

BORCEA–VOISIN MIRROR SYMMETRY FOR LANDAU–GINZBURG MODELS

ABSTRACT. FJRW theory is a formulation of physical Landau–Ginzburg models with a rich algebraic structure, rooted in enumerative geometry. By a major physical conjecture, the Landau–Ginzburg/Calabi–Yau correspondence, several birational morphisms of Calabi–Yau orbifolds should correspond to isomorphisms in FJRW theory. In this paper it is shown that not only does this claim prove to be the case, but is a special case of a wider FJRW isomorphism theorem.

Am: Should we say explicitly something about BV mirror symmetry here?

1. INTRODUCTION

Mirror symmetry is a phenomenon first described nearly thirty years ago by physicists. One can think of mirror symmetry as an exchange of information determined by Kähler structure of some physical model with information determined by complex structure of another physical model.

Mathematically, several different constructions have been described for various models, each in a different context. In this paper, we will focus on two particular versions of mirror symmetry. The first, which we call Borcea–Voisin (BV) mirror symmetry, relates two Calabi–Yau (CY) threefolds of a particular form, and the second, which has come to be known as Berglund–Hübsch–Krawitz (BHK) mirror symmetry, relates two Landau–Ginzburg (LG) models. We will extend BHK mirror symmetry to a new formulation of mirror symmetry for LG models based on the ideas of BV mirror symmetry and the Landau–Ginzburg/Calabi–Yau (LG/CY) correspondence. It is generally expected that computations should be less difficult on the LG side of the correspondence, thus we expect our new form of mirror symmetry to improve our ability to make calculations for Borcea–Voisin mirror symmetry.

Borcea–Voisin mirror symmetry was first described by Borcea in [4] and Voisin in [35] in the early 1990’s. In particular, they constructed a class of Calabi–Yau threefolds by first taking the product of an elliptic curve E (with involution σ_E), and a K3 surface S (with involution σ_S) and obtaining an involution $\sigma = (\sigma_E, \sigma_S)$ on the product. A *Borcea–Voisin threefold* is the crepant resolution $\tilde{\mathcal{Y}}$ of the quotient $(E \times S)/\sigma$. BV mirror symmetry is defined for this particular class of Calabi–Yau threefolds.

Borcea–Voisin mirror symmetry relies on mirror symmetry for K3 surfaces with involution, as described in the work of Nikulin (see [26]). That is, to each K3 surface S with involution, there is a family of K3 surfaces that are mirror dual to S .

With this in mind, the mirror symmetry posited by Borcea and Voisin is described as follows. Let S' be a K3 surface with involution $\sigma_{S'}$ mirror dual to S , and consider the crepant resolution $\tilde{\mathcal{Y}}'$ of the quotient

$$(E \times S')/\sigma'$$

where $\sigma' = (\sigma_E, \sigma_{S'})$. The pair $\tilde{\mathcal{Y}}$ and $\tilde{\mathcal{Y}}'$ are said to be Borcea–Voisin mirror pair. (Note that here we may treat elliptic curves as self-mirror.)

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On the other hand, for certain LG models, we can define BHK mirror symmetry. A LG model is a pair (W, G) of a quasihomogeneous polynomial W and group of symmetries G . Under certain conditions (see Section 1), we can associate to the pair (W, G) another LG model (W^T, G^T) , called the BHK mirror. It has been predicted and partially verified by mathematicians that the A-model construction for (W, T) will be equivalent in a certain way (even in non-CY cases) to the B-model construction for (W^T, G^T) .

N:Choose A-model or A-model.

Although mirror symmetry makes predictions for equivalence of the A-model and the B-model on many levels, in this paper we focus on the state space and Frobenius algebra structure of these models. In other words, we consider the A-model state space and algebra structure $\mathcal{A}_{W,G}$, constructed by Fan, Jarvis, and Ruan in [10]—and the B-model state space \mathcal{B}_{W^T, G^T} , given by Saito and Givental in [28, 29, 30, 31, 13, 14] and its associated algebra structure, given by Krawitz, following ideas of Kaufmann in [21, 19, 20, 23] (see Section 2 for a definition of both models). These are known to be isomorphic as vector spaces (cf. [23]), and in some cases as Frobenius algebras (cf. [12]), however in this article, we will consider the B-model state space for another pair, which we define later.

So far, we have described one form of mirror symmetry for CY threefolds, and one form of mirror symmetry for LG models. These two forms of mirror symmetry are related via the LG/CY correspondence. In particular, the LG/CY correspondence predicts a deep relationship between CY orbifolds and LG models in certain cases.

On the CY side of the correspondence, we have the orbifold quotient $[\{W = 0\}/\tilde{G}]$ in a quotient of weighted projective space, defined by the pair (W, G) , and on the LG side, we have the LG model defined by the same pair. The group \tilde{G} is simply the quotient of G by the exponential grading operator. On the simplest level, the LG/CY correspondence provides an isomorphism between the Chen-Ruan (orbifold) cohomology $H_{CR}^*([\{W = 0\}/\tilde{G}], \mathbb{C})$ (the CY state space) and the FJRW state space $\mathcal{A}_{W,G}$ mentioned above. This was proven for hypersurfaces by Chiodo–Ruan in [7].

Somewhat more generally, we may consider quotients of complete intersections $[\{W_1 = \dots = W_n = 0\}/\tilde{G}]$ in products of weighted projective spaces, and relate these to the FJRW theory of $(\sum_i W_i, G)$. It is expected that the correspondence will hold for these complete intersections as well.

The construction of Borcea–Voisin fits into this picture. First we construct the K3 surface. We have a list of the 95 weight systems $(v_0, \dots, v_3; d_2)$ given by Reid and Yonemura (cf. [27], [37]), where the resolution of a hypersurface of degree d_2 in the appropriate weighted projective space is a K3 surface. Each such weight system satisfies the *Calabi–Yau condition* $d_2 = \sum_{i=0}^3 v_i$. We consider a polynomial

$$(1) \quad W_2 = y_0^2 + f(y_1, y_2, y_3)$$

which is quasihomogeneous of degree d_2 with respect to this weight system and a group G_2 satisfying $J_2 \subset G_2 \subset \mathrm{SL}_{W_2}$. Exactly 44 of the 95 weight systems admit such a polynomial. We write X_{W_2, G_2} for the quotient by the group \tilde{G}_2 of the hypersurface defined by the vanishing of W_2 in the appropriate quotient of weighted projective space. There is an obvious nonsymplectic involution on X_{W_2, G_2} namely $\sigma_1 : y_0 \rightarrow -y_0$. The minimal resolution of X_{W_2, G_2} is a K3 surface with an involution induced by σ .

We construct an elliptic curve in a similar way. In other words, we can consider polynomials of the form $W_1 = x_0^2 + f_1(x_1, x_2)$ quasihomogeneous with respect to one of the weight systems $(u_0, u_1, u_2; d_1) = (3, 2, 1; 6)$ or $(2, 1, 1; 4)$, and a group of diagonal symmetries $J_1 \subset G_1 \subset \mathrm{SL}_{W_1}$. The resolution of X_{W_1, G_1} will be an elliptic curve and has an obvious corresponding involution.

Setting $\sigma = (\sigma_1, \sigma_2)$ we call the quotient

$$(X_{W_1, G_1} \times X_{W_2, G_2}) / \sigma$$

a *Borcea–Voisin orbifold*. As mentioned above, this is a complete intersection in the quotient of a product of weighted projective spaces. This orbifold has a crepant resolution, which is a Borcea–Voisin threefold. We will return to the question of whether the LG/CY correspondence holds for this case in Section 5.

Recall that BV mirror symmetry required a mirror K3 surface to construct the mirror manifold for a given BV threefold. In [1], Artebani–Boisierre–Sarti showed that the mirror symmetry described by Nikulin for K3 surfaces with involution agrees with what one would expect to obtain from BHK mirror symmetry for the pair (W_2, G_2) . Thus we can also describe a Borcea–Voisin mirror pair by constructing the K3 surface defined by (W_2^T, G_2^T) .

Inspired by the LG/CY correspondence, we will describe the corresponding LG model, given by the sum $W = W_1 + W_2$ and the group $\Sigma(G_1 \times G_2)$ generated by σ and $G_1 \times G_2$. Thus, our LG model is the pair $(W, \Sigma(G_1 \times G_2))$. In fact, we can generalize this construction on both sides of the LG/CY correspondence to more arbitrary symmetry groups and to any dimension.

Interestingly, the LG model provided by BHK mirror symmetry does not yield the same mirror as one would expect from Borcea–Voisin mirror symmetry. In this article we suggest another formulation of LG to LG mirror symmetry, which matches the Borcea–Voisin mirror symmetry construction, thus extending BHK mirror symmetry to this setting. We do this by defining the pair $(W^T, \Sigma(G_1^T \times G_2^T))$ as the Borcea–Voisin LG mirror dual to the pair $(W, \Sigma(G_1 \times G_2))$.

The main theorems of this work justify calling these a mirror pair. The first gives a state space correspondence.

Theorem 6. There is a mirror map $\mathcal{A}_{W, \Sigma G} \cong \mathcal{B}_{W^T, \Sigma G^T}$, which preserves the bi-degree.

Remark 0.1. If one views the bi-degree correctly, then the mirror map of Theorem 6 can be viewed as a rotation of the Hodge diamond by a right angle, as one expects in mirror symmetry.

On the way to proving this theorem, we obtain an isomorphism (called the twist map) within FJRW theory, which facilitates several computations of interest in the Landau–Ginzburg mirror symmetry arena. This is similar to the *twisting map* in [15].

The state space of FJRW theory itself possesses an interesting algebraic structure, namely it can be given the structure of a Frobenius algebra. The LG B-model state space can also be given such a structure. For LG to LG mirror symmetry, we also expect the isomorphism of state spaces to extend to an isomorphism of Frobenius algebras. This can be verified for many pairs (W, G) as we see in the second theorem.

Property ().* Let W be a non-degenerate, invertible singularity, and let G be an admissible group of symmetries of W . We say that *the pair (W, G) has Property (*)* if,

- (1) W can be decomposed as $W = \sum_{i=1}^M W_i$, where the W_i are themselves invertible polynomials having no variables in common with any other W_j .
- (2) For any element g of G whose associated sector $\mathcal{A}_g \subseteq \mathcal{A}_{W, G}$ is nonempty, and for each $i \in \{1, \dots, M\}$ the action of g fixes either all of the variables in W_i or none of them.
- (3) For any element g' of G^T whose associated sector of $\mathcal{B}_{g'} \subseteq \mathcal{B}_{W^T, G^T}$ is non-empty, and for each $i \in \{1, \dots, M\}$ the action of g' fixes either all of the variables in W_i^T or none of them.

Theorem 9. If W and ΣG satisfy Property (*) then $\mathcal{A}_{W, \Sigma G} \cong \mathcal{B}_{W^T, \Sigma(G^T)}$ as Frobenius algebras.

The Borcea–Voisin construction has been the topic of several recent articles, which are related to the current work. We briefly review those results.

In [2], Artebani–Boisier–Sarti prove a theorem in which they relate BHK mirror symmetry and BV mirror for the same polynomials considered here. There they consider only geometric models, using the LG model as a guide. They prove that the Calabi–Yau threefolds of Borcea–Voisin are birationally equivalent to the image of the twist map of Borcea (see [4] and Section 5.4).

Their result relies on the assumption that $\gcd(u_0, v_0) = 1$. In order to understand this restriction better, consider that there are two possible weight systems for W_1 yielding an elliptic curve and 44 possible weight systems for W_2 yielding a K3 surface with involution. Recall in this construction, we require our polynomials to be of the form (1). Only 48 of the 88 possible combinations of weight systems satisfy the gcd condition imposed in [2]. In this article we generalize this result in two ways. By considering the LG model more carefully, we remove the restriction on gcd’s, and extend the construction to all dimensions.

In [6] Chiodo–Kalashnikov–Veniani undertake to prove a mirror statement for the geometric model again using the LG model as a guide. There, however, the goal is to provide a mirror map for BV orbifolds. In other words, they prove an isomorphism of Chen–Ruan cohomology, whereas Borcea–Voisin considered the crepant resolution of these orbifolds. The authors consider two separate LG models, one for W_1 and one for W_2 , and then use a Künneth type formula to prove their result. Here we also generalize this result, removing the restriction that G be a product of groups, with a different method of proof. When G is not a product of groups, the crepant resolution of the quotient

$$X_{W_1} \times X_{W_2} / \widetilde{\Sigma G}$$

is no longer a Borcea–Voisin threefold, but rather is the Nakamura–Hilbert scheme of $\widetilde{\Sigma G}$ -orbits of $X_{W_1} \times X_{W_2}$.

