

# Gravgård

Gravgården er hvor tekster går for å dø. Ting som er fjernet fra main, men som jeg ikke helt klarer å slette enda

## 1 path integral

By Wick's theorem, an  $n$ -point correlated is given by the sum of all Feynman diagrams with  $n$  external vertices. The factor  $Z[0]^{-1}$  divides out all *vacuum bubbles*, that is diagrams without external vertices. We can show this by considering

Here we defined the *generating functional for connected diagrams*,  $W[J]$ . The reason for the name will become apparent later. (HUSK Å REFFERE TILBAKE) The expectation value of some function of the field-configuration,  $A = A[\varphi]$ , in the presence of the source  $J$  is

$$\langle A \rangle_J = \frac{1}{Z[J]} A \left( -i \frac{\delta}{\delta J} \right) Z[J]. \quad (1)$$

(DEFINE FUNCTIONAL DERIVATIVE) The expectation value of the field defines a functional,

$$\varphi[J](x) = \langle \varphi(x) \rangle_J = \frac{1}{Z[J]} \left( -i \frac{\delta}{\delta J} \right) Z[J] = \frac{\delta}{\delta J(x)} W[J], \quad (2)$$

and is sometimes called the *classical field*. The notation  $\mathcal{F}[f](x)$  means that  $\mathcal{F}$  is a functional which takes in a function  $f$ , and returns the new function  $(\mathcal{F}[f])(x)$ . One example is the Lagrangian density, which takes in a field, and returns a function which has a value for each point in space-time. We can reverse this relationship, by defining the functional  $J[\varphi](x)$  as *the current which causes the classical field*  $\varphi$ . That is, if  $\varphi[J_0](x) = \varphi_0(x)$  for some source  $J_0$ , then  $J[\varphi_0] = J_0$

## 2 free scalar

Comparing with the definitions of the thermal propagator in ??, we can write the free energy compactly as

$$\beta F = \frac{1}{2} \text{Tr} \{ \ln [D_0^{-1}(K, K')] \} = \frac{1}{2} \text{Tr} \{ \ln [\beta^2 D_0^{-1}(K)] \}. \quad (3)$$

## 3 interacting scalar

Notice that the constant factor from the Jacobian due to the change of variable  $\varphi \rightarrow \tilde{\varphi}$  does not affect the expectation value, as the same factor is in both the numerator and denominator. If the quantity  $A$  is a function of the momentum-space fields,  $A = A[\tilde{\varphi}(K)]$ , then this expectation value takes the form

$$\langle A \rangle_0 = \frac{\int_{\tilde{S}} \mathcal{D}\tilde{\varphi}(K) f[\tilde{\varphi}(K)] \exp \left\{ -\frac{1}{2} \langle \tilde{\varphi}^*, D \tilde{\varphi} \rangle \right\}}{\int_{\tilde{S}} \mathcal{D}\tilde{\varphi}(K) \exp \left\{ -\frac{1}{2} \langle \tilde{\varphi}^*, D \tilde{\varphi} \rangle \right\}}. \quad (4)$$

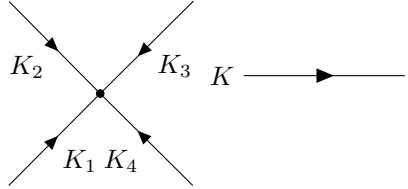
where, as before,

$$\langle \tilde{\varphi}^*, D\tilde{\varphi} \rangle = \int_{\Omega} \tilde{d}K [\beta^2(\omega_n^2 + \omega_n^2)] |\tilde{\varphi}(K)|^2 \quad (5)$$

The exponential form of  $Z[J]$  leads straight forwardly to Wick's theorem, which states that an expectation value of  $2n$  fields is a sum of *all possible, distinct* combination of  $n$  propagators. To write this in a formal way, we define the functions  $a$  and  $b$ , which define a way to pair up  $2m$  elements. The domain of the functions are the integers between 1 and  $m$ , the image a subset of the integers between 1 and  $2m$  of size  $m$ . A valid pairing is a set  $\{(a(1), b(1)), \dots (a(m), b(m))\}$ , where all elements  $a(i)$  and  $b(j)$  are different, such all integers up to and including  $2m$  are featured. A pair is not directed, so  $(a(i), b(i))$  is the same pair as  $(b(i), a(i))$ . Wick theorem states that,

$$\left\langle \prod_{i=1}^{2m} \varphi(X_i) \right\rangle_0 = \sum_{\{(a,b)\}} \langle \varphi(X_{a(i)}) \varphi(X_{b(i)}) \rangle, \quad (6)$$

where the sum is over all tuples  $(a, b)$  that define a valid and unique pairing.



$$(7)$$

The expression is the integrated over all *internal* momenta. The factor  $1/4!$  is removed as a general Feynman diagram represent all diagrams with the same form, but different pairing of the momenta. Some diagrams are more symmetric, such that an exchange of momenta still gives *the same pairing*.

