# 广义相对论基础

- 1、流形:物理体发生的背景 上有坐标系 $Q_{\alpha}$ 、 $Q_{\beta}$ ····
- 2、张量

定义:一点处的张量(k,l)型为 $n^{k+l}$ 个数,由 k 个上指标,l 个下指标区分。

$$T_{\nu_1\cdots\nu_l}^{\mu_1\ldots\mu_k}$$

坐标变换满足:

$$T_{\nu_1\cdots\nu_k}^{\prime\mu_1\cdots\mu_k} = \prod \left( \frac{\partial x^{\prime\mu_i}}{\partial x^{\sigma_i}} \frac{\partial x^{\rho_i}}{\partial x^{\prime\nu_j}} \right) \quad T_{\rho_1\ldots\rho_j}^{\sigma_1\cdots\sigma_k}$$

例子:

矢量:(1,0)型 V<sup>µ</sup>

坐标变化:

$$V^{\prime\mu} = \frac{\partial x^{\prime\mu}}{\partial x^{\sigma}} V^{\sigma}$$

对偶矢量:(0,1)型

坐标变化:

$$\omega_{\nu}' = \frac{\partial x^{\rho}}{\partial x'^{\nu}} \omega_{\rho}$$

直积:

$$C^{\mu}_{\nu} = A^{\mu}B_{\nu}$$

缩并:

$$V^{\sigma} = T_{\rho}^{\sigma\rho}$$

3、切矢

对于曲线 x=x(t),

其切矢:

$$\left(\frac{\partial}{\partial t}\right)^{\mu} = \frac{dx^{\mu}}{dt}$$

以上均谈及分量,若谈及矢量本身,可将希腊字母改为英文字母。

坐标线的切矢:坐标基矢: $\left(\frac{\partial}{\partial x^{\mu}}\right)^{a}$ 

可写为:

$$T^{a} \equiv \left(\frac{\partial}{\partial t}\right)^{a} = \frac{dx^{\mu}}{dt} \left(\frac{\partial}{\partial x^{\mu}}\right)^{a}$$

4、度规张量

 $g_{ab}$  (0,2)型对称张量,非退化(det  $(g_{\mu\nu}) \neq 0$ )

度规定义类似内积, 但不一定正定

定义矢量 , 边内积为:

$$g_{\mu\nu}u^{\mu}v^{\nu}$$

定义矢量 设长度:

$$|v| := \sqrt{|g(v, v)|}$$

正交归一基底 $(e_{\mu})^a$ :

$$g((e_{\mu})^a, (e_{\nu})^b) = \pm \delta_{\mu\nu}$$

不同基底下将 g 对角化,其号差相同。 全为+1:正定度规 1 个为-1:洛伦兹度规

洛伦兹度规:号差为2, 即 diag{-1,1,1,1}

g(v,v)<0: 类时矢量 =0 类光 >0 类空

曲线长度:

$$l := \int |T| dt = \int \sqrt{|g(T,T)|} dt$$

参数t有重参数化不变性。

$$l = \int \sqrt{|g_{\mu\nu} dx^{\mu} dx^{\nu}|}$$

记线元:

$$ds^2 \equiv g_{\mu\nu} dx^\mu dx^\nu$$

有:

$$l = \int \sqrt{|ds^2|}$$

欧氏度规

$$g_{\mu\nu} = \delta_{\mu\nu}$$

n=2:

$$ds^2 = \delta_{\mu\nu} dx^{\mu} dx^{\nu} = dx^2 + dy^2$$

闵氏度规:

$$g_{\mu\nu} = \eta_{\mu\nu}$$

$$\eta_{\mu\nu} = 0, \qquad \mu \neq \nu$$
 $\eta_{\mu\nu} = -1, \qquad \mu = \nu = 0$ 
 $\eta_{\mu\nu} = 1, \qquad else$ 

n=4:

$$ds^2 = -(dx^0)^2 + (dx^1)^2 + (dx^2)^2 + (dx^3)^2$$

度规张量 指标升降:

$$\det\left(g_{\mu\nu}\right)\neq0$$

度规有其逆: $g^{\mu\nu}$ 

矢量:

$$V^{\mu}g_{\sigma\mu} \equiv V_{\rho}$$

## 5、联络

张量的导数非张量

希望:

$$\tilde{V}\big|_q^\mu - V^\mu|_p = -\Gamma^\mu_{\sigma\nu} V^\mu\big|_p dx^\nu$$

为保证 $\tilde{V}|_{a}^{\mu}$ 为矢量,可用坐标变换来检查:

$$\left. \tilde{V} \right|_q^\mu = \left( \frac{\partial x'^\mu}{\partial x'^\nu} \right)_q \left. \tilde{V}^\nu \right. \mid_q$$

