

Exercise 1, TFY4235 Computational physics

Martin K. Johnsrud

Introduction

The goal of this exercise is to simulate particles as flat, hard disks in a square, 2D container. This is done with a event-driven simulation, as described in the exercise [1]. This is implemented in Python, using the built-in library `heapq`. The simulation is used to first tested with scenarios we know the outcome of, then used to demonstrate the Maxwell-Boltzmann distribution and to investigate the effect of a large, heavy disk hitting a large number of small, inert particles.

Implementation

The main engine of the code is the function `run_loop()` in `utilities.py`. It follows the algorithm, as laid out in [1], using the objects:

- `particles`, a numpy array with the position and velocity of all the particles.
- `collisions`, a priority queue containing the time of the collision, the index of the particle(s) involved, and the type of collision it is.
- `last_collided`, a list of when each particle was involved in a collision.

To be expanded ...

Tests

Several functions were developed to test the accuracy of the simulation. First, one particle, starting at in the middle of the box, all the way to the left, and with a velocity with at 45° to the x -axis should move in a titled rectangle. With $\xi = 1$ it should also conserve energy. Figure 1 shows that this is still the case after 10,000 events.

To test the validity of the particle collision, one small, light particle is sent towards a single, large and heavy particle, with varying impact parameters. The relationship between the impact parameter is is (ref

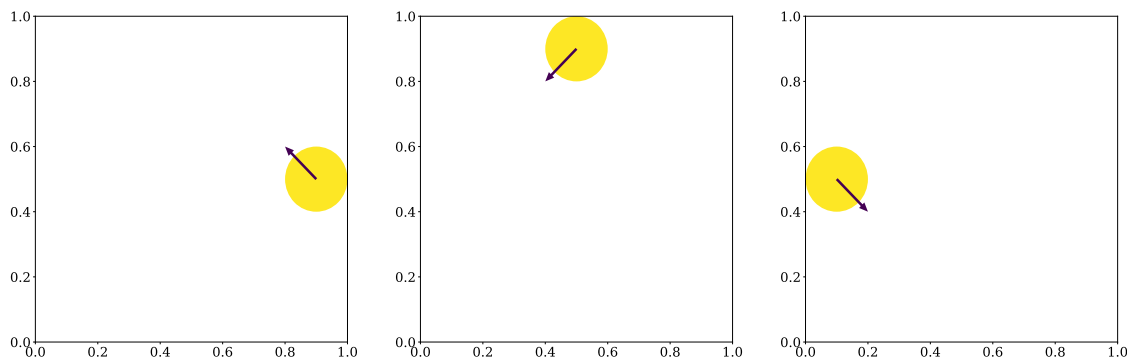


Figure 1: One ball being simulated. after 10000 steps, it still follows a regular pattern.