In [32], the last author has considered exactly the form of mirror symmetry we propose here with the restriction that the defining polynomials must be Fermat type. In fact, he was able to show that for the mirror pairs we consider here, there is a mirror map relating the FJRW invariants of the A-model to the Picard–Fuchs equations of the B model. In [33] he also gave an LG/CY correspondence relating the FJRW invariants of the pair $(W, \Sigma G)$ to the corresponding Gromov–Witten invariants of the corresponding Borcea–Voisin orbifold. Although these results are broader in scope, the restriction to Fermat type polynomials is significant, reducing the number of weight systems from which one can select a K3 surface to 10 (from the 48 mentioned above). Furthermore, there is no general method of proof for a state space isomorphism provided there. However, we expect results regarding the FJRW invariants, Picard–Fuchs equations and GW invariants to hold in general, and the state space isomorphism we establish here is the first step to such results. This will be the topic of future work.

Computational observations tell us that in some cases a pair (W, G) may have multiple mirrors in BHK mirror symmetry. For example, Shoemaker [34] and Kelly [22] investigated the birationality of multiple mirrors obtained via the so-called weights and groups theorem. But these do not account for all possible multiple mirrors. Our results give a new algebraic observation about possible multiple mirrors in the Landau–Ginzburg mirror symmetry arena.

Finally, an extension to automorphisms of higher order on the geometric side is considered by Goto–Livn’e–Yui in [15]. It will be interesting to see if the mirror symmetry of Borcea–Voisin can be extended to such examples as well on the LG side as on the geometric side.

N:Check this.

2. LANDAU-GINZBURG MODELS

We begin by reviewing the constructions of the LG model for a pair (W, G) .

A polynomial $W = \sum_{i=1}^m c_i \prod_{j=1}^n x_j^{a_{i,j}} \in \mathbb{C}[x_1, \dots, x_n]$ is called *quasihomogeneous* if there exist positive weights $q_j = \frac{w_j}{d}$ for all $j = 1, \dots, n$ such that every nonzero monomial of W has weighted degree one. We require $\gcd(w_1, \dots, w_n) = 1$, and we call $(w_1, \dots, w_n; d)$ a *weight system*.

Additionally, a polynomial $W \in \mathbb{C}[x_1, \dots, x_n]$ is called *nondegenerate* if it has an isolated singularity at the origin and the weights q_i are uniquely determined. A nondegenerate quasihomogeneous polynomial is called *invertible* if the number of variables is equal to the number of monomials in the polynomial.

The Landau-Ginzburg model also requires a choice of symmetry group. For an admissible polynomial W , the *maximal diagonal symmetry group* G_W^{max} is the group of elements of the form $g = (g_1, \dots, g_n) \in (\mathbb{Q}/\mathbb{Z})^n$ such that

$$W(e^{2\pi i g_1} x_1, e^{2\pi i g_2} x_2, \dots, e^{2\pi i g_n} x_n) = W(x_1, x_2, \dots, x_n).$$

Notice that if q_1, q_2, \dots, q_n are the weights of W , then the *exponential grading operator* $j_W = (q_1, \dots, q_n)$ is an element of G_W^{max} .

An *A-admissible symmetry group* for W is a subgroup $G \subset G_W^{max}$ which contains j_W . We denote $\langle j_W \rangle = J_W$ the subgroup generated by j_W , and thus any subgroup between J_W and G_W^{max} is A-admissible (cf. [23]).

We can also embed G_W^{max} into $\mathrm{GL}_n(\mathbb{C})$ via the diagonal matrices

$$(g_1, \dots, g_n) \mapsto \begin{pmatrix} e^{2\pi i g_1} & & 0 \\ & \ddots & \\ 0 & & e^{2\pi i g_n} \end{pmatrix}$$

and define the group $\mathrm{SL}_W = G_W^{max} \cap \mathrm{SL}_n(\mathbb{C})$. Equivalently, for $g = (g_1, \dots, g_n) \in G_W^{max}$, we can consider g_i simply as rational numbers with $0 \leq g_i < 1$. We define the *age* of g to be

$$\mathrm{age} \, g = \sum_{j=1}^n g_j.$$

Then SL_W is the subgroup of elements $g \in G_W^{max}$ with $\mathrm{age} \, g \in \mathbb{Z}$. A group $G \subseteq G_W^{max}$ is *B-admissible* if $G \subset \mathrm{SL}_W$.

We say that a nondegenerate quasihomogeneous polynomial W satisfies the Calabi-Yau condition if $\mathrm{age} \, j_W = 1$. Notice that this condition ensures that $J_W \subset \mathrm{SL}_W$. The following constructions of the A- and B-models do not require this condition. But we will require this condition beginning with Section 3.

2.1. Construction of A- and B-models. We now briefly review the construction of the A- and B-model state spaces. Each of these yields a graded vector space with a nondegenerate pairing.

Recall that we consider G_W^{max} as a subgroup of $\mathrm{GL}_n(\mathbb{C})$. For $g \in G_W^{max}$, we fix the following notation:

- $\mathrm{Fix}(g) = \{\mathbf{x} \in \mathbb{C}^n \mid g\mathbf{x} = \mathbf{x}\},$
- $W_g = W|_{\mathrm{Fix}(g)},$
- N_g the dimension of $\mathrm{Fix}(g),$
- $I_g = \{i \mid g \cdot x_i = x_i\}$

Note that I_g is the set of indices of those variables fixed by g , and that $N_g = |I_g|$.

The Landau–Ginzburg A-model. The A-model was first constructed by Fan–Jarvis–Ruan in [10] following ideas of Witten, and has since come to be known as FJRW theory. For W a nondegenerate quasihomogeneous polynomial and G an A-admissible group, the theory yields a state space, a moduli space of so-called W -curves, and a virtual cycle from which one can define numerical invariants. In this work, we are primarily concerned with the state space $\mathcal{A}_{W,G}$, which is defined in terms of Lefschetz thimbles.

We define $\mathcal{A}_{W,G}$, as

$$(2) \quad \mathcal{A}_{W,G} = \bigoplus_{g \in G} H^{mid}(\text{Fix}(g), (W)_g^{-1}(\infty))^G, .$$

where $(W)_g^{-1}(\infty)$ is a generic smooth fiber of the restriction of W to $\text{Fix}(g)$.

The middle-dimensional cohomology in Equation 2 is equipped with a natural non-degenerate pairing.

$$H^{mid}(\text{Fix}(g), (W)_g^{-1}(\infty))^G \times H^{mid}(\text{Fix}(g), (W)_g^{-1}(\infty))^G \rightarrow \mathbb{C}.$$

This pairing induces a non-degenerate pairing on $\mathcal{A}_{W,G}$. Because Lefschetz thimbles are more difficult to work with, we give an alternate definition of the A-model state space, which will facilitate computation.

The *Milnor Ring* \mathcal{Q}_W of the polynomial W is defined as

$$\mathcal{Q}_W = \mathbb{C}[x_1, \dots, x_n] / \langle \partial W / \partial x_i \rangle .$$

If W is nondegenerate, then \mathcal{Q}_W has finite dimension as a \mathbb{C} -vector space. We denote $\mu_W := \dim_{\mathbb{C}} \mathcal{Q}_W$.

The Milnor ring \mathcal{Q}_W has a natural residue pairing $\langle \cdot, \cdot \rangle_W$ determined by the equation

$$(3) \quad f \cdot g = \frac{\text{Hess}(W)}{\mu_W} \langle f, g \rangle_W + \text{lower order terms}.$$

The following theorem gives a connection between the middle dimensional cohomology and the Milnor ring (also cf. [5] where the isomorphism is given canonically).

Theorem 1 (Wall [36]). *Let $\omega = dx_1 \wedge \dots \wedge dx_n$, then*

$$H^{mid}(\mathbb{C}^n, (W)^{-1}(\infty)) \cong \frac{\Omega^n}{dW \wedge \Omega^{n-1}} \cong \mathcal{Q}_W \cdot \omega.$$

are isomorphic as G_W^{max} -spaces, and this isomorphism respects the pairing on both.

In this description of the state space G_W^{max} acts on both monomials in \mathcal{Q}_W as well as the volume form ω . The isomorphism in Theorem 1 certainly will hold for the restricted polynomials W_g as well. Therefore, setting $\mathcal{Q}_g = \mathcal{Q}_{W_g}$, we can rewrite the state space canonically as

$$(4) \quad \mathcal{A}_{W,G} = \bigoplus_{g \in G} (\mathcal{Q}_g \omega_g)^G ,$$

where $\omega_g = dx_{i_1} \wedge \dots \wedge dx_{i_s}$ for $i_j \in I_g$.

Notation 1. The state space $\mathcal{A}_{W,G}$ is a \mathbb{C} -vector space. We can construct a basis out of elements of the form $[m; g]$, where $I_g = \{i_1, \dots, i_r\}$ and m is a monomial in $\mathbb{C}[x_{i_1}, \dots, x_{i_r}]$, plus the volume form ω_g . We say that $[m; g]$ is *narrow* if $I_g = \emptyset$ (in which case $m = 1$), and *broad* otherwise.

We can now describe the A-model pairing.

$$(5) \quad \langle [m; g], [n; h] \rangle_A = \begin{cases} \langle m, n \rangle_{W_g} & \text{if } h = g^{-1} \\ 0 & \text{otherwise.} \end{cases}$$

Here the pairing on the right-hand side is the pairing from the Milnor ring. Notice that this pairing is well-defined, since $W_g = W_{g^{-1}}$, and so $(\mathcal{Q}_g \omega_g)^G \cong (\mathcal{Q}_{g^{-1}} \omega_{g^{-1}})^G$. This agrees with the original pairing on the state space described above. We will use η to denote the matrix associated to this pairing in the basis described above.

Remark 1.1. The above implies that $\langle [m; g], [n; h] \rangle \neq 0$ exactly when $g = h^{-1}$ and one of the monomials of $m \cdot n$ is a scalar multiple of $\text{Hess}(W_g)$.

The A-model also has a bigrading, defined as

$$(6) \quad (\deg_+^A([m; g]), \deg_-^A([m; g])) = (\deg m + \text{age } g - \text{age } j_W, N_g - \deg m + \text{age } g - \text{age } j_W).$$

The Landau–Ginzburg B-model. The B-model state space and its pairing are defined similarly to the A-model, with subtle but important differences. The input data is a pair (W, G) with W a nondegenerate quasihomogeneous polynomial and G a B-admissible symmetry group. The state space is defined as

$$\mathcal{B}_{W,G} = \bigoplus_{g \in G} (\mathcal{Q}_g \omega_g)^G.$$

The pairing is defined as in (3) and (5). The bidegree of $[m; g]$ in the B-model is

$$(\deg_+^B([m; g]), \deg_-^B([m; g])) = (\deg m + \text{age } g - \text{age } j_W, \deg m + \text{age } g^{-1} - \text{age } j_W).$$

Remark 1.2. As vector spaces (with pairing) we have an isomorphism $\mathcal{A}_{W,G} \cong \mathcal{B}_{W,G}$ (cf. Equation 4). This fact will be useful later. However, this isomorphism does not preserve the bigrading. Furthermore, both state spaces can be given the structure of \mathbb{C} -algebras (cf. Section 4.2), and, in general are not isomorphic as \mathbb{C} -algebras.

2.2. Berglund–Hubsch–Krawitz Mirror Symmetry. There is a more interesting relationship between A- and B-models which was first discovered by Berglund–Hubsch and later refined by Krawitz. This form of mirror symmetry relates the A-model of the pair (W, G) to the B-model of another pair (W^T, G^T) , which we will now define. We will hereafter refer to this formulation of mirror symmetry as BHK mirror symmetry.

Suppose $W = \sum_{i=1}^n \prod_{j=1}^n x_j^{a_{i,j}}$ is an invertible polynomial. We can suppress any coefficients via a change of variables. The *mirror potential* W^T is

$$W^T = \sum_{i=1}^n \prod_{j=1}^n x_j^{a_{j,i}}.$$

In other words, if A_W is the exponent matrix for W , then W^T is the polynomial corresponding to the exponent matrix A_W^T . One can check that W^T is also an invertible polynomial (cf. [7]).

If G is an admissible symmetry group for a polynomial W , then the *mirror group* G^T is defined by

$$G^T = \{ g \in G_{W^T}^{max} \mid g A_W h^T \in \mathbb{Z} \text{ for all } h \in G \}.$$

Remark 1.3. This is not the original definition of G^T . The first definition was given by Berglund–Henningson in [3], and then independently by Krawitz in [23]. In [1], it is shown that their definitions agree, and it is an exercise to check that the definition given here is equivalent to the one given by Krawitz.

In [1, Proposition 3] we learn that $(G^T)^T = G$ and that $J_W \subset G$ if and only if $G^T \subset \mathrm{SL}_W$. In other words, this duality exchanges A-admissible groups for B-admissible ones.

On the level of state spaces, we get an isomorphism between A- and B-models given by the theorem of Krawitz.

Theorem 2 (Krawitz [23]). *Let W be a non-degenerate invertible potential and G an A-admissible group of diagonal symmetries of W . There is a bigraded isomorphism of vector spaces*

$$\mathcal{A}_{W,G} \cong \mathcal{B}_{W^T,G^T},$$

where $\mathcal{A}_{W,G}$ is the FJRW A-model of (W, G) and \mathcal{B}_{W^T,G^T} is the orbifold B-model of (W^T, G^T) .

