

## Complex Inner Product Spaces

### The $\mathbb{C}^n$ spaces

The prototypical (and most important) real vector spaces are the Euclidean spaces  $\mathbb{R}^n$ . Any study of complex vector spaces will similar begin with  $\mathbb{C}^n$ . As a set,  $\mathbb{C}^n$  contains vectors of length  $n$  whose entries are complex numbers. Thus,

$$\begin{bmatrix} 2+i \\ 3-5i \\ i \end{bmatrix} \in \mathbb{C}^3,$$

$(5, -1)$  is an element found *both* in  $\mathbb{R}^2$  and  $\mathbb{C}^2$  (and, indeed, all of  $\mathbb{R}^n$  is found in  $\mathbb{C}^n$ ), and  $(0, 0, 0, 0)$  serves as the *zero* element in  $\mathbb{C}^4$ . Addition and scalar multiplication in  $\mathbb{C}^n$  is done in the analogous way to how they are performed in  $\mathbb{R}^n$ , except that now the scalars are allowed to be nonreal numbers. Thus, to rescale the vector  $(3+i, -2-3i)$  by  $1-3i$ , we have

$$(1-3i) \begin{bmatrix} 3+i \\ -2-3i \end{bmatrix} = \begin{bmatrix} (1-3i)(3+i) \\ (1-3i)(-2-3i) \end{bmatrix} = \begin{bmatrix} 6-8i \\ -11+3i \end{bmatrix}.$$

Given the notation  $\overline{3+2i}$  for the complex conjugate  $3-2i$  of  $3+2i$ , we adopt a similar notation when we want to take the complex conjugate simultaneously of all entries in a vector. Thus,

$$\text{if } \mathbf{z} = \begin{bmatrix} 3-4i \\ 2i \\ -2+5i \\ -1 \end{bmatrix}, \quad \text{then} \quad \bar{\mathbf{z}} = \begin{bmatrix} 3+4i \\ -2i \\ -2-5i \\ -1 \end{bmatrix}.$$

Both  $\mathbf{z}$  and  $\bar{\mathbf{z}}$  are vectors in  $\mathbb{C}^4$ . In general, if the entries of  $\mathbf{z}$  are all real numbers, then  $\bar{\mathbf{z}} = \mathbf{z}$ .

### The inner product in $\mathbb{C}^n$

In  $\mathbb{R}^n$ , the length of a vector  $\|\mathbf{x}\| = \sqrt{\mathbf{x} \cdot \mathbf{x}}$  is a real, nonnegative number. The modulus, or length, of a complex number  $z = a + ib$  is real and nonnegative as well:

$$|z| = \sqrt{\bar{z}z} = \sqrt{(a-ib)(a+ib)} = \sqrt{a^2+b^2}, \quad \text{or} \quad |z|^2 = \bar{z}z.$$

A natural idea, therefore, is to define an inner product between vectors  $\mathbf{z} = (z_1, \dots, z_n)$ ,  $\mathbf{w} = (w_1, \dots, w_n) \in \mathbb{C}^n$  in this manner:

$$\langle \mathbf{w}, \mathbf{z} \rangle = \sum_{j=1}^n \bar{w}_j z_j = \bar{w}_1 z_1 + \dots + \bar{w}_n z_n = \mathbf{w}^H \mathbf{z} = \bar{\mathbf{w}}^T \mathbf{z}. \quad (1)$$

Here,  $\mathbf{w}^H$  stands for  $\overline{\mathbf{w}}^T$ , the **conjugate transpose** of  $\mathbf{w}$ . For instance,

$$\text{if } \mathbf{z} = \begin{bmatrix} 3 - 4i \\ 2i \\ -2 + 5i \\ -1 \end{bmatrix}_{4 \times 1} \quad \text{then} \quad \mathbf{z}^H = \overline{\mathbf{z}}^T = \begin{bmatrix} 3 + 4i & -2i & -2 - 5i & -1 \end{bmatrix}_{1 \times 4}.$$

