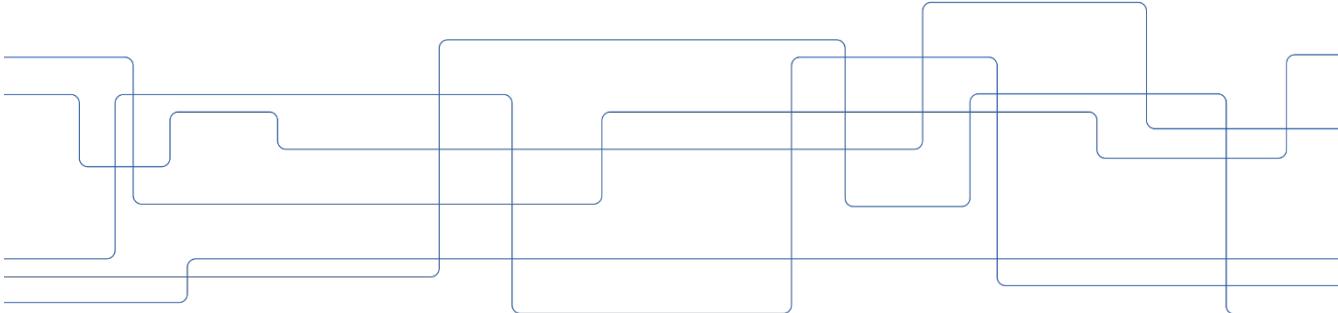


# PEDESTAL PHYSICS

## a phenomenological introduction

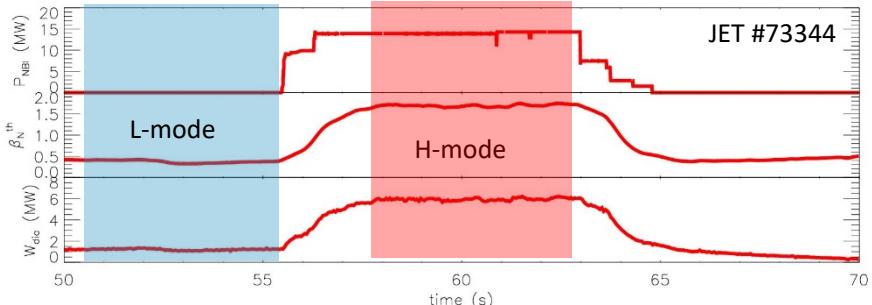
L. Frassinetti



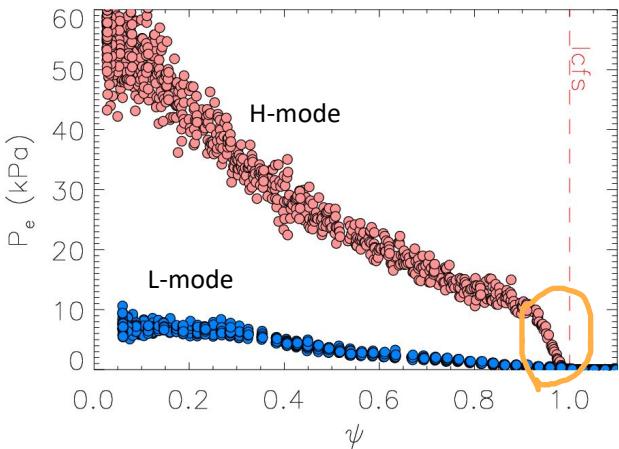
# H-mode plasma

Pedestal requires  
an L-H transition

- When the input power to the plasma is above a specific threshold, the plasma has a transition from a low confinement regime (L-mode) to a high confinement regimes (H-mode).



- The H-mode is characterized by:
  - steep gradients in the pressure "near" the edge of the plasma.  
This region is named "pedestal".
  - sudden releases of energy and particles from the pedestal region to the SOL and the divertor.  
These events are triggered by MHD instabilities and are named edge localized modes (ELMs)



# OUTLINE

- L-H transition
- Pedestal structure
- Edge localized modes (ELMs)
  - ELM energy losses
  - ELM types
- MHD stability of the pedestal
  - Role of MHD stability (and few words on transport)
  - The peeling-balloonning (PB) model
  - The ELM cycle within the PB model
  - Parameters that influences the pedestal
- Pedestal predictions
  - The EPED model:
    - The PB constraint
    - The KBM constraint
  - Non-linear MHD modelling
- Some of the most active research areas in pedestal physics

# L-H transition

- Above a specific threshold in power ( $P_{LH}$ ), the plasma enters the H-mode
- The  $P_{LH}$  threshold depends on several parameters.
- A scaling law based on results from several machines produces:

$$P_{LH} = 0.049 n_e^{0.72} B^{0.8} S^{0.94}$$

[Martin JPC2008]

Turbulence  
prevents H-mode

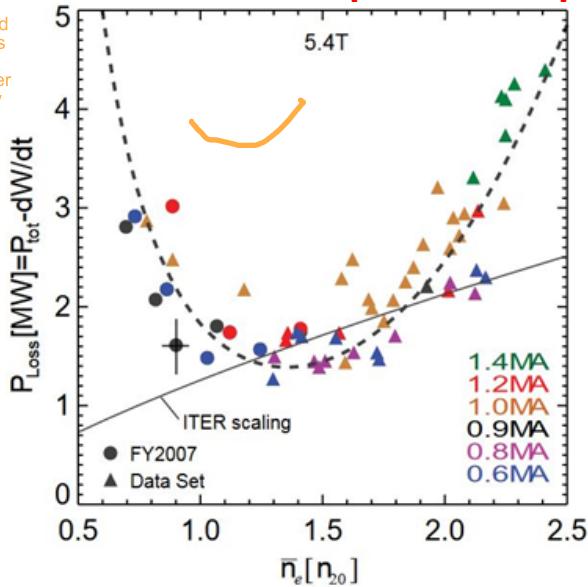
- However, the links between engineering/plasma parameters and  $P_{LH}$  is more complex. Some of the main parameters that affects  $P_{LH}$  are:
  - Magnetic field
  - Isotope mass ( $P_{LH}$  decreases with isotope mass) [Righi NF1999]
  - Divertor geometry [Delabie EPS2015]
  - Wall material ( $P_{LH}$  reduced from carbon to metal walls) [Neu JNM2013]
  - Plasma density [Martin JPC2008]
    - Minimum around  $0.2\text{-}0.4n_{GW}$
    - Non-monotonic behavior seem related to edge ion heating [Ryter NF2014]

To enter H-mode one must suppress turbulence, and the instabilities that generate turbulence differ based on density

inaccurate

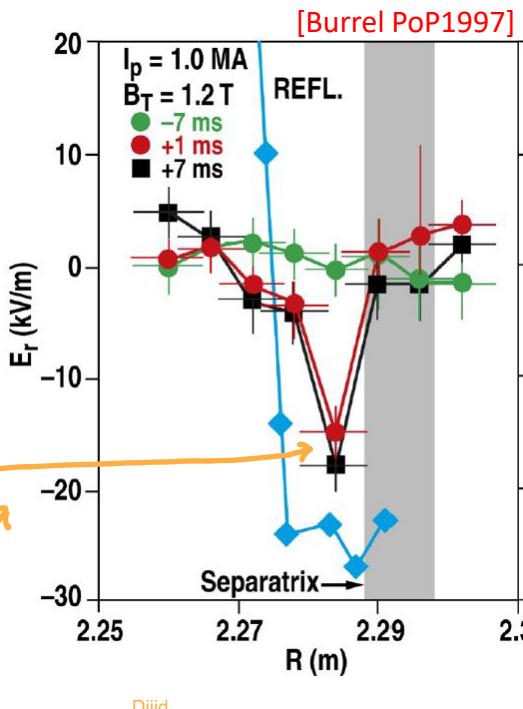
Physics not yet understood

[Gohil IAEA2013]



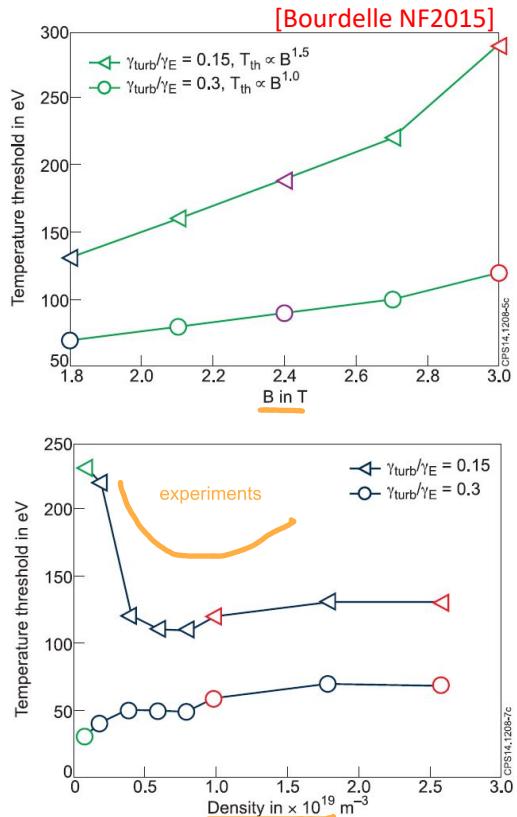
# L-H transition

- The physics of L-H transition is not yet fully understood
  - several models have been proposed to explain the experimental results
  - but a physics based model of the L-H transition with full predictive capabilities has not been developed yet.
- Some key experimental and theoretical concepts to explain the L-H transition are well established:
  - The L-H transition is due to stabilization of the turbulence near the plasma edge [Burrel PoP1997], [Terry RMP2000]
  - $\vec{E} \times \vec{B}$  shear stabilization plays a key role
    - higher  $\vec{E} \times \vec{B}$  in L-mode  $\rightarrow$  lower  $P_{LH}$ .
    - The formation of a  $E_r$  well, just inside the separatrix, occurs as the plasma enters H-mode
    - The well has to reach a certain depth to allow the transition



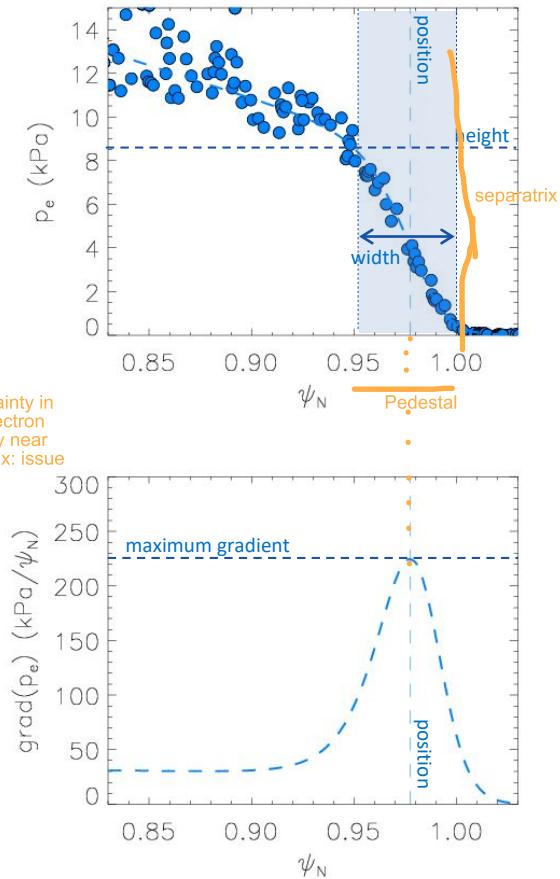
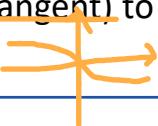
# L-H transition

- Many of the theoretical works are based on the interplay between the L-mode turbulence and  $E_r$  shearing. [Connor PPCF2000]
- A large part of other theoretical works are based on the stabilization of RBM via increased pressure gradient. [Rogers PRL1997]
- An example: [Bourdelle NF2015]
  - $\gamma_{\text{turb}}$  (growth rate of the turbulence) can be modeled from theory (either analytically or numerically)
  - $\gamma_E$  ( $E_r$  shear) can be obtained by modelling the  $E_r$  profiles. Stabilizing effect
  - $\gamma_{\text{turb}}/\gamma_E$  can be used to identify at which temperature the transition occurs  
→ Qualitative trends can be tested
- For a recent review on L-H transition: [Bourdelle NF2020]



# Pedestal structure

- To study the pedestal, it is necessary to quantify the parameters that identify its structure.
- The key parameters are
  - pedestal height
  - pedestal width
  - pedestal position (often defined as the position of the maximum gradient).
  - maximum gradient
- The pedestal parameters are determined for:
  - pressure
  - temperature
  - density
- These parameters are determined by fitting an analytical function (typically, a modified hyperbolic tangent) to the experimental data.



