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## Bioverse: The Habitable Zone Inner Edge Discontinuity as an Imprint of Runaway Greenhouse Climates on Exoplanet Demographics

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#### ABSTRACT

Long-term magma ocean phases on rocky exoplanets orbiting closer to their star than the runaway greenhouse threshold – the inner edge of the classical habitable zone – may offer insights into the physical and chemical processes that distinguish potentially habitable worlds from others. Thermal stratification of runaway planets is expected to significantly inflate their atmospheres, potentially providing observational access to the runaway greenhouse transition in the form of a "habitable zone inner edge discontinuity" in radius—density space. Here, we use Bioverse, a statistical framework combining contextual information from the overall planet population with a survey simulator, to assess the ability of ground- and space-based telescopes to test this hypothesis.

We find that the demographic imprint of the runaway greenhouse transition is likely detectable with high-precision transit photometry for sample sizes  $\gtrsim 100$  planets if at least  $\sim 10\,\%$  of those orbiting closer than the habitable zone inner edge harbor runaway climates. Our survey simulations suggest that in the near future, ESA's PLATO mission will be the most promising survey to probe the habitable zone inner edge discontinuity. We determine survey strategies that maximize the diagnostic power of the obtained data and identify as key mission design drivers: 1. A follow-up campaign of planetary mass measurements and 2. The fraction of low-mass stars in the target sample. Observational constraints on the runaway greenhouse transition will provide crucial insights into the distribution of atmospheric volatiles among rocky exoplanets, which may help to identify the nearest potentially habitable worlds.

## 1. INTRODUCTION

Despite recent advancements in observational tech-26 niques, our understanding of terrestrial-sized planets re-27 mains woefully limited, with fundamental aspects of 28 their nature, composition, and potential habitability 29 still largely unknown. Due to the inherent biases of 30 current exoplanet detection techniques, the best-studied 31 category of rocky exoplanets at present is that of hot or <sup>32</sup> warm, close-in planets (Jontof-Hutter 2019; Bean et al. 33 2021). These experience thermal states that are in some 34 aspects comparable to the ones of the inner solar sys-35 tem bodies at early stages of their evolution (Ikoma 36 et al. 2018; Chao et al. 2021), which likely profoundly 37 affected the distribution of volatiles between planetary 38 core, mantle, and atmosphere. Studying the geophysical 39 state of hot exoplanets can thus inform our understand-40 ing of the early evolutionary stages of Earth and other 41 habitable worlds (Lichtenberg et al. 2022; Krijt et al. 42 2022).

An example from the solar system for the potential significance of these early stages are the divergent atmo-

45 spheric evolutions of Venus and Earth (e.g., Kane et al. 46 2019, 2021; Salvador et al. 2023). While having accreted 47 from a similar mass reservoir (Raymond et al. 2020; 48 Kleine et al. 2020; Mezger et al. 2020; Zahnle & Carl-49 son 2020) and despite their similar bulk properties (Sm-50 rekar et al. 2018), they evolved into planets with very 51 different surface conditions (Donahue et al. 1982; Kast-52 ing 1988; Hamano et al. 2013; Kane et al. 2014; Way & 53 Del Genio 2020; Turbet et al. 2021). Both planets likely 54 underwent a giant impact phase (Raymond et al. 2020; 55 Gillmann et al. 2020; Liu et al. 2022) that melted their 56 mantles (Elkins-Tanton 2012; Schaefer & Elkins-Tanton 57 2018; Lichtenberg et al. 2022). Magma ocean states play 58 a substantial role in establishing the long-term geophys-59 ical and climatic regimes of rocky planets (Fegley et al. 60 2020), in particular owing to efficient heat and volatile 61 transfers between interior and atmosphere in the ab-62 sence of a stiff boundary separating them (Kite & Schae-63 fer 2021; Dorn & Lichtenberg 2021; Salvador et al. 2023). 64 Due to these similar formation sequences, it was com-65 monly assumed that the divergence of Venus and Earth <sub>66</sub> – in particular Venus' water loss – occurred late in their evolution (e.g., Way & Del Genio 2020).

Yet, Hamano et al. (2013) suggested that the present-69 day dry conditions on Venus may have been directly in-70 herited from the early magma ocean stage. If a strongly 71 infrared-absorbing, condensable species such as water 72 was dominant in the atmosphere, the resulting strong 73 thermal blanketing effect would prevent the planet to 74 efficiently radiate to space and maintain the surface 75 molten (Ingersoll 1969; Kasting 1988; Pierrehumbert 76 2010; Goldblatt et al. 2013; Leconte et al. 2015; Salvador 77 et al. 2017). This runaway greenhouse state can extend 78 the magma ocean stage to hundreds of Myr (Schaefer 79 et al. 2016; Barth et al. 2021), enough to remove the 80 entire water reservoir from a rocky planet by H<sub>2</sub>O pho-81 tolysis and subsequent hydrodynamic escape of hydro-82 gen (Wordsworth & Pierrehumbert 2013, 2014; Luger 83 & Barnes 2015). For Venus, the atmospheric composi-84 tion (Gillmann et al. 2020) and comprehensive analysis 85 of meteoritic samples across the Solar System (Alexan-86 der et al. 2018; Broadley et al. 2022) suggest that it 87 went through this phase. Although the past presence or 88 absence of a Venusian water ocean has not been defi-89 nitely established (Raymond et al. 2006, 2007; Hamano 90 et al. 2013; Way et al. 2016; Kane et al. 2019, 2021; 91 Turbet et al. 2021; Warren & Kite 2023), a transient 92 habitable phase cannot be conclusively ruled out (e.g., 93 Way et al. 2016; Salvador et al. 2017; Krissansen-Totton 94 et al. 2021).

The runaway greenhouse transition is a robust predic-<sup>96</sup> tion from climate models (Kasting 1988; Nakajima et al. 97 1992; Goldblatt & Watson 2012; Forget & Leconte 2014; 98 Boukrouche et al. 2021; Chaverot et al. 2022), and its 99 impact on planetary bulk properties has been shown to 100 be in the detectable range of current astronomical in-101 strumentation (Goldblatt 2015). In particular, planets in a runaway greenhouse state are expected to be thermally inflated (Turbet et al. 2019, 2020; Mousis et al. 104 2020), which directly increases their transit radii by an amount that is a function of the water content. However, dissolution of water (e.g., Elkins-Tanton & Seager 2008; 107 Hier-Majumder & Hirschmann 2017; Salvador & Samuel 108 2023) in the magma may decrease this effect (Dorn 109 & Lichtenberg 2021), and chemical exchange between 110 core and mantle material may influence the amount and 111 speciation of outgassed volatiles that are visible in the 112 atmosphere via transmission spectroscopy (Lichtenberg 113 2021; Schlichting & Young 2022). Astronomical ob-114 servations of planets that are currently in a runaway 115 greenhouse state may thus constrain properties of their 116 mantles and establish an observational connection between exoplanetary interiors and atmospheres (Lichtenberg et al. 2022; Wordsworth & Kreidberg 2022).