### Remarks

- Some authors define the inner product of complex vectors  $\mathbf{u}, \mathbf{v}$  to be the conjugate transpose of the second vector multiplied by the first—i.e.,  $\mathbf{v}^H \mathbf{u}$ . The two definitions suit the same purpose, but do not yield the same result. For example, if  $\mathbf{u} = (2 + i, 1 - 3i, 8)$  and  $\mathbf{v} = (-i, 3 + 2i, 1 - i)$ , then

$$\mathbf{v}^H \mathbf{u} = \begin{bmatrix} i & 3 - 2i & 1 + i \end{bmatrix} \begin{bmatrix} 2 + i \\ 1 - 3i \\ 8 \end{bmatrix} = 4 - i,$$

but

$$\mathbf{u}^H \mathbf{v} = \begin{bmatrix} 2 - i & 1 + 3i & 8 \end{bmatrix} \begin{bmatrix} -i \\ 3 + 2i \\ 1 - i \end{bmatrix} = 4 + i.$$

That is, the result of the one is always the complex conjugate of the other. It is a matter of preference which definition one uses, and I opt for this one just to be consistent with Strang's usage.

- When the entries of  $\mathbf{z}, \mathbf{w}$  are all real numbers (that is,  $\mathbf{z}, \mathbf{w} \in \mathbb{R}^n$ ), this new understanding for inner product exactly matches the dot product—that is,  $\langle \mathbf{z}, \mathbf{w} \rangle = \mathbf{z} \cdot \mathbf{w}$ .
- The inner product of vectors in  $\mathbb{C}^n$  no longer exclusively produces real numbers, as seen in the example above. However, when taking an inner product of  $\mathbf{z} = (z_1, z_2, \dots, z_n)$  with itself, the result

$$\langle \mathbf{z}, \mathbf{z} \rangle = \mathbf{z}^H \mathbf{z} = \sum_{j=1}^n \overline{z_j} z_j = \sum_{j=1}^n |z_j|^2,$$

is the sum of the moduli of the components of  $\mathbf{z}$ , guaranteed to be nonnegative. Thus, we define *length* for vectors  $\mathbf{z}$  in  $\mathbb{C}^n$  to be

$$\|\mathbf{z}\| = \sqrt{\langle \mathbf{z}, \mathbf{z} \rangle},$$

and note that the only instance in which  $\|\mathbf{z}\| = 0$  is when  $\mathbf{z}$  is, itself, the zero vector.

The essential list of properties that the inner product in  $\mathbb{C}^n$  has, for all vectors  $\mathbf{u}, \mathbf{v}, \mathbf{w} \in \mathbb{C}^n$  and all scalars  $c$ , is

- (i)  $\langle \mathbf{v}, \mathbf{v} \rangle \geq 0$ .
- (ii)  $\langle \mathbf{v}, \mathbf{v} \rangle = 0$  implies  $\mathbf{v} = \mathbf{0}$ .
- (iii)  $\langle \mathbf{u} + \mathbf{w}, \mathbf{v} \rangle = \langle \mathbf{u}, \mathbf{v} \rangle + \langle \mathbf{w}, \mathbf{v} \rangle$ .
- (iv)  $\langle c\mathbf{u}, \mathbf{v} \rangle = \bar{c} \langle \mathbf{u}, \mathbf{v} \rangle$ .
- (v)  $\langle \mathbf{u}, \mathbf{v} \rangle = \overline{\langle \mathbf{v}, \mathbf{u} \rangle}$ .

Note that (iii) and (v) together imply that

$$\langle \mathbf{u}, \mathbf{v} + \mathbf{w} \rangle = \langle \mathbf{u}, \mathbf{v} \rangle + \langle \mathbf{u}, \mathbf{w} \rangle,$$

while (iv) and (v) together give that

$$\langle \mathbf{u}, c\mathbf{v} \rangle = c \langle \mathbf{u}, \mathbf{v} \rangle.$$