## 4 effective action

In free theory, we may write

$$W[J] = \frac{1}{2} \int d^4x d^4y J(x) D_0(x-y) J(y), \quad (8)$$

where  $D_0$  is the free propagator. We may reverse the relation ?? to write the source in terms of the field,

$$J = D_0^{-1} \varphi(x) \quad (9)$$

This is the field equation for the free field with a source. For the scalar Klein-Gordon field,  $D_0^{-1} = \partial^2 + m^2$ . Inserting these two relation into the definition of the effective action, and assuming we can do partial integration with  $D_0^{-1}$ , we get

$$\Gamma[\varphi] = W[J] - \int d^4x J(x) \varphi(x) = \int d^4x \left( \frac{1}{2} \int d^4y (D_0^{-1} \varphi) D_0 (D_0^{-1} \varphi) - (D_0^{-1} \varphi) \varphi \right) = -\frac{1}{2} \int d^4x \varphi(x) D_0^{-1} \varphi(x) \quad (10)$$

This is the classical action. Thus, the effective action  $\Gamma$  and the classical action  $S$  are the same to first order in perturbation theory.

Let  $\varphi^*$  solve the quantum mechanical version of the equation of motion, i.e.

$$\frac{\delta \Gamma[\varphi^*]}{\delta \varphi} = 0. \quad (11)$$

We can Taylor-expand the classical action around this point, by setting  $\varphi(x) = \varphi^*(x) + \eta(x)$  for some function  $\eta$ . The generating functional becomes

$$Z[J] = \int \mathcal{D}(\varphi^* + \eta) \exp \left\{ iS[\varphi^* + \eta] + i \int d^4x J(\varphi^* + \eta) \right\} \quad (12)$$

The functional version of a Taylor expansion is

$$S[\varphi^* + \eta] = S[\varphi^*] + \int dx \frac{\delta S[\varphi^*]}{\delta \varphi(x)} \eta(x) + \frac{1}{2} \int dx dy \frac{\delta^2 S[\varphi^*]}{\delta \varphi(x) \delta \varphi(y)} \eta(x) \eta(y) + \dots \quad (13)$$

Inserting this into  $Z[J]$ , with  $S_I$  to denote the derivatives of higher order than 2, we get

$$Z[J] = \int \mathcal{D}\eta \exp \left\{ i \int d^4x (\mathcal{L}[\varphi^*] + J\varphi^*) + i \int dx \left( \frac{\delta S[\varphi^*]}{\delta \varphi(x)} + J(x) \right) \eta(x) + i \frac{1}{2} \int dx dy \frac{\delta^2 S[\varphi^*]}{\delta \varphi(x) \delta \varphi(y)} \eta(x) \eta(y) + i S_I[\eta] \right\}$$

In the first term we used the definition of the classical action. This term is constant with respect to  $\eta$ , and may therefore be taken outside the path integral. The next term is the classical equation of motion with a source,

$$\frac{\delta S[\varphi]}{\delta \varphi(x)} = -J(x), \quad (14)$$

evaluated at  $\varphi^*$ .

$$\frac{\delta S[\varphi^*]}{\delta \varphi(x)} + J(x) = \frac{\delta \Gamma[\varphi^*]}{\delta \varphi(x)} + J(x) + \left( \frac{\delta S[\varphi^*]}{\delta \varphi(x)} - \frac{\delta \Gamma[\varphi^*]}{\delta \varphi(x)} \right) = \left( \frac{\delta S[\varphi^*]}{\delta \varphi(x)} - \frac{\delta \Gamma[\varphi^*]}{\delta \varphi(x)} \right) := \delta J \quad (15)$$

The second to last term is a Gaussian integral, and may be evaluated as described in ??,

$$\int \mathcal{D}\eta \exp \left( i \frac{1}{2} \int dx dy \frac{\delta^2 S[\varphi^*]}{\delta \varphi(x) \delta \varphi(y)} \eta(x) \eta(y) \right) = C \det \left( \frac{\delta^2 S[\varphi^*]}{\delta \varphi^2} \right)^{-1/2} \quad (16)$$