# Edge Localized Modes (ELMs)

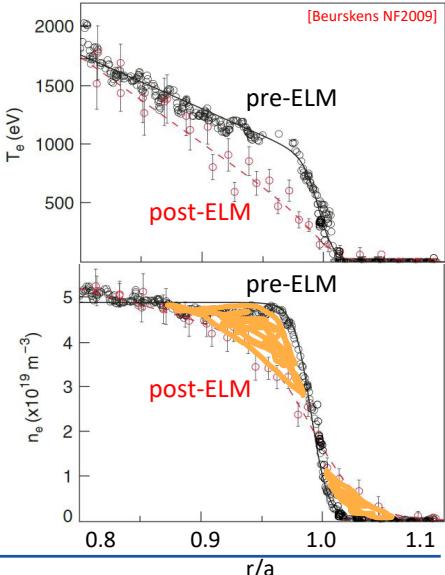
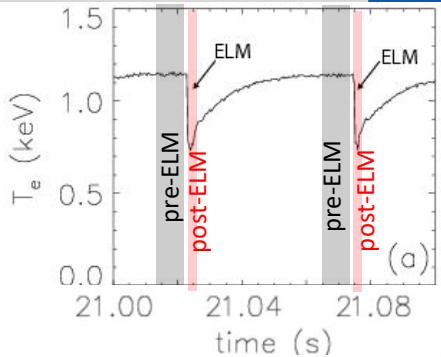
- The pedestal is characterized by sudden events, triggered by MHD instabilities, called edge localized modes (ELMs).
- The ELM triggers the collapse of the pedestal temperature and density, which in turn leads to the release of energy and particles to the divertor.
- The ELM collapse affects the kinetic profiles only in the pedestal region.
- The ELM losses can be calculated by integrating the profiles just before and soon after the ELMs:

[Beurskens NF2009]

$$\Delta W_{ELM} = W_{pre} - W_{post} = \\ = \frac{3}{2} k \int (n_{pre} T_{pre} - n_{post} T_{post}) dV \approx$$

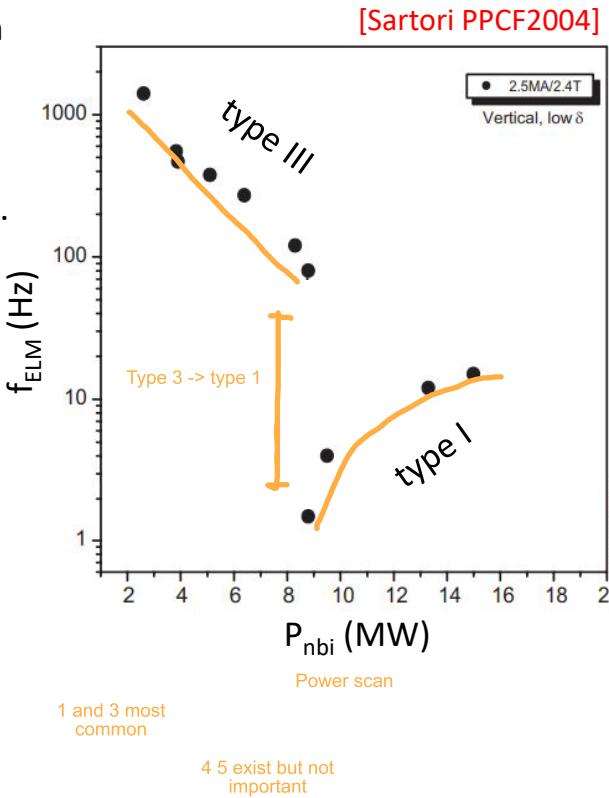
$$\approx \frac{3}{2} k \int \Delta n T dV + \frac{3}{2} k \int n \Delta T dV$$

convective losses      conductive losses



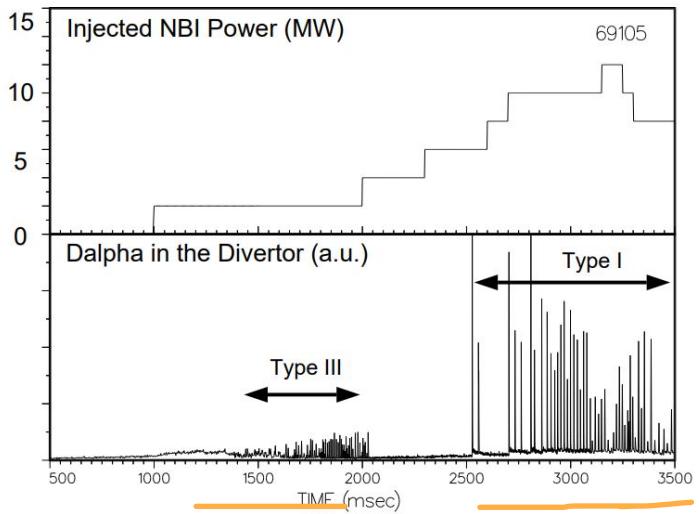
# ELM types: definitions

- H-mode plasma can be characterized by several types of ELMs. The ELM frequency ( $f_{\text{ELM}}$ ) is often used to identify the most common ELMs.
- The most common are:
  - **Type I ELMs.** Most common ones
    - $f_{\text{ELM}}$  increases with  $P_{\text{sep}} = P_{\text{in}} - P_{\text{rad}} - dW/dt$ .
    - typically occurs at  $P_{\text{sep}} \gg P_{\text{LH}}$ .
    - they are triggered by ideal MHD. Supposedly
    - they appear as sharp burst on the  $D_{\alpha}$ .
  - **Type III ELMs.**
    - $f_{\text{ELM}}$  decreases with  $P_{\text{sep}}$ . Negative correlation with power
    - typically occurs  $P_{\text{sep}} \approx P_{\text{LH}}$ .
    - they are not triggered by ideal MHD.
  - **Type II (or "grassy" ELMs).**
    - Not achieved in all machines.
    - Occurs at high confinement and high triangularity.
    - They lead to small but frequent energy losses.

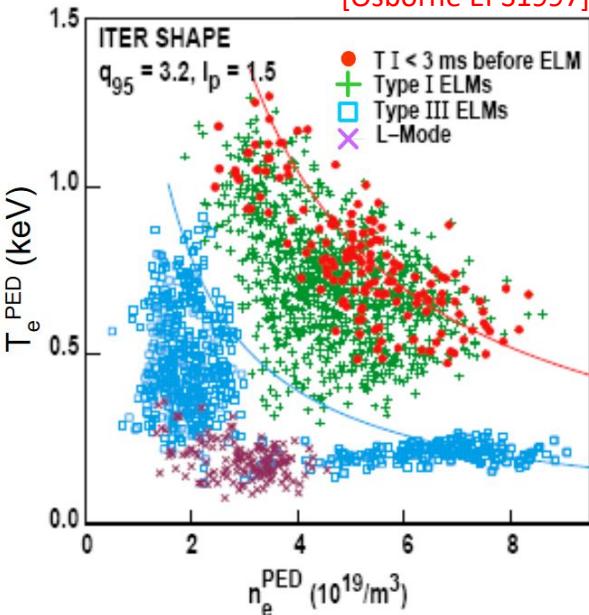


# ELM types: examples

[Zohm PPCF1996]



[Osborne EPS1997]



## Type I ELMs.

- $f_{\text{ELM}}$  increases with  $P_{\text{sep}} = P_{\text{in}} - P_{\text{rad}} - dW/dt$ .
- typically occurs at  $P_{\text{sep}} \gg P_{\text{LH}}$ .

## Type III ELMs.

- $f_{\text{ELM}}$  decreases with  $P_{\text{sep}}$ .
- typically occurs  $P_{\text{sep}} \approx P_{\text{LH}}$ .

Type one less freq but more power loss

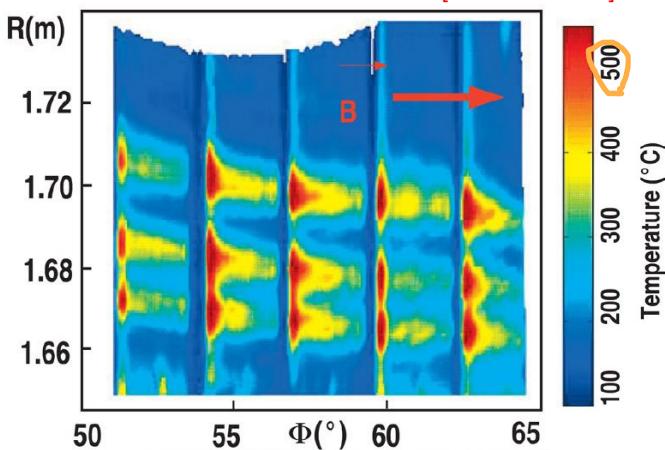
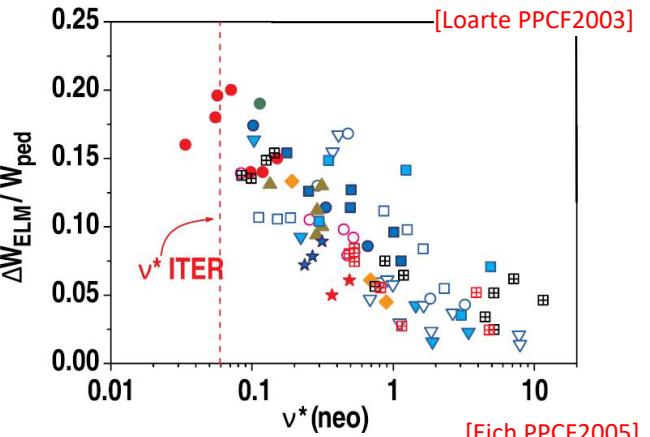
For reviews of ELM types:

- [Zohm PPCF1996]
- [Leonard PoP2014]

# ELMs: energy losses and heat loads

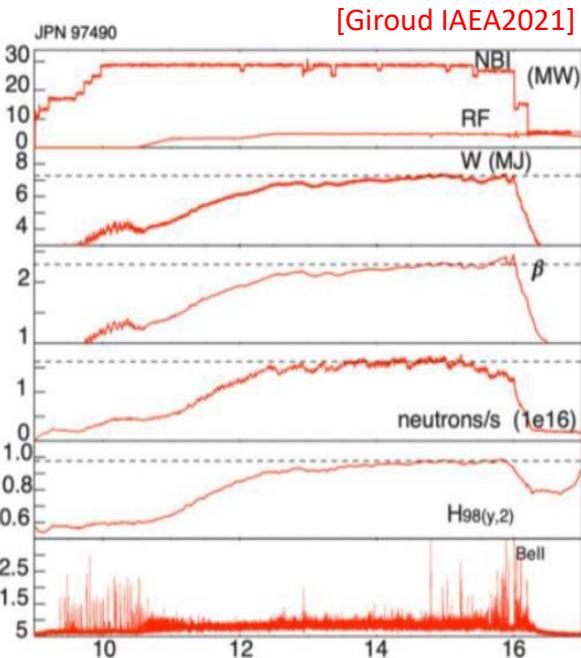
- ELM losses tend to increase with decreasing collisionality. Several machines
- At ITER collisionalities, the ELM energy losses might be 15%-20% of the pedestal stored energy.
- ELMs lead to fluxes of energy and particles to the divertor.
- The divertor can be damaged or could even melt. This could pose a problem for ITER. [Pitts JNM2013]
- It is essential to understand ELM pedestal physics to:
  - Minimize ELM energy losses
  - Develop techniques for ELM mitigation/suppressions. Some of the most developed techniques are:
  - RMPs for a review: [Evans JNM2013]
  - ELM pacing with pellets [Baylor NF2009]

+ small elms regimes  
+ vertical kicks: Moving plasma up-down with the control systems triggers elms



# ELM types: small ELMs scenarios

- Type I ELMs have been the most studies, so far
- In recent years, significant experimental efforts have been devoted to identify and study alternatives to type I ELMs that might be useful for a fusion reactor:
  - **Small ELMs (SE)**
    - at very low-gas and high performance baseline plasmas in JET [Garcia PoP2022]
  - **Quasy Continuous Exhaust (QCE)**
    - at high triangularity and high gas rate (type II ELMs) [Stober NF2001]
  - **Enhanced D-alpha (EDA) and quasi coherent mode (QCM)** [Greenwald PoP1999]
    - at high triangularity, low gas rate and power
  - **Quiescent H-mode (QH)** [Chen NF2020]
    - At high NBI torque, which excites a edge harmonic oscillation EHO which increase transport.
  - **Seeded small ELMs** [Giroud IAEA2021]
    - High power seeded plasmas.



