

# Formation of two-dimensional electron gas at AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructure and the derivation of its sheet density expression\*

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Models for calculating the sheet densities of two-dimensional electron gas (2DEG) induced by spontaneous and piezoelectric polarization in AlGa<sub>N</sub>/Ga<sub>N</sub>, AlGa<sub>N</sub>/Al<sub>N</sub>/Ga<sub>N</sub>, and Ga<sub>N</sub>/AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures are provided. The detailed derivation process of the expression of 2DEG sheet density is given. A longstanding confusion in a very widely cited formula is pointed out and its correct expression is analyzed in detail.

**Keywords:** high electron mobility transistors, Ga<sub>N</sub>, two-dimensional electron gas, polarization effect

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## 1. Introduction

High electron mobility transistors (HEMTs) based on AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructures have been widely studied and have exhibited outstanding performances in high-power and high-frequency devices.<sup>[1]</sup> Unlike conventional HEMTs made of modulation-doped semiconductor heterostructures, no intentional doping is required to generate two-dimensional electron gas (2DEG) at the AlGa<sub>N</sub>/Ga<sub>N</sub> interface.<sup>[2]</sup> The main problem that hinders the further development of AlGa<sub>N</sub>/Ga<sub>N</sub> HEMT is a lack of understanding the origin of 2DEG. There is no doubt that the large interface charge density in this system is correlated to the spontaneous and piezoelectric polarization and the resulting high electric field in the AlGa<sub>N</sub> barrier layer,<sup>[3]</sup> but the source of the large quantity of electrons is controversial. Among several different theoretical models to explain this issue, the most widely accepted one is that the 2DEG originates from the donor states on the AlGa<sub>N</sub> surface.<sup>[4]</sup> However, some details of the process of the formation of 2DEG are still somewhat unclear.

Also, an accurate expression of the sheet carrier concentration of 2DEG  $n_s$  is critical to the further development of AlGa<sub>N</sub>/Ga<sub>N</sub> HEMTs. There is an important and very widely cited equation reported for calculating the 2DEG concentration  $n_s$ <sup>[5]</sup> (see Eq. (4) in this article, which will be discussed later in more detail). However, according to our investigation, there is a serious confusion in the expression of this equation, especially about the term of polarization sheet charge density  $\sigma$ . In this widely cited equation, 2DEG is considered to be caused by the net polarization charge of AlGa<sub>N</sub> and Ga<sub>N</sub> (i.e.,  $\sigma_{\text{pol}} = \sigma_{\text{AlGa}_N} - \sigma_{\text{Ga}_N}$ ). However, according to our research, 2DEG is induced by the polarization of the AlGa<sub>N</sub> layer, and the term  $\sigma_{\text{Ga}_N}$  should not be added into the expres-

sion. Usually the term of polarization sheet charge density  $\sigma$  is regarded as a parameter to be determined, which means that the value of  $\sigma$  should be calculated by fitting to experiment results, rather than defined before fitting. So the different understanding of the meaning of  $\sigma$  will not affect the simulation results. However, a misconception or confusion of the expression may hamper the further understanding of the properties of AlGa<sub>N</sub>/Ga<sub>N</sub> structure.

In this work, we carefully analyze the process of 2DEG formation and provide a possible physical picture of the AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructure. Also, a derivation of the 2DEG density expression is given in detail to prove its validity that it is the right expression.

## 2. Modelling

### 2.1. Formation of the 2DEG in AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructure

In many studies concerning 2DEG, the existence of 2DEG is regarded as an *a priori* assumption.<sup>[3,6,7]</sup> There is no doubt about the existence of 2DEG in AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructure, but being aware of the existence of 2DEG is not helpful for understanding the mechanism of its formation. By splitting the whole heterostructure into two independent parts, analyzing separately and finally assembling together, we can see clearer physical pictures of AlGa<sub>N</sub>/Ga<sub>N</sub> heterostructure and the process of the formation of 2DEG. All structures discussed in this work are pseudomorphic heterostructures with Ga-face polarity.