One may notice that Theorem 6 is similar to the theorem of Krawitz, though they are significantly different in the choice of symmetry group. However, the theorem of Krawitz will be used in the proof of Theorem 6.

Remark 2.1. The proof of Theorem 2 in [23] does not hold in case W contains a summand of a so-called chain polynomial with weight $\frac{1}{2}$. However, such a polynomial will not appear in the current work.

2.3. Mirror Pairs inspired by Borcea–Voisin mirror pairs. Following ideas inspired by the LG–CY correspondence, we can define a Landau–Ginzburg model that will correspond to the Calabi–Yau threefolds constructed by Borcea and Voisin.

The LG-model corresponding to Borcea–Voisin threefolds can be described in many cases as follows. Consider two invertible polynomials of the form

$$(7) \quad W_1 = x_0^2 + f_1(x_1, \dots, x_n) \quad \text{and} \quad W_2 = y_0^2 + f_2(y_1, \dots, y_m),$$

with weights $(u_0, \dots, u_n; d_1)$ and $(v_0, \dots, v_m; d_2)$. Notice that because of the special form of the polynomials, $d_1 = 2u_0$ and $d_2 = 2v_0$. In order to ease notation slightly, in what follows we will denote $J_i = J_{W_i}$ generated by the exponential grading operator $j_i = j_{W_i}$ for $i = 1, 2$.

In the Borcea–Voisin construction we take the product of hypersurfaces defined by W_1 and W_2 , respectively, then quotient by a certain involution. While on the LG-side we consider the “twisted intersection” of these two polynomials. That is, let $W = W_1 + W_2$ and let G be any group satisfying $J_1 \times J_2 \leq G \leq \mathrm{SL}_{W_1} \times \mathrm{SL}_{W_2}$. Denote the involution $\sigma = (1/2, \mathbf{0}_x, 1/2, \mathbf{0}_y)$, and define the group $\Sigma G = G \cup \sigma G$ —that is, the smallest group that contains G and σ . This involution is the LG analogue of the involution in the Borcea–Voisin model. We call the LG model $(W, \Sigma G)$ a Borcea–Voisin Landau–Ginzburg (BVLG) model. We will justify this construction and the use of this terminology in Section 5.1.

Following BV, we call the pair $(W, \Sigma G)$ and $(W^T, \Sigma G^T)$ a BV mirror pair. As already mentioned, BV mirror symmetry does not actually match the picture given by BHK mirror symmetry, as one might expect. BHK mirror symmetry defines a group $(\Sigma G)^T$, which does not contain $J_1 \times J_2$, but ΣG^T does (e.g. since $\Sigma G \not\leq \mathrm{SL}_{W_1} \times \mathrm{SL}_{W_2}$). So the BHK mirror of a BVLG model is no longer a BVLG model.

Remark 2.2. In fact what we call a BVLG model is slightly more general.

Am: Than what? The LG model inspired by BV construction?

If we take for a group G a group that can be written as a product of subgroups $G = G_1 \times G_2$ with $G_1 \subset \mathrm{SL}_{W_1}$ and $G_2 \subset \mathrm{SL}_{W_2}$, then we recover the construction of Borcea–Voisin. If, on the other hand, one considers a group G that is not such a product, the corresponding geometry is no longer a variety of Borcea–Voisin type, but rather an orbifold, whose crepant resolution is the

Nakamura–Hilbert scheme of $\widetilde{\Sigma G}$ -regular orbits of $X_1 \times X_2$ (cf. [25]), as was also mentioned in the introduction.

3. THE TWISTING MAP

In this section we make some preparations needed for the proofs of Theorems 6 and 9. We will use ideas from the geometric model, namely the twist map from [4]. There is a similar notion on the LG side, which we will call the *LG twist map*.

Given a pair $(W, \Sigma G)$ as described above, namely $W = W_1 + W_2$ with W_1 and W_2 of the form (7), and G with $J_1 \times J_2 \subset G \subset \mathrm{SL}_{W_1} \times \mathrm{SL}_{W_2}$, we define another LG model (W_{tw}, G_{tw}) corresponding to the geometry of the image of the twist map (cf. Section 5.4 and [2]).

We first define $W_{tw} = f_1 - f_2$. If $\delta = \gcd(u_0, v_0)$, then this defines an invertible polynomial with weight system $(\frac{v_0 u_1}{\delta}, \frac{v_0 u_2}{\delta}, \dots, \frac{v_0 u_n}{\delta}, \frac{u_0 v_1}{\delta}, \dots, \frac{u_0 v_m}{\delta}; \mathrm{lcm}(d_1, d_2))$. We use the negative sign as a convention for clarity, in order to match the geometry. Since f_2 is invertible, this sign convention can be scaled away.

To see that W_{tw} is quasihomogeneous, note that $\mathrm{lcm}(d_1, d_2) = 2 \mathrm{lcm}(u_0, v_0)$, and the same for the gcd's. Further, $\mathrm{lcm}(d_1, d_2) \cdot \gcd(d_1, d_2) = d_1 d_2$, so $\delta \cdot \mathrm{lcm}(d_1, d_2) = \frac{d_1 d_2}{2}$ and $v_0 = \frac{d_1}{2}$. Hence

$$q_{x_i} = \frac{u_i}{d_1} = \frac{v_0 u_i}{\delta} \cdot \frac{1}{\mathrm{lcm}(d_1, d_2)}.$$

Furthermore, the gcd of all of the weights is equal to 1.

To define the group G_{tw} , notice that for any $g' \in G_{W_{tw}}^{\max}$, we can write $g' = (\alpha, \beta)$, where α acts on the variables x_1, \dots, x_n , and β on the variables y_1, \dots, y_m . There is an injective group homomorphism $\phi : G_{W_{tw}}^{\max} \hookrightarrow G_W^{\max}$, given by

$$\phi : (\alpha, \beta) \mapsto (0, \alpha, 0, \beta).$$

We define G_{tw} to be $\phi^{-1}(\Sigma G)$.

Remark 2.3. If $\delta = 1$, then this is the group considered by Artebani–Boisierre–Sarti in [2]. However, we will see that this condition is not necessary.

Now we consider the relationship between $\mathcal{A}_{W_{tw}, G_{tw}}$ and $\mathcal{A}_{W, \Sigma G}$. Geometrically, Borcea (and Artebani–Boisierre–Sarti) has shown that when $\delta = 1$ the corresponding geometric models are birational. We give a similar statement for the LG-models.

Theorem 3. *There is an isomorphism $tw_A : \mathcal{A}_{W, \Sigma G} \rightarrow \mathcal{A}_{W_{tw}, G_{tw}}$ which preserves bidegrees and pairings.*

This theorem is proved by giving a bijection on a basis of both spaces, and showing it preserves bidegrees and pairings. The map tw_A is really just a rescaled projection map. We first prove the following preparatory lemma.

Lemma 4. *If $(\mathcal{Q}_g)^{\Sigma G}$ is a non-empty sector of $\mathcal{A}_{W, \Sigma G}$, then g either fixes both of x_0 and y_0 or neither.*

Proof. Suppose $g = (1/2, \alpha, 0, \beta)$, for example. Then the volume form ω_g includes dy_0 , but not dx_0 . Then σ acts on any element of the Milnor ring \mathcal{Q}_{W_g} with weight $1/2$. Therefore there are no invariants. \square

Proof of Theorem 3. Consider $[p; g] \in \mathcal{A}_{W, \Sigma G}$. By Lemma 4, we can write $g = (\epsilon/2, \alpha, \epsilon/2, \beta)$, $\epsilon \in \{0, 1\}$, and $p = dx_0^{1-\epsilon} dy_0^{1-\epsilon} \prod x_j^{a_j} dx_j \prod y_j^{b_j} dy_j$. Here the products are taken over I_α and I_β , resp. (we can regard α and β as symmetries of f_1 and f_2 , resp.). We define the map via

$$tw_A([p; g]) = [p'; g'],$$

where

$$p' = 2^{(1-\epsilon)} \prod_{j \in I_\alpha} x_j^{a_j} dx_j \prod_{j \in I_\beta} y_j^{b_j} dy_j,$$

and $g' = (\alpha, \beta)$. We check that this map is well defined, injective and surjective, and that it preserves the bigradings and the pairings.

To see that the map is well defined, we check that $g' \in G_{tw}$, and that p' is fixed by G_{tw} . For the first, notice that either $g = (0, \alpha, 0, \beta)$ or $\sigma g = (0, \alpha, 0, \beta)$. Since $(0, \alpha, 0, \beta) \in \Sigma G \cap \phi(G_{W_{tw}}^{\max})$ we see that $g' \in G_{tw}$. For the second, note that p' is fixed by G_{tw} because ϕ restricts to an injection $\phi : G_{tw} \hookrightarrow \Sigma G$, and g' acts on p' exactly the same as $\phi(g') \in \Sigma G$ acts on p .

As for the surjectivity, in the notation already in use, suppose $[p'; g'] \in \mathcal{A}_{W_{tw}, G_{tw}}$. Let $p_x = \prod x_j^{a_j} dx_j$ and $p_y = \prod y_j^{b_j} dy_j$ (so, $p' = 2^{1-\epsilon} p_x p_y$). Again, the products are taken only over the respective sets of fixed variables. Let j_1, j_2 and j_{tw} be the exponential grading operators for W_1, W_2 , and W_{tw} , resp. When there is no danger of confusion, we will use the same notation for the elements in G_{tw} and $\phi(G_{tw})$. For example, if $j_{tw} = (j_x, j_y)$, then we may also write $j_{tw} = (0, j_x, 0, j_y)$. Notice that j_x and j_y are the exponential grading operators for f_1 and f_2 , resp. The context should make clear which group is meant. Notice we also have $j_1 = (1/2, j_x, 0, 0)$, and $j_2 = (0, 0, 1/2, j_y)$.

In the notation just described, $2j_{W_1} = (0, 2j_x, 0, 0)$, and so $(2j_x, 0) \in G_{tw}$. Since p' is fixed by G_{tw} , it is fixed by $(2j_x, 0)$, which acts trivially on the y variables, and so $(2j_x, 0)$ fixes p_x . Hence j_x either fixes or acts with weight $1/2$ on p_x . The same for j_y on p_y . And since $j_{W_{tw}}$ fixes p' , the action is the same for both.

If j_x fixes p' , then the preimage of $[p'; g']$ is $[p; g]$ with $g = (1/2, \alpha, 1/2, \beta)$ and $p = \prod x_j^{a_j} dx_j \prod y_j^{b_j} dy_j = p_x p_y$. Notice that the action of any $(\gamma_x, \gamma_y) \in G_{tw}$ on p' is the same as the action of $(\alpha_0, \gamma_x, \beta_0, \gamma_y) \in \Sigma G$ on p for any α_0 and β_0 . Thus, p is fixed by ΣG and $[p; g]$ is indeed an element of $\mathcal{A}_{W, \Sigma G}$.

If on the other hand, j_x acts with weight $1/2$ on p_x , then the preimage of $[p'; g']$ is $[p; g]$ with $g = (0, \alpha, 0, \beta)$ and $p = dx_0 dy_0 \prod x_j^{a_j} dx_j \prod y_j^{b_j} dy_j = dx_0 dy_0 p_x p_y$. Indeed, it is clear that $g \in \Sigma G$, we need only check that p is fixed by ΣG . Because p' is fixed by G_{tw} , we see that p is fixed by all elements of ΣG which have the form $(0, h_x, 0, h_y)$ or $(1/2, h_x, 1/2, h_y)$. We need only further consider the case when $h = (1/2, \eta_x, 0, \eta_y) \in \Sigma G$. Then we have $j_{W_1} + h = (0, j_x + \eta_x, 0, \eta_y) \in \Sigma G$, and therefore $j_{W_1} + h = (j_x + \eta_x, \eta_y) \in G_{tw}$ fixes p' . The action of j_x on p' has weight $1/2$, by assumption, which means the action of (η_x, η_y) on p' must also have weight $1/2$. Because h acting on dx_0 contributes another $1/2$, we see that h fixes p .

For the injectivity of the twist map, we need to show that if $[p; g] \in \mathcal{A}_{W, \Sigma G}$ with $g = (0, \alpha, 0, \beta)$, then $\left[\frac{p}{dx_0 dy_0}; \sigma g \right] \notin \mathcal{A}_{W, \Sigma G}$, and vice-versa. This follows by similar reasoning. For example, if j_{W_1} fixes p , then it acts on $\frac{p}{dx_0 dy_0}$ with weight $1/2$. Similarly, if j_{W_1} fixes $\frac{p}{dx_0 dy_0}$, then it acts on p with weight $1/2$. Thus both cannot belong to $\mathcal{A}_{W, \Sigma G}$.