Of particular relevance is that the radiation-induced 120 transition between a runaway greenhouse state and 121 a temperate climate is thought to occur at a rela-122 tively sharp instellation threshold (Goldblatt et al. 2013; 123 Leconte et al. 2013a; Kopparapu et al. 2013). Conse-124 quently, the instellation at which the runaway green-125 house transition occurs is aptly considered to be the 126 inner boundary of the habitable zone (e.g., Ramirez 127 2018; Salvador et al. 2023). Its prevalent definition 128 refers to the possibility of a sustained liquid water 129 body on the surface of an Earth-like planet with an 130 oxidized CO<sub>2</sub>/H<sub>2</sub>O/N<sub>2</sub>-rich atmosphere (Kasting et al. 131 1993; Kopparapu et al. 2013, 2014); its exact spatial 132 location and extent may be strongly influenced by the 133 interior and atmosphere oxidation state and resulting atmosphere composition (Pierrehumbert & Gaidos 2011; 135 Ramirez & Kaltenegger 2017, 2018; Katyal et al. 2019; 136 Graham & Pierrehumbert 2020; Graham et al. 2022; 137 Hakim et al. 2023). The fundamental concept of a hab-138 itable zone dates back centuries (Newton 1687; Whewell 139 1858; Shapley 1953; Huang 1959), and its modern form 140 has proven popular in the planetary literature<sup>1</sup>. How-141 ever, it should be emphasized that the habitable zone, 142 as it stands today, is merely a concept based on the-143 oretical predictions and geochemical evidence from one 144 planet – Earth (Catling & Zahnle 2020) – and its general 145 validity remains controversial (e.g., Cockell et al. 2016; 146 Moore et al. 2017; Tuchow & Wright 2023). The ques-147 tion naturally emerges if a planetary habitable zone – in 148 its common form with boundaries defined by stellar irra-149 diation – is a predictive theoretical concept and how the 150 habitable zone hypothesis can be tested observationally. Observational tests being considered include searches 152 for direct evidence of liquid water conveyed by ocean 153 glint (Williams & Gaidos 2008; Robinson et al. 2010; 154 Lustig-Yaeger et al. 2018) or water vapor in planetary 155 atmospheres (Suissa et al. 2020). A different approach 156 relies on comparative planetology: aiming for a statisti-157 cal detection in a planet population provides robustness 158 against ambiguity that could otherwise arise from indi-159 vidual variations in a planet's composition or geophys-160 ical history (Checlair et al. 2019; Apai et al. 2019b). 161 Tests suggested in the literature include determining, 162 for a range of orbital distances, atmospheric H<sub>2</sub>O and 163 CO<sub>2</sub> abundances, planetary albedos (Bean et al. 2017;

 $<sup>^1</sup>$  At the time of writing, a search of the term "habitable zone" in the titles and abstracts of refereed articles in the National Aeronautics and Space Administration (NASA) Astrophysics Data System returned  $\sim 1700$  results.

 $^{164}$  Bixel & Apai 2021), or colors (Crow et al. 2011; Bixel  $^{165}$  & Apai 2020), or testing the relationship between  $^{166}$  partial pressure and incident flux (Lehmer et al. 2020).  $^{167}$  All these tests require surveying a large enough sample  $^{168}$  of terrestrial-sized planets with next-generation instruments, rendering them out of reach in the immediate  $^{170}$  future.

Here, we explore the feasibility of a statistical test 171 172 of the habitable zone hypothesis by surveying plane-173 tary bulk properties close to its inner edge, the runaway 174 greenhouse transition. Our goal is to assess the ability of near-future transit surveys to test the hypothesis that 176 the runaway greenhouse effect causes a discontinuity of 177 planetary radii and bulk densities when ordered by re-178 ceiving instellation (Turbet et al. 2019). Our main tool 179 for this is Bioverse, a simulation framework for assess-180 ing the statistical power of exoplanet surveys (Bixel & 181 Apai 2021). It consists of a sample generator that populates stars from the Gaia catalog (Hardegree-Ullman et al. 2023; Prusti et al. 2016; Gaia Collaboration et al. <sup>184</sup> 2022a) with planetary systems based on state-of-the-art occurrence rates (Bergsten et al. 2022), a flexible survey 186 simulator that allows for a broad range of trade studies, and a hypothesis testing module that quantifies the 188 survey's ability to detect a previously injected trend. 189 Trying different instrumentation and survey designs, we 190 use Bioverse to recover runaway greenhouse-induced 191 effects based on model predictions (Turbet et al. 2020; 192 Dorn & Lichtenberg 2021) that we inject into a baseline 193 planet population.

In particular, we test the capability of the *PLATO* (PLAnetary Transits and Oscillation of stars, Rauer et al. 2016) mission, which will measure the radii of a large number of terrestrial-sized planets, to detect the radius/density discontinuity and determine its sensitivity ity to model assumptions and fundamental processes. We then perform a parameter study to explore which trades in the survey design of a *PLATO*-like mission maximize its diagnostic power to test runaway greenhouse climate models through the detection of the habitable zone inner edge discontinuity.

We organize the paper as follows: Section 2 introduces the baseline model we use to produce synthetic star and planet samples. In Section 3, we describe the model component that produces runaway greenhouse-induced transit radius changes. Section 4 explains our survey simulations and hypothesis tests. We present our results in Section 5 before interpreting them in Section 6. Finally, we summarize our findings in Section 7.

The goal of this study is to determine – for dif-215 ferent configurations of near-future exoplanet surveys <sub>216</sub> – the confidence level with which the runaway green-217 house threshold can be detected statistically. Our basic 218 methodology was as follows: We expanded the Bioverse 219 framework (Bixel & Apai 2020, 2021)<sup>2</sup> to generate syn-220 thetic samples of stars that host planets according to 221 the observed exoplanet demographics. We then adapted 222 planetary bulk properties as predicted from models of 223 runaway greenhouse atmospheres, simulated observa-224 tions of the planets, and computed Bayesian evidences 225 in favor of a habitable zone inner edge discontinuity 226 (see diagram in Figure 1). In this section, we review 228 the source of the stellar sample, the modeled luminosity 229 evolution, the generation of a synthetic planet sample, 230 and the orbital parameters of the planets. An overview 231 of our key assumptions and model parameters can be 232 found in Table 1.