可给出对于联络的约束:

$$\Gamma_{\mu\nu}^{\prime\tau} = \frac{\partial^{2}x^{\beta}}{\partial x^{\prime\mu} \partial x^{\prime\gamma}} \frac{\partial x^{\prime\tau}}{\partial x^{\beta}} + \Gamma_{\alpha\sigma}^{\beta} \frac{\partial x^{\alpha}}{\partial x^{\prime\mu}} \frac{\partial x^{\prime\tau} \partial x^{\sigma}}{\partial x^{\beta} \partial x^{\prime\gamma}}$$

注意到 $\Gamma_{\mu\nu}^{\tau}$ 并非张量。

### 6、协变导数∇μ

为使张量导数仍为张量,需使用协变导数:

$$\nabla_{\lambda}V^{\mu} = \lim_{q \to p} \frac{V^{\mu} |q - \tilde{V}^{\mu}| q}{\Delta x^{\lambda}} = \partial_{\lambda}V^{\mu} + \Gamma^{\mu}_{\sigma\lambda}V^{\sigma}$$

利用: $\nabla_{\mu}(f) := \partial_{\mu}f$ 

可推得对偶矢量协变导数:

$$\nabla_{\mu}w_{\nu} = \partial_{\mu}w_{\nu} - \Gamma^{\nu}_{\mu\sigma}w_{\nu}$$

张量的协变导数:

$$T^{\mu}_{\nu;\lambda} = T^{\mu}_{\nu,\lambda} + \Gamma^{\mu}_{\rho\lambda} T^{\rho}_{\nu} - \Gamma^{\rho}_{\nu\lambda} T^{\mu}_{\rho}$$

### 7、张量沿曲线的导数

标量场:

$$\frac{df}{dt} = \frac{dx^{\mu}}{dt} \frac{\partial f}{\partial x^{\mu}} = T^{\mu} \partial_{\mu} f = T^{\mu} \nabla_{\mu} f$$

定义矢量沿 $\zeta(t)$ 导数:

$$T^{b}\nabla_{b}v^{a} = \lim_{\Delta t \to 0} \frac{1}{\Delta t} (v^{q}|_{q} - \tilde{v}^{q}|_{q})$$

#### 8、测地线

曲线切矢沿曲线平移。

$$T^b \nabla_b T^a = 0$$

分量方程:

$$\frac{d^2x^{\mu}}{dt^2} + \Gamma^{\mu}_{\nu\sigma} \frac{dx^{\nu}}{dt} \frac{dx^{\sigma}}{dt} = 0$$

#### 9、适配导数算子

对于流形 M:

联络空间

$$(M, \nabla_a)$$

广义黎曼空间:

$$(M, g_{ab})$$

希望二者适配:

度规相容条件:

$$\begin{split} \nabla_c g_{ab} &= 0 \\ g_{\mu\nu,\lambda} - g_{\mu\gamma} \Gamma^{\gamma}_{\nu\lambda} - g_{\alpha\nu} \Gamma^{\alpha}_{\mu\lambda} &= 0 \end{split}$$

解出:

$$\Gamma^{\sigma}_{\mu\nu} = \frac{1}{2} g^{\sigma\rho} (g_{\nu\rho,\mu} + g_{\mu\rho,\nu} - g_{\mu\nu,\rho})$$

## 10、测量理论

定义质点4速度:

$$U^{a} \equiv \left(\frac{\partial}{\partial \tau}\right)^{a}$$
$$U^{a}U_{a} = -1$$

再取一观测者: $t,x^i$ 

$$U^a = \frac{dt}{d\tau} \left(\frac{\partial}{\partial \tau}\right)^a + \left(\frac{dx^i}{d\tau}\right) \left(\frac{\partial}{\partial x^i}\right)^a$$

观者4速度:

$$z^a = (1,0,0,0) = \left(\frac{\partial}{\partial t}\right)^a$$

定义4动量:

$$p^a := mU^a$$

 $z^a$ 测量到的动量与能量:

$$E = -p^a z_a = -mu^a z_a = -m\eta_{00} \frac{dt}{d\tau} \cdot 1 = \gamma m$$

$$p^{a} = mU^{a} - Ez^{a} = \gamma m(1, u^{a}) - \gamma m(1, \vec{0}) = \gamma mu^{a}$$

能动量关系:

$$p^{a}p_{a} = (Ez^{a} + p^{a})(Ez_{a} + p_{a}) = p^{2} - E^{2}$$
  
 $p^{a}p_{a} = m^{2}U^{a}U_{a} = -m^{2}$ 

得到:

$$E^2 = p^2 + m^2$$