Now we show that this map preserves degrees. We will consider the bigradings of the elements $[p; g]$ with $g = (\epsilon/2, \alpha, \epsilon/2, \beta)$ and $[p'; g']$ with $g = (\alpha, \beta)$. And indeed,

$$\begin{aligned} \deg_+^A([p; g]) &= \deg p + \text{age } g - 2 \\ &= \left(\frac{1-\epsilon}{2} + \frac{1-\epsilon}{2}\right) + \deg p' + \frac{\epsilon}{2} + \frac{\epsilon}{2} + \text{age } g' - 2 \\ &= \deg p' + \text{age } g' - 1 \\ &= \deg_+^A([p'; g']) \end{aligned}$$

and since for $g = (\epsilon/2, \alpha, \epsilon/2, \beta)$, we have $N_g = 2(1 - \epsilon) + N_{g'}$, we also have

$$\begin{aligned} \deg_-^A([p; g]) &= N_g - \deg p + \text{age } g - 2 \\ &= (2(1 - \epsilon) + N_{g'}) - \left(\frac{1-\epsilon}{2} + \frac{1-\epsilon}{2} + \deg p'\right) + \left(\frac{\epsilon}{2} + \frac{\epsilon}{2} + \text{age } g'\right) - 2 \\ &= N_{g'} - \deg p' + \text{age } g' - 1 \\ &= \deg_-^A([p'; g']) \end{aligned}$$

Finally, we show that tw_A preserves the pairing. First we will show that $\langle [p; g], [q; h] \rangle = 0$ if and only if $\langle [p'; g'], [q'; h'] \rangle = 0$. If $g = h^{-1}$ and mn is a scalar multiple of $\text{Hess}(W)$ (See Remark 1.1), then $g' = (h')^{-1}$ and $m'n'$ is a scalar multiple of $\text{Hess}(W_{tw})$. For the converse, suppose that $g' = (h')^{-1}$ and $m'n'$ is a scalar multiple of $\text{Hess}(W_{tw})$.

It is clear that mn is a scalar multiple of $\text{Hess}(W)$. We want to show that $g = h^{-1}$. This will be true whenever the action of j_x —which we know to be either trivial or with weight $1/2$ —is the same on m' and n' . The action of j_x on the Hessian is equivalent to its total weighted degree: $\hat{c}_{f_1} = \sum_{i=1}^n 1 - 2q_{x_i} \equiv 0 \pmod{\mathbb{Z}}$ (since W_1 satisfies the Calabi-Yau condition). Thus, the action of j_x on m' and n' must be the same.

Thus in the case where the pairings in $\mathcal{A}_{W, \Sigma G}$ and $\mathcal{A}_{W_{tw}, G_{tw}}$ are nonzero we have either $g = h^{-1} = (0, \alpha, 0, \beta)$ or $g = h^{-1} = (1/2, \alpha, 1/2, \beta)$. If $g = (0, \alpha, 0, \beta)$ and $g' = (\alpha, \beta)$ then

$$p \cdot q = \frac{\langle p, q \rangle_{W_g}}{\mu_g} \text{Hess}_g = \frac{\langle p, q \rangle_{W_g}}{\mu_{g'}} 4 \cdot \text{Hess}_{g'} = \frac{\langle 2p, 2q \rangle_{W_{g'}}}{\mu_{g'}} \text{Hess}_{g'}.$$

So, $\langle [p; g], [q; g^{-1}] \rangle_{W, \Sigma G} = \langle tw_A([p; g]), tw_A([q; g^{-1}]) \rangle_{W_{tw}, G_{tw}}$

For $g = (1/2, \alpha, 1/2, \beta)$, $W_g = W_{g'}$, and the equivalence of the pairing is automatic. \square

Corollary 5. *There is an isomorphism $tw_B : \mathcal{B}_{W, \Sigma G} \rightarrow \mathcal{B}_{W_{tw}, G_{tw}}$ which preserves the pairing.*

Proof. This isomorphism tw_B is defined as in Theorem 3. Recall that $\mathcal{B}_{W, G} \cong \mathcal{A}_{W, G}$ as vector spaces (with pairing) (see Remark 1.2). The only thing that remains to check is the bi-degree.

For $g = (\epsilon/2, \alpha, \epsilon/2, \beta)$,

$$\begin{aligned} \deg_+^B([p; g]) &= \deg p + \text{age } g - 2 \\ &= \left(\frac{1-\epsilon}{2} + \frac{1-\epsilon}{2}\right) + \deg p' + \frac{\epsilon}{2} + \frac{\epsilon}{2} + \text{age } g' - 2 \\ &= \deg p' + \text{age } g' - 1 \\ &= \deg_+^B([p'; g']) \end{aligned}$$

N: We switched somehow from p and q to m and n , and I think we mean just the polynomials and not the volume forms? Is that right?

and since $g^{-1} = (\epsilon/2, \alpha^{-1}, \epsilon/2, \beta^{-1})$, $(g')^{-1} = (\alpha^{-1}, \beta^{-1})$ we have $\text{age } g = \epsilon + \text{age}(g')^{-1}$. Thus

$$\begin{aligned} \deg_-^B(\lceil p; g \rceil) &= \deg p + \text{age } g^{-1} - 2 \\ &= \left(\frac{1-\epsilon}{2} + \frac{1-\epsilon}{2} + \deg p'\right) + (\epsilon + \text{age}(g')^{-1}) - 2 \\ &= \deg p' + \text{age}(g')^{-1} - 1 \\ &= \deg_-^B(\lceil p'; g' \rceil) \end{aligned}$$

□

Now we have two isomorphisms $tw_A : \mathcal{A}_{W, \Sigma G} \rightarrow \mathcal{A}_{W_{tw}, G_{tw}}$ and $tw_B : \mathcal{B}_{W, \Sigma G} \rightarrow \mathcal{B}_{W_{tw}, G_{tw}}$ for A-models and B-models, resp. We will use these isomorphisms to prove Theorems 6 and 9.

4. A- AND B-MODEL EQUIVALENCE

We now have all of the necessary tools to prove the mirror symmetry statements of the main theorems.

In some sense these theorems are LG analogues of results in both the work of Artebani–Boisier–Sarti in [2] and of Chiodo–Kalashnikov–Veniani in [6] simultaneously. Other than the examples computed by the last author in [32] this is the first time we are aware that this particular a LG model has been considered in the literature.

4.1. State Space isomorphism. In this section we will prove Theorem 6, which we restate here:

Theorem 6. *There is an isomorphism (the mirror map) $\mathcal{A}_{W, \Sigma G} \cong \mathcal{B}_{W^T, \Sigma G^T}$, which preserves the bi-degree.*

Proof. Consider the diagram

$$\begin{array}{ccc} \mathcal{B}_{(W^T)_{tw}, (G^T)_{tw}} & \xleftarrow{tw_B} & \mathcal{B}_{W^T, \Sigma G^T} \\ \updownarrow & & \updownarrow \\ \mathcal{A}_{W_{tw}, G_{tw}} & \xleftarrow{tw_A} & \mathcal{A}_{W, \Sigma G} \end{array}$$

The two horizontal arrows are the content of Theorem 3 and Corollary 5. The dashed arrow is the desired isomorphism of Theorem 6. We will establish this map by providing the leftmost vertical arrow in the diagram. This will follow from [23] of Krawitz (see Theorem 2). The only difference is the group of symmetries in the upper left corner is $(G^T)_{tw}$, instead of $(G^T)_{tw}^T$. Therefore all that remains is to show $(G^T)_{tw} = (G_{tw})^T$, which we do in the following lemma. □

Lemma 7. $(G^T)_{tw} = (G_{tw})^T$.

Proof. We first consider the inclusion $(G^T)_{tw} \subset (G_{tw})^T$. Recall $(G^T)_{tw} = (\phi^T)^{-1}(\Sigma G^T)$. Consider $g = (\alpha, \beta) \in (G^T)_{tw}$ and $g_\epsilon = ((1/2)^\epsilon, \alpha, (1/2)^\epsilon, \beta)$ for $\epsilon \in \{0, 1\}$. By definition, $g_0 \in G^T$ or $g_1 \in G^T$.

We want to show that $g \in (G_{tw})^T$, which means we need to show that for each $h = (\gamma_x, \gamma_y) \in G_{tw}$, $gA_{W_{tw}}h \in \mathbb{Z}$. Again, $h \in G_{tw}$ means that $h_{\epsilon'} = ((1/2)^{\epsilon'}, \gamma_x, (1/2)^{\epsilon'}, \gamma_y) \in G$ for either $\epsilon' = 0$ or $\epsilon' = 1$.

Let A_1 and A_2 be the exponent matrices for f_1 and f_2 , resp. Then for $\epsilon = 0$ or 1 , we know that

$$g_\epsilon A_W h_{\epsilon'} = 4(1/2)^{\epsilon+\epsilon'} + \alpha A_1 \gamma_x + \beta A_2 \gamma_y \in \mathbb{Z}.$$

Since the first summand is an integer, we have $g A_{W_{tw}} h = \alpha A_1 \gamma_x + \beta A_2 \gamma_y \in \mathbb{Z}$.

Now for the reverse inclusion $(G_{tw})^T \subset (G^T)_{tw}$. Suppose $g = (\alpha, \beta) \in (G_{tw})^T$, and again let $g_\epsilon = ((1/2)^\epsilon, \alpha, (1/2)^\epsilon, \beta)$. We need to show that either $g_0 \in G^T$ or $g_1 \in G^T$. In other words we need to show that $g_0 A_W h \equiv 0 \pmod{\mathbb{Z}}$ for every $h \in G$, or that $g_1 A_W h \equiv 0 \pmod{\mathbb{Z}}$ for every $h \in G$. For h of the form $((1/2)^{\epsilon'}, \gamma_x, (1/2)^{\epsilon'}, \gamma_y)$, the previous calculation shows that this indeed the case for both g_0 and g_1 . The difficulty comes with h is of the form $((1/2)^{\epsilon'}, \gamma_x, (1/2)^{1-\epsilon'}, \gamma_y)$. We consider two cases.

First suppose $\text{age } \alpha \equiv 0 \pmod{\mathbb{Z}}$. Then since $(G^T)_{tw} \subset (G_{tw})^T \subset \text{SL}_{W_{tw}}$, $\text{age } \beta \equiv 0 \pmod{\mathbb{Z}}$. We will show that, in fact, $g_0 \in G^T$. Notice that $\alpha A_1 j_x = \text{age } \alpha \equiv 0 \pmod{\mathbb{Z}}$. Suppose first that, $h = ((1/2), \gamma_x, 0, \gamma_y) \in G$. Then $j_1 + h \in \phi(G_{tw}) \cap G$, and so

$$g_0 A_W (j_1 + h) = \alpha A_1 j_x + \alpha A_1 \gamma_x + \beta A_2 \gamma_y = g A_{W_{tw}} \phi^{-1}(j_1 + h) \equiv 0 \pmod{\mathbb{Z}},$$

since $\phi^{-1}(j_1 + h) \in G_{tw}$ and, by assumption, $g \in (G_{tw})^T$. Since also

$$g_0 A_W (j_1 + h) = \text{age } \alpha + g_0 A_W h,$$

we have $g_0 A_W h \equiv 0 \pmod{\mathbb{Z}}$. If $h = (0, \gamma_x, 1/2, \gamma_y)$, the proof is the same, only using j_{W_2} instead.

On the other hand, suppose $\text{age } \alpha \equiv 1/2 \pmod{\mathbb{Z}}$. Then $\text{age } \beta \equiv 1/2 \pmod{\mathbb{Z}}$. We will show that, in fact, $g_1 \in G^T$. The proof is similar, but this time, we notice that $\alpha A_1 j_x$ is no longer an integer, but a half-integer. Again we first suppose that $h = ((1/2), \gamma_x, 0, \gamma_y)$. Then $j_1 + h \in \phi(G_{tw})$, and so

$$g_1 A_W (j_1 + h) = \alpha A_1 j_x + \alpha A_1 \gamma_x + \beta A_2 \gamma_y = g A_{W_{tw}} \phi^{-1}(j_1 + h) \equiv 0 \pmod{\mathbb{Z}}$$

since $\phi^{-1}(j_1 + h) \in G_{tw}$ and, by assumption, $g \in (G_{tw})^T$. Since also

$$g_1 A_W (j_1 + h) = \frac{1}{2} + \text{age } \alpha + g_1 A_W h,$$

we have $g_1 A_W h \equiv 0 \pmod{\mathbb{Z}}$. If $h = (0, \gamma_x, 1/2, \gamma_y)$, the proof is the same, only using j_{W_2} instead. \square

This concludes the proof of Theorem 6.

4.2. Landau–Ginzburg Algebra Isomorphism. We have shown that there is an isomorphism of bigraded vector spaces. But as we have remarked, each state space also has the structure of a Frobenius algebra over \mathbb{C} . We will now consider the algebra structure.

4.2.1. A-model Frobenius algebra. The product on the A-model is defined in [10] via the structure constants

$$a \star_{\mathcal{A}} b := \sum_{\xi, \xi'} \langle a, b, \xi \rangle_{0,3} \eta^{\xi, \xi'} \xi'$$

where the sum runs over a basis of $\mathcal{A}_{W,G}$, and $\eta^{\xi, \xi'}$ are the corresponding entries from the matrix inverse to the pairing matrix η .

