Poincaré, modified logarithmic Sobolev and isoperimetric inequalities for Markov chains with non-negative Ricci curvature

M. Erbar and M. Fathi, 2018

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Poincaré, modified logarithmic Sobolev and isoperimetric inequalities for Markov chains with non-negative Ricci curvature
as a follow up of
Ricci curvature of finite markov chains via convexity of the entropy

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   geometric structure in the study of Markov semigroups
   (Markov chains)

2. This is done via an interpratation of the semigroup as a
   gradient flow of the Shannon entropy w.r.t. a newly defined
   metric structure

3. Leading to a synthetic notion of Ricci curvature for discrete
   spaces to study contraction properties

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3. Leading to a synthetic notion of Ricci curvature for discrete spaces to study contraction properties

New ingredients (content of the new paper): functional inequalities and applications to bounds on the mixing time of the chain and to its convergence to equilibrium

Road from old paper to present

▶ Geometric structure on space of Markov chains

▶ Synthetic notion of Ricci curvature ($\mathrm{Ric} \geq 0$)

▶ Analogous version of functional inequalities for Markov chains

▶ Bounds on mixing time and rate of convergence to equilibrium
(among others)

Overview of the talk

The zero-range process I

Consider K interacting particles on the complete graph with L sites

The state space is $\mathcal{X}_{K,L} = \{\eta \in \mathbb{N}^L : \sum_i \eta_i = K\}$

The zero-range process with **constant rates** can be described as:
1. Choose a site i uniformly at random
   1.1 If $\eta_i = 0$, do nothing
   1.2 Else, choose another site j uniformly at random and move 1 particle from i to j

$\eta^{i,j}$ denotes the new configuration after such a move
For $\eta, \theta \in \mathcal{X}_{K,L}$,

$$Q_{K,L}(\eta, \theta) = \begin{cases} \frac{1}{L} & \theta = \eta^{i,j} \text{ for some } i, j \\ 0 & \text{else} \end{cases}$$

Stationary measure $\pi_{K,L}$ is the uniform measure

The zero-range process II

Good model for:

- Traffic-jam formation [Kaupužs et al., 2005]
- Population dynamics [Grange, 2020]
- Rewiring networks (e.g. Bianconi–Barabási model) [Grange, 2020], [Angel et al., 2005]
- ... any system in which deciding to move depends on the surroundings

Main result and interpretation

Convergence to equilibrium

The constant-rate zero range process with K particles and L sites has mixing time bounded by

$$\tau_{\mathrm{mix}}(\varepsilon) \leq K^2 L \log L \left(\frac{1}{8} - c \log \varepsilon \right) \qquad (1)$$

for some universal constant c

$\star$ This is exponential convergence to equilibrium! $\star$

Setup

Finite space of states $\mathcal{X}$, irreducible, reversible Markov kernel Q and stationary measure π

$$Q(x,y)\pi(x) = Q(y,x)\pi(y)$$

Operator L acting on functions $\psi : \mathcal{X} \to \mathbb{R}$ is

$$L\psi(x) = \sum_{y \in \mathcal{X}} (\psi(y) - \psi(x))Q(x,y)$$

is the generator of a continuous time Markov chain
Space of probability densities

$$\mathscr{P}(\mathcal{X}) := \left\{ \rho : \mathcal{X} \to \mathbb{R}_+ : \sum_{x \in \mathcal{X}} \rho(x)\,\pi(x) = 1 \right\}$$

Distance $\mathcal{W}$

We define a discrete transport distance, by analogy with the Wasserstein distance. For $\rho_0, \rho_1 \in \mathscr{P}(\mathcal{X})$,

$$\mathcal{W}(\rho_0, \rho_1)^2 := \inf_{\rho, \psi} \frac{1}{2} \int_0^1 \sum_{x,y \in X} (\psi_t(x) - \psi_t(y))^2 \, \hat{\rho}_t(x, y) \, Q(x, y) \, \pi(x) \, dt$$

where the infimum runs over all sufficiently regular curves satisfying a continuity equation

$$\begin{cases} \dfrac{d}{dt}\rho_t(x) + \sum_{y \in \mathcal{X}} \big(\psi_t(y) - \psi_t(x)\big)\hat{\rho}_t(x, y)Q(x, y) = 0, & \forall x \in \mathcal{X}, \\ \rho|_{t=0} = \rho_0, \quad \rho|_{t=1} = \rho_1 \end{cases}$$