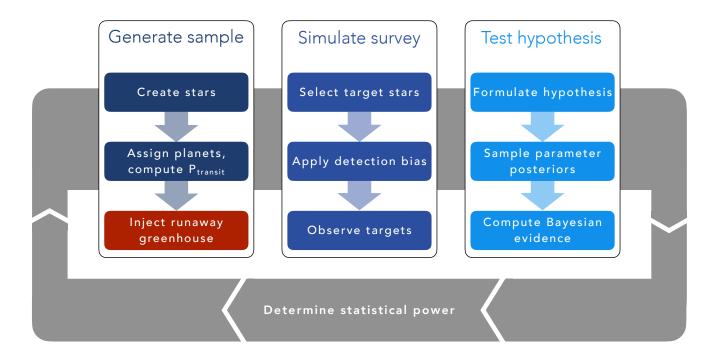
### 2.1. Stellar sample from Gaia DR3

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The original Bioverse stellar catalog was generated randomly from the Chabrier (2003) stellar mass function. Improved parallax and photometric data from the Gaia mission made it possible to generate a homogeneous and complete stellar catalog out to about log for Bioverse (Hardegree-Ullman et al. 2023). Here, we briefly describe how we derived the stellar effective temperature  $T_{\rm eff}$ , luminosity  $L_{\star}$ , stellar radius  $R_{\star}$ , and stellar mass  $M_{\star}$ .

Hardegree-Ullman et al. (2023) used the Gaia Catlateral alogue of Nearby Stars (hereafter GCNS, Smart et al.
lateral 2021) as the basis for deriving stellar parameters for
the Bioverse catalog. The GCNS identified stars out
to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and Gaia Collaboration et al. 2022b) and Gaia DR3 parallaxes and pholateral 2014 parameters from Gaia Collaboration et al. 2022b) and Gaia DR3 parallaxes and pholateral 2014 parameters from Gaia Collaboration et al. 2022b) and Gaia DR3 parallaxes and pholateral 2015 parallaxes and pholateral 2016 parallaxes and pholateral 2016 parallaxes and pholateral 2016 parallaxes and pholateral 2018 parallaxes and phola

<sup>&</sup>lt;sup>2</sup> Bioverse is actively maintained and documented open source software written in Python. Its latest version and documentation can be found at https://github.com/danielapai/bioverse.



**Figure 1.** Workflow of our hypothesis testing with **Bioverse**. First, we generate a sample of stars and populate them with planets based on *Kepler* demographics. A fraction of them are then assigned a runaway greenhouse climate based on the model described in Sect. 3. We then simulate an exoplanet survey, whereby selection effects and detection biases are introduced. Finally, we test the runaway greenhouse hypothesis based on data from the survey simulation. By iterating through these steps, we compute the statistical power of testing the hypothesis for different survey designs.

259 nitudes and a derived bolometric correction. We com-260 puted stellar radii with the effective temperatures and 261 luminosities using the Stefan-Boltzmann law or using 262 absolute  $K_S$ -band magnitudes and an empirical radiusluminosity relation from Mann et al. (2015) for targets within the absolute magnitude range of M dwarfs. Fi-265 nally, we derived masses from the mass-luminosity rela-266 tion of Torres et al. (2010) for stars with  $M_{\star} \gtrsim 0.7 \, M_{\odot}$ , 267 and from that of Mann et al. (2019) for targets within 268 the absolute magnitude range of M dwarfs. The derived 269 stellar parameters were compared to measured param-270 eters for all known exoplanet hosts from the literature 271 and were found to be consistent within 1%, 3%, and 272 5.5% for  $T_{\rm eff}$ ,  $R_{\star}$ , and  $M_{\star}$ , respectively, which are all 273 below the typical measurement uncertainties of 3.3%, 6.8%, and 7.9%, respectively (Hardegree-Ullman et al. 275 2023). From this catalog, Bioverse samples stars within 276 an isotropic distance from the solar system as required 277 by the planetary sample size.

#### 2.2. Stellar luminosity evolution

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Planetary systems are hosted by stars of a wide range of ages, and stellar luminosities evolve with time. Since the occurrence of a runaway greenhouse state is highly

<sup>282</sup> dependent on the amount of radiation received by the planet, and thus on the luminosity of the host star, we assigned age-dependent luminosities to our synthetic stars.

While stellar ages are notoriously poorly constrained (e.g., Adams et al. 2005), the age distribution of planet host stars in the Solar neighborhood was shown to be broadly consistent with uniform (Reid et al. 2007; Gaidos et al. 2023). For our synthetic stars, we thus drew random ages from a uniform distribution from OGyr to 10 Gyr. We then assigned each star a luminosity from the mass-dependent evolutionary models of Baraffe et al. (1998). Figure 2 shows the corresponding luminosity evolution as a function of stellar mass and

#### 2.3. Synthetic planet sample

Next, we assigned to the stellar sample planetary systems with frequencies, orbital parameters, and bulk properties derived from the *Kepler* mission. We adopted the model from Bergsten et al. (2022), which defines the occurrence rate of small planets in radius and orbital period. Following Youdin (2011), their inferred occurrence

Table 1. Key assumptions and model parameters used in our simulation setup

| Parameter                               | Value            | Unit              | Description                                     |
|---|------------------|-------------------|---|
| Stellar sample                          |                  |                   |   |
| $G_{ m max}$                            | 16               |                   | Maximum Gaia magnitude                          |
| $M_{\star,\mathrm{max}}$                | 1.5              | $M_{\odot}$       | Maximum stellar mass                            |
| Luminosity evolution                    |                  |                   | Baraffe et al. (1998)                           |
| Planetary parameters                    |                  |                   |   |
| $M_{ m P}$                              | 0.1 - 2.0        | ${ m M}_{\oplus}$ | Planetary mass range                            |
| $R_{ m P,min}$                          | 0.75             | $R_{\oplus}$      | Minimum planet radius                           |
| Baseline mass-radius relation           |                  |                   | Zeng et al. (2016) $100\% \text{ MgSiO}_3^a$    |
| $\delta_{\min}$                         | 80               | ppm               | Minimum transit depth                           |
| $P_{\max}$                              | 500              | d                 | Maximum orbital period [d]                      |
| S                                       | 10 - 2000        | ${ m Wm^{-2}}$    | Net instellation range                          |
| $S_{ m thresh}$                         | 280              | ${ m Wm^{-2}}$    | Threshold instellation for runaway greenhouse   |
| Runaway greenhouse model                |                  |                   |   |
| Runaway greenhouse atmospheric models   |                  |                   | Turbet et al. (2020); Dorn & Lichtenberg (2021) |
| $x_{H_2O}$                              | $10^{-5} - 0.1$  |                   | Bulk water mass fraction (fiducial case: 0.005) |
| $f_{ m rgh}$                            | 0-1              |                   | Dilution factor (fiducial case: 0.8)            |
| Priors                                  |                  |                   |   |
| $\Pi(S_{	ext{thresh}})$                 | [10, 1000]       | ${ m Wm^{-2}}$    | uniform   |
| $\Pi(x_{H_2O})$                         | $[10^{-5}, 0.1]$ |                   | log-uniform                                     |
| $\Pi(f_{\mathrm{rgh}})$                 | [0, 1]           |                   | uniform   |
| $\Pi(\langle R_{ m P}  angle_{ m out})$ | [0, 15]          | $R_{\oplus}$      | Mean radius of non-runaway planets, uniform     |

 $_{305}$  rate density can be expressed in the form

314

$$\frac{\mathrm{d}^2 n}{\mathrm{d}R \,\partial P} = F_0 C_n g(R, P, M_\star),\tag{1}$$

where  $F_0$  represents the average number of planets per star,  $C_n$  is a normalization constant, and the shape function  $g(R,P,M_\star)$  describes the distribution of planets in radius, orbital period, and stellar host mass. Bioverse generates planets based on the above occurrence rate density and assigns them to the previously generated stars.