The structure constants $\langle a, b, \xi \rangle_{0,3}$ are defined in FJRW theory via certain integrals over the moduli space of curves $\overline{\mathcal{M}}_{0,3}$, and there are corresponding numbers for higher genus and more marked points as well (see [10]). Explicit computations of some of these constants in certain cases are given in [10, 24, 8, 23] and most recently in [17], and methods for computations are given in [16, 11].

We will not need these numbers explicitly in this work. We rely instead on the B-model.

4.2.2. *B-model Frobenius algebra.* On the B-side, the product was defined by Intriligator, Vafa, and Kaufmann [18, 21, 19, 20], and written explicitly in the form we will use now by Krawitz [23]. It is defined on a basis, and then extended multilinearly to the entire state space. In what follows, we write

$$W_{g \cap h} := W|_{\text{Fix } g \cap \text{Fix } h}$$

and $\mu_{g \cap h}$ for the dimension of the corresponding Milnor algebra.

On the basis we have defined, the definition of the product is

$$[1; g] \star_{\mathcal{B}} [1; h] = \gamma_{g,h} [1; gh], \text{ where}$$

$$\gamma_{g,h} \frac{\text{Hess}(W_{g \cap h})}{\mu_{g \cap h}} = \begin{cases} \frac{\text{Hess}(W_{gh})}{\mu_{gh}} & \text{if } I_g \cup I_h \cup I_{gh} = \{1, \dots, N\} \\ 0 & \text{otherwise} \end{cases}$$

To prove our second theorem (Theorem 9), we will use the following lemma.

Lemma 8.

$$\gamma_{g,h} = \begin{cases} 4\gamma_{g,\sigma h} & \text{if } g = (1/2, \alpha, 1/2, \beta), \ h = (1/2, \delta, 1/2, \nu) \\ 1/4\gamma_{g,\sigma h} & \text{if } g = (1/2, \alpha, 1/2, \beta), \ h = (0, \delta, 0, \nu) \\ \gamma_{g,\sigma h} & \text{otherwise} \end{cases}$$

The straightforward proof is left to the reader.

In [12], it was shown that for LG models (W, G) satisfying Property (*), we have an isomorphism of Frobenius algebras

$$\mathcal{A}_{W,G} \cong \mathcal{B}_{W^T, G^T},$$

given by a rescaling of the Krawitz map. The conditions are satisfied for a majority of cases—in particular for all polynomials with no chains (see [12, Remark 1.1.1]).

4.2.3. *Proof of Theorem 9.* We will now exploit this isomorphism to prove Theorem 9.

We have the following diagram of Landau–Ginzburg models. The lower vertical (non-dashed) arrows come from Krawitz’s mirror map. As mentioned above, under the conditions stated in Theorem 9, there exist rescalings of these maps which are isomorphisms of Frobenius algebras. The horizontal arrows (and vertical double line) are the vector space isomorphisms of the previous section.

$$\begin{array}{ccc} \mathcal{B}_{(W^T)_{tw}, (G^T)_{tw}} & \xleftarrow{tw_B} & \mathcal{B}_{W^T, \Sigma G^T} \\ \parallel & & \updownarrow \\ \mathcal{B}_{(W_{tw})^T, (G_{tw})^T} & & \mathcal{B}_{W^T, (\Sigma G)^T} \\ \downarrow \kappa_2 & & \downarrow \kappa_1 \\ \mathcal{A}_{W_{tw}, G_{tw}} & \xleftarrow{tw_A} & \mathcal{A}_{W, \Sigma G} \end{array}$$

Using this description as a guide, we will define a map $\varphi : \mathcal{B}_{W^T, (\Sigma G)^T} \rightarrow \mathcal{B}_{W^T, \Sigma G^T}$, which is a rescaling of the composition $tw_B^{-1} \circ \kappa_2^{-1} \circ tw_A \circ \kappa_1$ of the solid arrows in the diagram, each of which is a degree-preserving isomorphism of graded vector spaces. We then show that φ preserves products and pairings. The isomorphism of Theorem 9 is the composition $\varphi \circ \kappa_1^{-1}$.

Theorem 9. *If W and ΣG satisfy Property (*), then $\mathcal{A}_{W,\Sigma G} \cong \mathcal{B}_{W^T,\Sigma(G^T)}$ as Frobenius algebras.*

Remark 9.1. It is also worth noting that in the diagram, there are two *different* B-models. This is related to the so-called multiple mirror phenomenon as mentioned in the Introduction. The B-model in the middle of the right-hand side is BHK mirror for $(W, \Sigma G)$. However, as we have seen, $\Sigma G^T \neq (\Sigma G)^T$ and so the corresponding geometry does not fit the Borcea–Voisin construction. Because of the LG/CY correspondence, we expect a B-model whose corresponding geometry is Borcea–Voisin type, which is what the upper right-hand B-model is. Also, since we have two B models, both mirror to the same A-model, it should be expected that the two B model state spaces are isomorphic.

For the proof of Theorem 9, we introduce the notation $\mathcal{B}_1 = \mathcal{B}_{W^T,(\Sigma G)^T}$ (the BHK mirror) and $\mathcal{B}_2 = \mathcal{B}_{W^T,\Sigma G^T}$ (the BV mirror).

It is straightforward to verify that for any $[m; g] \in \mathcal{B}_1$, where $g = (\epsilon/2, \alpha, \epsilon/2, \beta)$, $g' = (\alpha, \beta)$ and m' is the same as m except for the absence of the $dx_0 \wedge dy_0$ term (which may or may not be present in m).

$$\kappa_2^{-1} \circ tw_A \circ \kappa_1([m; g]) = k_{m,g}[m'; g'],$$

for some constant $k_{m,g}$. Thus, composing finally with tw_B^{-1} , gives the following map:

$$(8) \quad tw_B^{-1} \circ \kappa_2^{-1} \circ tw_A \circ \kappa_1([m; g]) = \begin{cases} k_{m,g}[m'; (1/2, \alpha, 1/2, \beta)] & \text{if } j_{x,B} \text{ acts with weight 0 on } m', \\ k_{m,g}[m'dx_0 \wedge dy_0; (0, \alpha, 0, \beta)] & \text{otherwise,} \end{cases}$$

where $j_{x,B}$ is the exponential grading operator associated with f_1^T .

Note that for $[m; g] \in \mathcal{B}_1$, g is always an element of ΣG^T , since $(\Sigma G)^T \subset \Sigma G^T$. But m may not be fixed by the action of all elements in the larger group. We shall use \tilde{m} to denote the monomial plus volume form which differs from m by the $dx_0 \wedge dy_0$ volume form (if m contains the term $dx_0 \wedge dy_0$, then \tilde{m} does not, and vice-versa). Lemma 4 shows that exactly one of $[m; g]$ or $[\tilde{m}; \sigma g]$ is in \mathcal{B}_2 . The following lemma further clarifies the map in Equation 8.

Lemma 10. *Let $[m; g] \in \mathcal{B}_1$, then*

$$tw_B^{-1} \circ \kappa_2^{-1} \circ tw_A \circ \kappa_1([m; g]) = \begin{cases} k_{m,g}[m; g] & \text{if } j_{W_1^T} \text{ fixes } m, \\ k_{m,g}[\tilde{m}; \sigma g] & \text{otherwise.} \end{cases}$$

Further, $[m; g] \in \mathcal{B}_2$ exactly when $j_{W_1^T}$ fixes m .

Proof. Suppose that $j_{W_1^T}$ fixes m . If, also, $j_{x,B}$ fixes m , then we know that $m = m_x m_y$ and $g = (1/2, \alpha, 1/2, \beta)$ for some α and β . In this case $tw_B([m'; g']) = [m; g]$, as desired. If, instead $j_{x,B}$ acts with weight 1/2 on m_x , then it must be that $m = m_x m_y dx_0 \wedge dy_0$ and $g = (0, \alpha, 0, \beta)$. Again, this yields $tw_B([m'; g']) = [m; g]$.

If $j_{W_1^T}$ doesn't fix m , a similar analysis of the cases where $j_{x,B}$ fixes/doesn't fix m shows that $tw_B([m'; g']) = [\tilde{m}; \sigma g]$ in both cases. \square

If φ is a rescaling of the map above then it will be a degree preserving bijection between B-models. It remains to define a rescaling and demonstrate that it will preserve products and pairings, which we do as follows:

$$\varphi([m; g]) = \begin{cases} [m; g] & \text{if } [m; g] \in \mathcal{B}_{W^T, \Sigma G^T} \\ \frac{1}{2}[\tilde{m}; \sigma g] & \text{if } [m; g] \notin \mathcal{B}_{W^T, \Sigma G^T}, g = (0, \alpha, 0, \beta) \\ 2[\tilde{m}; \sigma g] & \text{if } [m; g] \notin \mathcal{B}_{W^T, \Sigma G^T}, g = (1/2, \alpha, 1/2, \beta). \end{cases}$$

Notice that if $[m; g]$ and $[n; h]$ are elements in \mathcal{B}_1 with $I_g \cup I_h \cup I_{gh} = \{x_0, \dots, x_m, y_0, \dots, y_n\}$ and both are elements of \mathcal{B}_2 , then $[mn; gh]$ must also be an element of \mathcal{B}_2 . Also, if neither is an element of \mathcal{B}_2 , then $[mn; gh]$ is an element of \mathcal{B}_2 (Since $[\tilde{m}; \sigma g]$ and $[\tilde{n}; \sigma h]$ are in \mathcal{B}_2 and thus their product must be). Only if one of $[m; g]$ and $[n; h]$ is in \mathcal{B}_2 and the other is not, do we find that $[mn; gh] \notin \mathcal{B}_2$.

Lemma 11. *The map φ preserves products.*

Proof. Based on the definition of φ , there are several cases to check.

Case A: Suppose that $[m; g]$ and $[n; h]$ are each elements of both \mathcal{B}_1 and \mathcal{B}_2 . Based on the previous remarks, $[mn; gh]$ is in both as well, and so

$$\begin{aligned} \varphi([m; g] \star_{\mathcal{B}_1} [n; h]) &= \varphi(\gamma_{g,h} [mn; gh]) = \gamma_{g,h} [mn; gh] \\ &= [m; g] \star_{\mathcal{B}_2} [n; h] = \varphi([m; g]) \star_{\mathcal{B}_2} \varphi([n; h]) \end{aligned}$$

Case B: Next, suppose that $[m; g]$ and $[n; h]$ are each elements of \mathcal{B}_1 , but neither are elements of \mathcal{B}_2 . Notice that in this case

$$\varphi([m; g]) \star_{\mathcal{B}_2} \varphi([n; h]) = k_g [\tilde{m}; \sigma g] \star_{\mathcal{B}_2} k_h [\tilde{n}; \sigma h] = k_g k_h \gamma_{\sigma g, \sigma h} [mn; gh],$$

since $\sigma g \sigma h = gh$. Thus, to verify that $\varphi([m; g]) \star_{\mathcal{B}_2} \varphi([n; h]) = \varphi([m; g] \star_{\mathcal{B}_1} [n; h])$, we only need to check that

$$k_g k_h \gamma_{\sigma g, \sigma h} [mn; gh] = \gamma_{g,h} [mn; gh].$$

There are three relevant cases.

- (1) g and h are both of the form $(0, \alpha, 0, \beta)$
- (2) g and h are both of the form $(1/2, \alpha, 1/2, \beta)$
- (3) g is of one of the above forms and h is of the other.

It is straightforward to verify that in each case, $I_g \cup I_h \cup I_{gh} = \{x_0, \dots, x_m, y_0, \dots, y_n\}$ if and only if $I_{\sigma g} \cup I_{\sigma h} \cup I_{gh} = \{x_0, \dots, x_m, y_0, \dots, y_n\}$.

We verify the details of the preservation of the product in only the first case; the others follow by similar arguments.

- (1) g and h are both of the form $(0, \alpha, 0, \beta)$.

Here, $k_g = k_h = 1/2$. Suppose that $I_g \cup I_h \cup I_{gh} = \{x_0, \dots, x_m, y_0, \dots, y_n\}$. Lemma 8 gives

$$\gamma_{g,h} = \gamma_{g,\sigma h} = \frac{1}{4} \gamma_{\sigma g, \sigma h}$$

So, we have, $k_g k_h \gamma_{\sigma g, \sigma h} = 1/4 \cdot \gamma_{\sigma g, \sigma h} = \gamma_{g,h}$, as desired.

