

# Limitations on rotation and radius of NS

A presentation based on Glendenning's Compact Stars

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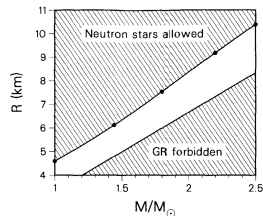
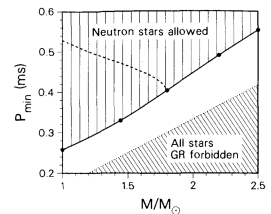
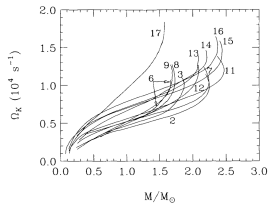
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The structure of compact stars

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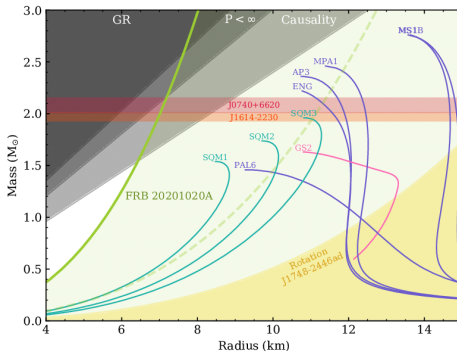
## Types of limiting factors of rotation and radius

- Limits for single stars (e.g. Kepler frequency, gravitational wave instabilities etc.)
- Absolute limits (e.g. General Relativity, Le Chateliér's principle, Causality etc.)



# Absolute limits

- General relativity: the Einstein's equations should hold for the stellar structure
- Le Chatelier's principle: for small distortions, an NS in equilibrium will bounce back to an equilibrium ( $dp/d\rho \geq 0$ )
- Causality: Speed of sound in the NS should not be larger than the speed of light ( $\sqrt{dp/d\epsilon} \leq 1$ )



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# Kepler frequency in GR

- The Kepler angular velocity or Kepler frequency for short is the absolute limit for a star to remain stable. A rotational velocity above the Kepler frequency causes the star to start shredding along its equator due to centrifugal forces.
- The Newtonian equilibrium of gravitational and centrifugal force is

$$\frac{Mm}{R^2} = m\Omega^2 R \quad (1)$$

- A self-consistency condition for a fluid element moving along the equator of the star with  $V$  velocity with respect to the ZAMO:

$$\Omega_K = e^{\nu-\mu} \frac{V}{R + \omega(R)} \quad (2)$$

# Kepler frequency in GR

- To only unknown in Eq. (2) is the  $V$  velocity of a fluid element along the equator. This can be determined by finding the extremum of

$$d\tau = \int_{t_1}^{t_2} dt \sqrt{e^{2\nu} - r^2 e^{2\mu} (\Omega - \omega(R))^2} \quad (3)$$

- After solving the variational problem and expanding  $V$ , we get the solution as

$$V_{\pm} = \frac{R\omega'}{2\psi'} e^{\mu-\nu} \pm \sqrt{\frac{\nu'}{\psi'} + \left( \frac{R\omega'}{2\psi'} e^{\mu-\nu} \right)^2}, \quad (4)$$

where  $\psi' = \mu' + 1/R$

- The physically interesting result is given by  $V_+$ , since that is the co-rotating case needed for Eq. (2).



# Approximations of the Kepler frequency

- Approximating the rotating object with the Schwarzschild metric, which gives the same result as the Hartle–Thorn perturbative solution:

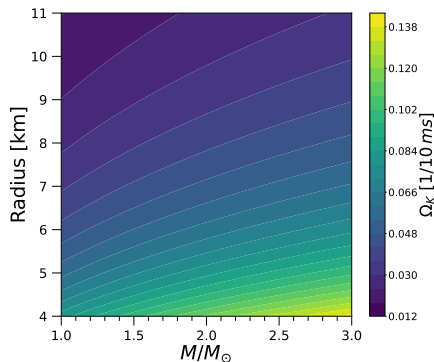
$$\Omega_K \approx \left[ 1 + \frac{\omega(R)}{\Omega_K} - 2 \left( \frac{\omega(R)}{\Omega_K} \right)^2 \right]^{-1/2} \sqrt{\frac{M}{R^3}} \quad (5)$$

- Approximating with a prefactor times the Newtonian formula:

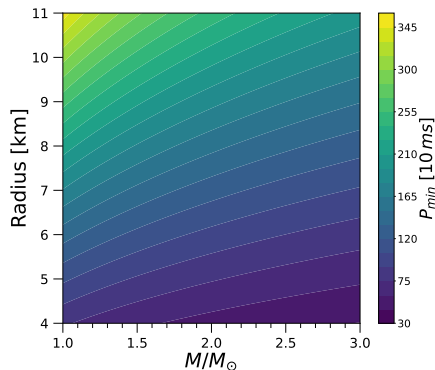
$$\Omega_K \approx 0.65 \cdot \left( \frac{M}{R^3} \right)^{1/2} \quad (6)$$



# Approximations of the Kepler frequency



**Figure 1:** Contour plot of the  $\Omega_K$  Kepler frequency calculated from the approximation in Eq. (6).



**Figure 2:** Contour plot of the  $P_{\min}$  minimal rotation period calculated from the approximation in Eq. (6).