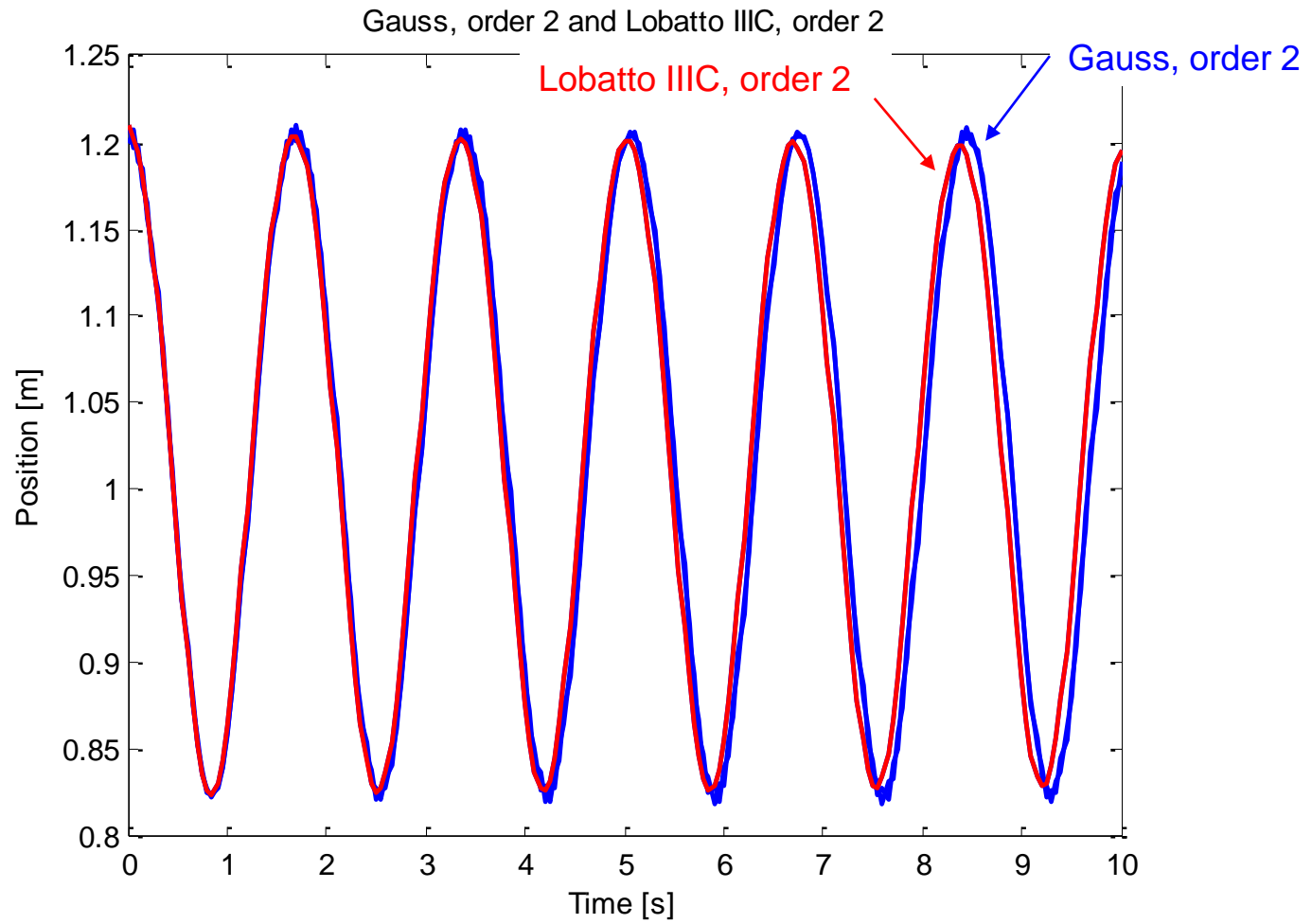
Energy of RK4 solution, $h = 0.0005$



- Energy related to fast dynamics slowly damped out

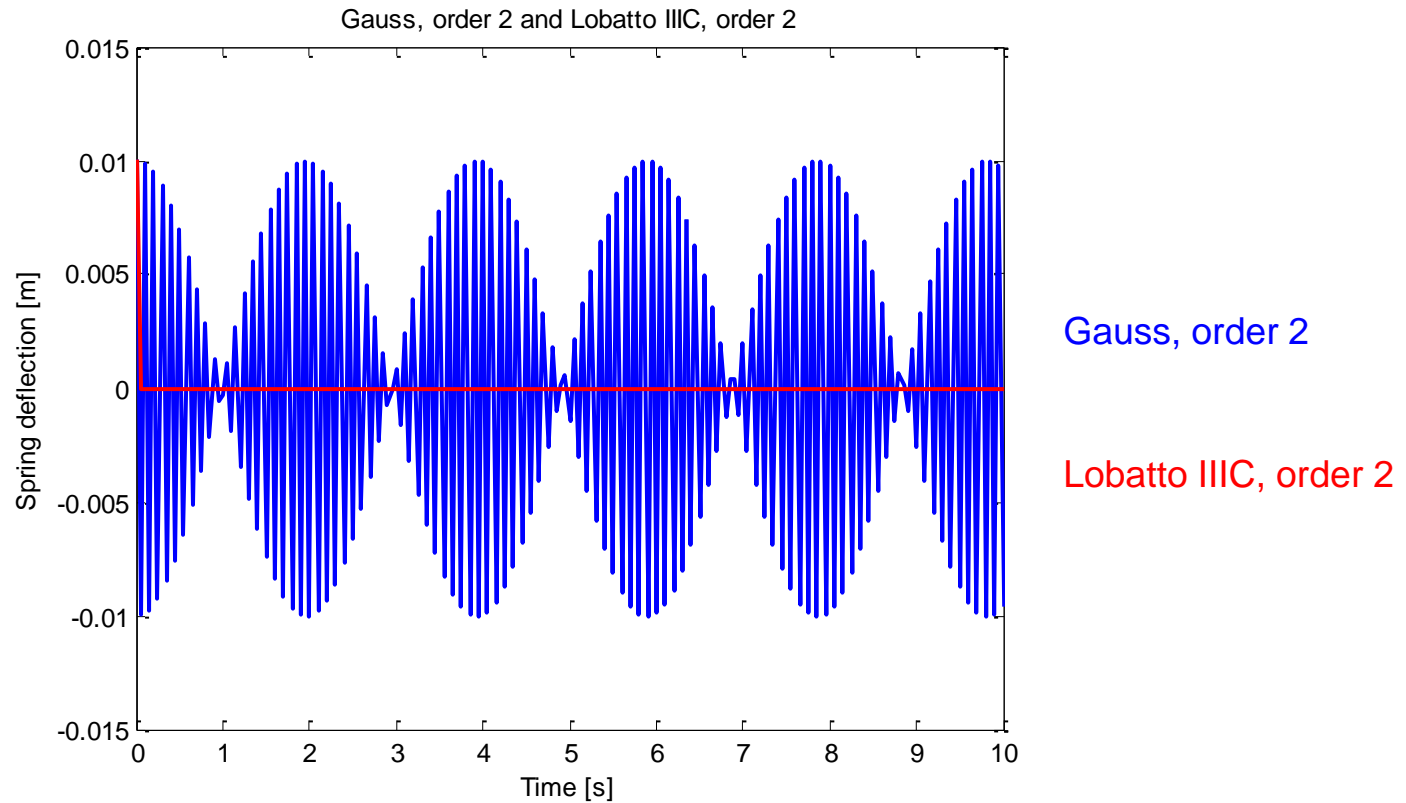
Position of the two masses

Gauss, order 2 and Lobatto IIIC, order 2, $h = 0.05$



Spring deflection

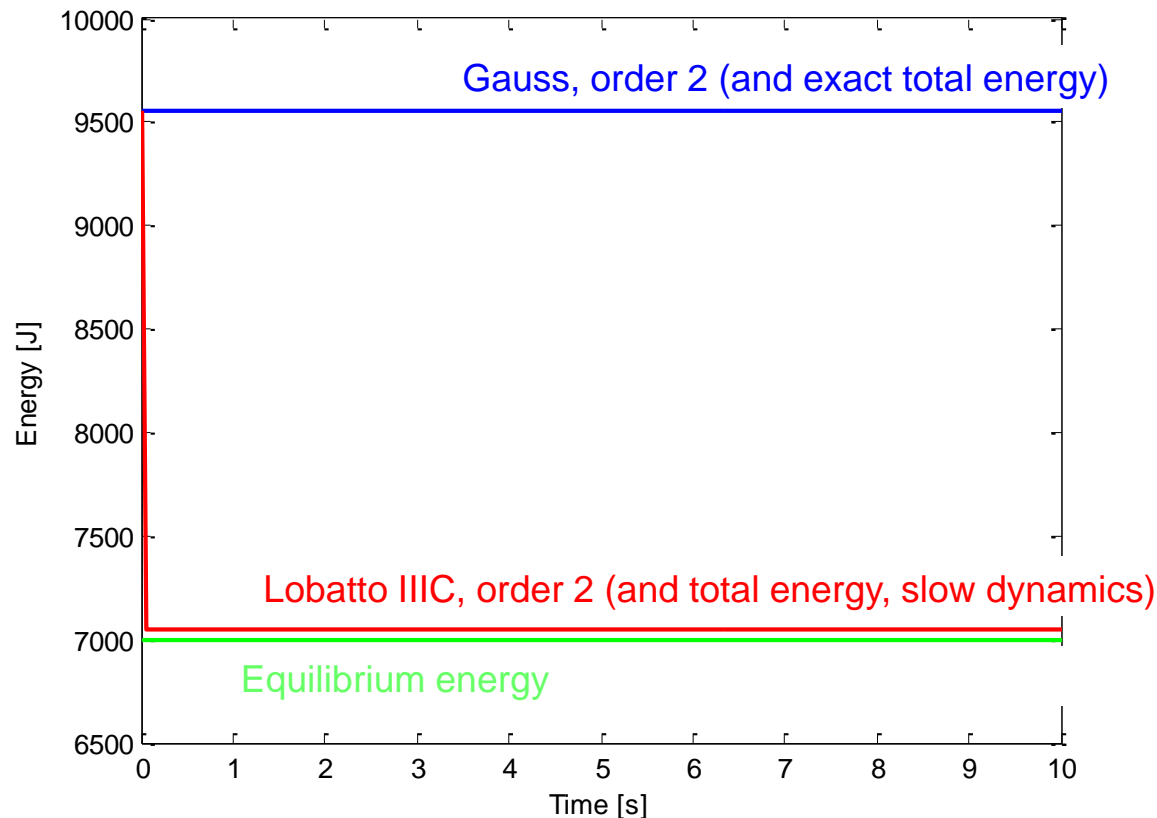
Gauss, order 2 and Lobatto IIIC, order 2, $h = 0.05$



- Gauss method gives no damping, but shifts fast dynamics and energy to frequencies below Nyquist frequency, $\omega_N = \frac{\pi}{h} = \frac{\pi}{0.05} = 62.8$
- Lobatto IIIC dampens out fast dynamics in one step

Total energies

Gauss, order 2 and Lobatto IIIC, order 2, $h = 0.05$



- Gauss does not dampen energies at all (same as exact total energy)
- Lobatto IIIC dampens out energy associated with fast dynamics in very few steps, to the energy of slow dynamics

Modelica

- Replacable / redeclare
- Choices
- Check out :
 - <http://book.xogeny.com/components/architectures/replaceable/>

Homework

- Check out the Modelica Circuit example (uploaded on Blackboard).
 - Look at the structure replaceable / redeclare
 - Read:
<http://book.xogeny.com/components/architectures/replaceable/>
- Read 14.7.