The polarization effect of an AlGa<sub>N</sub> layer grown on a Ga<sub>N</sub> buffer induces positive polarization charge at the AlGa<sub>N</sub>/Ga<sub>N</sub> interface and negative polarization charge at the top of the AlGa<sub>N</sub> layer. Thus an electric field is formed within

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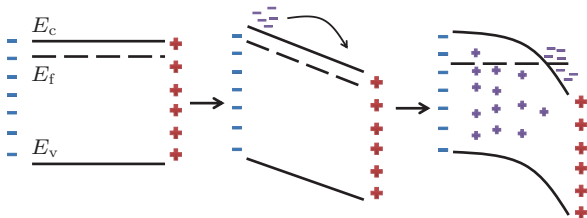
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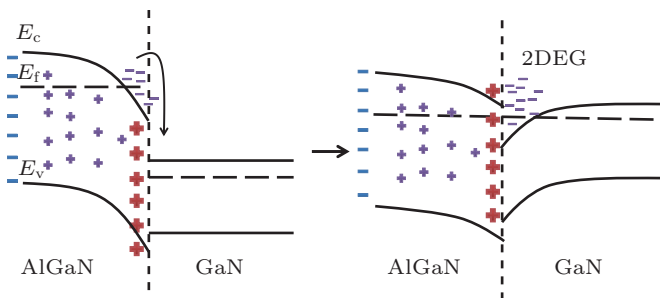
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the AlGaIn layer. It is reasonable to regard the AlGaIn surface and the AlGaIn/GaN interface as two infinite charged planes, as the thickness of the whole epilayer is just several  $\mu\text{m}$ , which is very small in comparison to the lateral size of the substrate.

The formation of 2DEG of an AlGaIn/GaN heterostructure with an n-doped AlGaIn layer is easy to understand. In the AlGaIn layer, the positive polarization charge at the AlGaIn/GaN interface and the negative polarization charge at the top of the AlGaIn layer induce an electric field and make the energy band (along with the Fermi level) tilt toward the interface. This situation is equal to applying a voltage to a piece of thin freestanding AlGaIn film, in other words, a planar plate capacitor with AlGaIn dielectric medium as shown in Fig. 1. Under the force of the electric field, conducting electrons (note that the AlGaIn layer is n-doped) in AlGaIn move to the positive electrode and accumulate there, leaving positive space charges in AlGaIn, bending the energy band and making the Fermi level flat again. Also, the built-in electric field will reduce the whole electric field in the AlGaIn layer and thus make the inclined energy band flatter. Once contacting GaN, electrons on the AlGaIn side will flow into the GaN and form 2DEG, because the Fermi level of GaN is lower than that of AlGaIn as shown in the left part of Fig. 2. This process will continue until the Fermi levels of AlGaIn and GaN become equal to each other (the right part of Fig. 2).



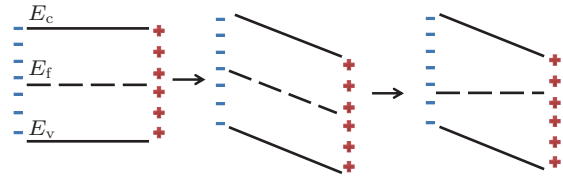
**Fig. 1.** (color online) Energy band of n-doped freestanding AlGaIn film (assuming the AlGaIn layer to be under the same tensile strain as that when growing on GaN). Conductive electrons move under the polarization-induced electric field from one side to the other, thereby bending the energy band, and leaving fixed positive charges in the AlGaIn layer.



**Fig. 2.** (color online) Energy band of an n-doped AlGaIn/GaN heterostructure. When contacting a GaN layer, electrons will flow into the GaN side, accumulate at the interface and form 2DEG.

