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# Homework #4

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February 12, 2016

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## Problem 1

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The following nonhomogeneous IBVP describes heat flow in a rod whose ends are held at temperatures  $u_0$ ,  $u_1$ :

$$\begin{aligned} u_t &= u_{xx} & 0 < x < 1, \quad t > 0 \\ u(0, t) &= u_0, \quad u(1, t) = u_1 \\ u(x, 0) &= f(x) \end{aligned} \tag{1}$$

(a) Find the steady state temperature  $U(x)$  that satisfies

$$\begin{aligned} U_{xx} &= 0 & 0 < x < 1 \\ U(0) &= u_0, \quad U(1) = u_1 \end{aligned}$$

$U_{xx} = 0$  implies  $U(x)$  is a linear function.

$$U(x) = a + bx$$

The boundary conditions imply

$$U(x) = u_0 + (u_1 - u_0)x$$

(b) Write  $u(x, t) = U(x) + v(x, t)$  and find the corresponding IBVP for  $v$ . Use separation of variables to solve for  $v$  and hence  $u$ .

If  $u(x, t) = U(x) + v(x, t)$ , then  $u(0, t) = u_0 = U(0) + v(0, t)$  implies  $v(0, t) = 0$ . Similarly,  $v(1, t) = 0$ . Also,  $u(x, 0) = f(x) = U(x) + v(x, 0)$  implies  $v(x, 0) = f(x) - U(x)$ . Lastly, partial differential equation  $u_t = u_{xx}$  implies  $v_t = U''(x) + v_{xx}$  but since  $U'' = 0$ , then  $v_t = v_{xx}$ . The problem becomes

$$\begin{aligned} v_t &= v_{xx} & 0 < x < 1, \quad t > 0 \\ v(0, t) &= 0, \quad v(1, t) = 0 \\ v(x, 0) &= f(x) - U(x) \end{aligned} \tag{2}$$

By our last homework, the solution of (2) is

$$v(x, t) = \sum_{n=1}^{\infty} c_n e^{-n^2 \pi^2 t} \sin(n\pi x)$$

where

$$c_n = 2 \int_0^1 (f(x) - U(x)) \sin(n\pi x) dx$$

Thus the solution to (1) is

$$u(x, t) = u_0 + (u_1 - u_0)x + \sum_{n=1}^{\infty} c_n e^{-n^2 \pi^2 t} \sin(n\pi x)$$

(c) How does  $u(x, t)$  behave as  $t \rightarrow \infty$ ?

As  $t \rightarrow \infty$ , the coefficients of the sin series decay hyper-exponentially, and thus the solution rapidly approaches the linear function  $U(x)$ , i.e.

$$\lim_{t \rightarrow \infty} u(x, t) = U(x)$$

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## Problem 2

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Define a first-order differential operator with complex coefficients acting  $L^2(0, 2\pi)$  by

$$A = -i \frac{d}{dx}.$$

(a) Show that  $A$  is formally self-adjoint.

$$\begin{aligned} \langle u, Av \rangle &= \int_0^{2\pi} -i \bar{u} v' dx \\ &= -i \left[ (\bar{u} v)_0^{2\pi} - \int_0^{2\pi} \bar{u}' v dx \right] \\ &= -i \left( \overline{u(2\pi)} v(2\pi) - \overline{u(0)} v(0) \right) + \int_0^{2\pi} i \bar{u}' v dx \\ &= -i \left( \overline{u(2\pi)} v(2\pi) - \overline{u(0)} v(0) \right) + \int_0^{2\pi} -i \bar{u}' v dx \\ &= -i \left( \overline{u(2\pi)} v(2\pi) - \overline{u(0)} v(0) \right) + \langle Au, v \rangle \end{aligned}$$

Thus, given adequate boundary conditions,  $A$  is self-adjoint, i.e.  $A$  is formally self-adjoint.

