基礎方程式:

$$\partial_t \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \boldsymbol{u} = -\nabla P - \frac{1}{\text{Re}_0} (\boldsymbol{u} - U_y \hat{\boldsymbol{y}}) + \frac{1}{\text{Re}} \nabla^2 \boldsymbol{u} + \sin(ny) \hat{\boldsymbol{x}},$$
$$\nabla \cdot \boldsymbol{u} = 0,$$
$$\boldsymbol{U} = \frac{a}{4\pi^2} \int dx dy \boldsymbol{u} = (0, U_y) \quad : \text{const.}$$

新しい変数:

$$\boldsymbol{v} = \boldsymbol{u} - \boldsymbol{U} = \nabla \times (\psi \hat{\boldsymbol{z}}) = (\partial_y \psi) \hat{\boldsymbol{x}} + (-\partial_x \psi) \hat{\boldsymbol{y}}$$

変換した基礎方程式:

$$\partial_t \boldsymbol{v} + \boldsymbol{v} \cdot \nabla \boldsymbol{v} + \boldsymbol{U} \cdot \nabla \boldsymbol{v} = -\nabla P - \frac{1}{\text{Re}_0} \boldsymbol{v} + \frac{1}{\text{Re}} \nabla^2 \boldsymbol{v} + \sin{(ny)} \hat{\boldsymbol{x}},$$

$$\nabla \cdot \boldsymbol{v} = 0,$$

$$\frac{a}{4\pi^2} \int dx dy \boldsymbol{v} = 0$$

∇× を掛けたい:

$$\nabla \times \boldsymbol{v} = \nabla \times \nabla \times (\psi \hat{\boldsymbol{z}}) = \nabla(\nabla \cdot \psi \hat{\boldsymbol{z}}) - \nabla^2(\psi \hat{\boldsymbol{z}}) = (-\nabla^2 \psi) \hat{\boldsymbol{z}}$$
$$\nabla \times (\boldsymbol{v} \cdot \nabla \boldsymbol{v}) = J(\psi, \nabla^2 \psi) \hat{\boldsymbol{z}},$$
$$\nabla \times (\boldsymbol{U} \cdot \nabla \boldsymbol{v}) = \boldsymbol{U} \cdot \nabla(\nabla \times \boldsymbol{v}) = (-U_u \partial_u \nabla^2 \psi) \hat{\boldsymbol{z}}$$

ただし

$$J(a,b) = (\partial_x a)(\partial_y b) - (\partial_y a)(\partial_x b)$$

基礎方程式に ∇× を掛ける:

$$-\partial_t \nabla^2 \psi + J(\psi, \nabla^2 \psi) - U_y \partial_y \nabla^2 \psi = \frac{1}{\text{Re}_0} \nabla^2 \psi - \frac{1}{\text{Re}} \nabla^2 \nabla^2 \psi - n \cos(ny)$$

整理:

$$\partial_t \nabla^2 \psi = n \cos(ny) - \frac{1}{\text{Re}_0} \nabla^2 \psi - U_y \partial_y \nabla^2 \psi + \frac{1}{\text{Re}} \nabla^2 \nabla^2 \psi + J(\psi, \nabla^2 \psi)$$

 $\omega = (-\nabla^2 \psi) \cdot \hat{\boldsymbol{z}} = (\nabla \times \boldsymbol{v}) \cdot \hat{\boldsymbol{z}}$ とおいて整理:

$$\partial_t \omega = -n\cos(ny) - \frac{1}{\text{Re}_0}\omega - U_y \partial_y \omega + \frac{1}{\text{Re}} \nabla^2 \omega + J(\psi, \omega)$$

線形部分 + cos 部分:

$$\partial_t \omega = -n\cos(ny) - \frac{1}{\text{Re}_0}\omega - U_y \partial_y \omega + \frac{1}{\text{Re}}\nabla^2 \omega$$

Fourier 変換 $F[\omega](k,\ell)$:

$$\partial_t F[\omega] = -F[n\cos(ny)] - \frac{1}{\text{Re}_0} F[\omega] - i\ell U_y F[\omega] - \frac{1}{\text{Re}} (k^2 + \ell^2) F[\omega]$$

整理:

$$\partial_t \{ F[\omega] e(t) \} = -F[n \cos(ny)] e(t),$$
$$e(t) = \exp(t\alpha),$$
$$\alpha = \frac{1}{\text{Re}_0} + i\ell U_y + \frac{1}{\text{Re}} (k^2 + \ell^2)$$

積分:

$$F[\omega](t) = F[\omega](t_0)e(-t+t_0) - e(-t)\int_{t_0}^t F[n\cos{(ny)}]e(s)ds$$

$$= F[\omega](t_0)e(-t+t_0) - F[n\cos{(ny)}] \cdot \frac{1 - e(-t+t_0)}{\alpha}$$

$$= \left(F[\omega](t_0) + \frac{F[n\cos{(ny)}]}{\alpha}\right)e(-t+t_0) - \frac{F[n\cos{(ny)}]}{\alpha}$$