This leaves us with

$$W[J] = -i \ln(Z) \quad (17)$$

$$= \int d^4x (\mathcal{L}[\varphi^*] + J\varphi^*) - \frac{1}{2} \text{Tr} \left\{ \ln \left( -\frac{\delta^2 S[\varphi^*]}{\delta \varphi^2} \right) \right\} + \int \mathcal{D}\eta \exp \left\{ i \int d^4x \delta J(x) \eta \right\} + \int \mathcal{D}\eta e^{i S_I} \quad (18)$$

$\delta J$  is ultimately dependent on our choice of  $J$  to define  $\varphi$ . It contributes to the expectation value of  $\eta$ , through tadpole diagrams

$$\langle \eta \rangle_{j=0} = \text{---} \longrightarrow \text{---} \bigcirc \text{---} \quad (19)$$

This can be removed by using the renormalization condition

$$\text{---} \longrightarrow \text{---} \bigcirc \text{---} = 0. \quad (20)$$

FULL LAGRANGIAN !!!!!

$\mathcal{L} =$

$$\begin{aligned}
& f^2 \left( 2B_0 m \cos \alpha + \frac{1}{2} \mu^2 \sin^2 \alpha \right), \\
& + f(\mu_I^2 \cos \alpha - 2B_0 m) \pi_1 \sin \alpha + f \mu_I \partial_0 \pi_2 \sin \alpha, \\
& + \frac{1}{2} \partial_\mu \pi_a \partial^\mu \pi_a + \mu_I \cos \alpha (\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) - B_0 m \pi_a \pi_a \cos \alpha + \frac{1}{2} \mu_I^2 \pi_a \pi_b k_{ab}, \\
& + \frac{\pi_a \pi_a \pi_1}{6f} (2B_0 m \sin \alpha - 2\mu_I^2 \sin 2\alpha) \\
& - \frac{2\mu_I}{3f} [\pi_1 (\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) + \pi_3 (\pi_3 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_3)] \sin \alpha, \\
& + \frac{1}{6f^2} \left\{ \frac{1}{2} B_0 m (\pi_a \pi_a)^2 \cos \alpha - [(\pi_a \pi_a) (\partial_\mu \pi_b \partial^\mu \pi_b) - (\pi_a \partial_\mu \pi_a) (\pi_b \partial^\mu \pi_b)] \right\} \\
& - \frac{\mu_I \pi_a \pi_a}{3f^2} \left[ (\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) \cos \alpha + \frac{1}{2} \mu_I \pi_a \pi_b k_{ab} \right], \\
& + \frac{l_1}{4} \left( \frac{8\mu_I^2}{f^2} (\partial_\mu \pi_a \partial^\mu \pi_a + 2\partial_\mu \pi_2 \partial^\mu \pi_2) \sin^2 \alpha + 16\mu_I^3 \left[ \frac{\partial_0 \pi_2}{f} \sin^3 \alpha + \frac{1}{f^2} (3\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) \cos \alpha \sin^2 \alpha \right] \right. \\
