

# A simplified unified propagator model for signed/unsigned order flow, concave price impact and rough volatility

14th Feb 2026

## Abstract

We simplify the setup in [MORS26] by exogenously modelling the signed order flow as an fBM-type process with  $H > \frac{1}{2}$ , which avoids the need for Hawkes processes, Mittag-Leffler functions and lengthy scaling limit arguments.

## 0.1 Positive order flow and the inversion formula

Similar to Theorem 3.1 in [MORS26]<sup>1</sup>, we consider a hyper-rough Heston process  $(F_t^+)_{t \geq 0}$  for cumulative positive order flow which satisfies

$$F_t^+ = g(t) + \int_0^t f(t-s)B_{F_s^+}ds$$

where  $B$  is a standard Brownian motion with  $B_0 = 0$  and  $f(t) = e^{-\theta t}t^{H-\frac{1}{2}}$  is the usual Gamma kernel with  $H \in (-\frac{1}{2}, \frac{1}{2})$ ,  $\theta > 0$  so  $f \in L^1$  and we assume  $g \in C^1$ .  $F_t^+$  is a.s.  $(2\alpha - \varepsilon) \wedge 1$ -Hölder continuous where  $\alpha = H + \frac{1}{2}$  (see e.g. Theorem 3.1 in [JM20] for details)<sup>2</sup> and  $F^+$  is an increasing process which is non-differentiable if  $H \leq 0$ . Then if  $h * f \equiv 1$ , we have the inversion formula:

$$\begin{aligned} \int_0^t h(t-s)(F_s^+ - g(s))ds &= \int_0^t h(t-s) \int_0^s f(s-u)B_{F_u^+}duds = \int_0^t \int_u^t h(t-s)f(s-u)dsB_{F_u^+}du \\ &= \int_0^t \int_0^{t-u} h(t-s-u)f(s)dsB_{F_u^+}du \\ &= \int_0^t (h * f)(t-u)B_{F_u^+}du = \int_0^t B_{F_u^+}du \end{aligned}$$

(where Fubini is justified since  $F_t^+ < \infty$  a.s. for finite  $t$  and  $B$  is a.s. continuous), and hence

$$\frac{d}{dt} \int_0^t h(t-s)(F_s^+ - g(s))ds = B_{F_t^+} \quad (1)$$

Lebesgue a.e. Taking Laplace transforms, the condition  $h * f \equiv 1$  becomes  $\hat{f}(\lambda)\hat{h}(\lambda) = \frac{1}{\lambda}$ , from which we find that

$$h(t) = \frac{\theta^\alpha}{\Gamma(\alpha)} \left(1 - \frac{\Gamma(-\alpha, t\theta)}{\Gamma(-\alpha)}\right)$$

where  $\Gamma(a, z) = \int_z^\infty s^{a-1}e^{-s}ds$  is the incomplete Gamma function, and in particular  $h(t) = O(t^{-H-\frac{1}{2}})$  as  $t \rightarrow 0$ .

## 0.2 Persistent signed order flow

Again following [MORS26], we now assume the *signed* order flow for an asset is

$$V_t = F_t^+ - F_t^-$$

where  $F_t^- = g(t) + \int_0^t f(t-s)W_{F_s^-}ds$  and  $W$  is another Brownian motion independent of  $B$  (so  $F^-$  is an i.i.d. copy of  $F^+$ ). Then we see that

$$\frac{d}{dt} \int_0^t h(t-s)V_sds = B_{F_t^+} - W_{F_t^-}$$

since the  $g$  terms cancel, and for  $0 \leq t \leq u$  we have

$$\mathbb{E}(F_u^+ | \mathcal{F}_{F_t^+}^B) = g(u) + \mathbb{E}\left(\int_0^u f(u-s)B_{F_s^+}ds | \mathcal{G}_t^+\right) = g(u) + \int_0^t f(u-s)B_{F_s^+}ds + B_{F_t^+} \int_t^u f(u-s)ds$$

(and similarly for  $\mathbb{E}(F_u^- | \mathcal{F}_{F_t^-}^W)$ ), where  $\mathcal{G}_t^+ = \mathcal{F}_{F_t^+}^B$ , assuming  $(B_{F_t^+})_{t \geq 0}$  is a  $\mathcal{G}_t^+$ -martingale (see e.g. Section 7 in [AJ21] for more on this).

<sup>1</sup>see also section 5.4 in [FGS21]

<sup>2</sup>we can simulate  $F^+$  using the Monte Carlo scheme in [AA25] using Normal Inverse Gaussian variates

### 0.3 Price dynamics and the propagator model

Making the usual assumption that the asset price  $P_t = \kappa \lim_{u \rightarrow \infty} \mathbb{E}(F_u | \mathcal{G}_t)$  where  $F_t = F_t^+ - F_t^-$  and  $\mathcal{G}_t = \sigma(\mathcal{G}_t^+, \mathcal{G}_t^-)$  with  $\mathcal{G}_t^- = \mathcal{F}_{F_t^-}^W$  (for some  $\kappa > 0$ )<sup>3</sup>, we find that  $f(u-s) \rightarrow 0$  as  $u \rightarrow \infty$  since  $\theta > 0$ , but  $\lim_{u \rightarrow \infty} \int_t^u f(u-s)ds = c_{H,\theta}$  where  $c_{H,\theta} = \theta^{-\frac{1}{2}-H} \Gamma(\frac{1}{2} + H)$ , so

$$P_t = \kappa \lim_{u \rightarrow \infty} \mathbb{E}(F_u^+ - F_u^- | \mathcal{G}_t) = \kappa c_{H,\theta} (B_{F_t^+} - W_{F_t^-})$$

i.e.  $P$  is the difference of two i.i.d. hyper-rough Heston models (each with correlation  $\rho = 1$  since each process only has one driving Brownian motion), and

$$P_t = \kappa c_{H,\theta} \frac{d}{dt} \int_0^t h(t-s) V_s ds$$

which is a weak formulation of the usual Propagator model, and the *unsigned* order flow is  $U_t = F_t^+ + F_t^-$ .