However, as for the case of undoped AlGaIn, the situation is different. As shown in Fig. 3, in an undoped AlGaIn/GaN

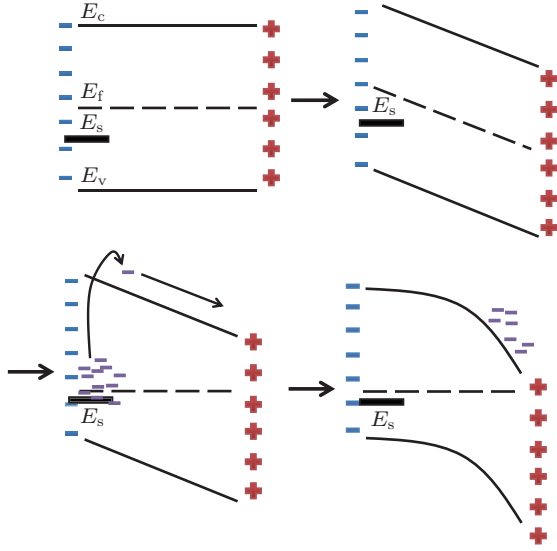
heterostructure, few conductive electrons will move under the force of the electric field and the energy band will almost remain unchanged because of the lack of conductive electrons. When electrons contact the GaN, Fermi levels on the sides are close to each other and even the Fermi level of GaN may be higher, so very few or even no electrons will flow into the GaN side and no 2DEG will form. However, this does not accord with the fact that no intentional doping is required to generate 2DEG at the AlGaIn/GaN interface.



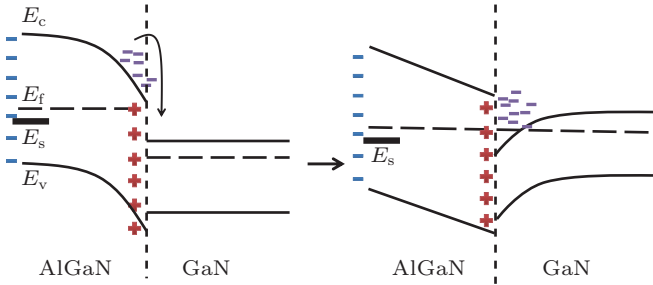
**Fig. 3.** (color online) Energy band of undoped freestanding AlGaIn on the assumption that the AlGaIn layer is under the same tensile strain as that when growing on GaN. The polarization-induced electric field will not cause curvature of the energy band because of the lack of conductive electrons.

Then where does the 2DEG come from? First, the 2DEG cannot come from the GaN buffer. The total polarization charge at the AlGaIn/GaN interface can be divided into two parts. One part comes from AlGaIn, and the other comes from GaN. The AlGaIn part will have no effect on the GaN side because the negative polarization charge at the top of the AlGaIn layer and the positive polarization charge at the interface forms a planar plate capacitor, which generates no electric field outside its two planes. Similarly, the GaN part, together with its counterpart at the bottom of the GaN layer, generates a uniform electric field pointing from the bottom of the GaN layer toward the AlGaIn/GaN interface. This electric field will reduce the net electron density at the AlGaIn/GaN interface, let alone form 2DEG. Second, the 2DEG cannot come from the body of the AlGaIn barrier because there are few conductive electrons in undoped AlGaIn as mentioned before.

The formation of the 2DEG at the AlGaIn/GaN interface can be well explained by assuming the existence of donor states on the AlGaIn surface.<sup>[4]</sup> Figure 4 shows the energy band of a piece of freestanding AlGaIn with surface donor states on the assumption that the AlGaIn layer is under the same tensile strain as grown on GaN. If the AlGaIn layer is thick enough, the Fermi level will reach the donor state level  $E_s$ , and the electrons will be stimulated into the conductive band and swiped toward the other side by the force of the polarization-induced electric field. Once contacting with a GaN layer, electrons will flow into the GaN side because of the drop of the Fermi level. The electrons will accumulate at the interface and form 2DEG as shown in Fig. 5. The 2DEG and the ionized surface donor states will generate an electric field pointing from the interface to the surface and reduce the polarization field in the AlGaIn layer.



**Fig. 4.** (color online) Energy band of undoped freestanding AlGaIn (taking surface donor states into account). Electrons are stimulated into the conductive band and move under the force of the polarization-induced electric field.



**Fig. 5.** (color online) Energy band of an undoped AlGaIn/GaN heterostructure. When contacting a GaN layer, electrons will flow into the GaN side, accumulate at the interface and form 2DEG.