Advantage compared to standard type I ELM scenarios:  
small ELMs, small heat loads on the divertor but good confinement

# MHD stability and transport

- What are the physical mechanisms that determines the pedestal structure and trigger the ELMs?
- Two main concepts
  - MHD stability
  - Heat and particle transport
- The time evolution is set by transport
  - Transport determines time evolution of
    - pedestal gradients
    - pedestal heights
- The pedestal grows till a critical threshold in pressure. Then, the MHD stability triggers an ELM.
  - MHD stability determines:
    - pedestal height
    - the maximum gradient.
  - In the pedestal, the main MHD instabilities are:
    - ballooning (B) modes
    - peeling (P) modes
    - coupled PB modes

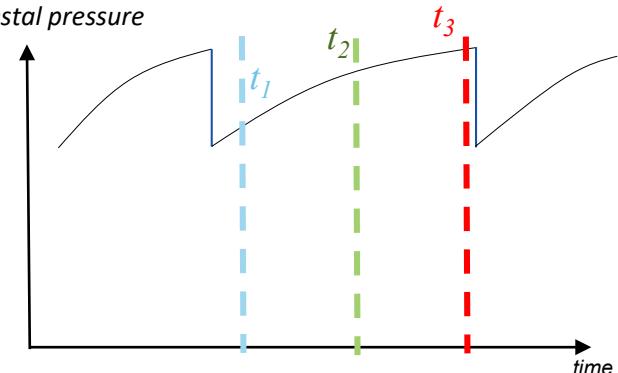
Gradient pressure triggers the mhd instability

Turbulent transport in pedestal

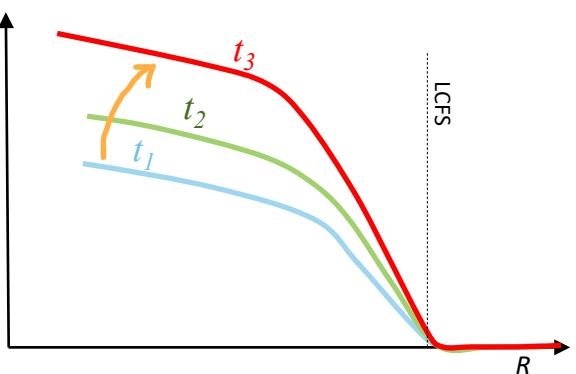
Transport control elm freq

3

Pedestal pressure



pressure



# The ballooning modes

- The ballooning instabilities are pressure driven: they are triggered when the pressure gradient exceeds a critical threshold.
- They arise from toroidicity
- B has an unfavourable curvature low field side → ballooning modes develop mainly on the LFS Low field side (?)

- Two key parameters define the ballooning stability

○ **the normalized pressure gradient  $\alpha$**

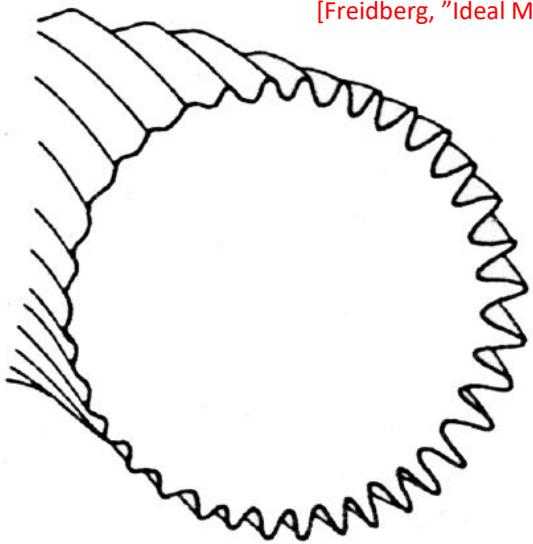
$$\alpha = -\frac{2\mu_0 R q^2}{B^2} \frac{dp}{dr}$$

has a destabilizing effect.

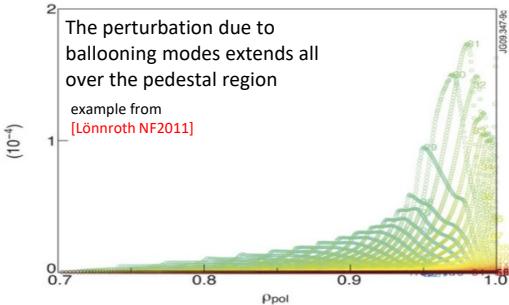
○ **the magnetic shear**

$$s = -\frac{r}{q} \frac{dq}{dr}$$

$s$  has a stabilizing effect.



[Freidberg, "Ideal MHD"]



# The ballooning modes

- the normalized pressure gradient  $\alpha$

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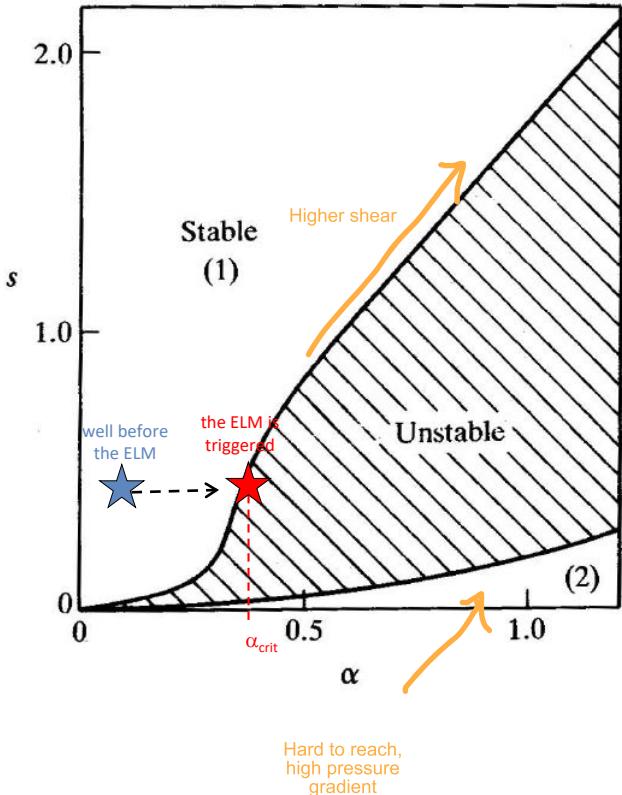
- the increase of  $\alpha$  destabilizes ballooning modes
- at a certain threshold in  $\alpha$  ( $\alpha_{\text{crit}}$ ) , the mode is unstable

- the magnetic shear

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- the shear has a stabilizing effect
- Increasing the shear leads to an increase in  $\alpha_{\text{crit}}$ .

[Wesson "tokamaks"]



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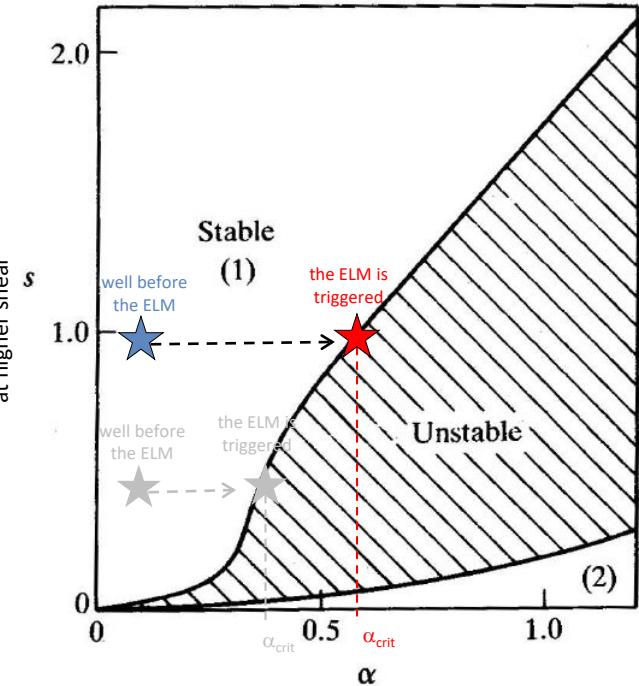
- the magnetic shear

$$s = -\frac{r dq}{q dr}$$

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- Most of the machines have a pedestal in region (1): the first stability region
- However, theory predicts a second stability region, at high  $\alpha$  and low shear

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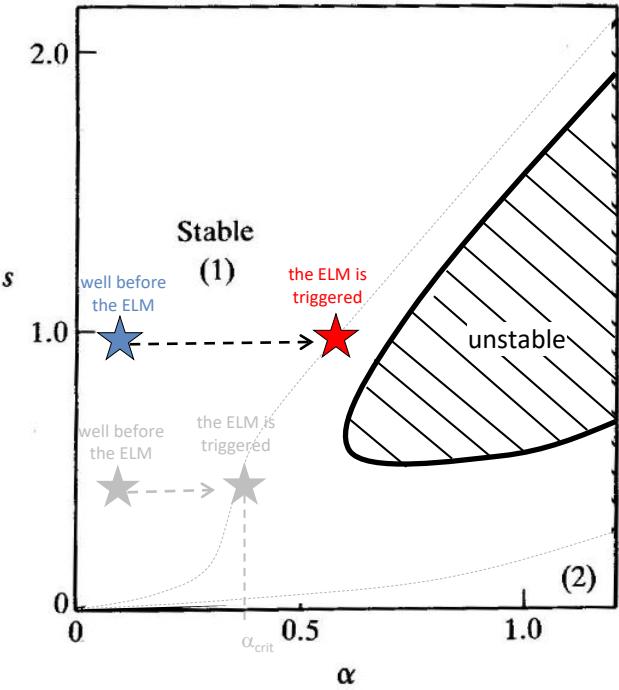
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adapted from  
[Wesson "tokamaks"]



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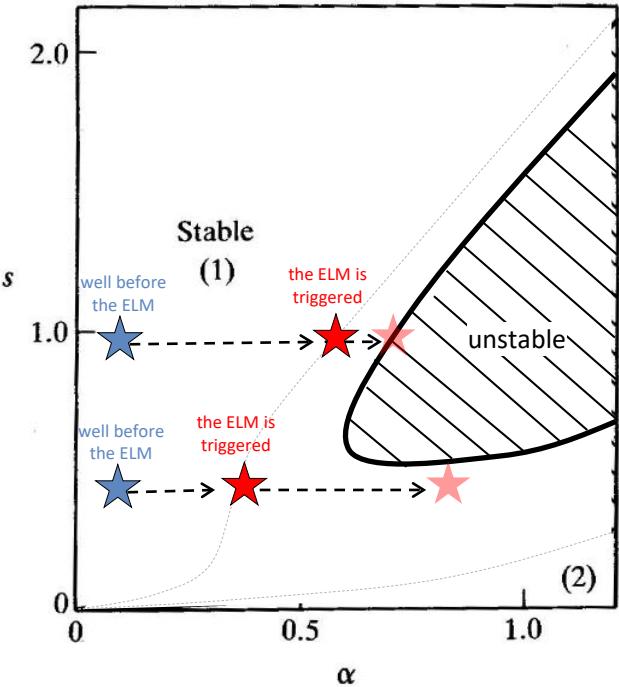
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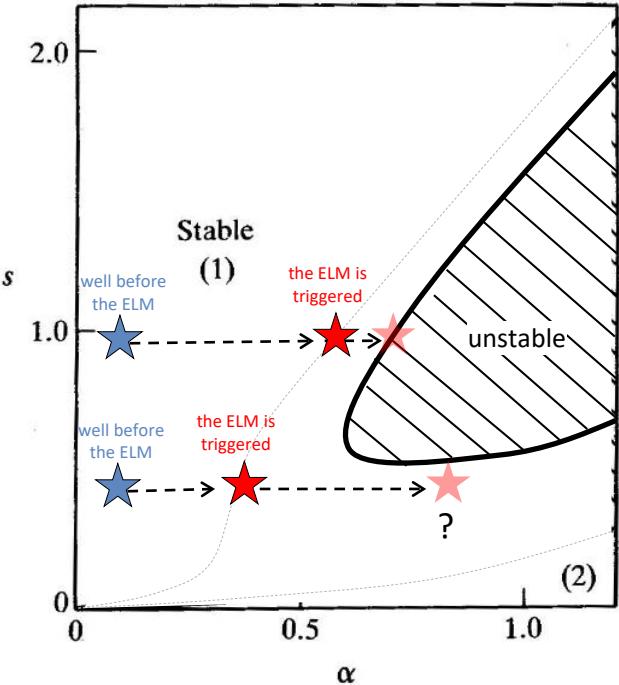
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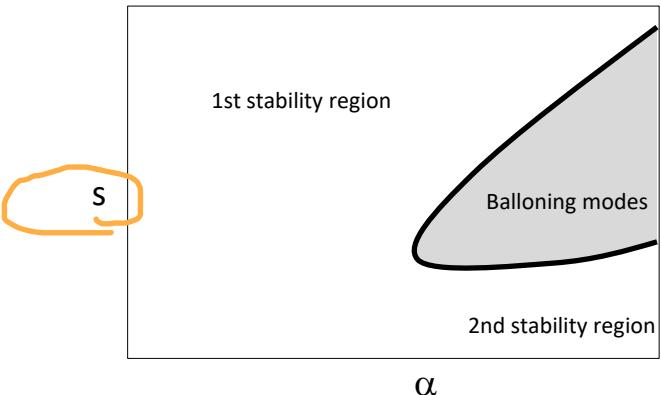
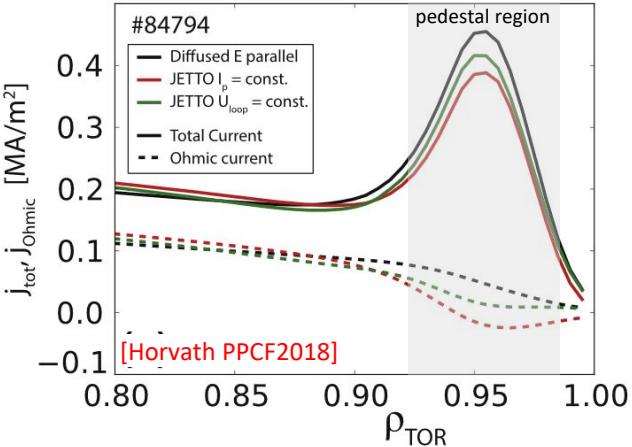


No!  
There are further instabilities → see later.  
But first...