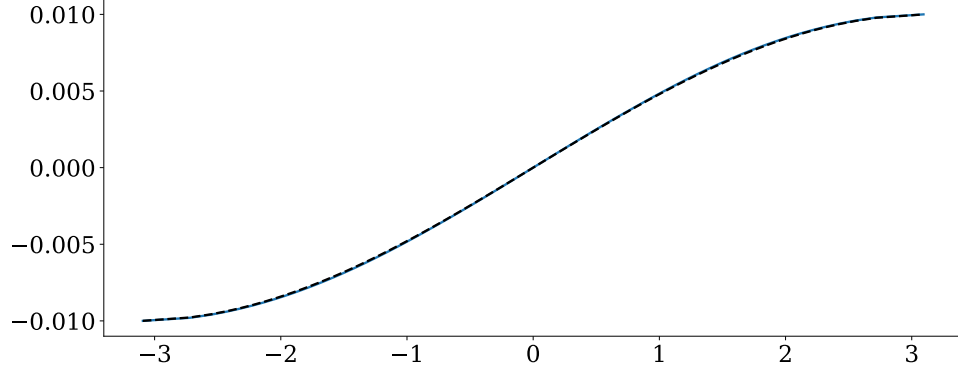


Figure 2: The

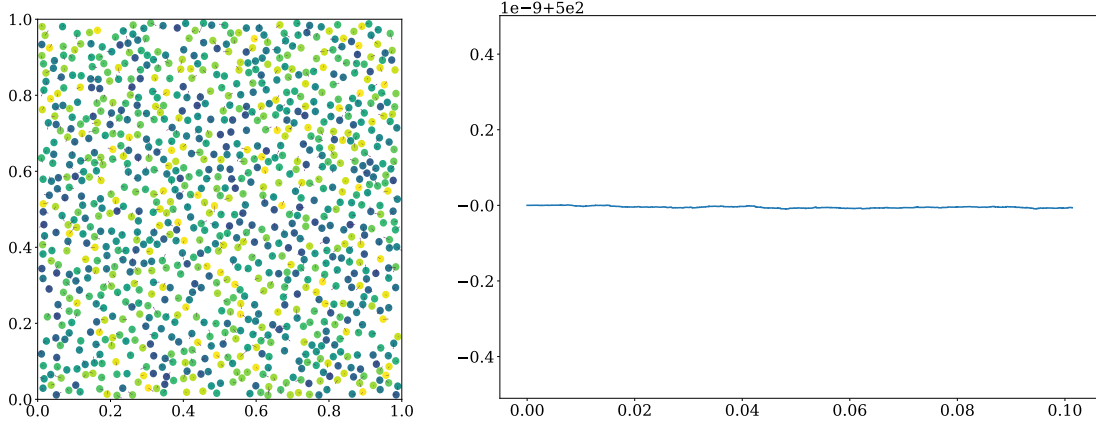


Figure 3: On the left is a snapshot of the particles. The arrows represent the velocities. On the right, the energy is plotted as a function of events. The energy loss can be seen to be negligible.

til goldstein, begunn hvorfor det blir $\cos \frac{ds}{d\theta} = a/2 \sin(\theta/2)$. The result is shown i Figure 2, and is in good agreement with the theory. Lastly, the energy from a simulation of N particles over T events is shown in

Results

As the system equilibrates, it should reach the Maxwell-Boltzmann distribution, which in 2D is

$$f(v) = \frac{mv}{T} \exp\left(-\frac{mv^2}{2T}\right),$$

when using units in which $k_b = 1$. The equipartition theorem gives the temperature $T = E$ in 2D. Figure 5 was used to find a good starting point for when the simulation has reach equilibrium. After that, the simulation is sampled every N event, where N is the number of particles. This ensures somewhat independent samples. Figure 4 shows the velocity distribution, compared to the Maxwell Boltzmann distribution.

Next, two different types of particles are simulated, one with a mass of 1, the other with a mass of 4. test

By setting the elasticity parameter (?) ξ to different values, the way the system approaches equilibrates changes.

To simulate the impact of a projectile, one disk is placed at $\vec{x} = (0.5, 0.75)$, and given a velocity of $v = (0, -20)$. The lower half of the square box is filled with small, light particles. The impact from the large

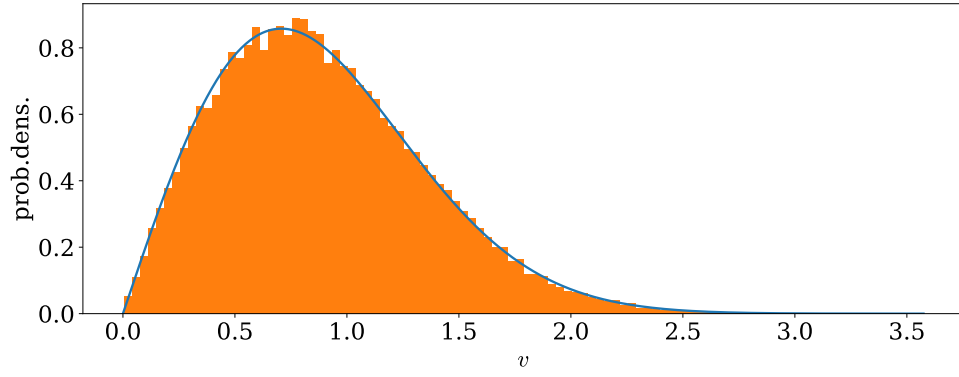


Figure 4: The velocity distribution is a good fit with the Maxwell-Boltzmann distribution, shown as a blue line.

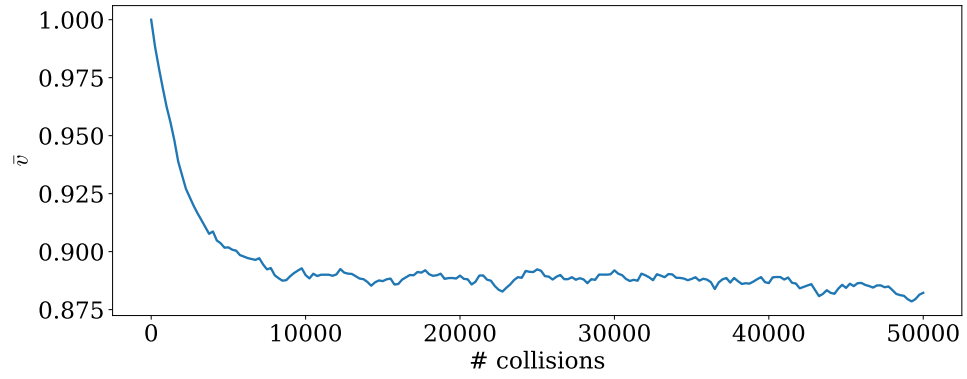


Figure 5: Average velocity, as a function of collisions. The distribution reaches equilibrium around 6000 collisions, or $3N$