## 2.2. Equations

### 2.2.1. The 2DEG density of AlGaIn/GaN heterostructure

Based on the above discussion, a profile of the charge distribution of an AlGaIn/GaN heterostructure is depicted in Fig. 6. There are three dipoles. First,  $\pm\sigma_{\text{AlGaIn}}$ , the AlGaIn polarization-induced charge; second,  $\pm\sigma_{\text{GaN}}$ , the GaN polarization-induced charge; third, the 2DEG and the ionized surface charge  $\sigma_s$ . These three dipoles are equivalent to three planar plate capacitors  $C_1$ ,  $C_2$ , and  $C_3$  respectively. Each of these capacitors only has a nonzero electric field between its two parallel planes and has no effect on electrons outside its inner body. The internal electric field intensity of a planar plate capacitor is given by  $E = \sigma/\epsilon$ , where  $\sigma$  is the quantity of electric charges on opposite planes,  $\epsilon$  is the dielectric constant of the dielectric medium of the capacitor. In the area where two or more pairs of charged planes overlap, the net electric field is the algebraic sum. Thus, the electric field intensity in the AlGaIn layer is determined by the charges of  $C_1$  and  $C_2$ , i.e., the AlGaIn polarization-induced charge and the 2DEG:

$$\frac{\Delta V_{\text{AlGaIn}}}{d} = \frac{\sigma_{\text{AlGaIn}}}{\epsilon} - \frac{qn_s}{\epsilon}, \quad (1)$$

where  $\Delta V_{\text{AlGaIn}}$  is the potential drop across the conductive band of the AlGaIn surface and the AlGaIn/GaN interface (see Fig. 7),  $d$  is the thickness of the AlGaIn layer,  $\sigma_{\text{AlGaIn}}$  is the sheet charge density that is induced by AlGaIn polarization,  $n_s$  is the sheet density of 2DEG,  $q$  is the electron charge, and  $\epsilon$  is the dielectric constant of AlGaIn. The potential drop  $\Delta V_{\text{AlGaIn}}$  can be expressed as

$$q\Delta V_{\text{AlGaIn}} = q\phi_b + E_f - \Delta E_c, \quad (2)$$

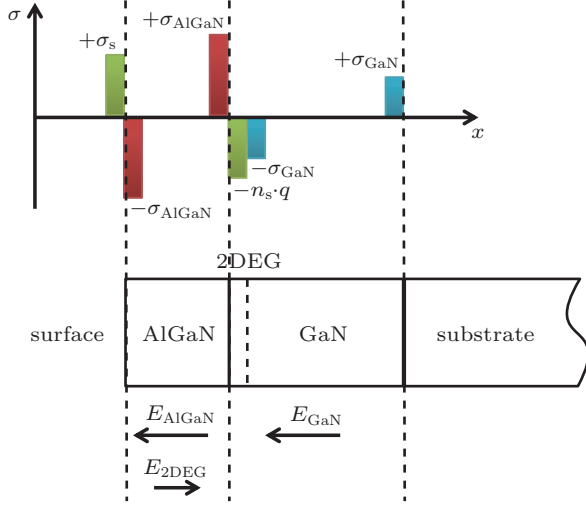
where  $\phi_b$  is the surface barrier height,  $E_f$  is the Fermi level position with respect to the GaN conduction-band edge at the AlGaIn/GaN interface, and  $\Delta E_c$  is the conduction band discontinuity between GaN and AlGaIn as shown in Fig. 7. Therefore, based on Eqs. (1) and (2), the sheet density of 2DEG can be expressed as

$$n_s = \frac{\sigma_{\text{AlGaIn}}}{q} - \frac{\epsilon}{q^2 d} (q\phi_b + E_f - \Delta E_c). \quad (3)$$