### 2.4. Orbit parameters and planet masses

Eccentric orbits alter the probability of a planet to transit (e.g., Barnes 2007). The distribution of eccentricities e of exoplanets has been found to resemble a Beta function (Kipping 2013), which we chose to draw synthetic eccentricities from. Following Kipping (2013), we used a Beta distribution with parameters a=0.867 and b=3.03, and truncated the distribution at e=0.8. <sup>322</sup> Assuming isotropic alignments of orbits, we assigned <sup>323</sup> each planet an inclination drawn from a distribution uni-<sup>324</sup> form in  $\cos(i)$ .

To assign masses to our planets, we use the semigraph assuming a pure  $^{326}$  empirical mass-radius relationship assuming a pure  $^{327}$  MgSiO $_3$  composition from Zeng et al. (2016) (see green  $^{328}$  line in Figure 3). This represents the baseline bulk den- $^{329}$  sity before any climate-related effects are applied.

## 2.5. Transit probability

We model the occurrence of transits by assuming isotropic orientations of planetary orbits and calculating the impact parameters  $b=a\cos(i)/R_{\star}$ . Following the approach in Bixel & Apai (2021), we further consider only planets with |b|<1. For these cases we calculate the transit depth

$$\delta = \left(\frac{R_{\rm P}}{R_{\star}}\right)^2,\tag{2}$$

<sup>&</sup>lt;sup>a</sup>For a comparison with alternative interior compositions, see Appendix A.3.

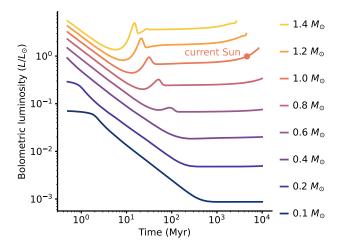


Figure 2. Bolometric luminosity tracks of stars with different masses, computed from stellar evolution models of Baraffe et al. (1998). Low-mass stars, which make up the majority of stars in the solar neighborhood, undergo an extended early phase of several magnitudes higher luminosity before entering a lifetime of relative faintness.

 $^{338}$  which is relevant for the detection probability of the  $^{339}$  respective planet (see Sect. 4.1). Excluding all non-  $^{340}$  transiting planets diminishes the sample to 1.5% of its  $^{341}$  original size.

## 3. RUNAWAY GREENHOUSE MODEL

The climate state of a planet has a direct influ-343 ence on its apparent size measured by transit photome-344 try (Turbet et al. 2019, 2020; Mousis et al. 2020; Aguichine et al. 2021). With even a fraction of the Earth's water inventory, a planet absorbing more flux than the 348 radiation limit of steam atmospheres was found to enter runaway greenhouse state resulting in a global magma ocean (Lichtenberg et al. 2021; Boukrouche et al. 2021, but see Selsis et al. (2023) for a contrasting viewpoint). We use predictions on transit atmospheric thickness 353 from geophysical models to derive the change in transit 354 radius and bulk density that planets with instellation-355 induced runaway greenhouse climates experience, depending on the distribution of water between planetary 357 interior and atmosphere, and on the resulting thermal atmospheric structure (Dorn & Lichtenberg 2021; Salvador & Samuel 2023).

The net absorbed stellar fluxes of planets are a funcintrinsic atmospheric properties such as their albedo, which are generally poorly constrained for planies outside the solar system (e.g., Angerhausen et al. intrinsic 2015; Parmentier & Crossfield 2018; Mansfield et al. intrinsic 2019). For instance, a planet with a high albedo may intrinsic sustain temperate conditions closer to the star than the intrinsic same planet with a lower albedo and located further <sup>368</sup> away from the star. Here, we assume global redistribu-<sup>369</sup> tion of incoming flux and a fixed Bond albedo of 0.3, <sup>370</sup> comparable to Earth's (Haar & Suomi 1971). We do <sup>371</sup> not take into account additional heating sources such as <sup>372</sup> tidal effects (e.g., Barnes et al. 2013).

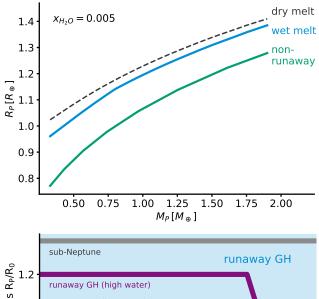
While we search for the signature of runaway greenhouse climates in demographic quantities such as avreage planet radii, the injected changes happen on the
planetary level: We changed each planet's transit radius based on its individual set of properties and the
associated predictions from steam atmosphere and water retention models. Relevant properties are a planet's
mass M, its net instellation S, and its bulk water inventory expressed as the total planetary water mass fraction  $x_{H_2O}$ . We consider the following cases (see Figure 3):

Non-runaway planets retain the radius assigned based on exoplanet occurrence rates (see Sect. 2.3). This case serves as our null hypothesis.

We consider as runaway planets those planets that ab-387 sorb a stellar flux higher than a dayside-averaged thresh- $_{388}$  old instellation  $S_{
m thresh}$ . For all planets absorbing an instellation exceeding  $S_{\rm thresh} = 280\,{\rm W\,m^{-2}}$ , we assume 390 an inflated transit radius due to a steam atmosphere. While the actual instellation threshold for a runaway 392 climate depends on planetary albedo, surface gravity, 393 and clouds (Pluriel et al. 2019; Turbet et al. 2021; Pier-394 rehumbert 2022), this value was found to be a typical 395 limit for the flux a planet can emit in a runaway green-396 house situation (Goldblatt et al. 2013; Kopparapu et al. 397 2013; Leconte et al. 2013a; Hamano et al. 2015; Sal-398 vador et al. 2017; Katyal et al. 2019; Boukrouche et al. 399 2021; Lichtenberg et al. 2021a). This choice translates 400 to about 1.18 times the instellation of present-day Earth 401 with a fixed albedo of 0.3, which compares favourably to 402 previous climate simulations (Leconte et al. 2013a; Wolf 403 & Toon 2015). We adopt the same threshold instellation 404 for all host star spectral types.

