

**Poissonovi oklepaji.** Ponovno imamo  $f(\underline{q}, \underline{p}, t)$ , kjer so posplošene koordinate  $\underline{q}$  in impulzi  $\underline{p}$  odvisni od časa. Spet je  $p_i = \frac{\partial L}{\partial \dot{q}_i}$ ,  $\dot{q}_i = \frac{\partial H}{\partial p_i}$ ,  $\dot{p}_i = -\frac{\partial H}{\partial q_i}$ .

$$\begin{aligned} \frac{df}{dt} &= \sum_i \left( \frac{\partial f}{\partial q_i} \dot{q}_i + \frac{\partial f}{\partial p_i} \dot{p}_i \right) + \frac{\partial f}{\partial t} \\ &= \sum_i \left( \frac{\partial f}{\partial q_i} \frac{\partial H}{\partial p_i} - \frac{\partial f}{\partial p_i} \frac{\partial H}{\partial q_i} \right) + \frac{\partial f}{\partial t} = [f, H] + \frac{\partial f}{\partial t} \end{aligned}$$

Definirali smo Poissonov oklepaj  $[f, g] = \sum_i \left( \frac{\partial f}{\partial q_i} \frac{\partial g}{\partial p_i} - \frac{\partial f}{\partial p_i} \frac{\partial g}{\partial q_i} \right)$ . Včasih se uporablja tudi oznako  $\{f, g\}$ .

**Lastnosti.** Poissonovi oklepaji imajo naslednje lastnosti:

1.  $[f, \lambda g + \eta h] = \lambda [f, g] + \eta [f, h]$
2.  $[f, g] = -[g, f]$
3.  $[f, gh] = [f, g]h + g[f, h]$
4.  $[f, f] = 0$
5.  $[x, p_x] = 1$
6.  $[x, p_y] = 0$

**Lagrangeov formalizem za zvezno sredstvo.** Na primer za struno:

$$\begin{aligned} dT &= \frac{1}{2} \dot{u}^2 \rho S_0 dx = \frac{1}{2} \rho S_0 \left( \frac{\partial u}{\partial t} \right)^2 dx \\ dV &= F(dl - dx) = F(\sqrt{dx^2 + du^2} - dx) \approx \frac{1}{2} F \left( \frac{\partial u}{\partial x} \right)^2 dx \\ dL &= \frac{1}{2} \rho S_0 \left( \frac{\partial u}{\partial t} \right)^2 dx \pm \frac{1}{2} F \left( \frac{\partial u}{\partial x} \right)^2 dx \pm \rho S_0 u dx \\ \mathcal{L} &= \frac{dL}{dx} = \dots \\ L &= \int_{x_1}^{x_2} \mathcal{L} dx = \int_{x_1}^{x_2} \mathcal{L}(u(x, t), u_x, u_t, x, t) \\ S &= \int_0^t L dt = \iint \mathcal{L} dx dt (dy, dz) \end{aligned}$$

Euler-Lagrangeova enačba za tak sistem je

$$\frac{d}{dt} \frac{\partial \mathcal{L}}{\partial u_t} + \frac{d}{dx} \frac{\partial \mathcal{L}}{\partial u_x} - \frac{\partial \mathcal{L}}{\partial x} = 0$$

Iz tega lahko izpeljamo na primer Maxwellove zakone, valovne enačbe itd.