and $\hat{\rho}(x, y) := \theta(\rho(x), \rho(y))$ is the logarithmic mean

Markov semigroup as gradient flow

$$\mathcal{W}(\rho_0, \rho_1)^2 := \inf_{\rho, \psi} \frac{1}{2} \int_0^1 \sum_{x,y \in X} \left(\psi_t(x) - \psi_t(y) \right)^2 \hat{\rho}_t(x,y) \, Q(x,y) \, \pi(x) \, dt$$

$$\text{s.t.} \quad \frac{d}{dt} \rho_t(x) + \nabla \cdot (\hat{\rho}(x,y) \nabla \psi(x,y)) = 0$$

Endowing $\mathcal{X}$ with this distance, we can interpret the Markov semigroup $P_t = e^{tL}$ as a gradient flow of Shannon's entropy

$$\mathcal{H}(\rho) := \sum_{x \in \mathcal{X}} \rho(x) \log \rho(x) \pi(x)$$

Ricci for discrete processes

Entropic Ricci curvature for discrete space

$(\mathcal{X}, Q, \pi)$ has entropic Ricci curvature bounded from below by $\kappa \in \mathbb{R}$ if for any geodesic $\{\rho_t\}_{t \in [0,1]}$ on $(\mathscr{P}(\mathcal{X}), \mathcal{W})$ we have

$$\mathcal{H}(\rho_t) \leq (1-t)\mathcal{H}(\rho_0) + t\mathcal{H}(\rho_1) - \frac{\kappa}{2}t(1-t)\mathcal{W}(\rho_0, \rho_1)^2$$

In this case, we write $\mathrm{Ric}(\mathcal{X}, Q, \pi) \geq \kappa$

Ricci for the zero-range process (in brief)

The zero range process with constant rates has $\mathrm{Ric} \geq 0$

In another contribution, M. Fathi and J. Maas provide explicit ways of computing bounds on Ricci for discrete processes

They show that for the ZRP with *increasing* rates, Ricci is positive

This easily extends to showing non-negative Ricci for *constant* rates

Overview of the talk

The modified log-Sobolev inequality I

Modified logarithmic Sobolev inequality

It is a **bound on the entropy** in terms of a norm of the gradient of our observable:

$$\mathcal{H}(\rho) \leq \frac{1}{2\lambda}\|\nabla \log \rho\|_\rho^2 \qquad \text{(mLSI}(\lambda))$$

where $\|\nabla \log \rho\|_\rho^2 =: \mathcal{I}(\rho)$ is the Fisher information

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mLSI for $\mathrm{Ric} \geq 0$

If $\mathrm{Ric}(\mathcal{X}, Q, \pi) \geq 0$ and the diameter of $(\mathcal{X}, d_{\mathcal{W}})$ is bounded by D, then the modified logarithmic Sobolev inequality $(\text{mLSI}(\lambda))$ holds with constant

$$\lambda = \frac{c}{D^2}$$

for some universal constant c

The modified log-Sobolev inequality II

The mLSI has many implications on convergence property, spectrum of the Laplacian and hypercontractivity

$\star$ Our focus: mLSI implies entropy decay, which allows to bound the mixing time $\star$

$$\mathcal{H}(P_t\rho) \leq e^{-2\lambda t}\mathcal{H}(\rho) \qquad \text{(Entropy decay)}$$

Overview of the talk

From mLSI to entropy decay

$$\mathcal{H}(\rho) \leq \frac{1}{2\lambda}\mathcal{I}(\rho) \quad \rightarrow \quad \mathcal{H}(P_t\rho) \leq e^{-2\lambda t}\mathcal{H}(\rho)$$

When P_t is the heat semigroup,

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}(P_t\rho) = -\mathcal{I}(P_t\rho) \overset{\mathrm{mLSI}}{\leq} -2\lambda\mathcal{H}(P_t\rho)$$

From mLSI to entropy decay

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When P_t is the heat semigroup,

$$\frac{\mathrm{d}}{\mathrm{d}t}\mathcal{H}(P_t\rho) = -\mathcal{I}(P_t\rho) \overset{\mathrm{mLSI}}{\leq} -2\lambda\mathcal{H}(P_t\rho)$$