### Conjugate transpose of a matrix

Suppose, now, that  $\mathbf{A}$  is an  $m$ -by- $n$  matrix whose entries are complex numbers. The idea of a *conjugate transpose*  $\mathbf{A}^H$  makes sense, as it did for vectors in  $\mathbb{C}^n$ . In fact, it is precisely what is computed when the *prime* symbol is invoked in OCTAVE.

```
octave:127> A = randi(7,3,2)-4*ones(3,2) + (randi(5,3,2)-4*ones(3,2))*i
A =

-2 - 3i   2 + 1i
 0 - 1i  -1 - 1i
 2 - 2i   1 + 1i

octave:128> A'
ans =

-2 + 3i   0 + 1i   2 + 2i
 2 - 1i  -1 + 1i   1 - 1i
```

Of course, taking the conjugate transpose of a matrix twice returns one to the original:  $(\mathbf{A}^H)^H = \mathbf{A}$ . For complex matrices  $\mathbf{A}$ ,  $\mathbf{B}$  of appropriate size, one can take the conjugate transpose of the product  $\mathbf{AB}$ . If we denote the matrix full of conjugates of entries found in  $\mathbf{A}$  by  $\overline{\mathbf{A}}$ , then we have

$$(\mathbf{AB})^H = (\overline{\mathbf{AB}})^T = (\overline{\mathbf{A}} \overline{\mathbf{B}})^T = \overline{\mathbf{B}}^T \overline{\mathbf{A}}^T = \mathbf{B}^H \mathbf{A}^H.$$

In particular, the conjugate transpose of a matrix-vector product  $\mathbf{Av}$  is

$$(\mathbf{Av})^H = \mathbf{v}^H \mathbf{A}^H,$$

and if we need to take an inner product between  $\mathbf{A}\mathbf{u}$  and  $\mathbf{v}$ , we have the convenient formula

$$\langle \mathbf{A}\mathbf{u}, \mathbf{v} \rangle = (\mathbf{A}\mathbf{u})^H \mathbf{v} = \mathbf{u}^H \mathbf{A}^H \mathbf{v} = \langle \mathbf{u}, \mathbf{A}^H \mathbf{v} \rangle. \quad (2)$$

Formula (2) holds whether or not  $\mathbf{A}$  is square.

Note that  $\mathbf{A}^H$  is sometimes called the **adjoint** matrix (as opposed to the *conjugate transpose*). If the entries in  $\mathbf{A}$  are real numbers only, then  $\mathbf{A}^H = \mathbf{A}^T$ . Any matrix which satisfies  $\mathbf{A}^H = \mathbf{A}$  (necessarily square) is said to be **self-adjoint**, or **Hermitian**.

## Symmetric Matrices

We use the term **symmetric** to describe a matrix  $\mathbf{A}$  which satisfies  $\mathbf{A}^T = \mathbf{A}$ . A complex symmetric matrix need not be self-adjoint, but a real symmetric matrix is. Symmetric real matrices are very important in applications. They have some very favorable properties, some of which hold for all self-adjoint matrices. Every fact we state for self-adjoint matrices is true of symmetric real matrices.

**Theorem 1:** Eigenvalues of a self-adjoint matrix are real.

**Proof:** To prove this, we note first that any complex number  $z$  can be expressed in the form  $z = a + ib$ ; here  $a, b$  are real numbers, called the *real* and *imaginary* parts of  $z$ , respectively. The number  $z$  is, in fact, *real* precisely when its imaginary part  $b = 0$ . Furthermore, the difference of  $z$  and its conjugate is

$$z - \bar{z} = (a + ib) - (a - ib) = i(2b),$$

which is zero if and only if  $z \in \mathbb{R}$ .

Now, suppose  $(\lambda, \mathbf{v})$  is an eigenpair (with  $\mathbf{v} \neq \mathbf{0}$ ) of a self-adjoint matrix  $\mathbf{A}$ . Consider the quantity  $\lambda \|\mathbf{v}\|^2$ , which may alternatively be expressed as

$$\lambda \langle \mathbf{v}, \mathbf{v} \rangle = \langle \mathbf{v}, \lambda \mathbf{v} \rangle = \langle \mathbf{v}, \mathbf{A}\mathbf{v} \rangle = \langle \mathbf{A}^H \mathbf{v}, \mathbf{v} \rangle = \langle \mathbf{A}\mathbf{v}, \mathbf{v} \rangle = \langle \lambda \mathbf{v}, \mathbf{v} \rangle = \bar{\lambda} \langle \mathbf{v}, \mathbf{v} \rangle.$$