Kahoot

- <https://play.kahoot.it/#/k/694ab821-e4e0-421a-a5d6-0d297fd2cf1c>

Self-study section

Example: "Lambert's problem"

- IVP:

$$\begin{aligned}\dot{u} &= \frac{1}{100} - \left(\frac{1}{100} + u + v\right)(1 + (u + 1000)(u + 1)), & u(0) &= 0 \\ \dot{v} &= \frac{1}{100} - \left(\frac{1}{100} + u + v\right)(1 + v^2), & v(0) &= 0\end{aligned}$$

- Task: Simulate from $t = 0$ s til $t = 100$ s

- Eigenvalues:

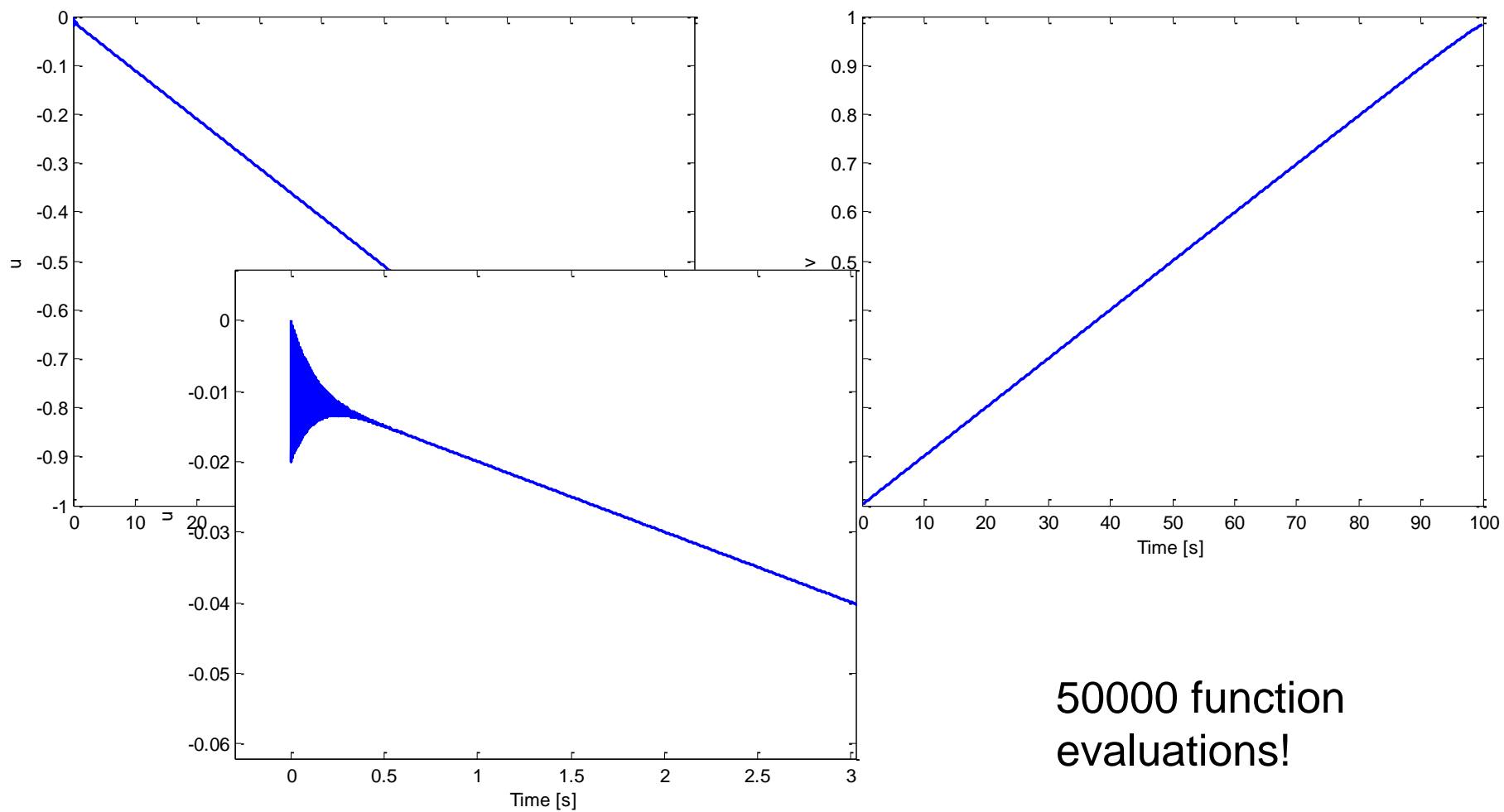
$$(u, v) = (0, 0) \Rightarrow \lambda_1 \approx -1000, \lambda_2 \approx -0.01$$