& + 4\mu_I^4 \left\{ \sin^4 \alpha + 2 \sin^2 \alpha \left[ \frac{\pi_1}{f} \sin 2\alpha + \frac{\pi_a \pi_b}{f^2} (k_{ab} + 2\delta_{a1} \delta_{a2} \cos^2 \alpha) \right] \right\} \Bigg) \\
& + \frac{l_2}{4} \left( \frac{4\mu_I^2}{f^2} (\partial_0 \pi_a \partial_0 \pi_a + \partial_0 \pi_2 \partial_0 \pi_2 + \partial_\mu \pi_2 \partial^\mu \pi_2) \sin^2 \alpha + 16\mu_I^3 \left[ \frac{\partial_0 \pi_2}{f} \sin^3 \alpha + \frac{1}{f^2} (3\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) \cos \alpha \sin^2 \alpha \right] \right. \\
& + 4\mu_I^4 \left\{ \sin^4 \alpha + 2 \sin^2 \alpha \left[ \frac{\pi_1}{f} \sin 2\alpha + \frac{\pi_a \pi_b}{f^2} (k_{ab} + 2\delta_{a1} \delta_{a2} \cos^2 \alpha) \right] \right\} \Bigg) \\
& + \frac{l_3 + h_1 - h_3}{16} \left( (8B_0 \bar{m})^2 \left[ \cos^2 \alpha - \frac{\pi_1}{f} \sin 2\alpha + \frac{1}{f^2} (\pi_1^2 \sin^2 \alpha - \pi_a \pi_a \cos^2 \alpha) \right] \right) \\
& + \frac{l_4}{4} \left( 8B_0 \bar{m} \left\{ 2 \frac{\partial_\mu \pi_a \partial^\mu \pi_a}{f^2} \cos \alpha + 4\mu_I \left[ \frac{\partial_0 \pi_2}{2f} \sin 2\alpha + \frac{1}{f^2} (\pi_1 \partial_0 \pi_2 \cos 2\alpha - \pi_2 \partial_0 \pi_1 \cos^2 \alpha) \right] \right. \right. \\
& \quad \left. \left. + \mu_I^2 \left[ 2 \cos \alpha \sin^2 \alpha - 2 \frac{\pi_1}{f} \sin \alpha (3 \sin^2 \alpha - 1) + \frac{1}{f^2} (\pi_1^2 [2 - 9 \sin^2 \alpha] + \pi_2^2 [2 - 3 \sin^2 \alpha] - 3\pi_3^2 \sin^2 \alpha) \cos \alpha \right] \right\} \right) \\
& + \frac{h_1 - h_3 - l_4 - l_7}{16} \left( -16 \left( \frac{2\Delta m B_0 \pi_3}{f} \right)^2 \right) + \frac{h_1 + h_2 - l_4}{4} (8B_0^2 (\bar{m}^2 + \Delta m^2)) \\
& - \frac{h_1 - h_3 - l_4}{8} \left( \left( \cos 2\alpha - 2 \frac{\pi_1}{f} \sin 2\alpha - 2 \frac{\pi_a \pi_a}{f^2} \cos^2 \alpha + 2 \frac{\pi_1^2}{f^2} \sin^2 \alpha \right) + 16B_0^2 \Delta m^2 \left( 1 - 2 \frac{\pi_3^2}{f^2} \right) \right) \\
& + \mathcal{O} \left[ t^6 \left( \frac{\pi}{f} \right)^5 \right]
\end{aligned}$$