### 0.4 Power-law price impact

In particular, the *market impact function* of an exogenous metaorder executed at constant trading speed 1 up to time  $t_0$  is given by

$$MI(t) = \frac{d}{dt} \int_0^t h(t-s) s ds = \frac{\theta^{-\alpha_-} (t\theta \Gamma(-\alpha) - t\theta \Gamma(-\alpha, t\theta) - \Gamma(\alpha_-, 0) + \Gamma(\alpha_-, t\theta))}{\Gamma(-\alpha) \Gamma(\alpha)}$$

for  $0 \leq t < t_0$  (where  $\alpha_- = \frac{1}{2} - H$ ) for  $0 \leq t \leq t_0$  which is  $O(t^{\frac{1}{2}-H})$  as  $t \rightarrow 0$  and globally concave in  $t$  for  $H \in (-\frac{1}{2}, \frac{1}{2})$ , consistent with empirical evidence (note the usual square-root impact law corresponds to  $H = 0$  here), and we can show that the asymptotic (i.e. permanent) impact as  $t \rightarrow \infty$  (given no trading after  $t_0$ ) is

$$MI(\infty) = \frac{t_0 \theta^\alpha}{\Gamma(\frac{1}{2} + H)}$$

(this is the asymptotic grey line in the plot below). For  $H = \frac{1}{2}$ ,  $h(t) = \theta$  which corresponds to pure permanent price impact; conversely  $\lim_{\theta \rightarrow 0} h(t) = \text{const.} \times t^{-\frac{1}{2}-H}$ . Recall from above that  $F_t^+$  has Hölder regularity  $(2\alpha - \varepsilon) \wedge 1$ , and empirical evidence in [MORS26] suggests that  $2\alpha \in (\frac{1}{2}, 1)$  which corresponds to  $H \in (-\frac{1}{4}, 0)$ .

### 0.5 Extensions

If we consider a mixed model where  $f(t) = e^{-\theta t} t^{H-\frac{1}{2}} + e^{-\theta_2 t} t^{H_2-\frac{1}{2}}$  with  $H \in (-\frac{1}{4}, 0)$  and  $H_2 \in (0, \frac{1}{2})$ , then we find that

$$h(t) = \frac{1}{\Gamma(\alpha)} \sum_{n=0}^{\infty} (-c)^n \sum_{k=0}^{\infty} \frac{(\beta n)_k}{k!} \delta^k \theta^{-(\mu_n+k)} \frac{\gamma(\mu_n+k, \theta t)}{\Gamma(\mu_n+k)}$$

where  $\beta = \frac{1}{2} + H_2$ ,  $c = \Gamma(\beta)/\Gamma(\alpha)$ ,  $\delta = \theta - \theta_2$ ,  $\mu_n = \beta n - \alpha(n+1)$ ,  $\gamma(a, z) = \int_0^z s^{a-1} e^{-s} ds$  is the lower incomplete Gamma function and  $(a)_k$  denotes the Pochhammer symbol. We can then use two (possibly different)  $f$  functions to drive  $F^\pm$ , and the presence of the positive  $H_2$  parameter means the price dynamics will now have a (non-hyper) rough component, which is arguably more realistic. We can have  $W$  and  $B$  correlated which gives greater flexibility in capturing skew for the price process  $P_t$ .

---

<sup>3</sup>note we are assuming  $P_0 = 0$  without loss of generality since we can easily add a  $P_0$  term

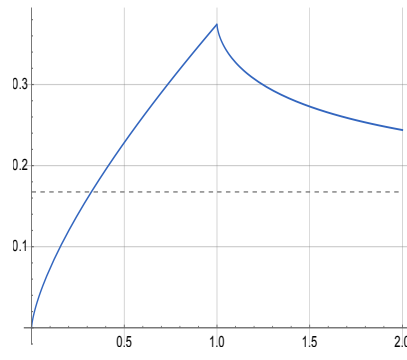


Figure 1: Concave price impact of an exogenous metaorder executed at constant trading speed 1 over  $[0, 1]$  with  $H = -0.2$ ,  $\theta = 0.1$  and  $\kappa = 1$ . The blue line asymptotes to the constant level  $MI(\infty)$  (grey line) as  $t \rightarrow \infty$  which represents the asymptotic permanent price impact.

## References

- [AJ21] Abi Jaber, E., “Weak existence and uniqueness for affine stochastic Volterra equations with  $L^1$ -kernels”, *Bernoulli* 27(3), 2021, 1583–1615
- [AA25] Abi Jaber, E. and E. Attal, “Simulating integrated Volterra square-root processes and Volterra Heston models via Inverse Gaussian”, Preprint, 2025
- [FGS21] Forde, M., S. Gerhold and B. Smith, “Small-time, large-time and  $H \rightarrow 0$  asymptotics for the rough Heston model”, *Mathematical Finance*, 31(1), 203–241, 2021.
- [JM20] P. Jusselin and M. Rosenbaum, “No-arbitrage implies power-law market impact and rough volatility”, 2018, to appear in *Mathematical Finance*, *Mathematical Finance*, Volume 30, Issue 4 pp. 1309–1336
- [MORS26] Muhle-Karbe, J., Y. Chahdi, M. Rosenbaum, and G. Szymanski, “A unified theory of order flow, market impact, and volatility”, preprint.