Subtracting the expression at one end from that on the other gives

$$0 = \lambda \langle \mathbf{v}, \mathbf{v} \rangle - \bar{\lambda} \langle \mathbf{v}, \mathbf{v} \rangle = (\lambda - \bar{\lambda}) \langle \mathbf{v}, \mathbf{v} \rangle = (\lambda - \bar{\lambda}) \|\mathbf{v}\|^2.$$

Since  $\|\mathbf{v}\| \neq 0$ , it follows that  $\lambda - \bar{\lambda} = 0$ , which implies  $\lambda \in \mathbb{R}$ . □

As well, self-adjoint matrices generate eigenvectors which are naturally orthogonal.

**Theorem 2:** Eigenvectors corresponding to distinct eigenvalues of a self-adjoint matrix are orthogonal.

Proof: Suppose  $(\mu, \mathbf{u}), (\lambda, \mathbf{v})$  are both eigenpairs of a self-adjoint matrix  $\mathbf{A}$  with  $\mu \neq \lambda$ . By the previous theorem,  $\mu$  and  $\lambda$  are real numbers, so  $\bar{\mu} = \mu$ . We have

$$\begin{aligned} (\mu - \lambda) \langle \mathbf{u}, \mathbf{v} \rangle &= \mu \langle \mathbf{u}, \mathbf{v} \rangle - \lambda \langle \mathbf{u}, \mathbf{v} \rangle = \langle \mu \mathbf{u}, \mathbf{v} \rangle - \langle \mathbf{u}, \lambda \mathbf{v} \rangle = \langle \mathbf{A} \mathbf{u}, \mathbf{v} \rangle - \langle \mathbf{u}, \mathbf{A} \mathbf{v} \rangle \\ &= \langle \mathbf{u}, \mathbf{A}^H \mathbf{v} \rangle - \langle \mathbf{u}, \mathbf{A} \mathbf{v} \rangle = \langle \mathbf{u}, \mathbf{A} \mathbf{v} \rangle - \langle \mathbf{u}, \mathbf{A} \mathbf{v} \rangle = 0. \end{aligned}$$

Since  $(\mu - \lambda) \langle \mathbf{u}, \mathbf{v} \rangle = 0$  and  $\mu - \lambda \neq 0$ , it follows that  $\langle \mathbf{u}, \mathbf{v} \rangle = 0$ . □

Now, if  $\mathbf{A}$  is self-adjoint (real and symmetric, for instance) and has  $n$  *distinct* eigenvalues (all real, of course)  $\lambda_1, \dots, \lambda_n$ , then

- the corresponding eigenspaces  $\text{Null}(\mathbf{A} - \lambda_j \mathbf{I})$  are all 1-dimensional (since  $\text{GM} = \text{AM} = 1$  for each eigenvalue) having a single basis eigenvector  $\mathbf{v}_j$ ,
- the vectors  $\{\mathbf{v}_1, \dots, \mathbf{v}_n\}$  are linearly independent (proved earlier), and form a basis of  $\mathbb{C}^n$ , and
- the basis  $\{\mathbf{v}_1, \dots, \mathbf{v}_n\}$  is *orthogonal*, not requiring a Gram-Schmidt process to make it so.

If we choose (or make) the lengths of the eigenvectors be 1, say, setting

$$\mathbf{q}_j = \frac{\mathbf{v}_j}{\|\mathbf{v}_j\|}, \quad j = 1, 2, \dots, n,$$

then the square matrix  $\mathbf{Q}$  having these vectors as its columns

$$\mathbf{Q} = \begin{bmatrix} \mathbf{q}_1 & \mathbf{q}_2 & \cdots & \mathbf{q}_n \\ \downarrow & \downarrow & & \downarrow \end{bmatrix},$$

satisfies the relationship  $\mathbf{Q}^H \mathbf{Q} = \mathbf{I}$ . Such a  $\mathbf{Q}$  is called a **unitary matrix**. Note that  $\mathbf{Q}^{-1} = \mathbf{Q}^H$ . Naturally,  $\mathbf{Q}$  serves to diagonalize  $\mathbf{A}$ , with

$$\mathbf{A} = \mathbf{Q} \mathbf{\Lambda} \mathbf{Q}^{-1} = \mathbf{Q} \mathbf{\Lambda} \mathbf{Q}^H.$$