# The bootstrap current

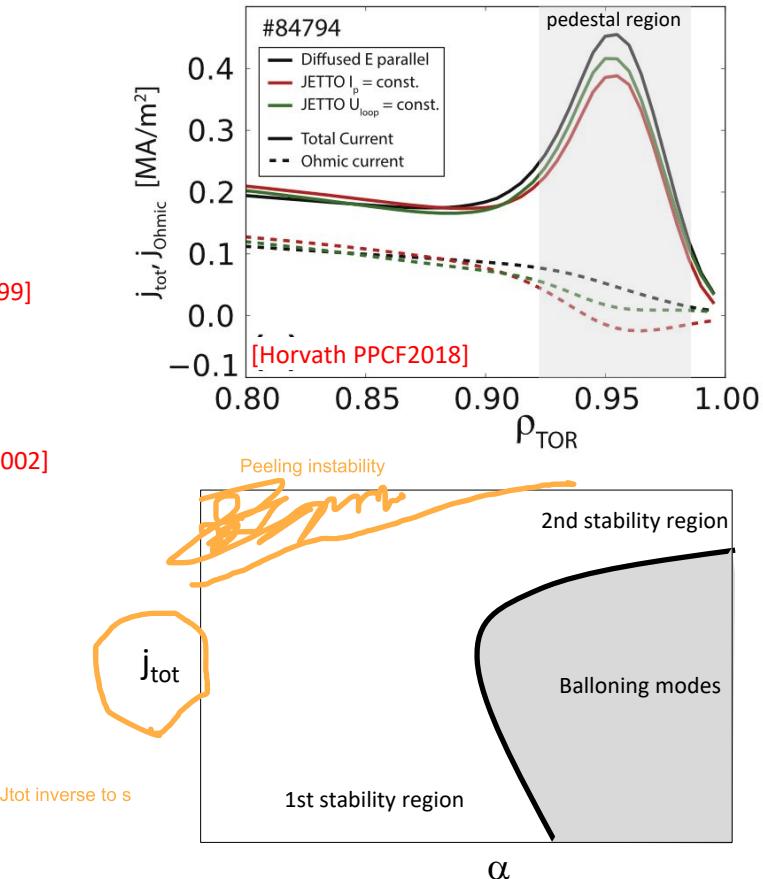
Related to  
pedestal stability

- Due to the steep gradients in the pedestal region, the bootstrap current ( $j_{bs}$ ) can give a significant contribution to the total current density.
  - For an expression of  $j_{bs}$ : [Sauter PoP1999]
  - The increase in the current density affects the shear [Miller PoP1999]
- $j_{bs}$  has an effect on the pedestal stability. [Snyder PoP2002]
- the parameters that affects  $j_{bs}$  will affect also the ballooning stability:
- collisionality
  - plasma shape
- It is common to use  $j_{tot}$  instead of the shear in the stability diagram



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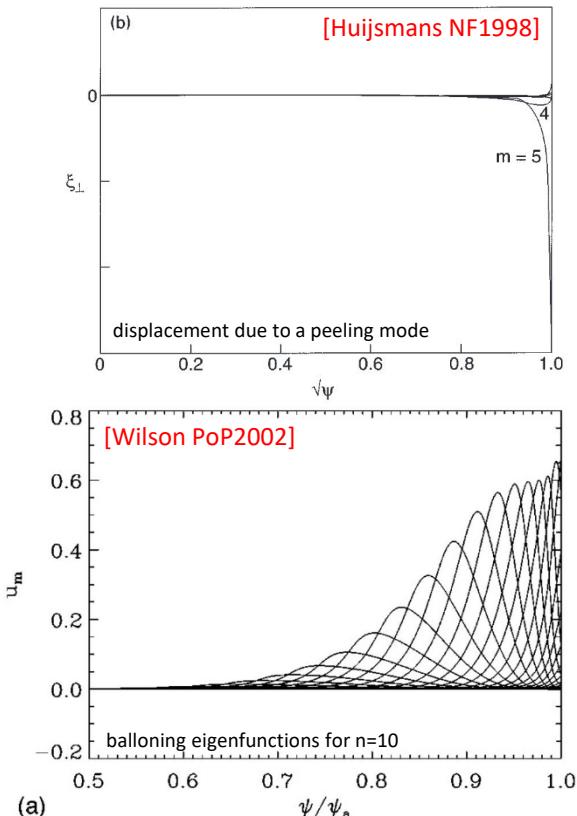
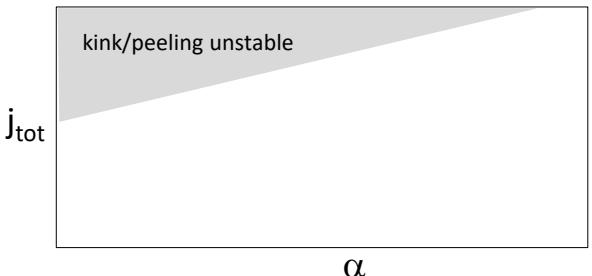


# The external kink / peeling mode

- The external kink mode is current driven
- The kink mode ( $m, n$ ) is destabilized when  $q$  at the plasma edge is low enough that  $q_{edge} < m/n$  and the resonance is very close to the plasma
  - the kink mode is resonant outside the plasma
  - the kink mode is strongly localized at the plasma edge.

For comparison, the ballooning modes have a more global structure.
- The kink mode depends on the edge current  
 $\rightarrow j_{bs}$  has a strong role

[Huijsmans NF1998]

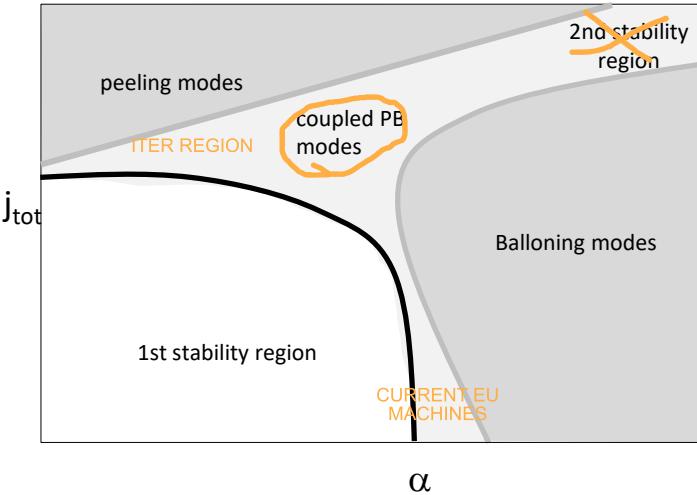


# The peeling-balloonning (PB) modes

- Toroidicity and shaping effects can couple peeling and ballooning (PB) modes.
- The coupled PB modes can be destabilized even if the single peeling mode and ballooning are stable.

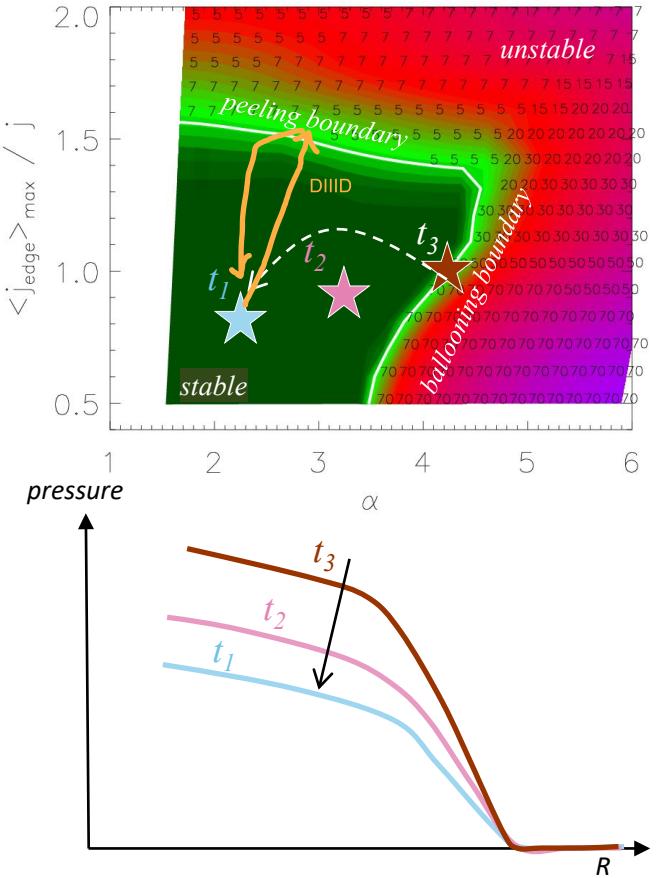
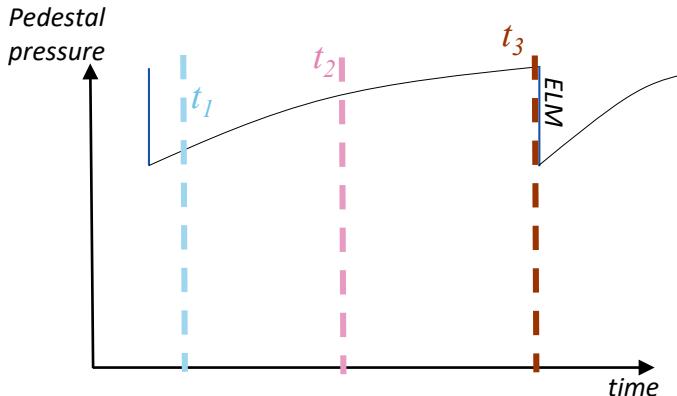
[Connor PoP1998]

- The PB stability are driven by both pressure gradient and current density.
- The PB stability is the leading theory to explain the pedestal behvaior in type I ELM My H-modes. [Snyder PoP2002]  
[Wilson PoP2002]
- The PB modes strongly limit the stable region.
- The access to the 2nd stability region is closed (most of the times).



# The PB model for the ELM cycle

- 1 Just after an ELM, the pedestal has low gradient and low  $j_{bs}$ .
- 2 During the ELM cycle, the pressure gradient (and hence  $j_{bs}$ ) increases
- 3 The process continues till the PB boundary is reached.
- 4 Then an ELM is triggered:
  - o the pressure gradient and the  $j_{bs}$  collapse.
  - o the process starts again.

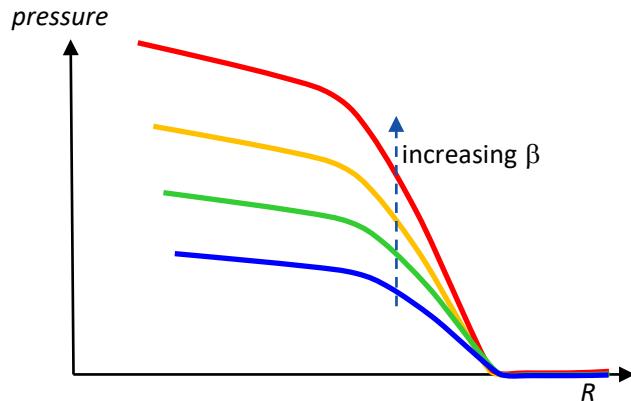
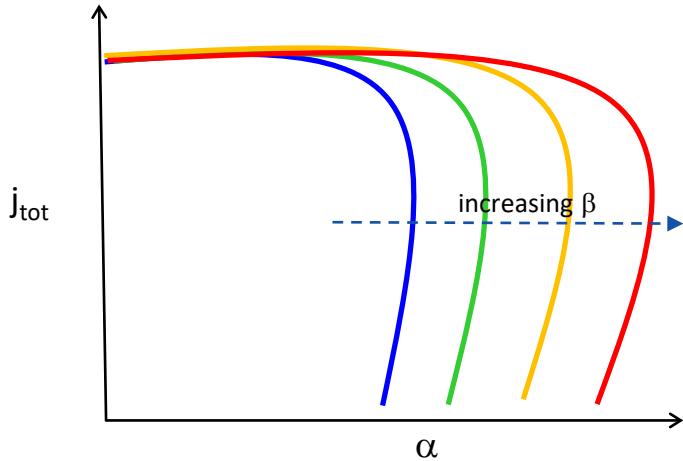


# Parameters that affect the pedestal: $\beta$

- $$\beta = \frac{\langle p \rangle}{B^2/(2\mu_0)}$$
 Volume averaged pressure
  - the increase of  $\beta$  leads to the increases of the Shafranov shift.
    - the Shafranov shift has a stabilizing effect on the ballooning modes.
    - the ballooning modes boundary moves to higher  $\alpha$
    - the pre-FIM pedestal pressure gradient increases
- $p_{\text{ped}}^{\text{ped}}$  increases with increasing  $\beta$ .