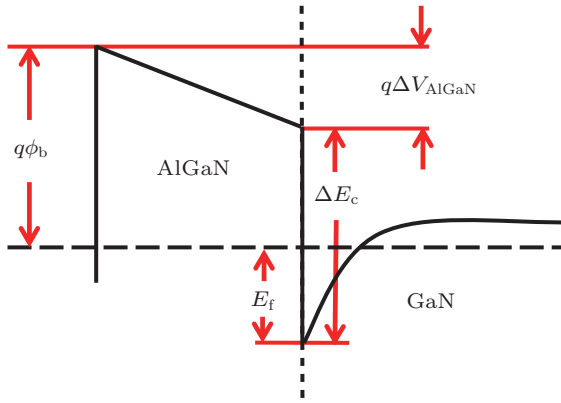
This equation is different from the one that has been widely cited<sup>[5]</sup>

$$n_s = \frac{\sigma_{\text{pol}}}{q} - \frac{\epsilon}{q^2 d} (q\phi_b + E_f - \Delta E_c), \quad (4)$$

where  $\sigma_{\text{pol}}$ , the net polarization-induced charge density at the surface, is given by  $\sigma_{\text{pol}} = \sigma_{\text{AlGaIn}} - \sigma_{\text{GaN}}$ . However, as discussed before,  $\sigma_{\text{GaN}}$ , the polarization charge induced by GaN has no effect on the potential drop of the AlGaIn layer, and  $\sigma_{\text{AlGaIn}}$  is the contributor to the formation of 2DEG. According to our investigation, there was confusion in expressing the polarized sheet charge density  $\sigma_{\text{pol}}$  in this equation, which was presented in a review article<sup>[5]</sup> and later widely cited by other research groups.<sup>[4,6–8]</sup> Actually, the expression in Ref. [5] was also cited from another paper,<sup>[9]</sup> which provides a 2DEG expression of the AlGaIn/GaN heterostructure for the first time, to the best of our knowledge. In this article the authors did not take  $\sigma_{\text{GaN}}$  into account and gave out an expression that was just the same as Eq. (3). For some unknown reason, in later citations,<sup>[4–8]</sup>  $\sigma_{\text{AlGaIn}}$  was replaced by  $\sigma_{\text{pol}}$  (equals  $\sigma_{\text{AlGaIn}} - \sigma_{\text{GaN}}$ ) in the  $n_s$  expression. It is also not known why in another article<sup>[10]</sup> written by the authors of the same research group as that in Ref. [9], the first term of the expression of the net polarization sheet charge density at the interface was also changed into the same one that was presented in Ref. [5], i.e.,  $\sigma_{\text{pol}}$ , rather than  $\sigma_{\text{AlGaIn}}$ , was used in the  $n_s$  expression. Unfortunately, there was no detailed derivation of this equation in these articles, so we actually do not know the real derivation process about how the authors obtained it. However, based on our analysis and discussion, we believe that the correct form of the expression should be the one as described by Eq. (3), i.e., the first term should be  $\sigma_{\text{AlGaIn}}$ , instead of  $\sigma_{\text{pol}}$ .



**Fig. 6.** (color online) Charge distribution profile of an AlGaIn/GaN heterostructure.



**Fig. 7.** (color online) Conductive band diagram of an AlGaIn/GaN heterostructure.

### 2.2.2. The 2DEG density of AlGaIn/AlN/GaN and GaN/AlGaIn/GaN heterostructure

Now an AlN interfacial layer<sup>[11]</sup> or a GaN cap layer<sup>[12]</sup> is often to be added into the AlGaIn/GaN HEMT structure to achieve a better performance. The expression of 2DEG sheet density  $n_s$  of these two structures is easy to obtain based on the above discussion.

The profile of the charge distribution of an AlGaIn/AlN/GaN heterostructure is shown in Fig. 8. The charges at the surface and interfaces can be regarded as four pairs of infinite charged planes with an equal quantity of heterogeneous charges. Each pair of planes only has a nonzero electric field within its two charged planes ( $E = \sigma/\epsilon$ ) and has no effect on the outside region. In the area where two or more pairs of charged planes overlap, the electric field is the algebraic sum. Thus we can obtain the expression of the potential energy drop of the AlGaIn layer and the AlN interfacial layer in a similar way

$$\frac{\Delta V_{\text{AlGaIn}}}{d_{\text{AlGaIn}}} = \frac{\sigma_{\text{AlGaIn}}}{\epsilon_{\text{AlGaIn}}} - \frac{qn_s}{\epsilon_{\text{AlGaIn}}}, \quad (5)$$

$$\frac{\Delta V_{\text{AlN}}}{d_{\text{AlN}}} = \frac{\sigma_{\text{AlN}}}{\epsilon_{\text{AlN}}} - \frac{qn_s}{\epsilon_{\text{AlN}}}. \quad (6)$$

