

Defining the Error in Solutions to Kepler's equation

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Abstract

This short derivation justifies the use of the error on the function's proximity to zero, ϵ , as a good approximation for the error on the root, δ , in the case of Kepler's equation. The region where this approximation might fail is also shortly discussed.

Let us define the following function as:

$$f(x; e, M) = x - e \sin x - M \quad (1)$$

where $e \in [0, 1)$ and $M \in [0, 2\pi]$ are parameters and $x \in [0, 2\pi]$ is the variable for which we want to find $f(x; e, M) = 0$, which corresponds to *Kepler's equation*.

There are many possible numerical methods for finding the roots, E , given the parameter values. But to be able to properly compare them, we should be concise in what we mean by the approximation error, δ , of the solution, x_0 .

$$\delta \equiv |E - x_0| \quad (2)$$

When a root is found, it is straightforward to obtain the error on the value of $f(x_0; e, M)$, we simply substitute x_0 back and see what we get, since:

$$\epsilon \equiv f(x_0; e, M) - 0 = f(x_0; e, M) \quad (3)$$

but equation 3 tells us very little about the error in the estimate of the actual root.

Suppose our algorithm has found a root $x_0 < E$, such that $E = x_0 + \delta$. If we expand $f(x_0; e, M)$ around $x = E$, we get:

$$f(x_0) = f(E) - \delta f'(E) + \frac{\delta^2}{2} f''(E) + \mathcal{O}(\delta^3) \quad (4)$$

Now, we know that the first term on the RHS of equation 4 is zero. We also defined the LHS as ϵ in equation 3, so we can write:

$$\epsilon = -\delta f'(E) + \frac{\delta^2}{2} f''(E) + \mathcal{O}(\delta^3) \quad (5)$$

Taking into account the formulas for the derivatives of $f(x; e, M)$, we are able to relate ϵ and δ by E :

$$\epsilon = -\delta (1 - e \cos E) + \frac{\delta^2}{2} e \sin E + \mathcal{O}(\delta^3) \quad (6)$$

The term of leading order in δ has a coefficient of magnitude ≈ 1 , thus the orders of magnitude of ϵ and δ are roughly the same. Knowing that in most of our uses we work with ϵ much smaller than 10^{-2} , we can confidently neglect the higher order terms in the expansion and claim that:

$$\epsilon \simeq \delta (-1 + e \cos E) \quad (7)$$

The function $f(x; e, M)$ is monotonous, since $f'(x; e, M) = 1 - e \cos x > 0$, $\forall x$, $e \in [0, 1)$. This means that:

$$\epsilon = \begin{cases} -|\epsilon|, & x_0 < E \\ +|\epsilon|, & x_0 > E \end{cases} \quad (8)$$

So, redefining $\epsilon \equiv |f(x_0; e, M)|$:

$$\delta \simeq \frac{\epsilon}{1 - e \cos E} \quad (9)$$

When treating a full orbit, i.e. an array of M values, the extremes of the δ values occur when $\cos E = -1$ and $\cos E = 1$, $\epsilon \approx (1 + e)\delta$ and $\epsilon \approx (1 - e)\delta$, so $\delta \in \left[\frac{\epsilon}{1+e}, \frac{\epsilon}{1-e} \right]$. This should not be a problem in most situations, but it might lead to larger error in the calculated eccentric anomaly near the pericenter of highly eccentric orbits. For $e = 0.9$, this would amount to a 10-fold increase in δ relative to ϵ , which is significantly higher.