Using Gronwall's inequality, this yields an estimate on entropy decay

$$\mathcal{H}(P_t\rho) \leq e^{-2\lambda t}\mathcal{H}(\rho)$$

that quantifies convergence to equilibrium

From entropy decay to bound on mixing time

$$\mathcal{H}(P_t\rho) \leq e^{-2\lambda t}\mathcal{H}(\rho) \quad \rightarrow \quad \tau_{\mathrm{mix}} \leq \frac{1}{8}D^2 - \frac{\log\varepsilon}{\lambda}$$

Total variation mixing time of a Markov chain

Let P_t be the Markov semigroup on the space of densities $\mathscr{P}(\mathcal{X})$.
For $\varepsilon > 0$,

$$\tau_{\mathrm{mix}}(\varepsilon) := \inf\left\{t > 0 : \|P_t\rho - 1\|_{TV} < \varepsilon \quad \forall x \in \mathcal{X}\right\}$$

From entropy decay to bound on mixing time

$$\mathcal{H}(P_t\rho) \le e^{-2\lambda t}\mathcal{H}(\rho) \quad \to \quad \tau_{\mathrm{mix}} \le \frac{1}{8}D^2 - \frac{\log\varepsilon}{\lambda}$$

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$\to$ Relation to entropy: Pinsker's inequality

Let ρ, σ be two probability densities over $\mathcal{X}$. Then

$$\|\rho - \sigma\|_{TV}^2 \le \frac{1}{2}\mathcal{H}(\rho|\sigma) \qquad \text{(Pinsker)}$$

From entropy decay to bound on mixing time I

Ingredients:

1. Bound on $\mathcal{H}$ in terms of the diameter: $\mathcal{H}(P_t\rho) \leq \frac{\mathcal{W}^2(\rho,1)}{4t}$
2. Pinsker's: $\|P_t\rho - 1\|_{TV}^2 \leq \frac{1}{2}\mathcal{H}(P_t\rho)$
3. Entropy decay: $\mathcal{H}(P_t\rho) \leq e^{-2\lambda t}\mathcal{H}(\rho)$

From (1) we obtain $\mathcal{H}(P_t\rho) \leq 2$ for

$$t \geq \frac{D^2}{8} \quad := t_0 \qquad \qquad \text{(bound \#1)}$$

Plug this into (2) and combine with (3) to obtain

$$\|P_t\rho - 1\|_{TV} \leq \sqrt{\frac{1}{2}\,\mathcal{H}(P_t\rho)} \leq \sqrt{\frac{1}{2}\,e^{-2\lambda t}\,\mathcal{H}(P_{t_0}\rho)} \leq e^{-\lambda t}$$

We get ε-closeness for

$$t \geq -\frac{\log \varepsilon}{\lambda} \qquad \qquad \text{(bound \#2)}$$

From entropy decay to bound on mixing time II
Ingredients:

1. Bound on $\mathcal{H}$ in terms of the diameter: $\mathcal{H}(P_t\rho) \leq \frac{\mathcal{W}^2(\rho,1)}{4t}$
2. Pinsker's: $\|P_t\rho - 1\|^2_{TV} \leq \frac{1}{2}\mathcal{H}(P_t\rho)$
3. Entropy decay: $\mathcal{H}(P_t\rho) \leq e^{-2\lambda t}\mathcal{H}(\rho)$

Sum the two bounds $t \geq D^2/8$ and $t \geq -\log\varepsilon/\lambda$ to obtain an upper bound on the mixing time

$$\tau_{\text{mix}} \leq \frac{1}{8}D^2 - \frac{\log\varepsilon}{\lambda}$$

Bound on the mixing time of the zero-range process

$$\tau_{\mathrm{mix}} \leq \frac{1}{8}D^2 - \frac{\log \varepsilon}{\lambda}$$

Recall: The ZRP has $\mathrm{Ric}(\mathcal{X}_{KL}, Q_{KL}, \pi_{KL}) \geq 0$, hence a mLSI with constant $\lambda = \frac{c}{D^2}$ holds

Moreover: There exists a constant $c > 0$ such that for any L, K

$$D_{KL} \leq cK\sqrt{L \log L}$$

Convergence of the ZRP

The total variation mixing time of the zero range process on $\mathcal{X}_{K,L}$ with constant uniform rates is

$$\tau_{\mathrm{mix}}(\varepsilon) \leq K^2 L \log L \left(\frac{1}{8} - c \log \varepsilon \right)$$

for some universal constant c

References I

📄 Andjel, E., Ferrari, P., Guiol, H., and Landim *, C. (2000).
Convergence to the maximal invariant measure for a
zero-range process with random rates.
Stochastic Processes and their Applications, 90(1):67–81.

