

Electronic Raman Dynamical Structure Factor

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The non-perturbative Hamiltonian is H_0 , and the effective perturbative Hamiltonian H' is

$$\begin{aligned} H' &= -e \frac{4\pi n(\vec{q}, \omega)}{q^2} e^{-i\omega t} \sum_{\vec{k}, \sigma} \gamma_{\vec{k}} c_{\vec{k}+\vec{q}\sigma}^\dagger c_{\vec{p}\sigma} \\ &= -e \phi(\vec{q}, \omega) e^{-i\omega t} \tilde{\rho}_{\vec{q}}^\dagger \end{aligned} \quad (1)$$

where $\phi(\vec{q}, \omega) = \frac{4\pi n_t(\vec{q}, \omega)}{q^2}$ represents the potential induced by the charge $n(\vec{q}, \omega)$.

When the total Hamiltonian is $H = H_0 + H'$, the following difference is obtained by electronic Raman scattering experiments: $\Delta \tilde{\rho}_{\vec{q}} = \langle \tilde{\rho}_{\vec{q}} \rangle_H - \langle \tilde{\rho}_{\vec{q}} \rangle_{H_0}$

This is expressed using the external force $F(t) = e\phi(\vec{q}, \omega)e^{-i\omega t}$:

$$\Delta \tilde{\rho}_{\vec{q}} = - \int_{-\infty}^t dt' F(t') \chi_{\tilde{\rho}\tilde{\rho}}^{(R)}(t - t') \quad (2)$$

Here $\chi_{\tilde{\rho}\tilde{\rho}}^{(R)}$, which is often called the response function (of the retarded part), was shown by R. Kubo in the linear response regime and is known as the Kubo formula.

$$\chi_{\tilde{\rho}\tilde{\rho}}^{(R)}(t) = -i\theta(t) \langle [\tilde{\rho}_{\vec{q}}(t), \tilde{\rho}_{\vec{q}}^\dagger] \rangle \quad (3)$$

where the thermal average $\langle \cdots \rangle$ can be expanded by using the thermodynamic potential Ω and orthonormal vectors $\{|n\rangle\}$ which satisfy $H|n\rangle = E_n|n\rangle$

$$\begin{aligned} \langle [\tilde{\rho}_{\vec{q}}(t), \tilde{\rho}_{\vec{q}}^\dagger] \rangle &= \text{Tr} \left[e^{\beta(\Omega - H)} (e^{iHt} \tilde{\rho}_{\vec{q}} e^{-iHt} \tilde{\rho}_{\vec{q}}^\dagger - \tilde{\rho}_{\vec{q}}^\dagger e^{iHt} \tilde{\rho}_{\vec{q}} e^{-iHt}) \right] \\ &= e^{\beta\Omega} \sum_{m,n} e^{i(\beta E_n - E_m)t} |\langle n | \tilde{\rho}_{\vec{q}} | m \rangle|^2 (e^{-\beta E_n} - e^{-\beta E_m}) \end{aligned} \quad (4)$$

After Fourier transformed, it can be also written as

$$\chi_{\tilde{\rho}\tilde{\rho}}^{(R)}(\vec{q}, \omega) = -i \int_0^\infty dt \langle [\tilde{\rho}_{\vec{q}}(t), \tilde{\rho}_{\vec{q}}^\dagger] \rangle \quad (5)$$

Generally, the scattering probability is proportional to the dynamical structure factor in the Born approximation.

$$S(\vec{q}, \omega) = e^{\beta\Omega} \sum_{m,n} e^{-\beta E_n} |\langle n | \rho_{\vec{q}} | m \rangle|^2 \delta(E_n - E_m + \omega) \quad (6)$$

From the correspondence between $\chi_{\tilde{\rho}\tilde{\rho}}^{(R)}$ and the dynamical structure factor, we should define the electronic Raman dynamical structure factor $\tilde{S}(\vec{q}, \omega)$ as follows.

$$\begin{aligned} \tilde{S}(\vec{q}, \omega) &= e^{\beta\Omega} \sum_{m,n} e^{-\beta E_n} |\langle n | \tilde{\rho}_{\vec{q}} | m \rangle|^2 \delta(E_n - E_m + \omega) \\ &= -\frac{1 + \coth(\beta\omega/2)}{2\pi} \text{Im} \chi_{\tilde{\rho}\tilde{\rho}}^{(R)}(\vec{q}, \omega) \end{aligned} \quad (7)$$