(b) Show that  $A$  with periodic boundary conditions  $u(0) = u(2\pi)$  is self-adjoint, and find the eigenvalues and eigenfunctions of the corresponding eigenvalue problem

$$-i u' = \lambda u, \quad u(0) = u(2\pi).$$

If  $u(0) = u(2\pi)$ , then  $\bar{0} = \overline{u(2\pi)}$ . Thus,

$$\langle u, Av \rangle = -i(u(0)[v(2\pi) - v(0)]) + \langle Au, v \rangle$$

and so the adjoint boundary condition is  $v(0) = v(2\pi)$ . In this case,  $A$  is self-adjoint.

The solution to  $-i u' = \lambda u$  is

$$u(x) = c \exp(\lambda i x)$$

for some constant  $c$ . The periodicity of  $u$  implies

$$c = c \exp(\lambda i 2\pi) \implies \lambda = n, \quad n \in \mathbb{Z}$$

- (c) What are the adjoint boundary conditions to the Dirichlet condition  $u(0) = 0$  at  $x = 0$ ? Is  $A$  with this Dirichlet boundary condition self-adjoint? Find all eigenvalues and eigenfunctions of the corresponding eigenvalue problem

$$-iu' = \lambda u, \quad u(0) = 0.$$

How does your result compare with the properties of finite-dimensional eigenvalue problems for matrices?

If  $u(0) = 0$  then  $\overline{u(0)} = 0$  and

$$\langle u, Av \rangle = -i\overline{u(2\pi)}v(2\pi) + \langle Au, v \rangle$$

Thus the adjoint condition is  $v(2\pi) = 0$ , and in this case,  $A$  is self-adjoint. The solution to  $-iu' = \lambda u$  is

$$u(x) = c \exp(\lambda i x)$$

for some constant  $c$ . The boundary condition  $u(0) = 0$  implies  $c = 0$ , and thus  $u \equiv 0$ . Since eigenvalues cannot have 0 eigenfunctions, there are no eigenvalues for this eigenvalue problem. This does not happen in finite-dimensional eigenvalue problems since all matrices (over algebraically closed fields, like  $\mathbb{C}$ ) have at least one eigenvalue.

### Problem 3

Let  $A$  be a Sturm-Liouville operator, given by

$$Au = -(pu')' + qu,$$

acting in  $L^2(a, b)$ . Verify that  $A$  with the Robin boundary conditions

$$\alpha u'(a) + u(a) = 0, \quad u'(b) + \beta u(b) = 0$$

is self adjoint.

$$\begin{aligned} \langle u, Av \rangle &= \int_a^b [u(-(pv')' + qv)] dx \\ &= - \int_a^b up'v' dx - \int_a^b upv'' dx + \int_a^b uqv dx \end{aligned}$$

We can integrate the middle integral by parts:

$$\begin{aligned} \int_a^b upv'' dx &= [upv']_a^b - \int_a^b (up'v' + u'pv') dx \\ \Rightarrow \langle u, Av \rangle &= -[upv']_a^b + \int_a^b u'pv' dx + \int_a^b uqv dx \end{aligned}$$

Again, we can integrate the leftmost integral by parts:

$$\int_a^b u'pv' dx = [u'pv]_a^b - \int_a^b (u''pv - u'p'v) dx$$

$$\begin{aligned}
\Rightarrow \langle u, Av \rangle &= [p(u'v - uv')]_a^b + \int_a^b (-(u''pv + u'p'v) + uqv) dx \\
&= [p(u'v - uv')]_a^b + \int_a^b [(-(pu')' + qu)v] dx \\
&= [p(u'v - uv')]_a^b + \langle Au, v \rangle
\end{aligned}$$

The Robin boundary conditions imply  $u'(a) = -\frac{u(a)}{\alpha}$  and  $u'(b) = -\beta u(b)$ . If we impose the adjoint condition

$$\alpha v'(a) + v(a) = 0, \quad v'(b) + \beta v(b) = 0,$$

then

$$\begin{aligned}
[p(u'v - uv')]_a^b &= p(b)(u'(b)v(b) - u(b)v'(b)) - p(a)(u'(a)v(a) - u(a)v'(a)) \\
&= p(b)(-\beta u(b)v(b) + \beta u(b)v(b)) - p(a)\left(-\frac{1}{\alpha}u(a)v(a) + \frac{1}{\alpha}u(a)v(a)\right) \\
&= 0
\end{aligned}$$

and thus  $A$  is self-adjoint.