Case C: Finally, we must consider the case when $[m; g], [n; h]$ are both elements of \mathcal{B}_1 , but only one of them is an element of \mathcal{B}_2 . Without loss of generality, we say $[m; g] \in \mathcal{B}_2$. Notice that in this case

$$\varphi([m; g] \star_{\mathcal{B}_1} [n; h]) = \varphi(\gamma_{g,h} [mn; gh]) = \gamma_{g,h} k_{gh} [\widetilde{mn}; \sigma gh],$$

and,

$$\varphi([m; g]) \star_{\mathcal{B}_2} \varphi([n; h]) = [m; g] \star_{\mathcal{B}_2} k_h [\tilde{n}; \sigma h] = \gamma_{g,\sigma h} k_h [\widetilde{mn}; \sigma gh].$$

We must prove $\gamma_{g,h} k_{gh} [mn; \sigma gh] = \gamma_{g,\sigma h} k_h [mn; \sigma gh]$ in the following four cases:

- (1) g and h are both of the form $(0, \alpha, 0, \beta)$ ($k_h = k_{gh} = 1/2$)
- (2) g and h are both of the form $(1/2, \alpha, 1/2, \beta)$ ($k_h = 2, k_{gh} = 1/2$).
- (3) g is of the form $(0, \alpha, 0, \beta)$, and h is of the form $(1/2, \alpha, 1/2, \beta)$ ($k_h = k_{gh} = 2$)
- (4) g is of the form $(1/2, \alpha, 1/2, \beta)$, and h is of the form $(0, \alpha, 0, \beta)$ ($k_h = 1/2, k_{gh} = 2$).

Again, we verify the preservation of the product in the first case only; the others follow similarly.

- (1) g and h are both of the form $(0, \alpha, 0, \beta)$ ($k_h = k_{gh} = 1/2$). From Lemma 8, $\gamma_{g,h} = \gamma_{g,\sigma h}$. So, $k_{gh}\gamma_{g,h} = \frac{1}{2}\gamma_{g,h} = \frac{1}{2}\gamma_{g,\sigma h} = k_h\gamma_{g,\sigma h}$, as desired. \square

Remark 11.1. Recall that for an element $\lceil m; g \rceil$ from an LG state space, m represents a monomial together with a volume form. In the following proof, we are sometimes interested in only the monomial portion of m (which by abuse of notation, we will call m). Further, we point out that the monomial portion of both m and \tilde{m} are the same.

Lemma 12. *The map φ preserves the pairing.*

Proof. Recall that if $\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1} \neq 0$, then we have

$$mn = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_g} \text{Hess}(W_g) + l.o.t$$

If both $\lceil m; g \rceil$ and $\lceil n; g^{-1} \rceil$ are in \mathcal{B}_2 , then the pairings are computed in exactly the same way in both \mathcal{B}_1 and \mathcal{B}_2

If neither $\lceil m; g \rceil$ nor $\lceil n; g^{-1} \rceil$ are in \mathcal{B}_2 , then we have two cases to consider.

- (1) g is of the form $(0, \alpha, 0, \beta)$.

$$m \cdot n = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_g} \text{Hess}(W_g) = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_{\sigma g}} 4 \text{Hess}(W_{\sigma g})$$

Thus, for $\varphi(\lceil m; g \rceil) = \frac{1}{2} \lceil \tilde{m}; \sigma g \rceil$, and $\varphi(\lceil n; g^{-1} \rceil) = \frac{1}{2} \lceil \tilde{n}; \sigma g^{-1} \rceil$

$$\frac{1}{2}m \cdot \frac{1}{2}n = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_{\sigma g}} \text{Hess}(W_{\sigma g}),$$

and $\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1} = \langle \varphi(\lceil m; g \rceil), \varphi(\lceil n; g^{-1} \rceil) \rangle_{\mathcal{B}_2}$.

- (2) g is of the form $(1/2, \alpha, 1/2, \beta)$.

$$m \cdot n = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_g} \text{Hess}(W_g) = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_{\sigma g}} \frac{1}{4} \text{Hess}(W_{\sigma g})$$

Thus, for $\varphi(\lceil m; g \rceil) = 2 \lceil \tilde{m}; \sigma g \rceil$, and $\varphi(\lceil n; g^{-1} \rceil) = 2 \lceil \tilde{n}; \sigma g^{-1} \rceil$

$$2m \cdot 2n = \frac{\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1}}{\mu_{\sigma g}} \text{Hess}(W_{\sigma g}).$$

So, $\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle_{\mathcal{B}_1} = \langle \varphi(\lceil m; g \rceil), \varphi(\lceil n; g^{-1} \rceil) \rangle_{\mathcal{B}_2}$.

Note that if $\lceil m; g \rceil$ is in \mathcal{B}_2 but $\lceil n; g^{-1} \rceil$ is not, then $\langle \lceil m; g \rceil, \lceil n; g^{-1} \rceil \rangle = 0$, since the j_x acts with weight zero on the Hessian. \square

We have verified that φ preserves both the product and the pairing, and therefore, the composition, $\varphi \circ \kappa_1^{-1}$ is an isomorphism of Frobenius algebras when Property $(*)$ is satisfied.

5. A GEOMETRIC VIEW

In this section, we consider the geometry of the Borcea–Voisin model. We begin by describing a construction called the GLSM, which gives some reason why we expect the LG/CY correspondence to hold. In Section 5.2 we will describe the LG/CY state space isomorphism, establishing an equivalence between the BVLG model state space and the state space of the BV orbifold. In Section 5.4 we describe the twist map on the geometric side, generalizing the twist map of [4] and [2]. In Section 5.5 we pull everything together and discuss some consequences in the corresponding geometry.

5.1. Gauged linear sigma models. We include this section, in order to demonstrate why the LG model described in Section 2.3 has the particular form we have described there. None of the ideas presented in this section will be used in any later proofs.

The LG/CY correspondence is part of a larger idea due to Ed Witten, in which he considers each theory in the correspondence (e.g. GW theory or FJRW theory) as different “phases” of some larger theory called the *gauged linear sigma model* (GLSM) that depends on a parameter. Variation of this parameter will then produce the various theories involved in the LG/CY correspondence. The different phases of the GLSM are conjectured to be equivalent to each other in some sense. The first evidence of this equivalence is an isomorphism of the respective state spaces in each phase.

These ideas have recently been mathematically formalised by Fan–Jarvis–Ruan in [9] and we begin this section by briefly describing their construction. We are intentionally vague about some aspects of this construction, since they are not necessary for our purposes. In this article, we are primarily interested in the state spaces, so we will focus our attention there.

A GLSM depends on a choice of (1) a finite dimensional vector space V over \mathbb{C} , (2) a reductive algebraic group G_V acting on V , (3) a G_V -character θ , and (4) a *superpotential* $\bar{W} : V \rightarrow \mathbb{C}$. From these ingredients, one obtains a *state space*, a *moduli space of LG-quasimaps* with a good virtual cycle, and numerical invariants defined as integrals over the virtual cycle.

If we vary the character θ , we get a different GIT quotient, and therefore a different theory. We can vary the character in the so-called phase space. This space of characters is partitioned into various chambers. Varying the character within a chamber, does not change the theory. However, if we cross into a different chamber, we obtain a different theory.

The idea behind the LG/CY correspondence is that if we choose the character θ in a certain chamber, we obtain GW theory of a particular orbifold, whereas if we choose the character in different chamber, we should expect to obtain FJRW theory. The moduli space, virtual cycle and numerical invariants from the GLSM will give rise to such structures for Gromov–Witten theory and for FJRW theory. It has been conjectured that these structures will agree for both theories, but as mentioned we will focus our attention on the state spaces.

The relevant GLSM’s for BV models are obtained using the following input. Let $V = \mathbb{C}^{n+m+2} \times \mathbb{C}^2$ with coordinates $x_0, \dots, x_n, y_0, \dots, y_m$ and p_1, p_2 . Next we define an action of $(\mathbb{C}^*)^2$ on V via the weights

$$\begin{pmatrix} u_0 & \dots & u_n & 0 & \dots & 0 & -d_1 & 0 \\ 0 & \dots & 0 & v_0 & \dots & v_m & 0 & -d_2 \end{pmatrix}$$

We can embed $(\mathbb{C}^*)^2$ as diagonal matrices into $\mathrm{GL}(V)$ via these weights.

We can similarly embed the group $G_{\max} \hookrightarrow \mathrm{GL}(V)$ acting only on the x and y coordinates. Our reductive group G_V we define as the subgroup of $\mathrm{GL}(V)$ generated by $(\mathbb{C}^*)^2$ and ΣG . For a superpotential, we take $\bar{W} = p_1 W_1 + p_2 W_2$.

N:I restructured this. The text is more dense now, but also a quicker read. What do you think? We could make list out of the input still, if it would make it look better there.

In order to describe the two relevant GIT quotients, we need to choose appropriate characters of G_V . In order to identify the characters, we use the method of symplectic reduction. For more on this perspective and the relationship between symplectic reduction and GIT quotients, see the original construction on GLSM by Fan–Jarvis–Ruan [9]. We take the standard Kähler metric on V . Since G_V is reductive, it is the complexification of a maximal compact Lie subgroup H acting on V via a faithful unitary representation. The Lie algebra in our case is \mathbb{R}^2 .

We consider the Hamiltonian action of H on V , which has moment map $\mu : V \rightarrow \mathbb{R}^2$ given by

$$\mu_1 = \sum_{i=0}^n u_i |x_i|^2 - d_1 |p_1|^2, \quad \mu_2 = \sum_{j=0}^m v_j |y_j|^2 - d_2 |p_2|^2.$$

The set of critical values for this moment map is $\{\tau_1 = 0\} \cup \{\tau_2 = 0\} \subset \mathbb{R}^2$.

The conditions for the appropriate GIT quotients translate to the requirement of having a regular value of the moment map in the symplectic setting. Once we have our regular values we can return to the algebraic setting to describe the GIT quotients. Notice the set of critical values divides \mathbb{R}^2 into 4 chambers. Each regular value within a given chamber will yield an isomorphic GIT quotient.

A derivation of a character defines a weight in the Lie algebra \mathbb{R}^2 . We will now discuss the two chambers which yield GW theory and FJRW theory. First we describe the regular values which yield the relevant characters; then we return to the GIT perspective, and describe the unstable locus and the corresponding GIT quotient for each character which has derivation in the given chamber.

GW phase. The first chamber we consider is defined by $\tau_1 > 0, \tau_2 > 0$. In this chamber, the unstable locus is the set of points in \mathbb{C}^{n+m+2} with $(x_0, \dots, x_n) \neq 0$, and $(y_0, \dots, y_m) \neq 0$. We get as a GIT quotient

$$[(\mathbb{C}^{n+1} \setminus 0) \times (\mathbb{C}^{m+1} \setminus 0) \times \mathbb{C}^2 / G_V] \cong [\mathcal{O}_{\mathbb{P}(u_0, \dots, u_n)}(-d_1) \times \mathcal{O}_{\mathbb{P}(v_0, \dots, v_m)}(-d_2) / \widetilde{\Sigma G}],$$

Here $\widetilde{\Sigma G} = \Sigma G / (J_1 \times J_2)$. The critical locus of \overline{W} is $\{W_1 = 0, W_2 = 0, p_1 = p_2 = 0\} / \widetilde{\Sigma G}$, which is the stack

$$[X_{W_1} \times X_{W_2} / \widetilde{\Sigma G}].$$

If we set $n = 2, m = 3$ this has a Borcea–Voisin variety as crepant resolution. Though it has not been verified in every case, the state space corresponding to this GIT quotient is expected to be the Chen–Ruan cohomology of the critical locus of \overline{W} . In the next section, we will show that the Chen–Ruan cohomology of the Borcea–Voisin orbifold is indeed isomorphic to the state space of FJRW theory.

FJRW phase. The other relevant chamber is defined by $\tau_1 < 0, \tau_2 < 0$. The unstable locus is the set of points in V with $p_1 \neq 0$, and $p_2 \neq 0$. We get the GIT quotient

$$[\mathbb{C}^{n+m+2} \times (\mathbb{C}^*)^2 / G_V] \cong [\mathbb{C}^{n+m+2} / \Sigma G]$$

With superpotential $\overline{W} = W$ (we scaled away the p 's).

The critical locus of \overline{W} is the origin. In this chamber, the state space is exactly the state space of FJRW theory.

Because both theories are merely different phases of the same GLSM, we expect the theories to be equivalent. The most basic manifestation of this equivalence is an isomorphism of state spaces, which we show in the next section.

5.2. LG/CY correspondence: state space isomorphism. In this section, we will prove that the two state spaces described above are isomorphic as graded vector spaces, i.e.

$$(9) \quad \mathcal{A}_{W, \Sigma G}^{p,q} \cong H_{CR}^{p,q} \left([X_{W_1} \times X_{W_2} / \widetilde{\Sigma G}]; \mathbb{C} \right)$$

This is similar to what was done by Chiodo–Ruan in [7]. However, there are some important differences, so we will give the details here. We begin by expressing each side of Equation (9) in a more useable form.

5.2.1. FJRW state space. We begin with FJRW theory. In order to simplify notation, we will write $J = J_1 \times J_2$. For $g \in \Sigma G$, we write $g = (g_1, g_2)$ with $g_1 \in G_{W_1}^{max}$ and $g_2 \in G_{W_2}^{max}$. Notice this differs slightly from previous sections, where we wrote $g = (\epsilon/2, \alpha, \epsilon'/2, \beta)$.