The different terms of the NLO Lagrangian is

$$\begin{aligned}
\mathcal{L}_4^{(0)} &= \frac{l_1}{4} 4\mu_I^4 \sin^4 \alpha + \frac{l_2}{4} 4\mu_I^4 \sin^4 \alpha + \frac{l_3 + h_1 - h_3}{16} (8B_0 \bar{m})^2 \cos^2 \alpha + \frac{l_4}{8} 8B_0 \bar{m} \mu_I^2 2 \cos \alpha \sin^2 \alpha \\
&\quad + \frac{h_1 + h_3 - l_4}{4} (8B_0^2 (\bar{m}^2 + \Delta m^2)) - \frac{h_1 - h_3 - l_4}{8} \left( 16B_0^2 \bar{m}^2 \cos 2\alpha + 16B_0^2 \Delta m^2 \right) \\
&= (l_1 + l_2) \mu_I^4 \sin^4 \alpha + l_3 (2B_0 \bar{m})^2 \cos^2 \alpha + l_4 [2B_0 \bar{m} \mu_I^2 \cos \alpha \sin^2 \alpha - 2B_0^2 (\bar{m}^2 (1 - \cos 2\alpha))] \\
&\quad + h_1 [2B_0^2 \bar{m}^2 (1 - \cos 2\alpha) + (2B_0 \bar{m})^2 \cos^2 \alpha] + h_3 [2B_0^2 \bar{m}^2 (1 + \cos 2\alpha) - (2B_0 \bar{m})^2 \cos^2 \alpha] + 4B_0 \Delta m^2 \\
&= (l_1 + l_2) \mu_I^4 \sin^4 \alpha + l_3 (2B_0 \bar{m})^2 \cos^2 \alpha + l_4 [2B_0 \bar{m} \mu_I^2 \cos \alpha \sin^2 \alpha - (2B_0 \bar{m})^2] \sin^2 \alpha + h_1 (2B_0 \bar{m})^2 + h_3 (2B_0 \Delta m)^2 \\
&= (l_1 + l_2) \mu_I^4 \sin^4 \alpha + (l_3 + l_4) (2B_0 \bar{m})^2 \cos^2 \alpha + l_4 (2B_0 \bar{m}) \mu_I^2 \cos \alpha \sin^2 \alpha - l_4 (2B_0 \bar{m})^2 + h_1 (2B_0 \bar{m})^2 + h_3 (2B_0 \Delta m)^2 \\
\mathcal{L}_4^{(1)} &= \frac{l_1}{4} \left( 16\mu_I^3 \left[ \frac{\partial_0 \pi_2}{f} \sin^3 \alpha \right] + 4\mu_I^4 \left\{ 2 \sin^2 \alpha \left[ \frac{\pi_1}{f} \sin 2\alpha \right] \right\} \right) + \frac{l_2}{4} \left( 16\mu_I^3 \left[ \frac{\partial_0 \pi_2}{f} \sin^3 \alpha \right] + 4\mu_I^4 \left\{ 2 \sin^2 \alpha \left[ \frac{\pi_1}{f} \sin 2\alpha \right] \right\} \right) \\
&\quad + \frac{l_3 + h_1 - h_3}{16} \left( (8B_0 \bar{m})^2 \left[ -\frac{\pi_1}{f} \sin 2\alpha \right] \right) + \frac{l_4}{4} \left( 8B_0 \bar{m} \left\{ 4\mu_I \left[ \frac{\partial_0 \pi_2}{2f} \sin 2\alpha \right] + \mu_I^2 \left[ -2\frac{\pi_1}{f} \sin \alpha (3 \sin^2 \alpha - 1) \right] \right\} \right) \\
&\quad - \frac{h_1 - h_3 - l_4}{8} \left( 16B_0^2 \bar{m}^2 \left( -2\frac{\pi_1}{f} \sin 2\alpha \right) \right) \\
&= (l_1 + l_2) \left( 4\mu_I^3 \frac{\partial_0 \pi_2}{f} \sin^3 \alpha + \mu_I^4 2 \sin^2 \alpha \frac{\pi_1}{f} \sin 2\alpha \right) \\
&\quad - (l_3 + h_1 - h_3) (2B_0 \bar{m})^2 \frac{\pi_1}{f} \sin 2\alpha + l_4 2B_0 \bar{m} \left\{ 4\mu_I \left[ \frac{\partial_0 \pi_2}{2f} \sin 2\alpha \right] - 2\mu_I^2 \frac{\pi_1}{f} \sin \alpha (3 \sin^2 \alpha - 1) \right\} \\
&\quad + (h_1 - h_3 - l_4) (2B_0 \bar{m})^2 \frac{\pi_1}{f} \sin 2\alpha \\
&= (l_1 + l_2) \frac{1}{f} (4\mu_I^3 \partial_0 \pi_2 \sin^3 \alpha + \mu_I^4 2 \sin^2 \alpha \pi_1 \sin 2\alpha) - (l_3 + l_4) \frac{1}{f} (2B_0 \bar{m})^2 \pi_1 \sin 2\alpha \\
&\quad + l_4 2B_0 \bar{m} \frac{1}{f} [2\mu_I \partial_0 \pi_2 \sin 2\alpha - 2\mu_I^2 \pi_1 \sin \alpha (3 \sin^2 \alpha - 1)] \\