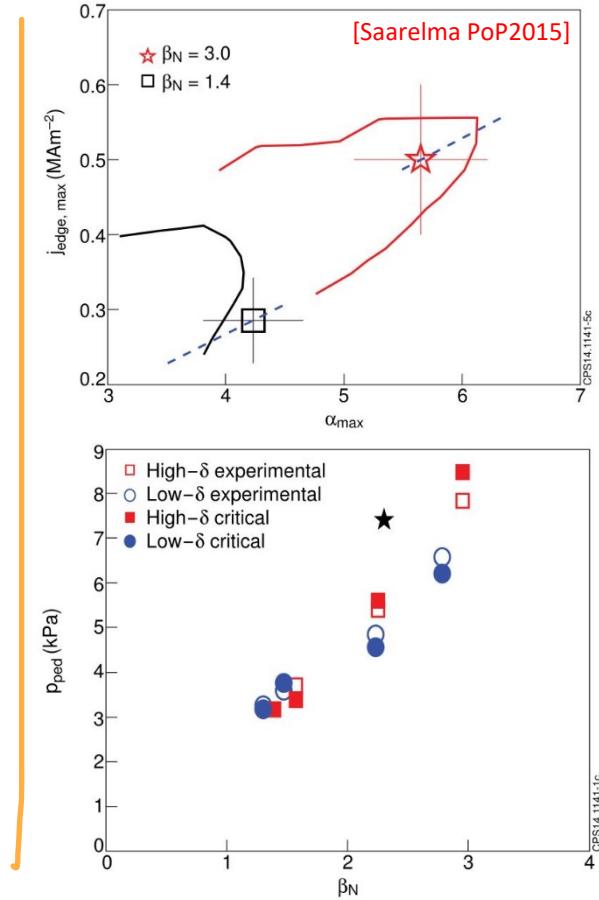
Pedestal pressure

To increase beta  
we increase power



# Parameters that affect the pedestal: $\beta$

- $\beta = \frac{\langle p \rangle}{B^2/(2\mu_0)}$
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    - the ballooning modes boundary moves to higher  $\alpha$
    - the pre-ELM pedestal pressure gradient increases
- $p_{\text{ped}}$  increases with increasing  $\beta$ .



# Parameters that affect the pedestal: $\nu^*$

- Collisionality

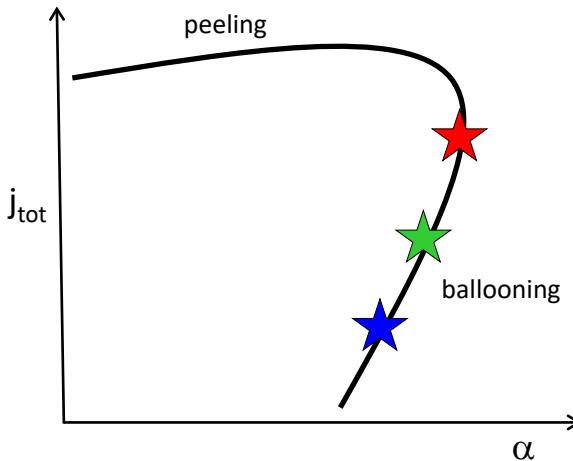
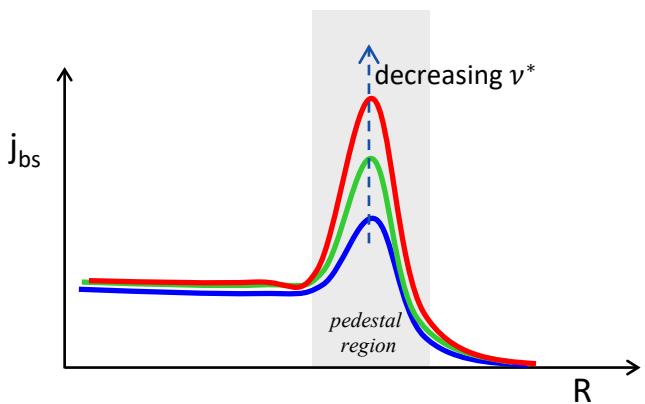
$$\nu^* = \text{cln} \Lambda \frac{R q n_e}{\varepsilon^{3/2} (T_e)^2}$$

- the collisionality has a major effect on  $j_{bs}$ .

[Sauter PoP1999] [Redl PoP2021]

- Approximately:

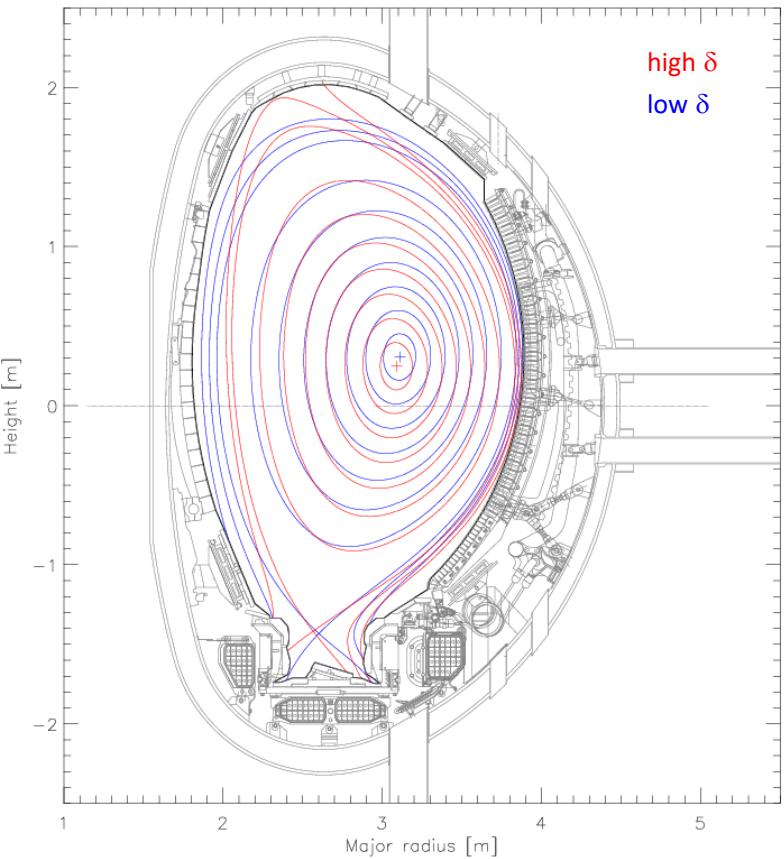
$$j_{bs} \approx \nu^{*-1}$$



- The reduction of collisionality tends to increase  $\nabla p$ , if the pedestal is near the ballooning boundary

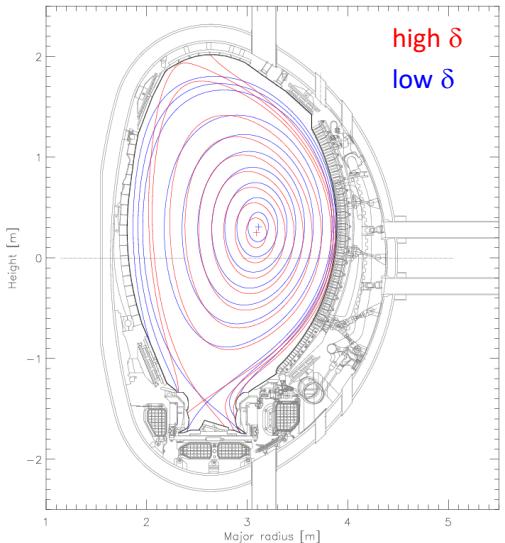
# Parameters that affect the pedestal: $\delta$

- $\delta$ : plasma triangularity

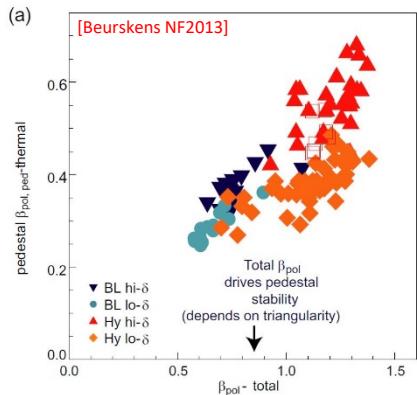
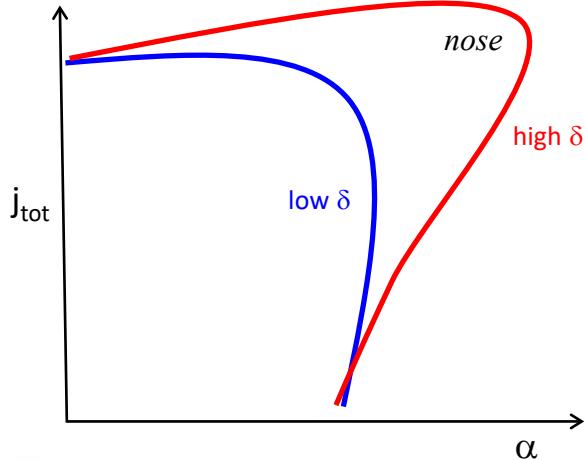


# Parameters that affect the pedestal: $\delta$

- $\delta$ : plasma triangularity
- the increase of  $\delta$  stabilizes part of the ballooning modes.
- the PB is strongly shaped at high  $\delta$  and a so called "nose" is formed:
  - high  $j_{bs}$   $\rightarrow \nabla p$  increases with increasing  $\delta$ .
  - low  $j_{bs}$   $\rightarrow \nabla p$  does not change much with  $\delta$ .



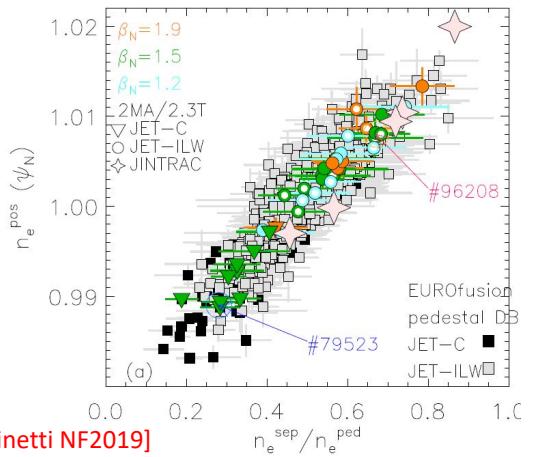
[Saibene PPCF2002]  
 [Beurskens NF2013]  
 [Urano NF2014]



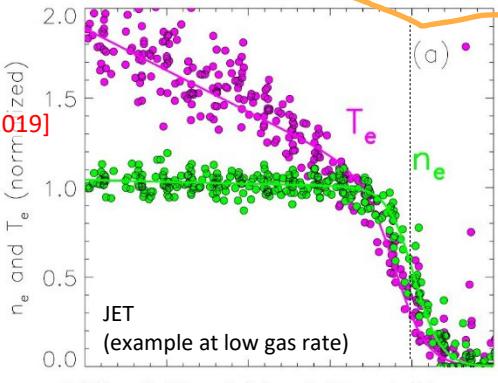
# Parameters that affect the pedestal: $n_e^{\text{sep}}/n_e^{\text{ped}}$

- Separatrix density and pedestal position:

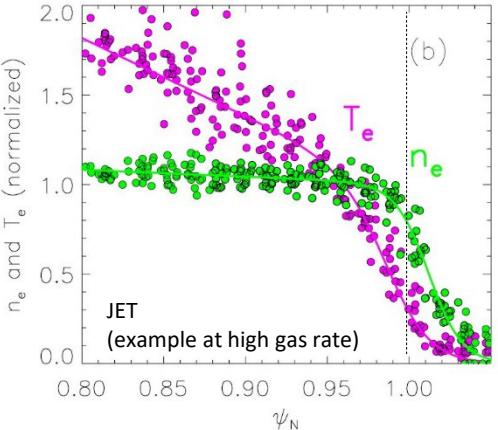
- The pedestal position can vary significantly depending on engineering parameters [Stefanikova NF2019]
- $n_e$  pedestal position and separatrix density are strongly correlated. [Frassinetti NF2021]



