

Parerga and paralipomena

1 p_T^2 integration

In the phase space decomposition employed for the rapidity distribution, the integration on p_T^2 is computed with a Dirac delta that leaves out a factor of $|\tilde{p}_T^2|^\epsilon$. This depends on the integration variables l^2 and k'^2 . Explicitly, the Dirac delta in the frame of reference where l is at rest gives

$$p_H^0 + k'^0 - \sqrt{l^2} = 0. \quad (1)$$

Because initial gluon emissions are soft in the threshold limit, this frame of reference is also the frame of reference of the partonic CoM. The equation above can be rewritten using the usual parametrization of p_H . Therefore we have the equation

$$\sqrt{m_H^2 + p_T^2} \cosh y + \sqrt{k'^2 + p_T^2 \cosh^2 y + m_H^2 \sinh^2 y} - \sqrt{l^2} = 0. \quad (2)$$

Solving with Mathematica with respect to p_T^2 gives

$$\tilde{p}_T^2 = \frac{[(k'^2)^2 - 2k'^2 m_H^2 + m_H^4 - 2k'^2 l^2 + l^4 - 2m_H^2 l^2 \cosh 2y] \operatorname{sech}^2 y}{4l^2} \quad (3)$$

Using $\cosh 2y = 2 \cosh^2 y - 1$ this reduces to

$$\tilde{p}_T^2 = \frac{(k'^2 - m_H^2 - l^2)^2 \operatorname{sech}^2 y}{4l^2} - m_H^2. \quad (4)$$

The expression is different than De Ros (eq. 4.68 pag. 46) because at that point he is still using as variable p_z instead of y .

In each passage of this calculation, one can verify that in the threshold limit, be it for fixed x_1 or x_2 , $\tilde{p}_T^2 \rightarrow 0$ using $k'^2 \rightarrow 0, l^2 \rightarrow s$ and $s \rightarrow m_H^2/x_2$. Therefore $|\tilde{p}_T^2|^\epsilon$ carries a factor of $[(1 - x_1)(1 - x_2)]^\epsilon$ that **contributes to the collinear scale**. Possibly, as shown explicitly at NLO, it also carries factor that regularize the integrals in l^2 e k'^2 .

2 Rewriting Glosser and Schmidt

We shortly revise and carry on the rewriting of the GS result. Equation 3.17 of GS is

$$\begin{aligned} C_{sing}^{(2)} = & \delta(Q^2) \left\{ (11 + \delta + N_c U) g_{gg} \right. \\ & + (N_c - n_f) \frac{N_c}{3} \left[\frac{m_H^4}{s} + \frac{m_H^4}{t} + \frac{m_H^4}{u} + m_H^2 \right] \Big\} \\ & + \left\{ \left(\frac{1}{-t} \right) \left[-P_{gg}(z_t) \log \frac{\mu_F^2 z_t}{-t} + p_{gg}(z_t) \left(\frac{\log(1 - z_t)}{1 - z_t} \right)_+ \right] g_{gg,t}(z_t) \right. \\ & + \left(\frac{1}{-t} \right) \left[-2n_f P_{qg}(z_t) \log \frac{\mu_F^2}{Q_{max}^2} + 2n_f z_t (1 - z_t) \right] g_{qg,t}(z_t) \\ & + \left(\frac{z_t}{-t} \right) \left[\left(\frac{\log(1 - z_t)}{1 - z_t} \right)_+ - \log \frac{Q_T^2 z_t}{-t} \left(\frac{1}{1 - z_t} \right)_+ \right] \\ & \cdot \frac{N_c^2}{2} \left[\frac{m_H^8 + s^4 + t^4 + u^4 + Q^8 + z_t z_u (m_H^8 + s^4 + Q^8 + (u/z_u)^4 + (t/z_t)^4)}{sut} \right] \\ & - \left(\frac{z_t}{-t} \right) \left(\frac{1}{1 - z_t} \right)_+ \frac{\beta_0}{2} N_c \left(\frac{m_H^8 + s^4 + z_t z_u ((u/z_u)^4 + (t/z_t)^4)}{sut} \right) \\ & \left. + (t \leftrightarrow u) \right\} \end{aligned}$$

$$\begin{aligned}
& + N_c^2 \left[\frac{(m_H^8 + s^4 + Q^8 + (u/z_u)^4 + (t/z_t)^4)(Q^2 + Q_T^2)}{s^2 Q^2 Q_T^2} + \right. \\
& \left. + \frac{2m_H^4((m_H^2 - u)^4 + (m_H^2 - t)^4 + t^4 + u^4)}{sut(m_H^2 - t)(m_H^2 - u)} \right] \frac{1}{p_T^2} \log \frac{p_T^2}{Q_T^2}
\end{aligned}$$