$$(u, v) = (-.5, .5) \Rightarrow \lambda_1 \approx -500, \lambda_2 \approx -0.03$$

$$(u, v) = (-1, 1) \Rightarrow \lambda_1 \approx -11, \lambda_2 \approx -1$$

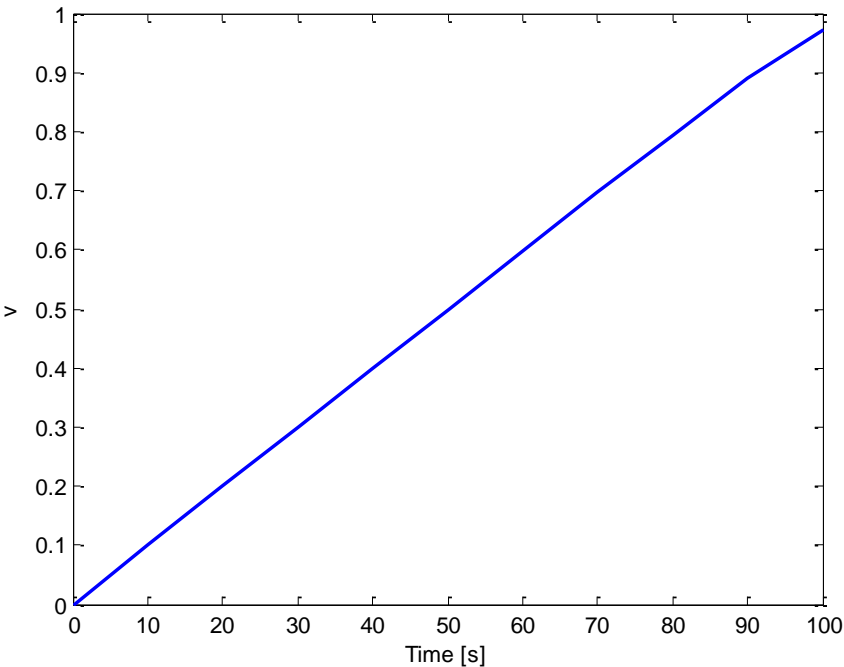
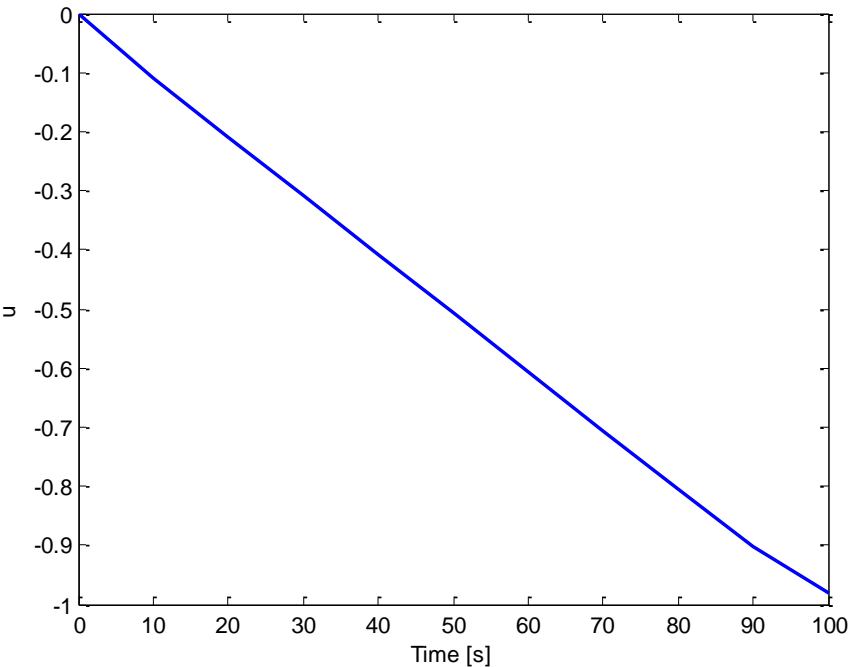
Using Euler (explicit), $h = 0.002$