[Frassinetti NF2019]

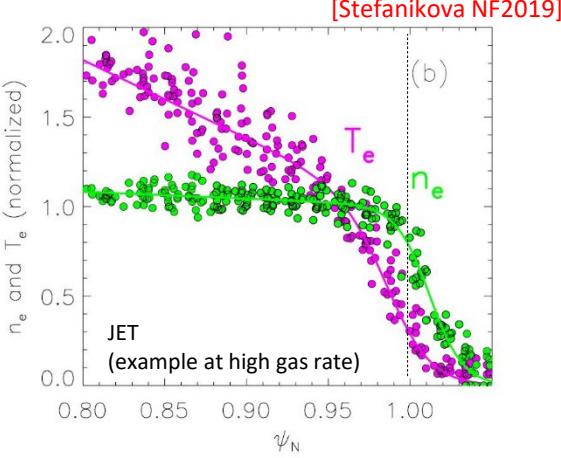
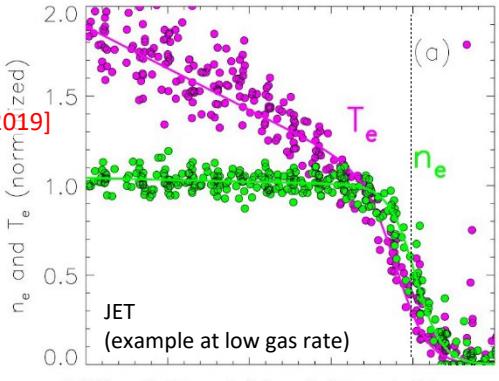
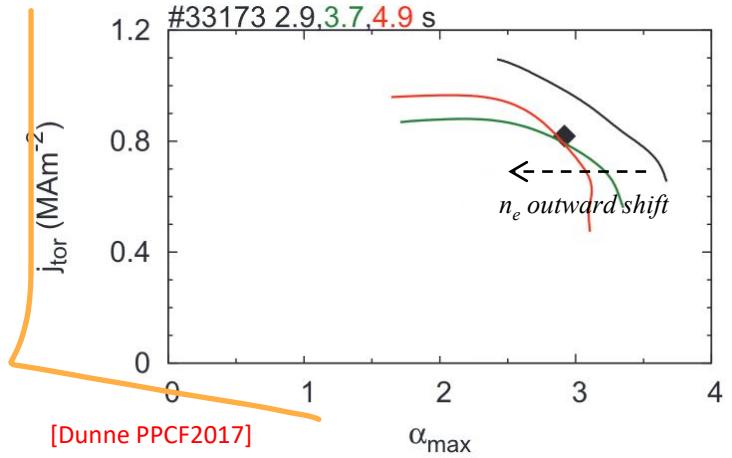


[Stefanikova NF2019]



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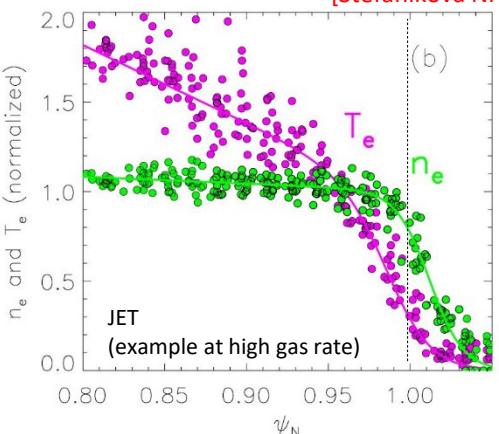
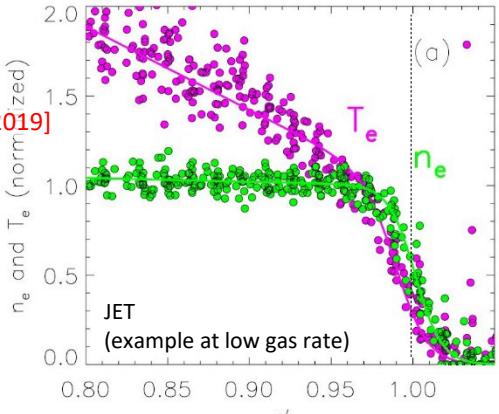
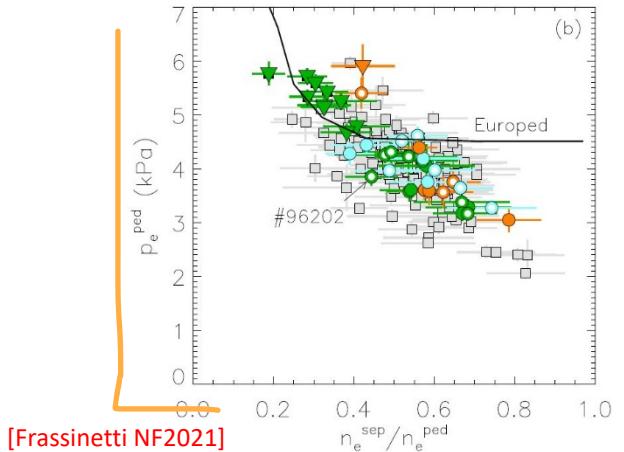
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  - The increase of  $n_e^{\text{sep}}/n_e^{\text{ped}}$  and the outward shift of the pedestal position destabilizes ballooning modes and reduces the pedestal height [Dunne PPCF2017][Frassinetti NF2019]



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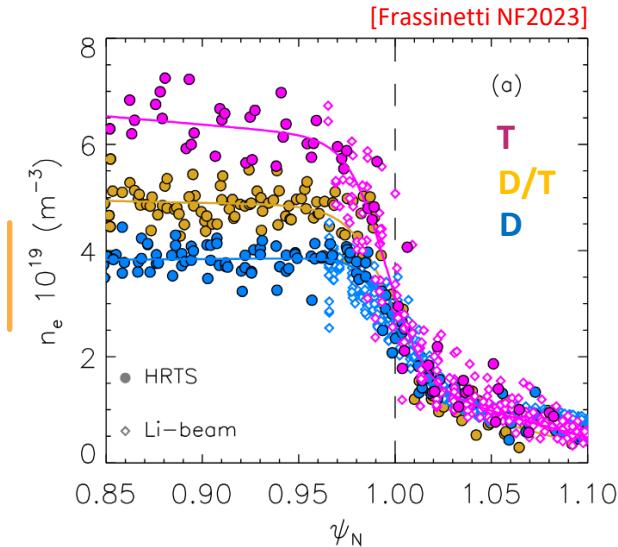
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# Parameters that affect the pedestal: isotope

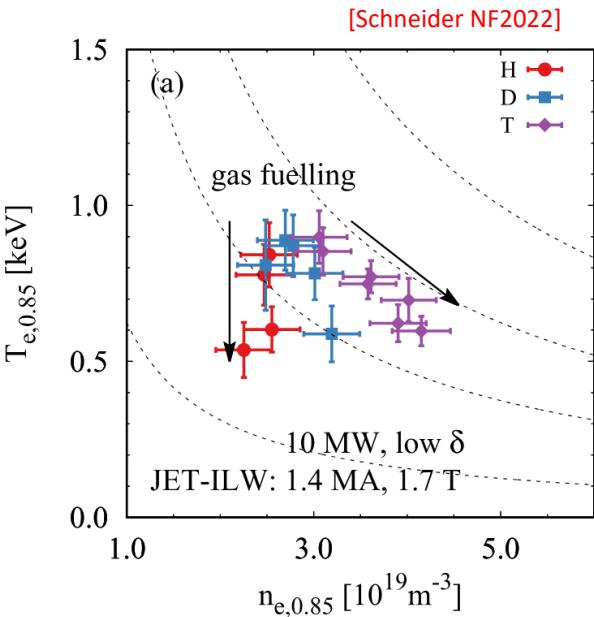
- Isotope mass ( $A$ ) can affect the pedestal:
  - pedestal density
    - $n_e^{\text{ped}}$  increases with increasing  $A$  (JET, DIII-D, AUG, JT60-U...)
    - Effect observed both from  $H \rightarrow D$  and from  $D \rightarrow T$

[Urano NF2013] [Horvath NF2021]  
[Gohil IAEA2008] [Schneider PPCF2021]  
[Maggi PPCF2018] [Schneider NF2022]  
[Frassinetti NF2023]



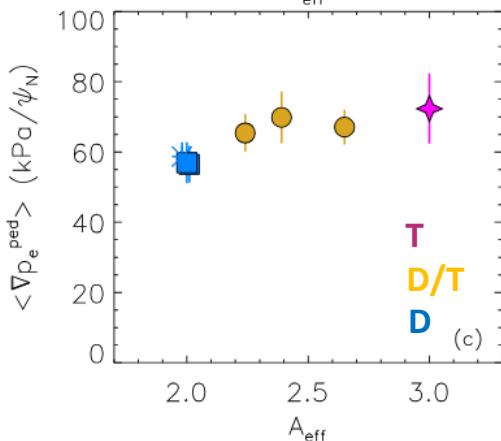
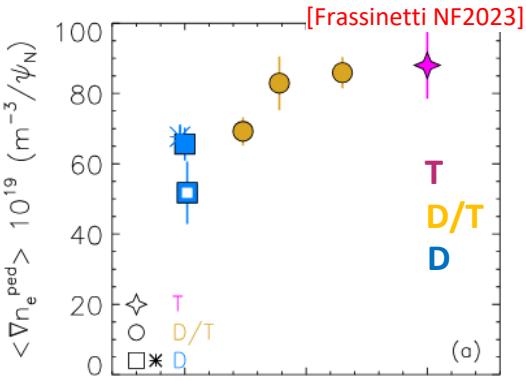
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    - In JET and AUG:  $p_e^{\text{ped}}$  increases with increasing  $A$ . [Maggi PPCF2018, Schneider NF2023]
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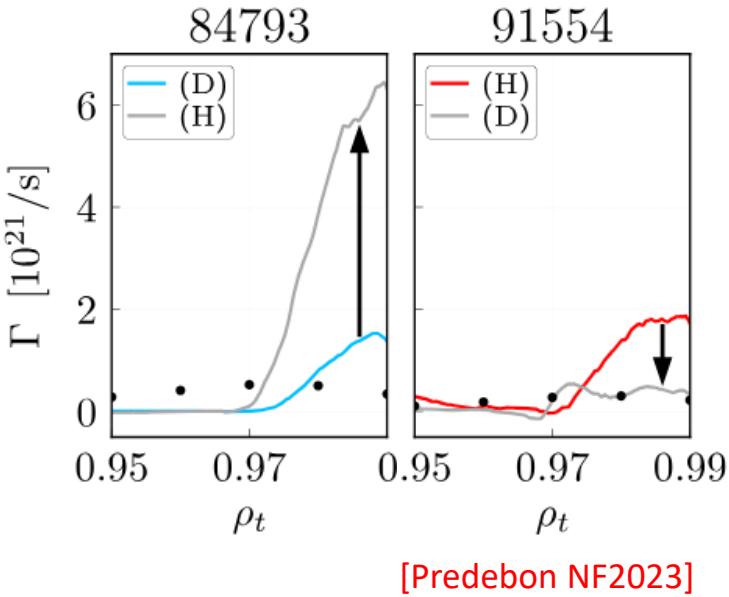
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    - In DIII-D and JT60-U: no significant effect related to A [Urano NF2013] [Gohill IAEA2008]
  - Pedestal gradients:
    - The increase in  $n_e^{\text{ped}}$  is due to an increase in  $\nabla n_e \rightarrow$  likely, A affects the transport (but other reasons cannot be excluded yet) [Frassinetti NF2023]
    - The increase in  $p_e^{\text{ped}}$  is due to an increase in  $\nabla p_e \rightarrow$  A can affects the stability [Frassinetti NF2023]
  - Pedestal width: no significant variation



# Parameters that affect the pedestal: isotope

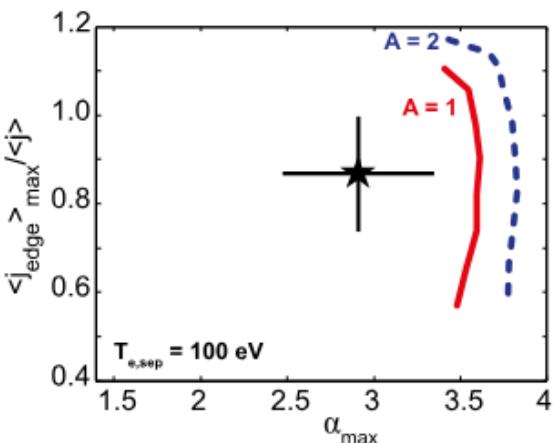
- Understandings of the effects of A
  - pedestal density
    - Sources: experimental estimate very challenging
    - GK results show in H vs D plasmas in JET-ILW show reduction of particle transport with increasing mass [Predebon NF2023] (due to reduced ITG turbulence)
    - Further work on-going before conclusive claims