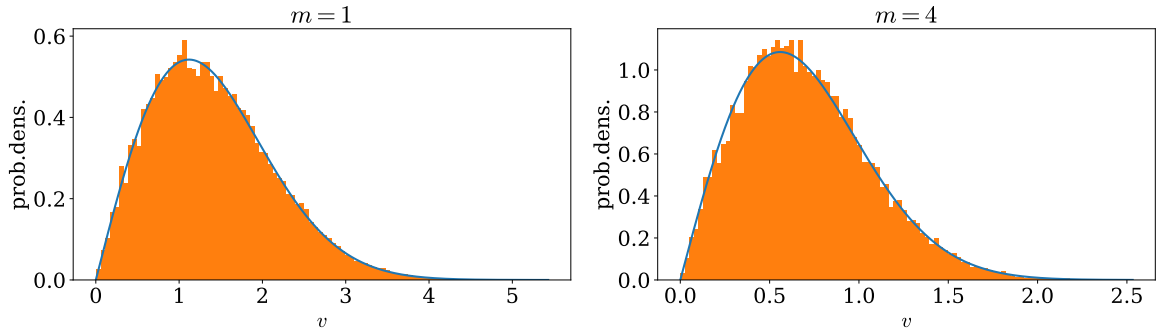


Figure 6: The velocity distribution of the particles with $m = 1$ and $m = 4$ is show left and right, resp. and compared to the Maxwell-Boltzmann distribution.

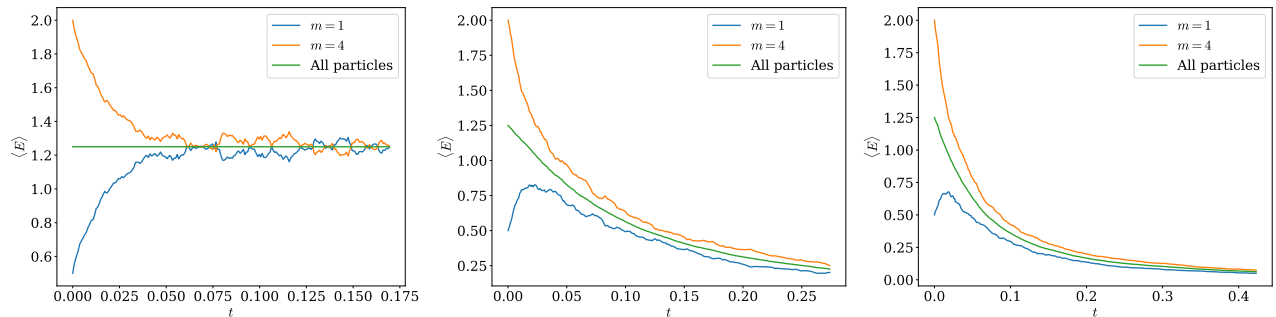


Figure 7: The average energy, as a function of time, of resp. $\xi = 1$, $\xi = 0.9$ and $\xi = 0.8$.

disk into the densely packed disks will create a crater. To investigate the effect of the mass and radius on the size of the crater, the size of the crater must be measured. This is done by laying a grid with a spacing Δx , and searching through each cell if there is a disk inside it. The process for checking if a disk is inside each square cell is illustrated in figure 8. Then, the size of the crater is given by $s = m\Delta x^2$, where m is the number of unoccupied cells. This method relies on choosing a cell size Δx large enough so that only cells within the crater are empty. This is done by inspecting the result, and comparing it to a plot of the crater.