$$h|\lambda| \leq 2$$



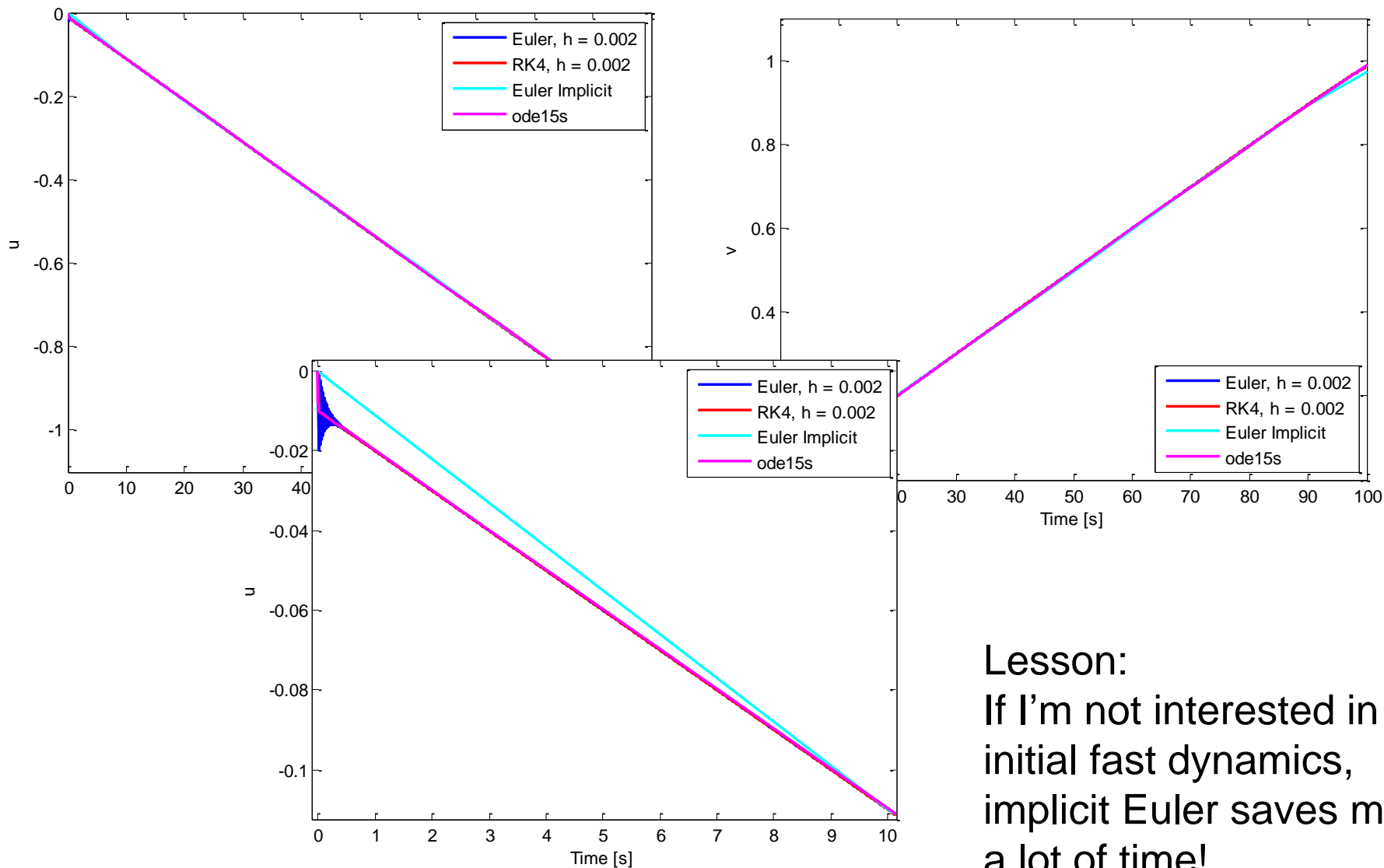
50000 function evaluations!

Attempt 3: Euler implicit, $h = 10$



149 function evaluations!
(dependent on solution algorithm)

Comparisons



Lesson:
If I'm not interested in initial fast dynamics, implicit Euler saves me a lot of time!