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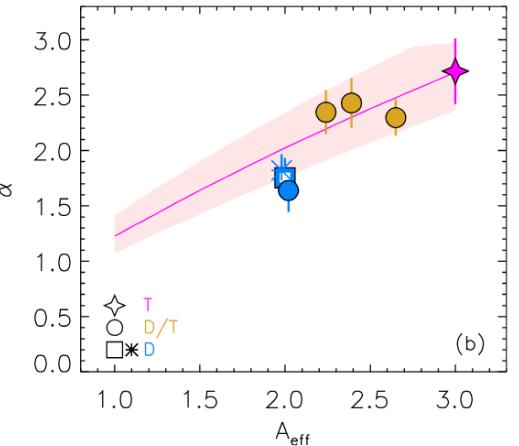
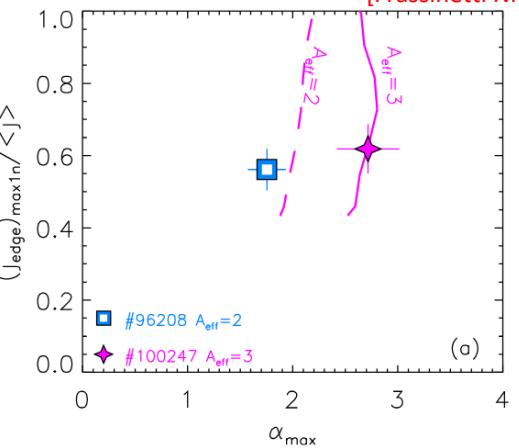


[Horvath NF2021]

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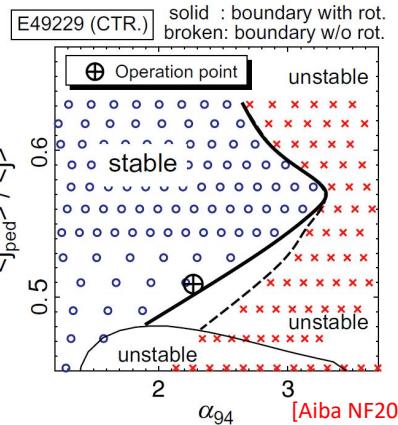
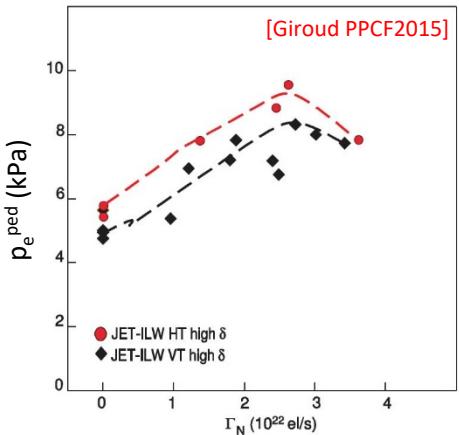
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  - pedestal pressure
    - Ideal MHD cannot explain the increase of  $\nabla p_e$  [Horvath NF2021]
    - Resistive MHD shows a reasonable agreement with experimental results [Frassinetti NF2023]



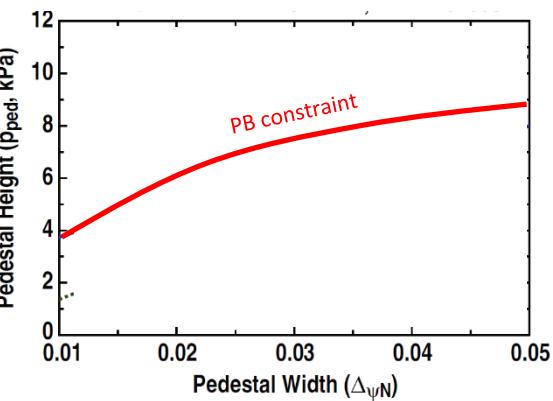
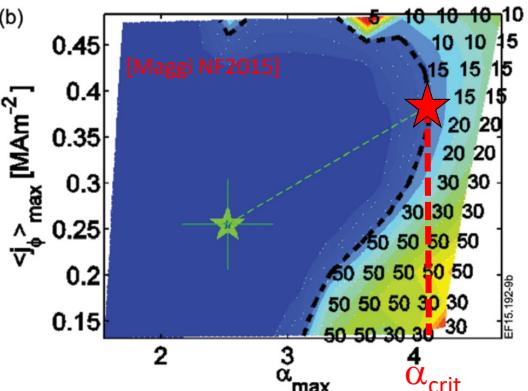
# Parameters that affect the pedestal

- Other parameters that affect the pedestal stability are:
  - **Impurities and seeding.**  $Z_{\text{eff}}$  affects collisionality and  $j_{\text{bs}}$ . It affects the electron pressure via the dilution effect. It can affect turbulent transport. Not fully understood yet.  
 [Giroud NF2013, PPCF2015, IAEA2018, Dunne PPCF2017]  
 [Saarela PoP2015]
  - **q-profile.** A change in q-profiles affects the shear. [Snyder PoP2002]
  - **Pedestal width.** A wider pedestal can contain more ballooning modes, so it is more unstable  
 [Snyder PoP2002]
  - **Plasma rotation.** [Aiba NF2018]
  - **Density at the pedestal top.** Not trivial effects, see later  
 [Snyder NF2011]



# Pedestal predictions: the PB constraint

- Can we use the PB model to predict the pedestal pressure height before the ELM?
- The PB model identify the critical normalized pressure gradient ( $\alpha_{\text{crit}}$ ) above which the PB modes are destabilized.
  - It can be used to determine  $\nabla p$ .
- For a specific pedestal width, the PB model can determine the critical  $\nabla p$  at which the PB modes are destabilized.
  - for this specific width, the critical pressure height can determined from  $(\nabla p)_{\text{crit}}$ .
  - A correlation between width and critical pressure can be obtained. This is often called "PB constraint"
- More information is necessary to predict pedestal height and width.



# Pedestal predictions: the KBM constraint

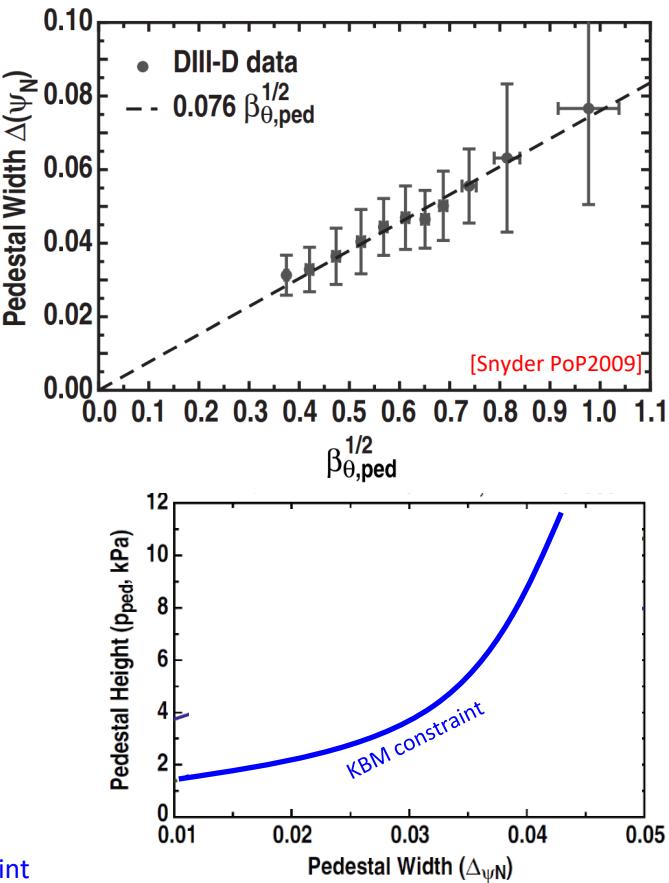
- The other constraint can come from pedestal transport
- The problem is that the pedestal transport is (often) driven by turbulence. Turbulence studies are not trivial and very time consuming
- The most successful approach, so far, has been developed in DIII-D [Snyder PoP2009]
  - experimental results suggest that DIII-D pedestal transport is driven by kinetic ballooning modes (KBMs)
  - from the theoretical arguments, it can be derived that for pedestals limited by the KBM turbulence:

$$w_{ped} = c \sqrt{\beta_{\theta}^{ped}}$$

- an experimental fit from DIII-D data gives:

$$w_{ped} = 0.076 \sqrt{\beta_{\theta}^{ped}}$$

KBM constraint



# Pedestal predictions: the KBM constraint

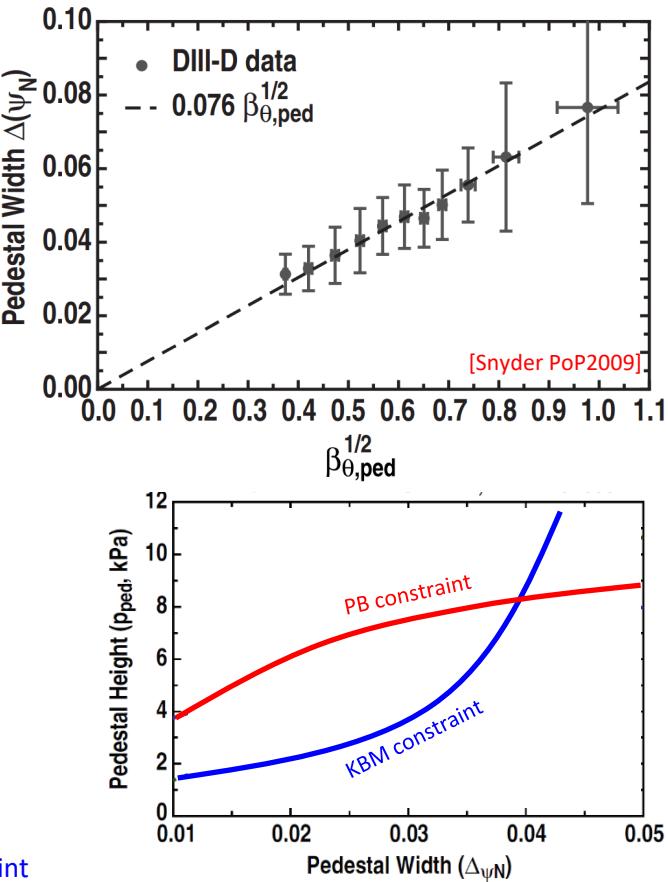
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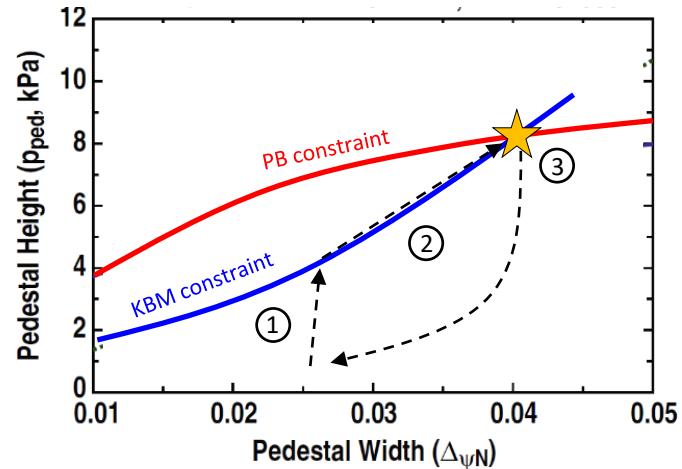


# The EPED1 model

- The EPED1 model predicts pedestal pressure height and pedestal pressure width using the
  - [Snyder PoP2009]
  - [Snyder NF2011]
- KBM constraint:  
local KBM stability → "clamps"  $\nabla p$
- PB constraint:  
global PB stability → triggers the ELM

## THE ELM CYCLE ACCORDING TO EPED1:

- ①  $\nabla p$  grows unconstrained
- ② KBM boundary is reached:
  - $\nabla p$  is "clamped"
  - The pedestal height grows via the increase of the pedestal width:
- ③ PB boundary is reached
  - ELM triggered



# The EPED1 model

- EPED1 tends to predict the pedestal pressure height rather well, for a large of parameters and in many machines.