Note if  $\mathbf{A}$  is real and symmetric, then the entries of  $\mathbf{Q}$  can be taken to be real as well, making  $\mathbf{Q}$  an orthogonal matrix whose columns, simple eigenvectors of  $\mathbf{A}$ , form an orthonormal basis of  $\mathbb{R}^n$ . In that case, the above can be written as

$$\mathbf{A} = \mathbf{Q} \mathbf{\Lambda} \mathbf{Q}^T.$$

Now, if degenerate matrices (those with at least one eigenvalue for which  $GM < AM$ ) are undesirable, here is the really great news. Even though a self-adjoint matrix can have repeated eigenvalues—eigenvalues with  $AM > 1$ —no eigenvalue will have  $GM < AM$ . This fact, which we do not prove at this time, is worthy of a gray box.

**Fact 1:** If  $\mu$  is an eigenvalue of a self-adjoint matrix  $\mathbf{A}$ , then the algebraic multiplicity of  $\mu$  (i.e., the maximum power  $m$  for which  $(\lambda - \mu)^m$  is a factor of the characteristic polynomial  $\det(\mathbf{A} - \lambda\mathbf{I})$ ) is equal to its geometric multiplicity (i.e., the dimension of the eigenspace  $E_\mu = \text{Null}(\mathbf{A} - \mu\mathbf{I})$ ).

For a given eigenvalue  $\lambda$  whose  $AM = k > 1$ , it is still true that  $\text{Null}(\mathbf{A} - \lambda\mathbf{I})$  has many bases. But, using Gram-Schmidt, it is always possible to choose an orthonormal basis  $\{\mathbf{q}_1, \dots, \mathbf{q}_k\}$  of eigenvectors for  $\text{Null}(\mathbf{A} - \lambda\mathbf{I})$ . This set of vectors, put together with orthonormal bases of the other eigenspaces, generate one complete orthonormal basis of  $\mathbb{R}^n$ . These various results are summarized in the Spectral Theorem.

**Theorem 3 (Spectral Theorem for Self-Adjoint Matrices):** Suppose  $\mathbf{A}$  is a self-adjoint matrix. Then there exists an orthonormal basis of  $\mathbb{C}^n$  consisting of eigenvectors  $\{\mathbf{q}_1, \dots, \mathbf{q}_n\}$  of  $\mathbf{A}$ .  $\mathbf{A}$  is, hence, diagonalizable, and the matrix  $\mathbf{Q}$  whose  $j^{\text{th}}$  column is  $\mathbf{q}_j$ ,  $j = 1, \dots, n$ , is a unitary matrix which serves to diagonalize  $\mathbf{A}$ , so that  $\mathbf{A} = \mathbf{Q}\mathbf{\Lambda}\mathbf{Q}^T$ . As a special case, if  $\mathbf{A}$  is real and symmetric, then there exists an orthonormal basis of  $\mathbb{R}^n$  consisting of eigenvectors  $\{\mathbf{q}_1, \dots, \mathbf{q}_n\}$  of  $\mathbf{A}$ . The matrix  $\mathbf{Q}$  formed from these eigenvectors serves to diagonalize  $\mathbf{A}$ , so that  $\mathbf{A} = \mathbf{Q}\mathbf{\Lambda}\mathbf{Q}^T$ .

The result of the spectral theorem can be realized in steps like those outlined in the following example.

### Example 1:

Consider the matrix

$$\mathbf{A} = \begin{bmatrix} 1 & 0 & \sqrt{2} \\ 0 & 2 & 0 \\ \sqrt{2} & 0 & 0 \end{bmatrix}.$$

Since  $\mathbf{A}$  is a symmetric matrix, the Spectral Theorem guarantees  $\mathbb{R}^3$  has a basis which consists of eigenvectors of  $\mathbf{A}$ . The following steps lead to a realization of such a basis.