To quantify the radius change, we applied the massradius relationships derived by Turbet et al. (2020) using and a 1D inverse radiative-convective model (Turbet et al. 2019). Their calculations rely on the same mass-radius relations for rocky interiors that we apply for our nonrunaway planets (Zeng et al. 2016). For each planet above the instellation threshold, we assigned the predicted radius for the given water mass fraction and planet mass.

Nominally, the above models assume a *dry melt* without dissolved volatiles. Here, however, we consider a \*\*wet melt magma ocean and take into account a radius decrease from retention of water in the melt (Dorn & \*Lichtenberg 2021). The impact of the water distribution between melt and atmosphere on the change of the



Time [Myr]

Top: Mass-radius relationships for a bulk water mass fraction  $x_{H_2O} = 0.005$  and different planet states. Green: planets with a solid mantle and no steam atmosphere. Dashed: planets with steam atmospheres. Blue: planets with steam atmospheres and including the effect of water retention in the melt. Steam atmospheres cause a significant radius increase, which is slightly reduced when water retention in the melt is considered. *Bottom*: Radius evolution of different planet types, illustrating degeneracies and potential for confusion among planet classes. Shown is a schematic time evolution of the transit radius normalized to the atmosphere-free radius  $R_0$  for different scenarios. Planets can move between planet classes through processes such as atmospheric loss and desiccation, which ultimately ends a runaway greenhouse phase on a timescale dependent on a plaret's water content.

transit radius depends on the planet's mass and water content and is generally small compared to the radius inflation from the steam atmosphere. We computed radius deviations between a wet magma ocean and a solid mantle following Dorn & Lichtenberg (2021) and assuming a tropopause pressure  $P_{\rm iso}=0.1\,{\rm bar}$ . We then added the (in almost all cases negative) radius deviations to the inflated planet radii computed for the dry melt case.

We illustrate the mass-radius relations of the three cases in Figure 3 where a fiducial bulk water mass fraction of  $x_{H_2O}=0.005$  is assumed. In the following, we only distinguish between the non-runaway greenhouse and wet melt scenarios.

To account for planets unable to sustain a steam atmo-434 sphere over extended time spans, as well as evolutionary 435 effects such as desiccation through water photodissocia-436 tion and H escape (see bottom panel of Figure 3), we in-437 troduce a dilution parameter  $f_{\rm rgh}$ . It represents the frac-438 tion of planets above the instellation threshold whose 439 atmospheres are currently inflated due to a runaway 440 greenhouse climate. Our simulation setup is such that 441 all planets receiving a net instellation  $S < S_{\rm thresh}$  follow 442 the non-runaway greenhouse relation, and a fraction  $f_{
m rgh}$ 443 of the planets with  $S>S_{
m thresh}$  follow the wet melt rela-444 tion. The choice  $f_{\rm rgh} < 1$  reflects the fact that planets – 445 despite a high irradiation – can evade a runaway green-446 house climate. This situation may, for instance, arise in 447 the absence of an atmosphere or of volatiles that could 448 form a steam atmosphere. In the following, we test if 449 and under what conditions this parametrization causes 450 a demographic trend that is large enough to be detected 451 with high significance.

## 4. EXOPLANET SURVEY SIMULATIONS AND HYPOTHESIS TESTING

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The survey module of Bioverse converts the synthetic planet sample into a set of uncertainty-laden measurements on a subset of that sample. This task includes
selection of the targets, application of detection biases,
and conducting simulated measurements, all of which
are specific to the particular survey. For each planetlevel measurement such as transit radius or instellation,
we draw the measured value from a normal distribution
centered on the true value with a standard deviation set
by the survey's precision. We then follow a Bayesian hypothesis testing approach to assess various realizations
of simulated surveys in terms of their ability to detect
and characterize the runaway greenhouse transition.

#### 4.1. Detection bias, target selection, and sensitivity

Not all transiting planets are detectable with the same likelihood and detection biases have an impact on the demographic measurements we are interested in. A detailed characterization of the detection biases of individual missions would not be justifiable given the uncertainties of the theoretical predictions. Instead, we derived generic observing limits that reflect the limitations of state-of-the-art transit surveys.

A successful transit detection requires a sufficient signal-to-noise ratio, which is sensitive to the achieved

photometric precision. PLATO (PLAnetary Transits and Oscillation of stars) is an ESA mission designed to characterize terrestrial planets in the habitable zones of Sun-like stars via long-term high-precision photometric monitoring of a sample of bright stars (Rauer et al. 2016). In line with this requirement, PLATO is designed to enable the detection of a 80 ppm transit signal (Matuszewski et al. 2023; ESA 2017). To reflect its sensitivity, we chose a minimum transit depth of 80 ppm as a detection limit and consider only measurements of planets ets exceeding this threshold. We further exclude target stars with Gaia magnitudes  $M_{\rm G} > 16$ .

The runaway greenhouse effect becomes obsolete both for very small instellations and where no atmosphere can be maintained due to proximity to the host star and resulting atmospheric erosion. Ensuring to stay well clear of such regions, we clear our sample from all planets with a net instellation  $S < 10\,\mathrm{W\,m^{-2}}$  or  $S > 2000\,\mathrm{W\,m^{-2}}$ . We further consider only rocky planets with masses below  $2\,\mathrm{M_{\oplus}}$ .

#### 4.2. Measurements and their uncertainties

Under real-world conditions, the planetary properties in question can only be probed with a finite precision that is specific to each exoplanet mission. PLATO's definition study report (ESA 2017) states precision requirements for planet radii (3%), planet masses through radial velocity (RV) follow-up (10%), and stellar masses, radii, and ages (10%). We adopted these estimates and assumed a 10% error on instellation measurements.

Since planetary bulk density  $\rho \propto R_{\rm P}^{-3}$ , we expect a stronger runaway greenhouse signal when measured through bulk density instead of transit radius. We thus simulated measurements of planetary densities assuming the mass-radius relation defined above. For uncertainties in bulk density measurements, we propagated the errors of the mass measurements assuming  $\sigma_{\rm M_P} = 10\,\%$ .