Finally, define the following notation:

$$\begin{aligned} \mathbb{C}_{g_1}^{n+1} &= \{x \in \mathbb{C}^{n+1} \mid g_1 x = x\} \\ W_{1,g_1} &= W|_{\mathbb{C}_{g_1}^{n+1}} \end{aligned}$$

with similar definitions for g_2 .

Recall the definition of the state space:

$$(10) \quad \mathcal{A}_{W, \Sigma G}^{p,q} = \bigoplus_{g \in \Sigma G} (\mathcal{Q}_{W_g}^{p,q} \omega_g)^{\Sigma G}$$

Since $W_g = W_{1,g_1} + W_{2,g_2}$, we obtain $\mathcal{Q}_{W_g} = (\mathcal{Q}_{W_{1,g_1}} \otimes \mathcal{Q}_{W_{2,g_2}})$. Thus in the g -sector, we can write

$$(11) \quad \mathcal{Q}_{W_g}^{p,q} \omega_g = \bigoplus_{\substack{h_1+h_2=p \\ k_1+k_2=q}} (\mathcal{Q}_{W_{1,g_1}}^{h_1,k_1} \omega_{g_1} \otimes \mathcal{Q}_{W_{2,g_2}}^{h_2,k_2} \omega_{g_2})^{\Sigma G}.$$

This isomorphism sends

$$[m; g] \rightarrow [m_1; g_1] \otimes [m_2; g_2]$$

where $m = m_x m_y$. Recall the definition of bidegree for FJRW theory in equation (6). Notice $\text{age } g = \text{age } g_1 + \text{age } g_2$ and $N_g = N_{g_1} + N_{g_2}$. Therefore the tensor product preserves the bigrading.

Furthermore we have the bidegree on the left equal to the sum of the bidegrees on the right hand side.

We now decompose ΣG into $2M = 2|G|/(d_1 d_2)$ cosets of J . We denote by \mathcal{C} the set of cosets. We can choose these representatives for each coset so that the first M representatives fix x_0 and y_0 , and the last M coset representatives are simply the first ones multiplied by σ .

Now we can write degree (p, q) part of the state space as a sum over these cosets. Using (10) and (11) this becomes

$$(12) \quad \mathcal{A}_{W, G}^{p,q} = \bigoplus_{g \in \mathcal{C}} \bigoplus_{\substack{h_1+h_2=p \\ k_1+k_2=q}} \left[\left(\bigoplus_{k=1}^{d_1} \mathcal{Q}_{g_1 j_1^k}^{h_1, k_1} \right)^{J_1} \otimes \left(\bigoplus_{k=1}^{d_2} \mathcal{Q}_{g_2 j_2^k}^{h_2, k_2} \right)^{J_2} \right]^{\Sigma G}.$$

Recall that j_i is the generator of J_i and for $i = 1, 2$. Notice we have taken J_1 and J_2 invariants in each of the factors of the tensor product in the second line. We do this simply to make the isomorphism more clear.

5.2.2. *Chen–Ruan cohomology.* The Chen–Ruan cohomology is slightly more subtle. On the Calabi–Yau side, the state space is

$$H_{CR}^* \left(\left[X_{W_1} \times X_{W_2} / \widetilde{\Sigma G} \right] \right).$$

The *Chen–Ruan orbifold cohomology* is defined via the *inertia orbifold* (see [7]). If $\mathcal{X} = [X/H]$ is a global quotient of a nonsingular variety X by a finite group H , the inertia orbifold $I\mathcal{X}$ takes a particularly simple form. Let S_H denote the set of conjugacy classes (h) in H , then

$$I[V/H] = \coprod_{(h) \in S_H} [V^h / C(h)].$$

As a vector space, the Chen–Ruan cohomology groups $H_{CR}^*(\mathcal{X})$ of an orbifold \mathcal{X} are the cohomology groups of its inertia orbifold:

$$H_{CR}^*(\mathcal{X}) := H^*(I\mathcal{X}).$$

The bidegree on the Chen–Ruan cohomology is the normal bidegree with an *age shift*, which we will describe below.

In order to compute the Chen–Ruan cohomology, we need to describe the orbifold $[X_{W_1} \times X_{W_2} / \widetilde{\Sigma G}]$ in a different form. We write T^2 for the torus $(\mathbb{C}^*)^2$, acting on \mathbb{C}^{m+n+2} via the weights

$$\begin{pmatrix} u_0 & \dots & u_n & 0 & \dots & 0 \\ 0 & \dots & 0 & v_0 & \dots & v_m \end{pmatrix}$$

and the group ΣGT^2 for the product of the two groups. Notice that $\Sigma G \cap T^2 = J$. We will denote by μ_{d_1} the cyclic group of order d_1 in the first copy of \mathbb{C}^* and μ_{d_2} the corresponding one in the second factor. Notice that μ_{d_i} corresponds with the cyclic group J_i .

Define

$$\begin{aligned} V_{W_1} &= \{W_1 = 0\} \subset \mathbb{C}^{n+1} \setminus \{0\} \\ V_{W_2} &= \{W_2 = 0\} \subset \mathbb{C}^{m+1} \setminus \{0\}. \end{aligned}$$

Notice that, since $\Sigma G \cap T^2 = J$, we have $\Sigma GT^2 / T^2 \cong \widetilde{\Sigma G}$. Using this description, we can express the orbifold (see [[Romagny]] as

$$\left[X_{W_1} \times X_{W_2} / \widetilde{\Sigma G} \right] \cong [V_{W_1} \times V_{W_2} / \Sigma GT^2].$$

The Chen–Ruan cohomolgy can therefore be written as

$$(13) \quad H_{CR}^{p,q} \left(\left[X_{W_1} \times X_{W_2} / \widetilde{\Sigma G} \right] \right) = \bigoplus_{g \in \Sigma GT^2} H^{p-a(g), q-a(g)}((V_{W_1} \times V_{W_2})_g / \Sigma GT^2; \mathbb{C}).$$

Here $a(g)$ denotes the age shift mentioned earlier. To define this we look at the action of g on the tangent space at a given point. The action can be diagonalized in a suitable basis as

$$\begin{pmatrix} e^{2\pi i a_1} & & 0 \\ & \ddots & \\ 0 & & e^{2\pi i a_{n+m}} \end{pmatrix}.$$

with $0 \leq a_i < 1$. The age $a(g) = \sum_{l=1}^{m+n} a_l$. Notice this differs from the age in FJRW theory, because of g acting on the tangent space to $V_{W_1} \times V_{W_2}$ instead of the action on \mathbb{C}^{n+m+2} . The difference in the two age contributions is discussed at some length in [7].

With this description of the Chen–Ruan cohomology, we can now write the Chen–Ruan cohomology of the Borcea–Voisin orbifold in a more suitable manner. Given $g \in \Sigma GT^2$, we can write

$g = (g_1, g_2)$ with g_1 acting only on the x 's and g_2 only on the y 's, as we did with FJRW theory. We can define $\mathbb{C}_{g_1}^{n+1}$ and W_{1,g_1} similar to what was in the previous section, and we define the hypersurfaces

$$\begin{aligned} V_{W_1, g_1} &= \{W_{1, g_1} = 0\} \subset \mathbb{C}_{g_1}^{n+1} \setminus \{0\} \\ V_{W_2, g_2} &= \{W_{2, g_2} = 0\} \subset \mathbb{C}_{g_2}^{m+1} \setminus \{0\}. \end{aligned}$$

Now consider a single summand corresponding to $g \in \Sigma GT^2$. We can write $(V_{W_1} \times V_{W_2})_g = V_{W_1, g_1} \times V_{W_2, g_2}$

$$(V_{W_1} \times V_{W_2})_g / \Sigma GT^2 \cong (V_{W_1, g_1} \times V_{W_2, g_2} / T^2) / \widetilde{\Sigma G}.$$

But

$$V_{W_1, g_1} \times V_{W_2, g_2} / T^2 \cong X_{W_1, g_1} \times X_{W_2, g_2}.$$

So for a fixed $g \in \Sigma GT^2$, we can write the cohomology as

$$\begin{aligned} H^{p-a(g), q-a(g)}(V_{W_1, g_1} \times V_{W_2, g_2} / \Sigma GT^2; \mathbb{C}) &= H^{p-a(g), q-a(g)}(X_{W_1, g_1} \times X_{W_2, g_2} / \widetilde{\Sigma G}; \mathbb{C}) \\ &= H^{p-a(g), q-a(g)}(X_{W_1, g_1} \times X_{W_2, g_2}; \mathbb{C})^{\widetilde{\Sigma G}} \\ &= \bigoplus_{\substack{h_1+h_2=p \\ k_1+k_2=q}} \left(H^{h_1-a(g_1), k_1-a(g_1)}(X_{W_1, g_1}; \mathbb{C}) \otimes H^{h_2-a(g_2), k_2-a(g_2)}(X_{W_2, g_2}; \mathbb{C}) \right)^{\widetilde{\Sigma G}}. \end{aligned}$$

Notice that the summand corresponding to g in Equation (13) is empty if $(V_{W_1} \times V_{W_2})_g$ is empty. As with FJRW theory, we will rewrite this sum as a sum over cosets. To this end, if we write $g_1 = (g_{10}, g_{11}, \dots, g_{1n}) \in G_{W_1}$ and $g_2 = (g_{20}, \dots, g_{2m})$ we can define

$$\Lambda_g^1 = \bigcup_{k=0}^n \{ \lambda \in \mathbb{C}^* \mid \lambda^{-w_k} = g_{1k} \}.$$

This set denotes the complex numbers $\lambda \in \mathbb{C}^*$ where $g_1 \lambda$ has non-trivial fixed locus. The set Λ_g^2 is defined similarly.

Notice that in this description, since the tangent space of a product is the product of the respective tangent spaces, we have $a(g) = a(g_1) + a(g_2)$. So the degree shift agrees on both sides of the equation. In other words, we can think of the summands as tensor product of factors of some Chen–Ruan cohomology.

As in the FJRW state space, we can choose $2M$ cosets of T^2 , using the same set \mathcal{C} of coset representatives as before. We can then write the Chen–Ruan cohomology as

$$\begin{aligned} H_{CR}^{p,q}([X_{W_1} \times X_{W_2} / \Sigma G]; \mathbb{C}) &= \\ (14) \quad \bigoplus_{g \in \mathcal{C}} \bigoplus_{\substack{h_1+h_2=p \\ k_1+k_2=q}} &\left[\left(\bigoplus_{\lambda \in \Lambda_g^1} H^{h_1, k_1}(X_{W_1, g\lambda}; \mathbb{C}) \right) \otimes \left(\bigoplus_{\lambda \in \Lambda_g^2} H^{h_2, k_2}(X_{W_2, g\lambda}; \mathbb{C}) \right) \right]^{\Sigma G}. \end{aligned}$$

On both sides of this equation, we have written bidegrees *with* the age shift. This is standard on the left hand side of the equation. On the right hand side, however, we mean the $h_1 + a(g_1), k_1 + a(g_2)$ part of the ordinary cohomology.

5.2.3. *The Isomorphism.* Comparing expressions (12) and (14), we see that the isomorphism of state spaces will follow, once we establish the isomorphism

$$(15) \quad \left(\bigoplus_{k=1}^{d_i} \mathcal{Q}_{g_i j_i^k}^{h_i, k_i} \right)^{J_i} \cong \bigoplus_{\lambda \in \Lambda_g^i} H^{h_i, k_i}(X_{W_i, g_i \lambda}; \mathbb{C}).$$

as an isomorphism of bigraded vector spaces. Again here, we take the convention that the right hand side of (15) is shifted by the age shift and the left hand side has the bidegree defined by FJRW theory. We are summing over J_1 and the elements of \mathbb{C}^* contributing a nonzero sector to the Chen–Ruan cohomology, resp. Furthermore, because this isomorphism is $G_{W_i}^{max}$ -equivariant, the action of $G_{W_i}^{max}$ is the same on both sides of the isomorphism (see [7]), and hence the action of ΣG is the same on the summands of both (12) and (14).

Equation (15) is proven by Chiodo–Ruan in [7], but we give a brief outline here. For ease of exposition, we will focus on the isomorphism with $i = 1$. The same will be true for $i = 2$. On the left hand side of (9) we have

$$(16) \quad \left(\bigoplus_{k=1}^{d_1} \mathcal{Q}_{g_1 j_1^k}^{h_1, k_1} \right)^{J_1} = \bigoplus_{\lambda \in \mu_{d_1} \cap \Lambda_g^1} H^{N_{g_1 \lambda}}(\mathbb{C}_{g_1 \lambda}^{n+1}, W_{g_1 \lambda}^{+\infty}; \mathbb{C})^{J_1} \oplus \bigoplus_{\lambda \in \mu_d \setminus \Lambda_g^1} 1_{g_1 \lambda} \mathbb{C}.$$

This is the definition of the various sectors of FJRW theory. Notice we have decomposed the left hand side into a sum of broad sectors and narrow sectors. The action of $G_{W_1}^{max}$ is trivial on narrow sectors.