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## Problem 4

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Show that the eigenvalues of the Sturm-Liouville problem

$$\begin{aligned}
-u'' &= \lambda u & 0 < x < 1 \\
u(0) &= 0, & u'(1) + \beta u(1) &= 0
\end{aligned}$$

are given by  $\lambda = k^2$  where  $k > 0$  satisfies the equation

$$\beta \tan k + k = 0.$$

Show graphically that there is an infinite sequence of simple eigenvalues  $\lambda_1 < \lambda_2 < \dots < \lambda_n < \dots$  with  $\lambda_n \rightarrow \infty$  as  $n \rightarrow \infty$ . What is the asymptotic behavior of  $\lambda_n$  as  $n \rightarrow \infty$ ?

If  $\lambda = 0$ , the solutions to  $-u'' = \lambda u$  are

$$u(x) = c_1 + c_2 x$$

Then  $u(0) = 0 \Rightarrow c_1 = 0$ , i.e.

$$\begin{aligned}
u(x) &= c_2 x \\
u'(x) &= c_2
\end{aligned}$$

The Robin boundary condition  $u'(1) + \beta u(1) = 0$  implies

$$0 = c_2 + \beta c_2 = c_2(1 + \beta)$$

If  $\beta \neq -1$ , then  $c_2 = 0$  and thus  $u \equiv 0$ . If  $\lambda \neq 0$ , the solutions to  $-u'' = \lambda u$  are

$$u(x) = c_1 \exp(\sqrt{\lambda}x) + c_2 \exp(-\sqrt{\lambda}x)$$

If  $\lambda > 0$ , then  $u(0) = 0$  implies  $c_1 = -c_2$ , i.e.

$$\begin{aligned} u(x) &= c_1 \left[ \exp(\sqrt{\lambda}x) - \exp(-\sqrt{\lambda}x) \right] \\ u'(x) &= c_1 \sqrt{\lambda} \left[ \exp(\sqrt{\lambda}x) + \exp(-\sqrt{\lambda}x) \right] \end{aligned}$$

The Robin boundary condition  $u'(1) + \beta u(1) = 0$  implies

$$\begin{aligned} 0 &= c_1 \sqrt{\lambda} \left[ \exp(\sqrt{\lambda}) + \exp(-\sqrt{\lambda}) \right] + \beta c_1 \left[ \exp(\sqrt{\lambda}) - \exp(-\sqrt{\lambda}) \right] \\ &= (\sqrt{\lambda} + \beta) \exp(\sqrt{\lambda}) + (\sqrt{\lambda} - \beta) \exp(-\sqrt{\lambda}) \end{aligned}$$

If  $\lambda < 0$ , then the solution is given by

$$u(x) = c_1 \cos(\sqrt{-\lambda}x) + c_2 \sin(\sqrt{-\lambda}x)$$

The boundary condition  $u(0) = 0$  implies  $c_1 = 0$ , i.e.

$$\begin{aligned} u(x) &= c_2 \sin(\sqrt{-\lambda}x) \\ u'(x) &= c_2 \sqrt{-\lambda} \cos(\sqrt{-\lambda}x) \end{aligned}$$