#### 4.3. Hypothesis tests

We now turn to quantifying the ability of the simulated surveys to detect the habitable zone inner edge discontinuity and to constrain parameters associated to the runaway greenhouse transition. To do this, we rely on a Bayesian hypothesis testing approach where we quantify the evidence of a hypothesis over another based on the (simulated) data. For our specific problem, this implies comparing evidences for a demographic imprint of the runaway greenhouse effect to its absence. As a null hypothesis, we consider the case where the planetary radius distribution is independent of the instellation,

$$H_0(\theta, S) = \theta, \tag{3}$$

where  $\theta$  is the set of parameters defining the radius distribution. We further define an alternative hypothesis that describes radius changes due to runaway greenhouse climates and inflated steam atmospheres. As motivated above, this hypothesis takes the form of a step function in net instellation S, where the step occurs at the outer edge of the runaway greenhouse region. Our main observable shall be the average transit radius in the planet population on either side of this threshold. The runaway greenhouse hypothesis is then defined as

$$H_{\rm rgh}(\theta, S) = \begin{cases} H_0, & S \le S_{\rm thresh} \\ \langle R_{\rm P} \rangle (f_{\rm rgh}, \Delta R_{\rm stm}, \Delta R_{\rm wtr}), & S > S_{\rm thresh}. \end{cases}$$
(4)

Here,  $f_{\rm rgh}$  is the fraction of planets above the instellation threshold experiencing a runaway greenhouse effect.  $\Delta R_{\rm stm}$  and  $\Delta R_{\rm wtr}$  are predicted radius changes from the steam atmosphere and water retention models, respectively. They are assumed to act additively on the planet radii and thus on their average  $\langle R_{\rm P} \rangle$ .

The only free parameter of the null hypothesis, which assumes the average transit radius to be independent of instellation, is the predicted mean radius  $\langle R_{\rm P} \rangle$ . The functional form of the runaway greenhouse hypothesis is more complex: Besides the mean radius of planets outside the threshold  $\langle R_{\rm P} \rangle_{\rm out}$ , which is a nuisance parameter necessary to define the hypothesis, it relies on the threshold instellation for the "step"  $S_{\rm thresh}$ , the planetary bulk water mass fraction  $x_{H_2O}$ , and the dilution factor  $f_{\rm rgh}$ . For hypothesis tests based on bulk density instead of radius, we proceeded in the same way and substituted  $R_{\rm P}$  by the bulk density  $\rho$ .

A sensible choice of priors is central for evidence estimation via nested sampling. As the parameters of intersest are poorly constrained by previous data, we used relatively uninformative priors to sample the entire physically plausible parameter space. For  $S_{\rm thresh}$ , we chose a uniform prior in  $[10,1000]~{\rm W\,m^{-2}}$ . We sampled  $x_{H_2O}$  from a log-uniform distribution to imply scale-invariant ignorance. Its boundaries  $[10^{-5},0.1]$  are motivated by the water mass fractions covered by the geophysical models (Sect. 3). For  $f_{\rm rgh}$ , we chose a uniform prior for  $\langle R_{\rm P} \rangle_{\rm out}$  bound by  $[0.1,15]~{\rm R}_{\oplus}$ . In the case of measuring bulk densities instead of transit radii, we drew uniformly from  $[1,6]~{\rm g\,cm^{-3}}$ .

The measured radii  $R_{\rm P,i}$  or bulk densities  $\rho_{\rm i}$  cannot be directly used for the hypothesis tests as they include intrinsic scatter that is not caused by measurement errors.  $H_{\rm rgh}$  and  $H_0$  should thus be tested against a statistical estimator that represents the population mean. To avoid binning and the artificial patterns it may introduce, we

576 chose to test our hypotheses against a simple moving SMA along the instellation axis with a window 578 of size 25 centered around each measurement. We fur-579 ther computed the uncertainty of this moving average 580 by propagating the individual measurement errors and applying a rolling standard error of the mean.

As our procedure involves random sampling of the model parameters  $\theta$ , we need to define the probability 584 of obtaining a data set given the model parameters, i.e., 585 a likelihood function  $\mathcal{L}$ . We assumed here that the in-586 dividual moving averages  $SMA_i$  are measured with a normally distributed uncertainty  $\sigma_{SMA_i}$  and adopted a 588 normal distribution

$$\mathcal{L}(SMA \mid \boldsymbol{\theta}) = \prod_{i}^{N} \frac{1}{\sqrt{2\pi\sigma_{SMA_{i}}^{2}}}$$

$$\times \exp\left(-\frac{(SMA_{i} - H(\boldsymbol{\theta}, S_{i}))^{2}}{2\sigma_{SMA_{i}}^{2}}\right).$$
 (6)

$$\times \exp\left(-\frac{(SMA_i - H(\boldsymbol{\theta}, S_i))^2}{2\sigma_{SMA_i}^2}\right).$$
 (6)

Here,  $H(\theta, S_i)$  corresponds to the functional form of the runaway greenhouse or null hypothesis.

## 4.4. Bayesian model comparison

We can now assess the relative plausibility of  $H_{rgh}$ 594 595 and  $H_0$  given the synthetic data we have generated, 596 assigning equal a priori probabilities to these models. This is done by comparing the Bayesian evidence  $\mathcal{Z}$  of the models, which we estimated with the nested sampling (Skilling 2004) algorithm dynesty (Speagle 2020). We initialized the sampler with the priors defined above 601 to let it estimate the evidence and sample the posterior 602 distributions. Our criterion to reject the null hypothesis 603 is

$$\Delta \ln \mathcal{Z} = \ln \mathcal{Z}_{rgh} - \ln \mathcal{Z}_0 > 3. \tag{7}$$

## 5. RESULTS

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## 5.1. Statistical signature of the runaway greenhouse threshold

To characterize the population-level imprint of indi-608 vidual radius changes, we generated a generic planet 610 population with an injected runaway greenhouse effect assuming a water fraction  $x_{\rm H_2O} = 0.005$ . Figure 4 shows 613 the resulting planetary radii. When ordered in orbital 614 period space, the different planet types overlap, dilut-615 ing the demographic imprint. With net instellation as 616 an independent variable, planets above and below the 617 runaway greenhouse threshold separate: The runaway 618 greenhouse-induced radius inflation introduces a discon-619 tinuity of average planet radii and bulk densities as a 620 function of stellar irradiation. We also show the pre-621 dictions of observable average planet radii from the sta-622 tistical hypotheses defined above. Within the runaway  $_{623}$  greenhouse regime, an average radius change of 15 % oc-624 curs. This pattern is consistent with the injected radius 625 inflation as predicted from the atmospheric models (see 626 Appendix A.2 for an investigation of the interplay be-627 tween model predictions and our synthetic planet pop-628 ulation).

#### 5.2. Testability of the runaway greenhouse hypothesis

Figure 5 shows a prototypical statistical detection of 631 the habitable zone inner edge discontinuity. Before we 632 study limiting cases of such a detection below, we first 633 demonstrate the interpretation process based on an op-634 timistic scenario where the sample of characterized plan-635 ets is large (N = 500) and the measurement uncer-636 tainties are small ( $\sigma_R = 2\%, \sigma_S = 5\%$ ). Here, we 637 assumed that the fraction of those planets irradiated  $_{638}$  stronger than  $S_{
m thresh}$  that have runaway greenhouse cli- $_{639}$  mates is  $f_{\mathrm{rgh}}=0.8,$  and we chose a water mass fraction of  $x_{H_2O} = 0.005$  for each planet. In this case, the hab-641 itable zone inner edge discontinuity was detected with <sub>642</sub> high significance ( $\Delta \ln \mathcal{Z} \approx 100$ ).