\mathcal{L}_4^{(2)} &= \frac{l_1}{4} \left( \frac{8\mu_I^2}{f^2} (\partial_\mu \pi_a \partial^\mu \pi_a + 2\partial_\mu \pi_2 \partial^\mu \pi_2) \sin^2 \alpha + 16\mu_I^3 [(3\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) \cos \alpha \sin^2 \alpha] \right. \\
&\quad \left. + 4\mu_I^4 \left\{ 2 \sin^2 \alpha \left[ \frac{\pi_a \pi_b}{f^2} (k_{ab} + 2\delta_{a1} \delta_{a2} \cos^2 \alpha) \right] \right\} \right) \\
&\quad + \frac{l_2}{4} \left( \frac{4\mu_I^2}{f^2} (\partial_0 \pi_a \partial_0 \pi_a + \partial_0 \pi_2 \partial_0 \pi_2 + \partial_\mu \pi_2 \partial^\mu \pi_2) \sin^2 \alpha + 16\mu_I^3 \left[ \frac{1}{f^2} (3\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) \cos \alpha \sin^2 \alpha \right] \right. \\
&\quad \left. + 4\mu_I^4 \left\{ 2 \sin^2 \alpha \left[ +\frac{\pi_a \pi_b}{f^2} (k_{ab} + 2\delta_{a1} \delta_{a2} \cos^2 \alpha) \right] \right\} \right) \\
&\quad + \frac{l_3 + h_1 - h_3}{16} \left( (8B_0 \bar{m})^2 \left[ +\frac{1}{f^2} (\pi_1^2 \sin^2 \alpha - \pi_a \pi_a \cos^2 \alpha) \right] \right) \\
&\quad + \frac{l_4}{4} \left( 8B_0 \bar{m} \left\{ 2\frac{\partial_\mu \pi_a \partial^\mu \pi_a}{f^2} \cos \alpha + 4\mu_I \left[ +\frac{1}{f^2} (\pi_1 \partial_0 \pi_2 \cos 2\alpha - \pi_2 \partial_0 \pi_1 \cos^2 \alpha) \right] \right. \right. \\
&\quad \left. \left. + \mu_I^2 \left[ +\frac{1}{f^2} (\pi_1^2 [2 - 9 \sin^2 \alpha] + \pi_2^2 [2 - 3 \sin^2 \alpha] - 3\pi_3^2 \sin^2 \alpha) \cos \alpha \right] \right\} \right) \\
&\quad + \frac{h_1 - h_3 - l_4 - l_7}{16} \left( -16 \left( \frac{2\Delta m B_0 \pi_3}{f} \right)^2 \right) \\
&\quad - \frac{h_1 - h_3 - l_4}{8} \left( 16B_0^2 \bar{m}^2 \left( -2\frac{\pi_a \pi_a}{f^2} \cos^2 \alpha + 2\frac{\pi_1^2}{f^2} \sin^2 \alpha \right) + 16B_0^2 \Delta m^2 \left( -2\frac{\pi_3^2}{f^2} \right) \right) \\
&= l_1 \frac{2\mu_I^2}{f^2} (\partial_\mu \pi_a \partial^\mu \pi_a + 2\partial_\mu \pi_2 \partial^\mu \pi_2) \sin^2 \alpha + l_2 \frac{4\mu_I^2}{f^2} (\partial_0 \pi_a \partial_0 \pi_a + \partial_0 \pi_2 \partial_0 \pi_2 + \partial_\mu \pi_2 \partial^\mu \pi_2) \sin^2 \alpha \\
&\quad + \frac{l_1 + l_2}{f^2} [ > 4\mu_I^3 (3\pi_1 \partial_0 \pi_2 - \pi_2 \partial_0 \pi_1) \cos \alpha \sin^2 \alpha + 2\mu_I^4 \sin^2 \alpha \pi_a \pi_b (k_{ab} + 2\delta_{a1} \delta_{a2} \cos^2 \alpha) ] \\
&\quad + \frac{l_3 + l_4}{f^2} (2B_0 \bar{m})^2 (\pi_1^2 \sin^2 \alpha - \pi_a \pi_a \cos^2 \alpha) + \frac{l_4}{f^2} 2B_0 \bar{m} \left[ 2\partial_\mu \pi_a \partial^\mu \pi_a \cos \alpha + 4\mu_I (\pi_1 \partial_0 \pi_2 \cos 2\alpha - \pi_2 \partial_0 \pi_1 \cos^2 \alpha) \right. \\
&\quad \left. + \mu_I^2 (\pi_1^2 [2 - 9 \sin^2 \alpha] + \pi_2^2 [2 - 3 \sin^2 \alpha] - 3\pi_3^2 \sin^2 \alpha) \cos \alpha \right] + \frac{l_7}{f^2} (2\Delta m B_0)^2 \pi_3^2
\end{aligned}$$