📄 Angel, A., Evans, M., Levine, E., and Mukamel, D. (2005).
Critical phase in nonconserving zero-range processes and
rewiring networks.
*Physical review. E, Statistical, nonlinear, and soft matter
physics*, 72:046132.

📄 Brannan, M., Gao, L., and Junge, M. (2020).
Complete logarithmic sobolev inequalities via ricci curvature
bounded below.

📄 Caputo, P., Pra, P. D., and Posta, G. (2007).
Convex entropy decay via the bochner-bakry-emery approach.

References II

Erbar, M., Henderson, C., Menz, G., and Tetali, P. (2017).
Ricci curvature bounds for weakly interacting Markov chains.
Electronic Journal of Probability, 22(none):1 – 23.

Grange, P. (2020).
Non-conserving zero-range processes with extensive rates
under resetting.

Kaupužs, J., Mahnke, R., and Harris, R. J. (2005).
Zero-range model of traffic flow.
Phys. Rev. E, 72:056125.

Münch, F. (2023).
Ollivier curvature, isoperimetry, concentration, and log-sobolev
inequalitiy.

Discussion

▶ The main contribution of this paper for the ZRP is a *quantification* of its convergence ([Andjel et al., 2000])

▶ The dependence on the diameter is optimal, since it is sharp (up to the values of the constant) for the random walk on the one-dimensional discrete torus

▶ However, the mLSI constant for ZRP is not believed to be optimal: dependence of curvature on D in high dimensions is not optimal (Gaussian concentration instead)

▶ Developments in this sense: [Münch, 2023], [Brannan et al., 2020]

Some examples of other applications

mLSI constant (and Poincaré) and implied convergence rate for:

- ▶ Bernoulli-Laplace model (diffusion of two incompressible gases)
- ▶ Random transposition model (random permutations) [Caputo et al., 2007]
- ▶ Hard-core models (nearest-neighbour exclusion)
- ▶ Glauber dynamics (Ising model) [Erbar et al., 2017]

Additional formulas I

Distance $d_{\mathcal{W}}$ on $\mathcal{X}$

The distance $\mathcal{W}$ on $\mathscr{P}(\mathcal{X})$ induces a distance on $\mathcal{X}$ by restricting to Dirac masses, i.e. for $x, y \in \mathcal{X}$

$$d_{\mathcal{W}}(x, y) := \mathcal{W}(\delta_x, \delta_y)$$

Induced Riemannian structure on $\mathcal{X}$

$$(\Psi, \Phi)_\rho := \frac{1}{2} \sum_{x,y} \Psi(x, y) \Phi(x, y) \hat{\rho}(x, y) Q(x, y) \pi(x)$$

Additional formulas II

Discrete gradient and divergence

$$\nabla\psi(x,y) = \psi(y) - \psi(x)$$

$$\nabla \cdot (\nabla\psi)(x) = \tfrac{1}{2} \sum_y \big(\nabla\psi(x,y) - \nabla\psi(y,x)\big)Q(x,y)$$

$$= \tfrac{1}{2} \sum_y \big(\psi(y) - \psi(x) - \psi(x) + \psi(y)\big)Q(x,y)$$

$$= \sum_y (\psi(y) - \psi(x))Q(x,y)$$

Additional formulas III

Heat flow as gradient flow of entropy

Heat equation $\rho_t' = \Delta \rho_t = \nabla \cdot \nabla \rho_t$

Continuity equation $\rho_t' + \nabla \cdot (\hat{\rho} \nabla \psi) = 0$

The heat equation can be re-written as a continuity equation if
$\nabla \psi_t = -\frac{\nabla \rho_t}{\hat{\rho}_t}$

The gradient of the entropy is $\text{grad}_{\mathcal{W}} \mathcal{H}(\rho_t) = \nabla \log \rho_t$

Hence, we get that the heat flow is the gradient flow of the entropy if

$$\nabla \psi_t = -\frac{\nabla \rho_t}{\hat{\rho}_t} = -\nabla \log \rho_t;$$

which gives

$$\frac{\nabla \rho_t}{\hat{\rho}_t} = \nabla \log \rho_t \tag{2}$$

i.e. precisely that $\hat{\rho}_t$ is the logarithmic mean