# Hawkes processes

Consider a time-inhomogenous Poisson process  $(N_t)_{t\geq 0}\in\mathbb{N}$  whose intensity is itself a random process  $\lambda_t$  which evolves as

$$\lambda_t = \mu + \int_{[0,t]} \phi(t-s) dN_s$$

i.e.  $\lambda$  depends on the history of N itself. where  $\mu$  is a positive constant and  $\phi$  a positive function. For this reason we say that N is **self-exciting**, and this is a special type of **Stochastic Volterra equation** with no Brownian motion. The meaning of the  $\lambda_t$  is that  $\lim_{h\to 0} \frac{1}{h} \mathbb{P}(N_{t+h} - N_t > 0 | \mathcal{F}_t^{\lambda,N}) = \lambda_t$ , and note we can re-write  $\lambda$  as

$$\lambda_t = \lambda_0 + \sum_{0 \le s_i \le t} \phi(t - s_i)$$

where  $s_1, s_2, ...$  are the random **jump times** of N, which we can also take as the definition of  $\lambda$ .

If we let  $M_t = N_t - \int_0^t \lambda_u du$ , then M is a martingale and we can re-write the  $\lambda$  equation as

$$\lambda_t = \mu + \int_{[0,t]} \phi(t-s)(dM_t + \lambda_t dt) = \mu + (\phi * dM)_t + (\phi * \lambda)_t.$$

Note that

$$\|\phi * \lambda\|_{\infty} \leq \|\lambda\|_{\infty} \int_{[0,t]} \phi(t-s) ds = \|\lambda\|_{\infty} \int_{[0,t]} \phi(u) du < \|\lambda\|_{\infty} \int_{0}^{\infty} \phi(u) du < \|\lambda\|_{\infty}$$

if  $\|\phi\| = \int_0^\infty \phi(u) du < 1$ .

So  $\phi *$  is a contraction on  $C_b[0,\infty)$  under the sup norm, so its inverse is well defined, and also on  $C_b[0,T]$ . We can re-write this in operator notation as

$$(I - \phi *)\lambda = \mu + \phi * dM$$
.

To make sense of  $(I - \phi *)^{-1}$ , we look for a function  $\psi$  such that

$$(I - \phi *)^{-1} f = (I + \psi *) f.$$

for any test function f, so

$$\begin{array}{rcl} f & = & (I - \phi *)(I + \psi *)f & = & (I - \phi *)(f + \psi * f) \\ & = & f - \phi * f + \psi * f - \phi * \psi * f \end{array}$$

which we can re-write in operator form (i.e. without the f) as  $\phi * \psi = \psi - \phi$ ,.  $\psi$  is known as the **resolvent of**  $\phi$ , **note definition here is opposite way round to chap 3 in FM14**. Applying this to our Hawkes process i.e. setting  $f(t) = \lambda_t$ , we see that

$$\lambda = (I - \phi *)^{-1} (\mu + \phi * dM) = (1 + \psi) * \mu + (I + \psi) * (\phi * dM)$$

$$= \mu + \psi * \mu + (\phi * + (\phi *)^2 + ...) * dM$$

$$= \mu + \psi * \mu + \psi * dM$$

where  $\psi = \sum_{k=1}^{\infty} (\phi *)^k$ , which is shorthand for

$$\lambda_t = \mu + \mu \int_{[0,t]} \psi(t-s)ds + \int_{[0,t]} \psi(t-s)dM_s$$
 (1)

#### The propagator model - concave price impact from Hawkes order flow

Consider two independent Hawkes processes  $N_t^{\pm}$  with associated intensities  $\lambda_t^{\pm}$  which evolve as with

$$\lambda_t^{\pm} = \mu + \mu \int_0^t \psi(t-s)ds + \int_0^t \psi(t-s)dM_s^{\pm}$$

where  $dM_t^{\pm} = dN_t^{\pm} - \lambda_t^{\pm} dt$ . Then

$$\mathbb{E}_{t}(N_{u}^{+}) = g(t) + \mathbb{E}_{t}(\int_{0}^{u} \int_{0}^{s} \psi(s-v)dM_{v}^{+}ds) = g(t) + \mathbb{E}_{t}(\int_{0}^{u} \int_{v}^{u} \psi(s-v)dsdM_{v}^{+}) \\
= g(t) + \int_{0}^{t} \int_{v}^{u} \psi(s-v)dsdM_{v}^{+}$$

for some function g(t), and note that  $\int_v^\infty \psi(s-v)ds = \int_0^\infty \psi(s)ds = \|\psi\|_1$ . Then if we assume the current price  $P_t = \kappa \lim_{u \to \infty} \mathbb{E}(N_u^+ - N_u^- | \mathcal{F}_t)$  for some constant  $\kappa > 0$ , then

$$\frac{1}{\kappa}P_t = \lim_{u \to \infty} \mathbb{E}(N_u^+ - N_u^- | \mathcal{F}_t) = \int_0^t \sigma(dM_v^+ - dM_v^-)$$