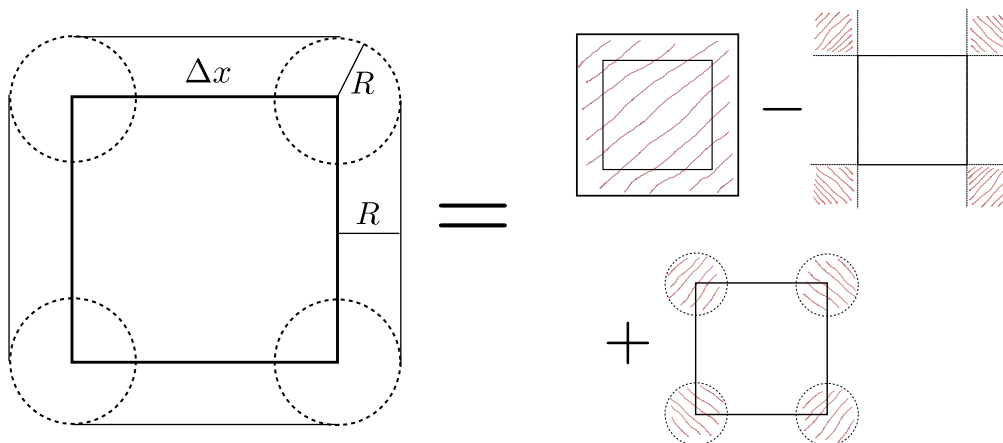


Figure 8: Checking if a disk of radius R is inside of a square of side length Δx , is equivalent to checking if the center of the disk is inside the shape on the left side. Thus, the task is reduced to checking if the center is inside the larger square of side lengths $R + 2\Delta x$, but not at the corners, or inside one of the circles of radius R centered at the corners of the square.

To avoid inelastic, which was first observed when the projectile was approaching the bottom of the box, the simulation was implemented as at TC-model as described in [2]. Good results, without inelastic collapse, were found with $t_c = 10^{-8}$.

To investigate the effect of the size of projectile on the crater, the simulation is run 10 times, using projectiles with different radius. The mass of the projectile is proportional to R^2 , so that the density is constant. Figure 9 shows the size of the crater,

References

- [1] TFY4235 Computational Physics Exercise 1, 2021. Institutt for fysikk.
- [2] Stefan Luding. How to handle the inelastic collapse of a dissipative hard-sphere gas with the tc model. *Granular Matter*, 1:113–128, 1998.

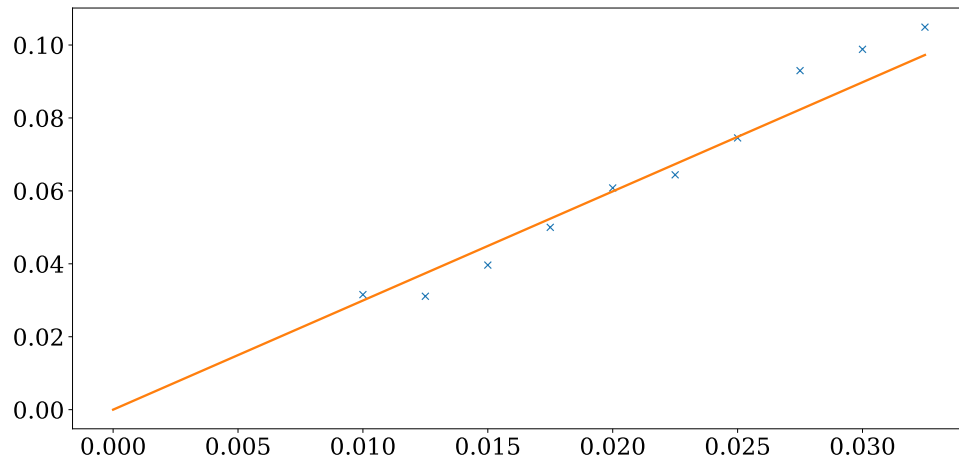
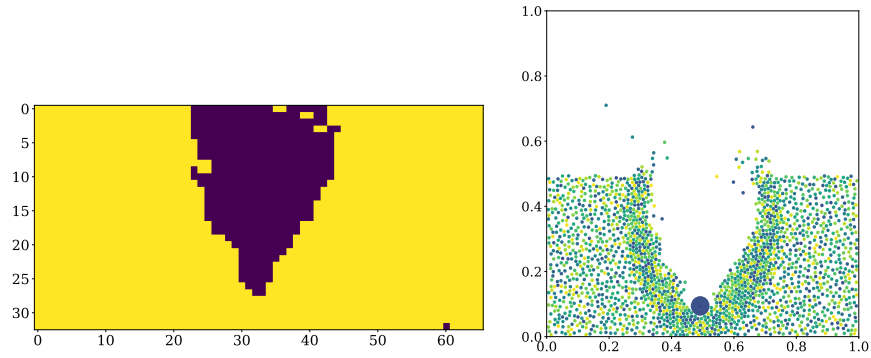


Figure 9: The size of the crater, as a function of the radius of the incoming projectile.