Calculating the free energy density:

$$\mathcal{F} = -f^2 \left( m_\pi^2 \cos \alpha + \frac{1}{2} \mu_I^2 \sin^2 \alpha \right) + \mathcal{F}_{\text{fin}, \pi_\pm}^{(1)} \quad (21)$$

$$- \frac{1}{2} \frac{1}{(4\pi)^2} \left[ \frac{3}{2} \left( \frac{3}{2} m_\pi^4 \cos^4 \alpha + m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha + \mu_I^4 \sin^4 \alpha \right) + \frac{1}{3} (\bar{l}_1 + 2\bar{l}_2 - 3) \mu_I^4 \sin^4 \alpha + \right. \quad (22)$$

$$\left. \frac{1}{2} (-\bar{l}_3 + 4\bar{l}_4 - 3) m_\pi^4 \cos^2 \alpha + 2 (\bar{l}_4 - 1) m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha \right] \quad (23)$$

$$- \frac{1}{2} \frac{1}{(4\pi)^2} \left[ \frac{1}{\epsilon} \left( \frac{3}{2} m_\pi^4 \cos^4 \alpha + m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha + \mu_I^4 \sin^4 \alpha \right) + \left( \ln \frac{\mu^2}{m_3^2} + \frac{1}{2} \ln \frac{\mu^2}{\tilde{m}_2^2} \right) m_\pi^3 \cos^2 \alpha \right. \quad (24)$$

$$\left. + \ln \frac{\mu}{m_3^2} (\mu_I^4 \sin^4 \alpha + 2m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha) - \left( \frac{1}{\epsilon} + \ln \frac{\mu^2}{M^2} \right) \left( \mu_I^4 \sin^4 \alpha + \frac{3}{2} m_\pi^4 \cos^2 \alpha + 2m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha \right) \right] \quad (25)$$

$$= -f^2 \left( m_\pi^2 \cos \alpha + \frac{1}{2} \mu_I^2 \sin^2 \alpha \right) + \mathcal{F}_{\text{fin}, \pi_\pm}^{(1)} \quad (26)$$

$$- \frac{1}{2} \frac{1}{(4\pi)^2} \left[ \frac{1}{3} \left( \bar{l}_1 + 2\bar{l}_2 - 3 + \frac{3^2}{2} \right) \mu_I^4 \sin^4 \alpha + \frac{1}{2} \left( -\bar{l}_3 + 4\bar{l}_4 - 3 + \frac{3^2}{2} \right) m_\pi^4 \cos^2 \alpha + 2 \left( \bar{l}_4 - 1 + \frac{3}{4} \right) m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha \right. \quad (27)$$

$$\left. (\mu_I^4 \sin^4 \alpha + m_\pi^4 \cos^2 \alpha + 2m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha) \ln \frac{M}{m_3^2} + \frac{1}{2} m_\pi^2 \cos^2 \alpha \ln \frac{M^2}{\tilde{m}_2^2} \right] \quad (28)$$

$$= -f^2 \left( m_\pi^2 \cos \alpha + \frac{1}{2} \mu_I^2 \sin^2 \alpha \right) + \mathcal{F}_{\text{fin}, \pi_\pm}^{(1)} - \frac{1}{2} \frac{1}{(4\pi)^2} \left[ \frac{1}{3} \left( \bar{l}_1 + 2\bar{l}_2 + \frac{3}{2} + 3 \ln \frac{M^2}{m_3^2} \right) \mu_I^4 \sin^4 \alpha \right. \quad (29)$$

$$\left. + \frac{1}{2} \left( -\bar{l}_3 + 4\bar{l}_4 + \frac{3}{2} + 2 \ln \frac{M^2}{m_3^2} + \ln \frac{M^2}{\tilde{m}_2^2} \right) m_\pi^4 \cos^2 \alpha + 2 \left( \bar{l}_4 - \frac{1}{4} + \frac{1}{2} \ln \frac{M^2}{m_3^2} \right) m_\pi^2 \mu_I^2 \cos \alpha \sin^2 \alpha \right]. \quad (30)$$

## 5 Forsøk på CCWZ

Goldstone's theorem tells us that if a theory is invariant under the actions of a group  $G$ , while the ground state of that theory i.e. the symmetry is broken, then there will appear massless modes. The low energy dynamics of the theory will be dominated by these modes, as they can be excited by arbitrarily small perturbation to the ground state. If we want to treat the theory perturbatively, then, the original degrees of freedom might not be the best way to treat the theory. For example, in QCD, the approximate chiral symmetries are apparent when describing the theory using the quark spinors,  $\psi_{i,\alpha}$ . However, low energy QCD is notoriously non-perturbative, due to the strong coupling of the strong force. Thus, we seek a way to be able to expand the theory in powers of the momenta of the Goldstone modes.

For concreteness, consider a theory consisting of  $N$  real fields  $\varphi_i(x)$ . The underlying fields of the theory might be complex, or Grassmann-number valued. However, in the case of  $N$  complex fields, they can be described as  $2N$  real fields, and the transformation of real Grassmann-numbers, for our purposes, follows the same rules as real numbers. We can then assume that  $G$  is a real, compact Lie Group. The action of  $g \in G$  can then be represented as a matrix  $M_{ij}(g)$  acting on  $\varphi_j$ , and infinitesimal transformations has the form  $M_{ij} = \delta_{ij} + i\epsilon n_\alpha t_{ij}^\alpha$ . Here,  $t_{ij}^\alpha$  is the generators of  $g$ , and a basis for the Lie Group corresponding to  $G$ , and  $n^\alpha$  is a normal vector.