The Robin boundary condition  $u'(1) + \beta u(1) = 0$  implies

$$\begin{aligned} 0 &= c_2 \sqrt{-\lambda} \cos(\sqrt{-\lambda}) + \beta c_2 \sin(\sqrt{-\lambda}) \\ &= \sqrt{-\lambda} + \beta \tan(\sqrt{-\lambda}) \quad \text{provided } c_2 \neq 0 \\ &= \beta \tan k + k \end{aligned}$$

where  $k^2 = -\lambda$ . Let the solutions to  $0 = \beta \tan k + k$  be  $k_n$  where  $\frac{\pi}{2} < |k_{n+1} - k_n| < \pi$  for all  $n$ . Also,  $k_n \rightarrow \infty$ , and thus  $\lambda_n \rightarrow \infty$  as  $n \rightarrow \infty$ .

## Problem 5

The following IBVP describes heat flow in a rod whose left end is held at temperature 0 and whose right end loses heat to the surroundings according to Newton's law of cooling (heat flux is proportional to the temperature difference):

$$\begin{aligned} u_t &= u_{xx} & 0 < x < 1, \quad t > 0 \\ u(0, t) &= 0, & u_x(1, t) &= -\beta u(1, t) \\ u(x, 0) &= f(x) \end{aligned}$$

Solve this IBVP by the method of separation of variables.

Assume the solution  $u(x, t) = F(x)G(t)$ . Then  $FG' = F''G$ , which implies  $\frac{G'}{G} = \frac{F''}{F} = \lambda$ , i.e.

$$G' - \lambda G = 0 \quad \text{and} \quad F'' - \lambda F = 0$$

Then

$$G(t) = c \exp[\lambda t]$$

and, if  $\lambda \neq 0$ , then

$$F(x) = c_1 \exp\left[\sqrt{\lambda}x\right] + c_2 \exp\left[-\sqrt{\lambda}x\right]$$

If  $\lambda = 0$ , then

$$F(x) = c_1 + c_2 x$$

The left boundary condition implies  $F(0) = 0$ , and so if  $\lambda = 0$  then  $c_1 = 0$ . Then  $F(x) = c_2 x$ . The right boundary condition implies  $F'(1) = -\beta F(1)$ . Thus  $c_2 = -\beta c_2 \iff c_2(1 + \beta) = 0 \iff \beta = -1$  for nontrivial solutions  $F$ . If  $\lambda > 0$ , then the left boundary condition implies  $c_1 = -c_2$ , and thus

$$\begin{aligned} F(x) &= c_1 \left( \exp\left[\sqrt{\lambda}x\right] - \exp\left[-\sqrt{\lambda}x\right] \right) \\ F'(x) &= c_1 \sqrt{\lambda} \left( \exp\left[\sqrt{\lambda}x\right] + \exp\left[-\sqrt{\lambda}x\right] \right) \end{aligned}$$

The right boundary condition implies

$$\begin{aligned} -c_1 \beta \left( \exp\left[\sqrt{\lambda}\right] - \exp\left[-\sqrt{\lambda}\right] \right) &= c_1 \sqrt{\lambda} \left( \exp\left[\sqrt{\lambda}\right] + \exp\left[-\sqrt{\lambda}\right] \right) \\ \left( \sqrt{\lambda} + \beta \right) \exp\left[\sqrt{\lambda}\right] + \left( \sqrt{\lambda} - \beta \right) \exp\left[-\sqrt{\lambda}\right] &= 0 \end{aligned}$$

If  $\lambda < 0$ , then

$$F(x) = c_1 \sin\left(\sqrt{-\lambda}x\right) + c_2 \cos\left(\sqrt{-\lambda}x\right)$$

By number 4,

$$F(x) = c_1 \sin\left(\sqrt{-\lambda}x\right)$$

Thus the general solution can be written as

$$u(x, t) = \sum_{n=1}^{\infty} c_n \exp[\lambda_n t] \sin\left(\sqrt{-\lambda_n}x\right)$$

where  $\lambda_n < 0$  for all  $n$  and  $\lambda_n \rightarrow -\infty$  as  $n \rightarrow \infty$ .