where  $\sigma = ||\psi||$ . Then

$$\frac{1}{\sigma}P_{t} = N_{t}^{+} - N_{t}^{-} - \int_{0}^{t} \lambda_{s}^{+} ds + \int_{0}^{t} \lambda_{s}^{-} ds = N_{t}^{+} - N_{t}^{-} - \int_{0}^{t} \int_{u}^{t} \phi(s-u)ds(dN_{u}^{+} - dN_{u}^{-})$$

$$= \int_{0}^{t} (1 - \int_{0}^{t-u} \phi(s)ds)(dN_{u}^{+} - dN_{u}^{-})$$

so

$$P_t = \int_0^t \zeta(t - u)(dN_u^+ - dN_u^-)$$
 (2)

where  $\zeta(t) = \kappa \sigma(1 - \int_0^t \phi(s) ds)$ , so

$$\zeta'(t) = -\kappa \sigma \phi(t) = -\zeta(t)\phi(t)$$
.

Hence from the Volterra form in (2), we see that P is a **propagator model** (e.g. like transient price impact) but also retains the martingale property.

### Impact of a metaorder executed at constant rate before and after completion

Consider the additional contribution from an additional agent who buys at a fixed rate v for duration  $\tau$ . Then impact the cumulative impact time t is

$$P_t = \int_0^t \zeta(t-u)(dN_u^+ - dN_u^-) + v \int_0^{t \wedge \tau} \zeta(t-s)ds$$

whose expectation is  $MI(t) = v \int_0^{t \wedge \tau} \zeta(t-s) ds$ , which we call the **market impact function**, see plot below where we see **concave price impact** up to  $\tau$ , and then decay thereafter (which is broadly consistent with empirical findings where MI(t) is often found to be  $const. \times t^{\frac{1}{2}}$  for  $t < \tau$  (the so-called **square root impact law**.)

#### Example $\phi$ and $\psi$ functions

One can easily check that  $\|\phi\| = \frac{\|\psi\|}{1+\|\psi\|}$  so  $\|\psi\| = \frac{\|\phi\|}{1-\|\phi\|}$ . A common choice for  $\phi$  is

$$\phi(t) = \nu t^{\alpha - 1} E_{\alpha, \alpha}(-\lambda t^{\alpha})$$

where  $E_{\alpha,\alpha}$  is the **Mittag-Leffler** function (which is **heavy-tailed** since  $\phi(t) \sim \frac{const.}{t^{1+\alpha}}$  as  $t \to \infty$ ), and  $\int_0^\infty \phi(t)dt = \frac{\nu}{\lambda}$  so we choose  $\nu < \lambda$ . For this choice of  $\phi$ , the resolvent is

$$\psi(t) = \nu t^{\alpha - 1} E_{\alpha, \alpha}(-(\lambda - \nu)t^{\alpha})$$

and as  $\nu \nearrow \lambda$ ,  $\|\phi\| \nearrow 1$  and  $\|\psi\| \nearrow \infty$  (see also table below).

	$\psi(t)$	$\phi(t)$
Constant	c	$c e^{-ct}$
Fractional	$\frac{c  t^{\alpha - 1}}{\Gamma(\alpha)}$	$ct^{\alpha-1}E_{\alpha,\alpha}(-ct^{\alpha})$
Exponential	$c e^{-\lambda t}$	$c e^{-(\lambda+c)t}$
Gamma	$\frac{c  e^{-\lambda t} t^{\alpha - 1}}{\Gamma(\alpha)}$	$c e^{-\lambda t} t^{\alpha - 1} E_{\alpha, \alpha}(-c t^{\alpha})$

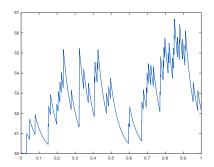


Figure 1: Here we we have simulated the intensity process of the form  $\lambda_t = \lambda_0 + \int_0^t k(t-s)dN_s$  for  $k(t) = e^{-10t}$  and  $\lambda_0 = 50$ .

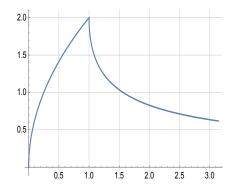


Figure 2: Here we see a typical concave impact function  $\int_0^{t\wedge\tau} \zeta(t-s)ds$  with  $\tau=1$ .

## Interpretation of a Hawkes process in terms of population dynamics

Let us define a population model: At time zero, there are no individuals. Some individuals (migrants) arrive as a uniform Poisson process with intensity  $\mu$ . If a migrant arrives at time s, the birth dates of its children form a Poisson process of intensity  $\phi(t-s)$  at time t, with  $\int_0^\infty \phi(t)dt < 1$ . In the same way, if a child is born at s', the birth dates of its children form a Poisson process of intensity  $\phi(\cdot - s')$ . Let  $N_t$  be the number of individuals who were born or migrated until time t. Then N is a Poisson-type process with intensity

$$\lambda_t = \mu + \int_0^t \phi(t-s) \, dN_s \tag{3}$$

i.e. N is a Hawkes process. This captures the notion of the process being self-exciting.

### References

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