On the other hand, we can express the cohomology of a hypersurface in weighted projective space as a direct sum of the ambient cohomology (coming from projective space) and the primitive cohomology. In e.g. [7] we see that the primitive cohomology can be expressed as the cohomology of the Milnor fiber invariant under the monodromy action. Thus we can write:

$$H^*(X_{W_1, g_1}; \mathbb{C}) \cong H^{N_{g_1 \lambda}}(\mathbb{C}_{g_1 \lambda}^{n+1}, W_{g_1 \lambda}^{+\infty}; \mathbb{C})^{J_1} \oplus H^{amb}(X_{W_1, g_1})$$

In this description, $G_{W_1}^{max}$ acts trivially on the ambient classes, and the group action on the primitive cohomology is the same as the action on the FJRW state space.

One last thing to note is that the Chen Ruan cohomology contains a sum over elements $\lambda \in \mathbb{C}^*$, or rather a sum over $\lambda \in \Lambda_g^1$, since only these contribute to the cohomology. It may happen that for some of these the corresponding diagonal symmetry does not lie in J_1 , i.e. when $\lambda \notin \mu_{d_1}$ in the notation of the decompositions written above. In this case, $W_{1, g_1 \lambda}$ vanishes on all of $\mathbb{C}_{g_1 \lambda}^{n+1}$ (see [7, Theorem ??]), and therefore the primitive cohomology vanishes. For these summands, we obtain simply the cohomology of weighted projectives space (the ambient classes).

Thus we have

$$(17) \quad \bigoplus_{\lambda \in \mathbb{C}^*} H^\bullet(X_{W_1, g_1 \lambda}; \mathbb{C}) \cong \bigoplus_{\lambda \in \mu_{d_1} \cap \Lambda_g^1} H^{N_{g_1 \lambda}}(\mathbb{C}_{g_1 \lambda}^{n+1}, W_{g_1 \lambda}^{+\infty}; \mathbb{C})^{J_1} \oplus \bigoplus_{\lambda \in \Lambda_g^1} H^{amb}(X_{W_1, g}).$$

In this expression, $G_{W_1}^{max}$ acts trivially on the ambient classes.

Comparing expressions (16) and (17), we see that we only need to compare the narrow sectors and the ambient classes. The degree shift for the broad sectors and for the primitive classes in the Chen–Ruan cohomology agree because the two age shifts agree (see [7, Lemma 22]). The key observation in proving (15) is that the number of narrow sectors in FJRW theory is equal to the sum of the dimension of all of the primitive cohomology, when one sums over $\lambda \in \mathbb{C}^*$. Furthermore, the bidegrees of the classes also agree, after the degree shift on both sides. As mentioned before,

$G_{W_1}^{max}$ acts trivially on both the narrow sectors of FJRW theory and the ambient classes of the Chen–Ruan cohomology.

This establishes (15) for $i = 1$. We can do the same with W_2 obtaining a similar isomorphism. Putting these together with (12) and (14), we obtain the isomorphism of state spaces.

5.3. Borcea–Voisin mirror symmetry. One of the first predictions of mirror symmetry is the rotation of the Hodge diamond. Such a statement is one of the main results in [6], for the Borcea–Voisin orbifolds of the form we consider here. In other words, they prove that if $G = G_1 \times G_2$ with $G_i \in G_{W_i}^{max}$ and we let $N = n + m - 2$, then

$$H_{CR}^{p,q}([X_{W_1} \times X_{W_2}/\widetilde{\Sigma G}]; \mathbb{C}) \cong H_{CR}^{N-p,q}([X_{W_1^T} \times X_{W_2^T}/\widetilde{\Sigma G^T}]; \mathbb{C}).$$

Their method of proof uses what they call 1/2-Calabi–Yau manifolds. In this section, we provide an alternative proof of this fact, while dropping the condition on G . If we relax the condition on G , we no longer obtain Borcea–Voisin orbifolds, as we have seen previously, but as we see from the BVLG mirror symmetry, we expect mirror symmetry to hold nonetheless. We describe this isomorphism now.

From Section 4 we have

$$\mathcal{A}_{W,\Sigma G}^{p,q} \cong \mathcal{B}_{W^T,\Sigma G^T}^{p,q}.$$

On the other hand, from Section 5.2, we have

$$\mathcal{A}_{W,\Sigma G}^{p,q} \cong H_{CR}^{p,q}(X_{W_1} \times X_{W_2}/\widetilde{\Sigma G}; \mathbb{C}).$$

Putting these together, we get the following corollary generalizing the result of Chiodo–Kalashnikov–Veniani.

Corollary 13. *With $N = n + m - 2$*

$$H_{CR}^{p,q}([X_{W_1} \times X_{W_2}/\widetilde{\Sigma G}]; \mathbb{C}) \cong H_{CR}^{N-p,q}([X_{W_1^T} \times X_{W_2^T}/\widetilde{\Sigma G^T}]; \mathbb{C})$$

In case $n = 2$, $m = 3$ and $G = G_1 \times G_2$ is a product of groups $G_1 \subset G_{W_1}^{max}$, $G_2 \subset G_{W_2}^{max}$, then a crepant resolution of both orbifolds exists, and we obtain the mirror symmetry of Borcea–Voisin at the level of state spaces.

5.4. The twisting map. Now we have established the state space isomorphism, we turn our attention to the twist map of Borcea [4] and Artebani–Boisierre–Sarti [2]. Geometrically, the twist map relates the orbifold $[X_{W_1} \times X_{W_2}/\widetilde{\Sigma G}]$ to a hypersurface in a quotient of weighted projective space. What follows is an instance where the LG side of the LG/CY correspondence informs the CY geometry.

In [2] the authors describe the twist map for those pairs of polynomials W_1, W_2 with the property that $\gcd(u_0, v_0) = 1$. As mentioned before, there are 44 of the 95 weight systems for K3 surfaces that admit such a polynomial, and two weight systems for an elliptic curve. Thus we have 88 total combinations. The restriction $\gcd(u_0, v_0) = 1$ limits us to 48 of these combinations. However, as mentioned in Section 3, the LG/CY correspondence allows us to understand the twist map more clearly. For example, the restriction on gcd’s can be lifted so that the twisting map is valid for all choices of polynomials that have form (7), as soon as we understand the group G_{tw} .

In order to define the twist map, let $\delta = \gcd(u_0, v_0)$. We define s_0 and t_0 via the equations $s_0 u_0 + \delta \equiv 0 \pmod{v_0}$ and $t_0 v_0 + \delta \equiv 0 \pmod{u_0}$. Let σ be the involution on $\mathbb{P}(u_1, \dots, u_m) \times$

$\mathbb{P}(v_0, v_1, \dots, v_n)$ given by $(x_0, y_0) \mapsto (-x_0, -y_0)$, Finally, define

$$s = \frac{s_0 u_0 + \delta}{v_0}, \quad t = \frac{t_0 v_0 + \delta}{u_0}.$$

Recall that the LG twist map relates the LG model $(W_1 + W_2, \Sigma G)$ to the LG model $(f_1 - f_2, G_{tw})$. One can check that $(J_1 \times J_2)_{tw} \subset G_{tw}$ and contains $J_{f_1 - f_2}$ and that $(J_1 \times J_2)_{tw} / J_{f_1 - f_2}$ is cyclic of order δ . Thus this group acts on $\mathbb{P}(\frac{v_0}{\delta} u_1, \dots, \frac{v_0}{\delta} u_m, \frac{u_0}{\delta} v_1, \dots, \frac{u_0}{\delta} v_n)$ (cf. Section 3). We define the map

$$\tilde{\tau} : \mathbb{P}(u_0, u_1, \dots, u_m) \times \mathbb{P}(v_0, v_1, \dots, v_n) \rightarrow \left[\mathbb{P}\left(\frac{v_0}{\delta} u_1, \dots, \frac{v_0}{\delta} u_m, \frac{u_0}{\delta} v_1, \dots, \frac{u_0}{\delta} v_n\right) / \widetilde{G_{tw}} \right]$$

by

$$((x_0, x_1, \dots, x_n), (y_0, y_1, \dots, y_m)) \mapsto \left((x_0^{s_0} y_0^t)^{\frac{u_1}{\delta}} x_1, \dots, (x_0^{s_0} y_0^t)^{\frac{u_m}{\delta}} x_m, (x_0^s y_0^{t_0})^{\frac{v_1}{\delta}} y_1, \dots, (x_0^s y_0^{t_0})^{\frac{v_n}{\delta}} y_n \right).$$

This map depends on a choice of δ -th root of unity, and on the choice of s_0 and t_0 . However, one can check that with a different choice of any of these, the image differs exactly by the action of an element of $(J_1 \times J_2)_{tw}$. Since the image lands in the quotient by G_{tw} , the map is well-defined.

Furthermore, the image written above is equivalent to

$$\left(\left(\frac{y_0}{x_0} \right)^{\frac{u_1}{u_0}} x_1, \dots, \left(\frac{y_0}{x_0} \right)^{\frac{u_m}{u_0}} x_m, y_1, \dots, y_n \right).$$

from which we see that τ descends to a well-defined map on the orbits of σ .

The twisting map is defined as the restriction of τ to the product $X_{W_1} \times X_{W_2}$. We need to check the image of τ is contained in $\{f - g = 0\}$. Recall that the weights of x_0 and y_0 are $1/2$, so we have $2u_0 = d_1$ (the degree of the first polynomial), and $2v_0 = d_2$. From the definition of s , we obtain

$$\frac{s}{\delta} d_2 = \frac{s_0}{\delta} d_1 + 2$$

$$\frac{t}{\delta} = \frac{t_0}{\delta} + 2.$$

Evaluating $f - g$ at any point in the image of τ , we get

$$x_0^{s_0 d_1 / \delta} y_0^{t_0 d_2 / \delta} (y_0^2 f(x) - x_0^2 g(y)) = 0$$

And we obtain the map

$$\tau : [X_{W_1} \times X_{W_2} / \widetilde{\Sigma G}] \rightarrow [X_{f-g} / \widetilde{G_{tw}}]$$

For a given choice of root, this map is smooth and a diffeomorphism almost everywhere. Furthermore, one can see from the definitions that if $\delta = 1$, we obtain the twist map of [2], which is in fact a birational morphism. So in case both domain and image are Calabi-Yau threefolds, and thus related by a sequence of simple flops, the two orbifolds related by the twist map have isomorphic genus-zero Gromov-Witten invariants. In general, it is more difficult to give any relationship between Gromov-Witten invariants, or even between the Chen-Ruan cohomologies. However, using the LG/CY correspondence, we will see that the Chen-Ruan cohomology of the two objects related by the twist map are indeed isomorphic.

N:citation?

5.5. Final applications. As mentioned in the previous section, we obtain an isomorphism of state spaces on the CY side. This relates the Chen–Ruan cohomology of a product with the cohomology of a hypersurface.

Corollary 14. *There is an isomorphism of bigraded vector spaces*

$$H_{CR}^{p,q}([X_{W_1} \times X_{W_2}/\widetilde{\Sigma G}]; \mathbb{C}) \cong H_{CR}^{p,q}([X_{f-g}/\widetilde{G_{tw}}]; \mathbb{C})$$

Proof. From the Section 5.2 we have

$$\mathcal{A}_{W,\Sigma G}^{p,q} \cong H_{CR}^{p,q}([X_{W_1} \times X_{W_2}/\widetilde{\Sigma G}]; \mathbb{C}).$$

From [7] we have

$$\mathcal{A}_{W_{tw},G_{tw}}^{p,q} \cong H_{CR}^{p,q}([X_{f-g}/\widetilde{G_{tw}}]; \mathbb{C}).$$

Finally from Section 4.1, we have

$$\mathcal{A}_{W,\Sigma G}^{p,q} \cong \mathcal{A}_{W_{tw},G_{tw}}^{p,q}$$

□

N: We may want to just remove this last bit, but I'm not sure. There was also some comment about a paper by Yui. I want to look at that one, but I think it was done in higher characteristic.

We expect this twist map to generalise naturally to complete intersections for polynomials f_i ,

$$\left(\prod_{i=1}^r \mathbb{P} \left(w_0^{(i)}, \dots, w_{k_1}^{(i)} \right) \right) / \mathbb{Z}_2 \rightarrow \mathbb{P} \left(\left(\prod_{i \neq 1} w_0^{(i)} \right) w_1^{(1)}, \dots, \left(\prod_{i \neq j} w_0^{(i)} \right) w_1^{(j)}, \dots, \left(\prod_{i \neq r} w_0^{(i)} \right) w_{k_r}^{(r)} \right),$$

where \mathbb{Z}_2 acts by negating the first coordinate of each factor, and sends

$$\left(\prod_{i=1}^r \{f_i = 0\} \right) / \mathbb{Z}_2$$

to

$$\{f_1 = f_2 = \dots = f_r\}.$$

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