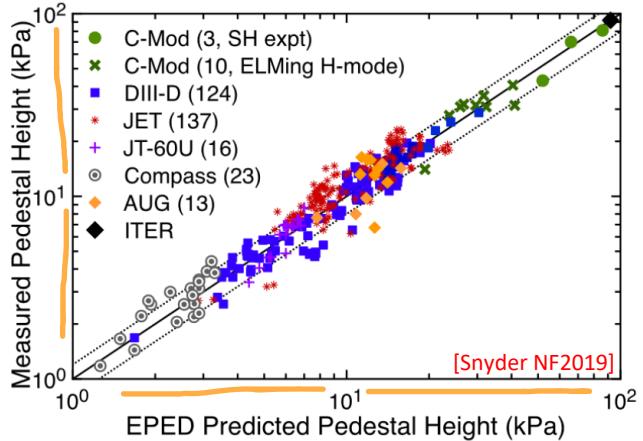
[Snyder NF2019]

- EPED1 is a useful tool to test the PB model.

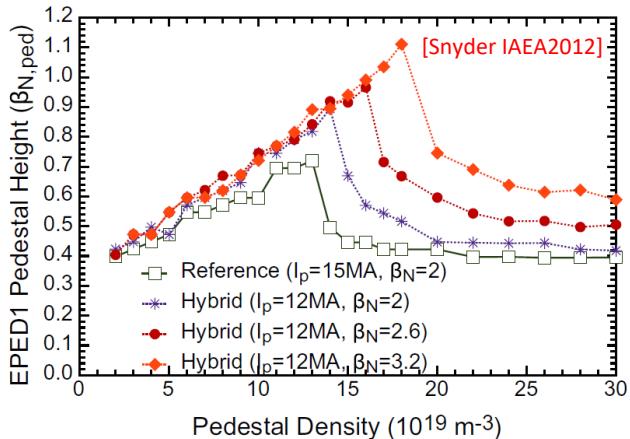
- EPED1 is widely used to predict the pedestal height (also in ITER).

- Example: prediction of pedestal pressure dependence with:

- density
- $\beta$



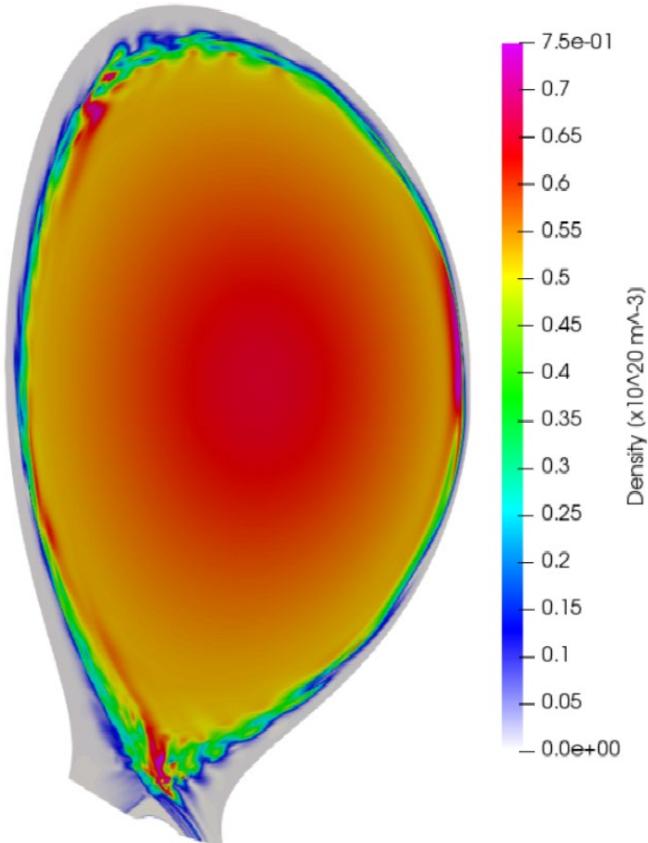
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[Snyder IAEA2012]

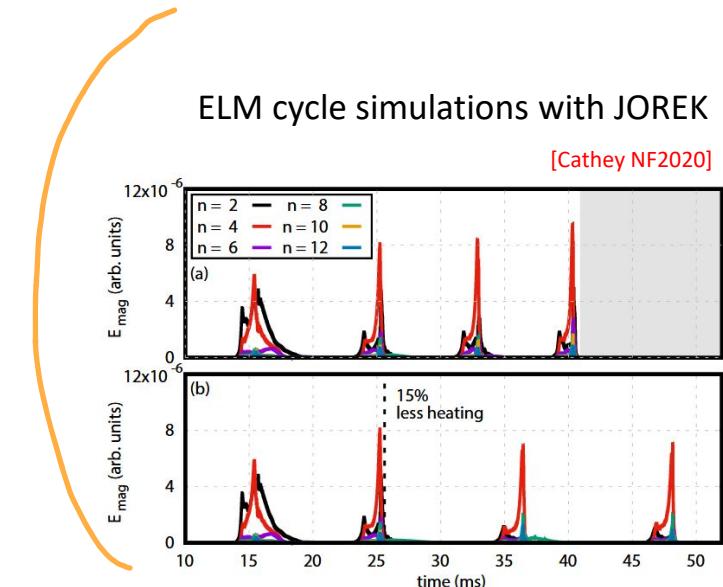
# Non-linear MHD modelling

- EPED1 works relatively well, but it is a linear model:
  - it does not predict time evolutions
  - cannot predict ELM energy losses
- Non-linear codes are necessary for modelling the details of the ELMs.
- Recent results with the JOREK code are very promising: [Huijsmans NF2007]



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- Recent results with the JOREK code are very promising: [Huijsmans NF2007]
  - type I ELMs start to be modeled rather well  
[Cathey NF2021]
  - ELMs similar to those in small ELMs scenarios have also been modelled.  
[Cathey PPCF2022]



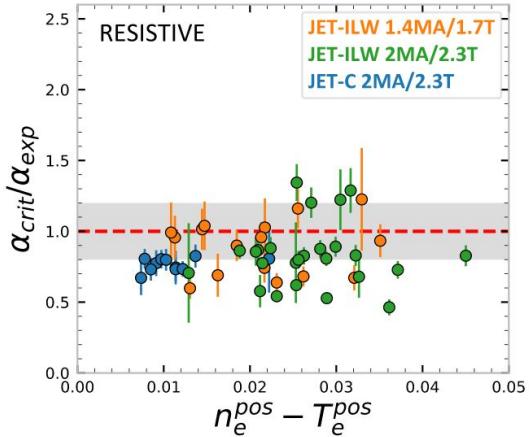
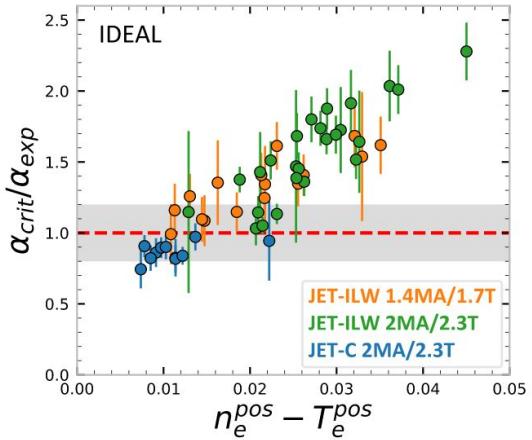
# Some active research areas

- Discrepancies between PB stability and experimental results, especially in JET-ILW, have been observed.
  - [Frassinetti NF2019], [Frassinetti NF2021], [Nyström NF2022]
  - Resistive MHD might play a role
- Super H-mode: DIII-D results show that at high  $\delta$  the 2nd stability region can be accessed. [Snyder NF2015]
  - can other experiments reach this region?
- Peeling limited pedestals
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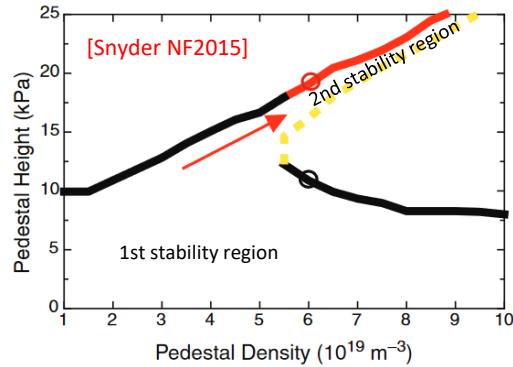
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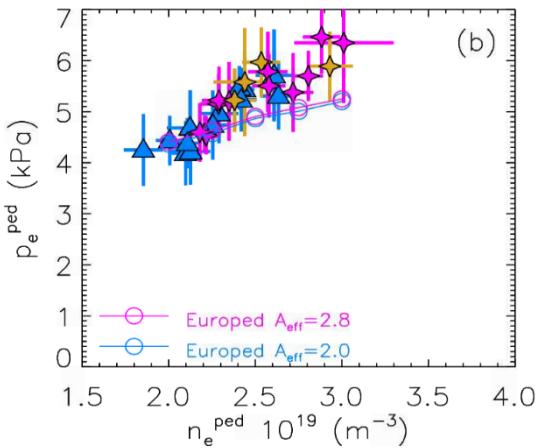
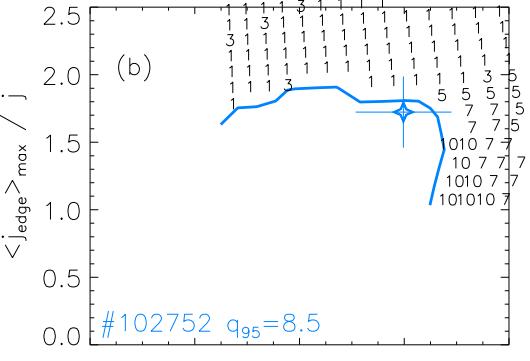
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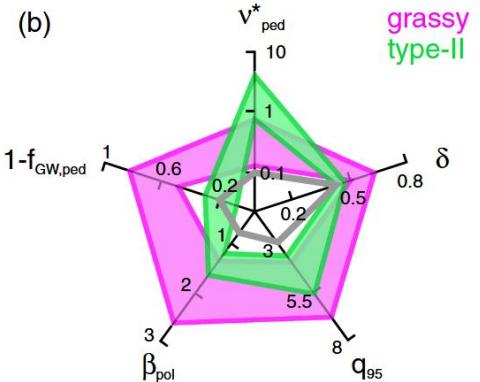
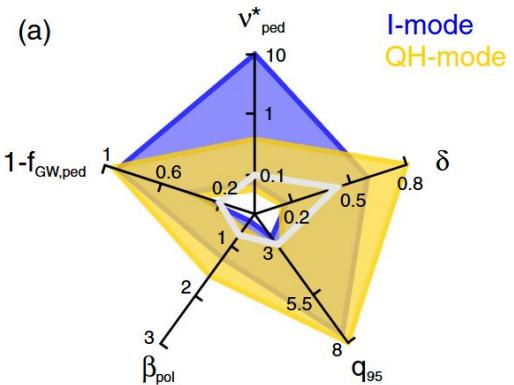
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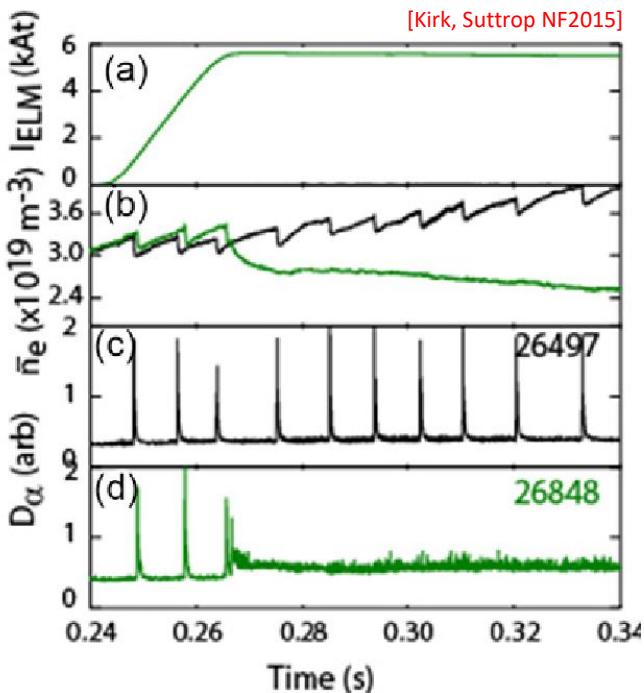
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# Some useful references

The choice of the following papers is based on two criteria:

- overview papers, when possible.
- most recent papers.

This list does not necessarily cite the original papers on the topic.  
Many excellent papers have not been included.

- Pedestal physics: [Groebner PPCF2023]  
[Urano NF2014]  
[Leonard PoP2014]
- LH transition: [Bourdelle NF2020]
- Pedestal structure: [Frassinetti NF2021]
- Isotope effect: [Maggi PPCF2018]
- ELMs: [Zohm PPCF1996]  
[Leonard PoP2014]
- PB model: [Wilson PoP1999]  
[Snyder PoP2002]  
[Snyder NF2004]
- EPED model: [Snyder PoP2009]  
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