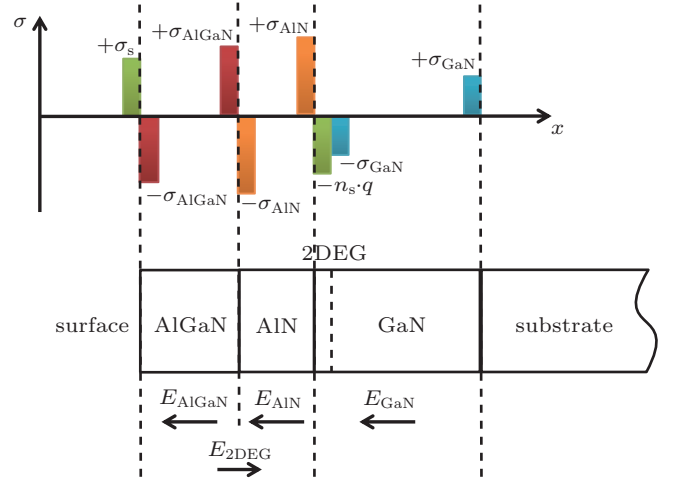
As shown in Fig. 9, the potential drop of the AlGaIn layer across the surface and the AlGaIn/AlN interface can be expressed as

$$q\Delta V_{\text{AlGaIn}} = q\phi_b + E_f - \Delta E_{c,\text{eff}}, \quad (7)$$

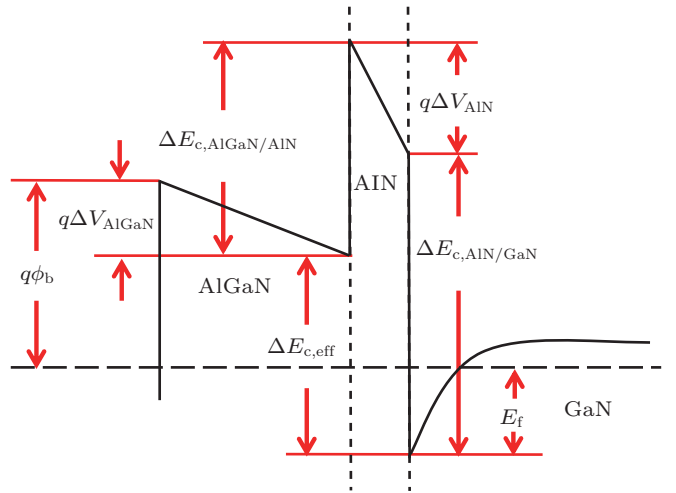
where the effective conductive band gap  $\Delta E_{c,\text{eff}}$  can be given as

$$\Delta E_{c,\text{eff}} = q\Delta V_{\text{AlN}} + \Delta E_{c,\text{AlN/GaN}} - \Delta V_{c,\text{AlGaIn/AlN}}, \quad (8)$$

where  $\Delta E_{c,a/b}$  is the conductive band gap between material a and b.



**Fig. 8.** (color online) Charge distribution profile of an AlGaIn/AlN/GaN heterostructure.



**Fig. 9.** (color online) Conductive band diagram of an AlGaIn/AlN/GaN heterostructure.

From Eqs. (5)–(8), we can obtain the expression of the sheet density of 2DEG of AlGaIn/AlN/GaN as

$$n_s = \frac{1}{q(d_{\text{AlGaIn}} + d_{\text{AlN}})} \left[ d_{\text{AlGaIn}} \sigma_{\text{AlGaIn}} + d_{\text{AlN}} \sigma_{\text{AlN}} - \frac{\epsilon}{q} (q\phi_b + E_f - \Delta E_{c,\text{AlN/GaN}} + \Delta E_{c,\text{AlGaIn/AlN}}) \right], \quad (9)$$

where  $\varepsilon$  is the dielectric constant of GaN, AlN, and AlGaIn on the assumption that their dielectric constants are equal to each other,  $d_{\text{AlGaIn}}$  and  $d_{\text{AlN}}$  are the thickness values of the AlGaIn and AlN layer respectively,  $\sigma_{\text{AlGaIn}}$  and  $\sigma_{\text{AlN}}$  are the polarization (including spontaneous and piezoelectric) charges induced by AlGaIn and AlN layers, respectively.