With such a strong signal, we can attempt an inference of the parameters defining the injected effect. Figure 6 shows the posterior distributions of  $S_{\text{thresh}}$ ,  $x_{H_2O}$ , and  $f_{\rm rgh}$  as determined by the nested sampler. The thresh-647 old instellation can be accurately constrained. Both a 648 higher water mass fraction and a higher dilution factor 649 lead to larger average radii, thus these parameters are 650 strongly correlated.

Figure 7 explores the statistical power of the hypothe-652 sis test achieved in the above scenario for different combinations of the poorly constrained parameters  $x_{H_2O}$ <sub>654</sub> and  $f_{\rm rgh}$ . It is highest for large water inventories and 655 large dilution factors. For all but very low water frac $f_{\rm rgh}$  dominates this trend: It enters linearly into 657 the average planet radius, whereas the contribution of  $x_{H_2O}$  - as predicted by the geophysical models - is sub-659 linear with a power-law exponent of  $\sim 0.3$ . Within the framework of our model and as long as  $f_{
m rgh}$  is larger than  $\sim 0.2$ , a sample size of 500 is sufficient for a 50 % detection rate even for water ratios as low as  $10^{-3}$ .

## 5.3. Detecting the runaway greenhouse transition with PLATO

To simulate the transit survey of ESA's PLATO mis-666 sion, we considered a volume-limited sample with a size 667 according to projections and including all stellar spec-668 tral types. Given PLATO's expected radius precision, we find that a yield of  $\sim 300$  is needed for a significant

<sup>&</sup>lt;sup>3</sup> See Appendix A.1 for a robustness test using a different estima-

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## Synthetic planets

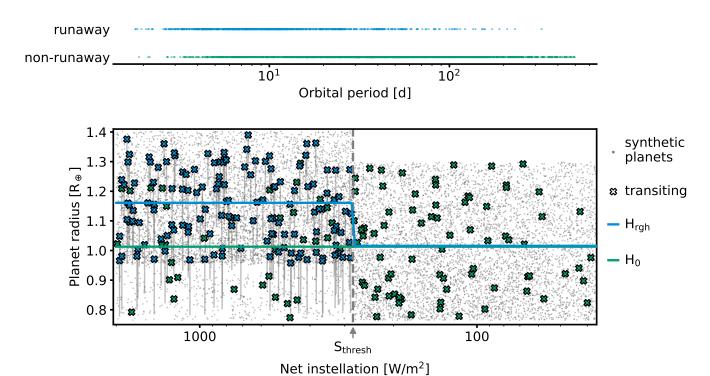


Figure 4. Synthetic planets above and below the runaway greenhouse threshold. Top: Planet state as a function of orbital period. Planets with and without a runaway greenhouse climate mix and are not distinguishable in orbital period space. Bottom: Transit radii of synthetic planets with injected radius deviation as a function of net instellation. Only the planets marked as transiting are observable. Above the runaway greenhouse threshold  $S_{\rm thresh} = 280 \, {\rm W \, m^{-2}}$ , some planets maintain their original radii (green crosses) whereas some have their transit radius inflated (blue crosses) by the amount indicated with gray lines. The sharp boundary at  $S_{\rm thresh}$  causes a discontinuity in the average planet radius (blue line). This runaway greenhouse hypothesis can be tested against the null hypothesis  $H_0$  (green line), where average radii are independent of instellation.

670 detection if the fraction of runaway greenhouse planets 671  $f_{\rm rgh} = 0.1$  (see Figure 8). The minimum needed fraction 672 rises to 0.2 for N = 100. For much smaller samples, only 673 an optimistically strong signal is likely to be detected.

# 5.4. Statistical power of different mission designs 5.4.1. Additional planet mass measurements

Comparing a measurement of the habitable zone informal ner edge discontinuity in radius space with a measureformal ment in density space (which requires planetary mass formal measurements), a stronger detection occurs in the latter formal case: With an optimistic choice of geophysical parameformal ters (see Sect. 5.2), the average measured radius change formal is 15% whereas the average density change is -33%.

Figure 8 shows that investigations of the discontinuity are more constraining when radius measurements can be augmented with mass measurements: At unchanged sample size, the difference in Bayesian evidence can be up to an order of magnitude larger. Consequently, we achieve a statistically significant detection with smaller samples or lower dilution factors  $f_{\rm rgh}$ . A density-based

survey of 100 targets is roughly equivalent to a radiusbased survey of 300 targets. At N=100, pure radius measurements require  $f_{\rm rgh}\gtrsim 0.2$  whereas bulk density measurements enable a detection from  $f_{\rm rgh}\gtrsim 0.1$ .

#### 5.4.2. Dependence on host star spectral type

Since the incident radiation at a given orbital distance depends on the spectral type of the host star, the relative number of planets on either side of  $S_{\rm thresh}$  is different for FGK and M dwarfs. We tested the detectability of the habitable zone inner edge discontinuity when only FGK or only M dwarfs are considered (see Figure 8). All other parameters of the climate models are kept the same. The samples are volume and magnitude-limited to reflect the target counts of PLATO's provisional Long-duration Observation Phase fields (15996 FGK stars in the P1 and P2 samples, 33948 M stars in the P4 sample, Nascimbeni et al. 2022). The resulting M dwarf planet sample is significantly larger with  $228 \pm 14$  planets compared to  $40 \pm 6$  planets in the FGK sample. No significant detection is possible in the pure FGK sample,

#### Optimistic survey (500 planets)

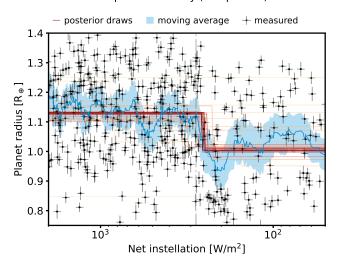


Figure 5. Detection of the runaway greenhouse threshold. From simulated radius and instellation measurements of a large (N=500) survey, we compute the moving average (blue confidence intervals) and fit the runaway greenhouse hypothesis to it (Eqn. 4, random draws from the posterior in  $\mathbb{R}^2$ d). The pattern is detected with high significance.

independent of the assumed geophysical parameters. In the M dwarf sample, the evidence threshold is reached around  $f_{\rm rgh} \sim 0.2$ , similar to the case above where all spectral types are considered.

