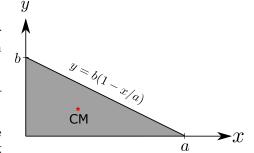
# Exercise 8 - TFY4345 Classical Mechanics

#### 2020

#### 1 Principal moments of inertia of a triangular slab

(Exam Aug. 2019)

(a) Compute the center-of-mass (CM) for the planar triangle in the figure, assuming it to be of uniform two-dimensional mass density  $\rho$ .



- (b) Compute the inertia tensor with respect to the origin, for the same triangle.
- (c) (Optional) If the origin is shifted to the CM, the inertia tensor becomes (this can be show by using the Steiner's parallel axis theorem)

$$I_{CM} = \begin{pmatrix} I_{11} & I_{12} & 0 \\ I_{21} & I_{22} & 0 \\ 0 & 0 & I_{33} \end{pmatrix} = \frac{M}{18} \begin{pmatrix} a^2 & \frac{1}{2}ab & 0 \\ \frac{1}{2}ab & b^2 & 0 \\ 0 & 0 & a^2 + b^2 \end{pmatrix}$$

where  $I_{xy}=I_{yx}$  and  $I_{xx}+I_{yy}=I_{zz}$  in the general form show first. Define next

$$A = \frac{1}{2}(I_{xx} + I_{yy}), \quad B = \sqrt{\frac{1}{2}(I_{xx} - I_{yy})^2 + I_{xy}^2}, \quad \vartheta = \tan^{-1}\left(\frac{2I_{xy}}{I_{xx} - I_{yy}}\right).$$

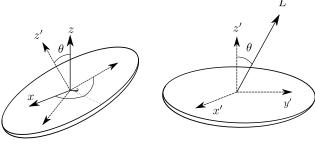
Derive the principal moments of inertia and the principal axes by using the general form of the inertia tensor, and these new variables.

[Hint] The last equations comprises a relationship that can be described by a right triangle.

### 2 Precession of a frisbee

(Exam Aug. 2016)

Consider an axial-symmetric body with the principal moments of inertia  $I_1 = I_2 \neq I_3$  for rotation around the axis of symme-



try. The angular momentum in the laboratory frame is  $\mathbf{L} = L\mathbf{e}_z$ , where z is the symmetry axis.

- (a) Derive the equations of motion for the body, using the Euler equations and the angles  $\theta, \psi$  and  $\phi$ . Define also the components of  $\omega$ . (See lecture notes, we derived this already!)
- (b) Find the expression for the Euler angles  $\theta, \psi, \phi$  as a function of time.
- (c) Assume  $I_3 = 2I_1$ . The precession (wobble) of the frisbee is given by  $\dot{\phi}$ . Show that the precession is twice as fast as the rotation frequency of the frisbee, assuming that  $\theta$  is small (i.e. that  $\cos(\theta) \approx 1$ ).

## 3 Precession of a heavy spinning top

(Based on example p. 208-223 in Goldstein 3rd. ed., p. 70-74 in the compendium)

In the example we define the shifted energy as

$$E' = \frac{1}{2}I_1\dot{\theta}^2 + V(\theta), \quad V(\theta) = \frac{(p_\phi - p_\psi \cos(\theta)^2)}{2I_1\sin(\theta)^2} + Mgh\cos(\theta),$$

which is a constant of motion. We also defined  $\theta_0$  to be the constant angle of inclination of spinning top with regular precession, e.g. that the symmetry axis rotates around the z at a fixed angle. Consider the shape of the effective potential  $V(\theta)$  for  $\theta_0$ . What is the condition in this case? The following change of variables will come in handy for the result:

$$\beta = p_{\phi} - p_{\psi} \cos(\theta_0).$$

You will encounter a quadratic equation for  $\beta$ . Show that for the equilibrium precession inclination angle  $\theta_0$ , the following must hold true:

$$\omega_3 \le \frac{2}{I_3} \sqrt{MghI_1\cos(\theta_0)}.$$

What can you say about the corresponding precession angular velocity  $\dot{\phi}_0$ ?