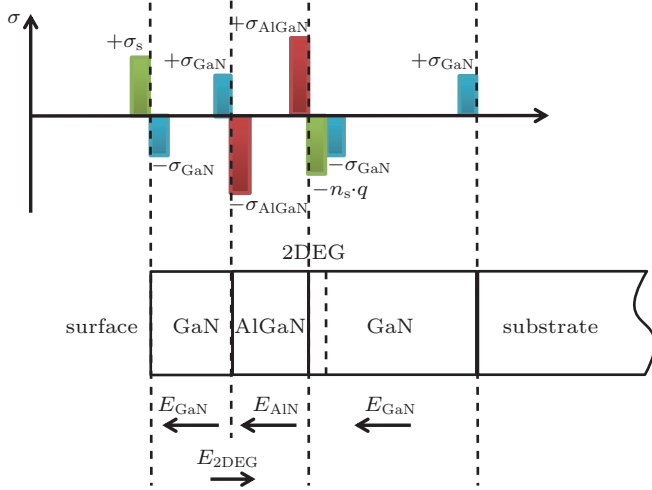


Fig. 10. (color online) Charge distribution profile of a GaN/AlGaIn/GaN heterostructure.

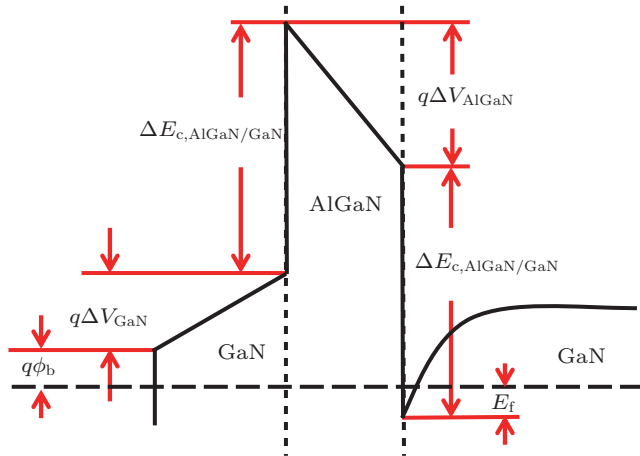


Fig. 11. (color online) Conductive band diagram of a GaN/AlGaIn/GaN heterostructure.

Similarly, we can obtain the expression of the sheet density of 2DEG of GaN/AlGaIn/GaN as follows:

$$n_s = \frac{1}{q(d_{\text{GaN}} + d_{\text{AlGaIn}})}$$

$$\times \left[ d_{\text{GaN}} \sigma_{\text{GaN}} + d_{\text{AlGaIn}} \sigma_{\text{AlGaIn}} - \frac{\varepsilon}{q} (q\phi_b + E_f) \right]. \quad (10)$$

The charge distribution and the energy band figure are shown in Figs. 10 and 11.

### 3. Summary and conclusions

In summary, according to the surface state model, we analyzed the formation of 2DEG in an AlGaIn/GaN heterostructure. A detailed derivation of the expression of 2DEG sheet density of the AlGaIn/GaN heterostructure is provided. A longstanding confusion about the expression of  $n_s$  is analyzed in detail. It is proven that the polarization charge in the 2DEG formula should be  $\sigma_{\text{AlGaIn}}$ , the spontaneous and piezoelectric polarization-induced sheet charge of AlGaIn, i.e., the net polarization charge at the AlGaIn surface ( $\sigma_{\text{AlGaIn}} - \sigma_{\text{GaN}}$ ), instead of  $\sigma_{\text{pol}}$ . The exact expressions of  $n_s$  of AlGaIn/AlN/GaN and GaN/AlGaIn/GaN heterostructures are also derived in a corresponding way.

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