Now defining new coefficients this is rewritten as

$$\begin{aligned}
C_{sing,rap}^{(2)} = \int_0^1 dq J(x_1, x_2, q) Q_{max}^2 & \left\{ \frac{\delta(q)}{Q_{max}^2} \mathcal{V} + \right. \\
& \frac{1}{-t} \mathcal{A}_t \left[-p_{gg} \left(- \left(\frac{\ln(1-z_t)}{1-z_t} \right)_+ + \ln \frac{\mu^2 z_t}{-t} \left(\frac{1}{1-z_t} \right)_+ \right) - \beta_0 \ln \frac{\mu^2 z_t}{-t} \frac{\delta(q)}{Q_{max}^2} \right] + \\
& \frac{1}{-t} \left[\mathcal{B}_{1,t} \ln \frac{1}{q Q_{max}^2} + \mathcal{B}_{2,t} \right] + \\
& \frac{z_t}{-t} \mathcal{C}_t \left[\left(\frac{\ln(1-z_t)}{1-z_t} \right)_+ - \ln \frac{Q_T^2 z_t}{-t} \left(\frac{1}{1-z_t} \right)_+ \right] - \\
& \frac{z_t}{-t} \mathcal{D}_t \left(\frac{1}{1-z_t} \right)_+ + (t \leftrightarrow u) \\
& \left. + \mathcal{E} \right\}
\end{aligned} \tag{5}$$

where

$$\begin{aligned}
\mathcal{A}_t &= g_{gg,t}(z_t) \\
\mathcal{A}_u &= g_{gg,u}(z_u) \\
\mathcal{B}_{1,t} &= -2n_f P_{qg}(z_t) g_{qg,t}(z_t) \\
\mathcal{B}_{2,t} &= 2n_f z_t (1-z_t) g_{qg,t}(z_t) \\
\mathcal{B}_{1,u} &= -2n_f P_{qg}(z_u) g_{qg,t}(z_t) \\
\mathcal{B}_{2,t} &= 2n_f z_u (1-z_u) g_{qg,t}(z_t), \\
\mathcal{C}_t = \mathcal{C}_u &= \frac{N_c^2}{2} \left[\frac{m_H^8 + s^4 + t^4 + u^4 + Q^8 + z_t z_u (m_H^8 + s^4 + Q^8 + (u/z_u)^4 + (t/z_t)^4)}{sut} \right] \\
\mathcal{D}_t = \mathcal{D}_u &= \frac{\beta_0}{2} N_c \left(\frac{m_H^8 + s^4 + z_t z_u ((u/z_u)^4 + (t/z_t)^4)}{sut} \right) \\
\mathcal{E} &= N_c^2 \left[\frac{(m_H^8 + s^4 + Q^8 + (u/z_u)^4 + (t/z_t)^4)(Q^2 + Q_T^2)}{s^2 Q^2 Q_T^2} + \frac{2m_H^4((m_H^2 - u)^4 + (m_H^2 - t)^4 + t^4 + u^4)}{sut(m_H^2 - t)(m_H^2 - u)} \right] \frac{1}{p_T^2} \log \frac{p_T^2}{Q_T^2}
\end{aligned}$$

notice that in the definition of the \mathcal{B} coefficients, quarks and gluons functions are exchanged too. Rearranging

$$\begin{aligned}
C_{sing,rap}^{(2)} = \int_0^1 dq J(x_1, x_2, q) & \left\{ \right. \\
& Q_{max}^2 \left[\mathcal{E} + \frac{\mathcal{B}_{2,t}}{-t} + \frac{\mathcal{B}_{1,t}}{-t} \ln \frac{1}{q Q_{max}^2} \right] \\
& + \delta(q) \left[\mathcal{A}_t \beta_0 \ln t + N_c \mathcal{A}_t \ln^2 Q_{max}^2 - N_c \mathcal{A}_t \ln^2 t - \mathcal{C}_t \ln Q_{max}^2 \ln p_T^2 + \frac{1}{2} \mathcal{C}_t \ln^2 Q_{max}^2 \right. \\
& \left. - \frac{1}{2} \mathcal{C}_t \ln^2 t + \mathcal{C}_t \ln p_T^2 - \mathcal{D}_t \ln Q_{max}^2 + \mathcal{D}_t \ln t + \mathcal{C}_t \ln Q_T^2 \ln t + \mathcal{V} \right] \\
& + \left(\frac{1}{q} \right)_+ \left[\frac{\mathcal{A}_t p_{gg}}{z_t} \ln Q_{max}^2 + \mathcal{C}_t \ln Q_{max}^2 - \mathcal{C}_t \ln Q_T^2 - \mathcal{D}_t \right] \\
& \left. + \left(\frac{\ln q}{q} \right)_+ \left[\frac{\mathcal{A}_t p_{gg}}{z_t} + \mathcal{C}_t \right] + (t \leftrightarrow u) \right\},
\end{aligned} \tag{6}$$

where in terms proportional to $\delta(q)$ we have set $z_t = 1$ and therefore $p_{gg}(1) = 2N_c$.