1. **Find the eigenvalues.** The process is the same as for finding eigenvalues of any square matrix, except we know the results will be real numbers. They are, in fact  $\lambda = -1, 2, 2$  (i.e., 2 is an eigenvalue with  $\text{AM} = 2$ ).
2. **Find bases of the various eigenspaces.** The eigenspace  $E_{-1}$  consists of solutions to  $(\mathbf{A} + \mathbf{I})\mathbf{v} = \mathbf{0}$ . The augmented matrix

$$\left[ (\mathbf{A} + \mathbf{I}) \mid \mathbf{0} \right] \sim \left[ \begin{array}{ccc|c} 1 & 0 & 1/\sqrt{2} & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

has, as expected, one free column, leading to eigenvector  $\mathbf{v}_1 = (-1, 0, \sqrt{2})$ .

The eigenspace  $E_2$  consists of solutions to  $(\mathbf{A} - 2\mathbf{I})\mathbf{v} = \mathbf{0}$ . The augmented matrix

$$\left[ (\mathbf{A} - 2\mathbf{I}) \mid \mathbf{0} \right] \sim \left[ \begin{array}{ccc|c} 1 & 0 & -\sqrt{2} & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{array} \right]$$

has two free columns, matching the algebraic multiplicity of the eigenvalue 2. One can verify that  $\mathbf{v}_2 = (\sqrt{2}, 0, 1)$  and  $\mathbf{v}_3 = (0, 1, 0)$  are independent eigenvectors in  $E_2$ .

3. **For those eigenspaces of dimension  $> 1$ , find orthogonal bases using Gram-Schmidt.** In this instance,  $E_2$  is the only eigenspace of dimension higher than 1. As luck would have it, the basis  $\{\mathbf{v}_2, \mathbf{v}_3\}$  is already orthogonal.
4. **Amass the bases of the various eigenspaces and normalize.** It is already the case that  $\{\mathbf{v}_1, \mathbf{v}_2, \mathbf{v}_3\}$  is a basis of  $\mathbb{R}^3$  consisting of eigenvectors of  $\mathbf{A}$ . The steps prior to this one ensure they form, in fact, an orthogonal basis. We now set

$$\begin{aligned} \mathbf{q}_1 &= \frac{\mathbf{v}_1}{\|\mathbf{v}_1\|} = \frac{1}{\sqrt{3}} \begin{bmatrix} -1 \\ 0 \\ \sqrt{2} \end{bmatrix} = \begin{bmatrix} -1/\sqrt{3} \\ 0 \\ \sqrt{2/3} \end{bmatrix} \\ \mathbf{q}_2 &= \frac{\mathbf{v}_2}{\|\mathbf{v}_2\|} = \frac{1}{\sqrt{3}} \begin{bmatrix} \sqrt{2} \\ 0 \\ 1 \end{bmatrix} = \begin{bmatrix} \sqrt{2/3} \\ 0 \\ 1/\sqrt{3} \end{bmatrix} \\ \mathbf{q}_3 &= \frac{\mathbf{v}_3}{\|\mathbf{v}_3\|} = \mathbf{v}_3 = \begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}. \end{aligned}$$

Then

$$\mathbf{Q} = \begin{bmatrix} \mathbf{q}_1 & \mathbf{q}_2 & \mathbf{q}_3 \\ \downarrow & \downarrow & \downarrow \end{bmatrix} = \begin{bmatrix} -1/\sqrt{3} & \sqrt{2/3} & 0 \\ 0 & 0 & 1 \\ \sqrt{2/3} & 1/\sqrt{3} & 0 \end{bmatrix}$$

is an orthogonal matrix, and

$$\mathbf{A} = \mathbf{Q} \begin{bmatrix} -1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 2 \end{bmatrix} \mathbf{Q}^T = -\mathbf{q}_1 \mathbf{q}_1^T + 2\mathbf{q}_2 \mathbf{q}_2^T + 2\mathbf{q}_3 \mathbf{q}_3^T.$$

■