Our goal is to remove the Goldstone-modes from the original degrees of freedom,  $\varphi_i(x)$ , by transforming them into a restricted set  $\tilde{\varphi}_i(x)$ . The condition that  $\tilde{\varphi}_i$  is

$$\tilde{\varphi}_i(x) t_{ij}^\alpha \varphi_j^* = 0, \quad (31)$$

where  $\varphi^*$  is the vacuum. If  $t^\alpha$  is not a broken generator, then this is trivially fulfilled, as  $t_{ij}^\alpha \varphi_j = 0$ . Thus, this is one constraint per broken generator.

Consider now the quantity

$$V_\varphi(g) = \varphi_i g_{ij} \varphi_j^*. \quad (32)$$

This is a continuous bounded function of  $g$ , as  $G$  is compact. This means that it has a maximum. Given an arbitrary function  $\varphi(x)$ , there is a function  $g(x)$  that maximizes  $V$  for each  $x$ . At this maximum,  $V$  is stationary, and thus invariant under a small change in  $g(x)$ ,  $\delta g(x) = i\epsilon n_\alpha g(x) t^\alpha$ . Thus,

$$\delta V_{\varphi(x)}(g(x)) = i n_\alpha \epsilon \varphi_i(x) g_{ij}(x) t_{jk}^\alpha \varphi_k = 0. \quad (33)$$

As  $n_\alpha$  is arbitrary, this gives us our transformed field,

$$\tilde{\varphi}_i(x) = \varphi_j(x) g_{ij}(x) \quad (34)$$

Let  $H \subset G$  be the non-broken group left after the broken symmetry. If this set is non-empty, then the choice of  $g_{ij}(x)$  is highly non-unique. This is because  $h \in H$ ,  $h_{ij} \varphi_j^* = \varphi_i^*$ , by definition. Thus,  $V_\varphi(gh) = V_\varphi(g)$ , and if  $g(x)$  maximizes  $V_{\varphi(x)}$ , so does  $g(x)h$ . We therefore consider  $g$  and  $gh$  equivalent. This is an equivalence relation, in the sense that it is reflective, symmetric and transitive. This partitions  $G$  into equivalence classes, where all the elements of the right coset,

$$gH = \{gh | \forall h \in H\} \quad (35)$$

are equivalent to  $g$ . The set of cosets, called the quotient group  $G/H$ , is a new group. We only need one representative element for each coset, i.e. we have a bijective function from the equivalence classes of  $g$  and  $G/H$ .

We now insert  $\varphi_i(x)$  into the Lagrangian. The original theory was invariant under global transformations  $g \in G$ . This means that any terms in the Lagrangian  $f(\varphi)$  that does not depend on derivatives of  $\varphi(x)$  only depend on  $\tilde{\varphi}(x)$ , as  $f(\varphi(x)) = f(\tilde{\varphi}(x))$ . The derivative of  $\varphi(x)$  is

$$\partial_\mu \varphi(x) = [\partial_\mu \tilde{\varphi}_i(x) + \tilde{\varphi}(x)(\partial_\mu g^{-1}(x))g(x)]g^{-1}(x), \quad (36)$$

As terms that depend on the derivatives of  $\varphi(x)$  also are invariant under a global transformation, all terms that depend on  $g$  will be proportional to at least one derivative of  $g(x)$ . Thus, there will be no mass terms, and all terms can be ordered in terms of powers of the momenta of the Goldstone bosons. The field  $g(x)$  takes on values in  $G/H$ , and can therefore be parametrized as

$$g(x) = \exp\{i\xi_a(x)t^a\}, \quad (37)$$

where  $t^a$  are the generators of  $G/H$ . These new fields  $\xi_a$  are identified with the Goldstone bosons. We see that there are one field per broken generator, as  $|G/H| = |G| - |H|$ . These fields might in general transform non-linearly under  $G$ . We may deduce their new transformation rule by the fact that we might write any transformation  $g'g(\xi(x))$  as an element in  $G/H$ , which we can write as  $g(\xi'(x))$  for some  $\xi$ , and an element in  $H$ . Thus, a transformation  $\varphi(x) \rightarrow g'\varphi(x)$  induces a transformation  $\xi \rightarrow \xi'$  defined by

$$g'g(\xi) = g(\xi')h(\xi, g) \quad (38)$$