## 5.4.3. Constraining the threshold instellation

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Although the exact location of the runaway green-716 house transition depends on the stellar energy actually 717 absorbed by the planet, a measurement of  $S_{\rm thresh}$  would 718 be a key constraint resulting from a detection of the 719 habitable zone inner edge discontinuity. Here, we as-720 sess the ability of different mission concepts explored above to constrain this parameter. Figure 9 shows pos-722 terior distributions from a grid of inferences together 723 with the true values of the injected signal for different fractions  $f_{\rm rgh}$ . We consider three cases: only radius 725 measurements, radius and mass measurements, and ra-726 dius and mass measurements of only planets orbiting 727 M dwarfs. The simulations otherwise represent the sim-728 ulated PLATO example described above. We chose a planetary sample size of  $N=100\pm10$ , as this was found 730 to be a threshold case in Sect. 5.3.

We find that retrievals from radius measurements alone require high fractions of greenhouse climate-bearing planets to achieve an accurate constraint on  $S_{\rm thresh}$ : dilution factors  $f_{\rm rgh} \gtrsim 0.5$  yield posterior probability distributions that are condensed at the order of magnitude of the true value; accurate constraints to

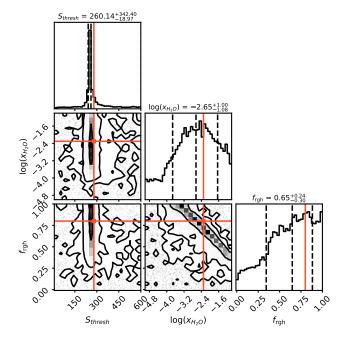


Figure 6. Retrieved posterior distribution of key parameters in the optimistic scenario. The density maps in each panel show relationships between and marginalized distributions of the threshold instellation  $S_{\rm thresh}$ , the bulk water mass fraction  $x_{H_2O}$ , and the dilution factor  $f_{\rm rgh}$  as they could be retrieved with a high-precision transit survey and a sample of 500 planets. True values of the parameters for the injected effect are shown in orange. The threshold instellation can be reasonably constrained; the predominant water fraction and the fraction of planets with runaway greenhouse climates are degenerate.

737 within  $\pm 0.25\,\mathrm{dex}$  of the truth are reached only from 738  $f_\mathrm{rgh} \approx 0.8.$ 

In contrast, if planet masses are available and the hy-740 pothesis test is conducted in bulk density space, useful 741 constraints emerge already from about  $f_{\rm rgh} \approx 0.35$ . Ac-742 curacy and precision of the retrievals are improved.

A sample containing only planets around M dwarf yields still better performance and result in the overall best accuracy and precision. The reason is that these planets contribute most to the statistical power of the retrieval.

## 6. DISCUSSION

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#### 6.1. Statistical imprint of exoplanet climates

We showed in Sect. 5.1 that injecting the theoretically predicted radius inflation effect into a synthetic planet population following the currently known demographics leaves a distinct pattern in the radius and density distribution as a function of instellation. The transition radius inflation was suggested before as an observational diagnostic to probe the runaway greenhouse transity.

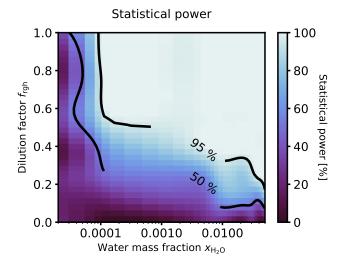


Figure 7. Statistical power of the runaway greenhouse hypothesis test as a function of model parameters. For a sample size N=500, the color code shows the fraction of simulations resulting in a sound detection  $(\Delta \ln \mathcal{Z} > 3)$  for different combinations of bulk water mass fraction and dilution factor. Higher values in either parameter result in a more reliable detection. For water mass fractions  $\gtrsim 10^{-4}$ , the statistical power largely depends on the fraction of greenhouse climate placets in the sample.

757 sition (Turbet et al. 2019), and our simulations show quantitatively how the contributions of the individual r59 radius changes combine to create a significant demographic signal. Its strength depends on largely unknown factors such as the planetary volatile content, but turned out to be well in the detectable range under reasonable assumptions of these factors.

Our finding that the expected habitable zone inner 765 edge discontinuity is strong except for extremely low water mass fractions gives reason for optimism regard-767 ing its detection. It also presents a potential for con-768 straints on the water inventory of terrestrial planets in 769 the case of a non-detection. Such constraints will pro-770 vide insight into whether the initial water content of 771 rocky planets is varied by systematic inter-system effects (e.g., Raymond et al. 2004; Mulders et al. 2015; 773 Sato et al. 2016; Lichtenberg et al. 2019; Lichtenberg 774 & Krijt 2021; Lichtenberg & Clement 2022) or if there 775 is a predominant pattern of volatile-enrichment across 776 planetary systems that is only modified by intra-system 777 effects such as planet migration (Schlecker et al. 2021a) 778 or atmospheric escape (Owen & Wu 2016). Ultimately, 779 such measurements will shed light on how delivery and 780 loss effects during rocky planet formation and evolution 781 shape the diversity of exoplanetary climates. Overall, 782 probing the runaway greenhouse discontinuity appears

783 to be the most promising approach toward a first em-784 pirical test of the habitable zone concept.

## 6.2. Detectability of the habitable zone inner edge discontinuity

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We showed in Sect. 5.2 that under favorable conditions, a sufficiently large (500 planets) photometric survey is likely to detect the demographic imprint of the
runaway greenhouse transition and accurately constrain
the associated threshold instellation. Of course, even assuming capable instrumentation and an optimal survey
design, our current models may not capture fully the
complexity of trends in the geophysics and demographics of small planets. The detectability of the runaway
greenhouse threshold is thus a function of both survey
design and the diverse outcomes of rocky planet formation and evolution. In the following, we will explore
the key factors that may influence the emerging demographic signal.

## 6.2.1. Key factors influencing tests of the runaway greenhouse hypothesis

What influences the probability of correctly rejecting a false null hypothesis? We revidentified six drivers of the diagnostic power for detecting the runaway greenhouse transition with a transit survey. Many of these factors directly influence the fraction of planets currently in runaway greenhouse states, which in our model is represented by the dilution factor  $f_{\rm rgh}$ .

- Occurrence rate of planets forming steam atmospheres
- Planetary evolution and duration of the steam atmosphere phase
- Prevalent water inventory
- Size and composition of the planetary sample
- Radius measurement precision
  - Availability and precision of mass measurements

818 We will now briefly explore the above drivers.

Occurrence rate of planets forming steam atmospheres:
The runaway greenhouse climate relies on sufficient amounts of atmospheric water vapor that can act as a greenhouse gas. It was shown that already about ~10–20 bar of water vapor – corresponding to a mi-nor fraction of one Earth ocean and thus the lower limit of water on Earth – is enough to sustain sufficiently high surface temperatures to keep the planet in a magma ocean stage (Boukrouche et al. 2021; Lichtenberg et al. 2021). However, recent simulations accounting for radiative—convective profiles of near-surface

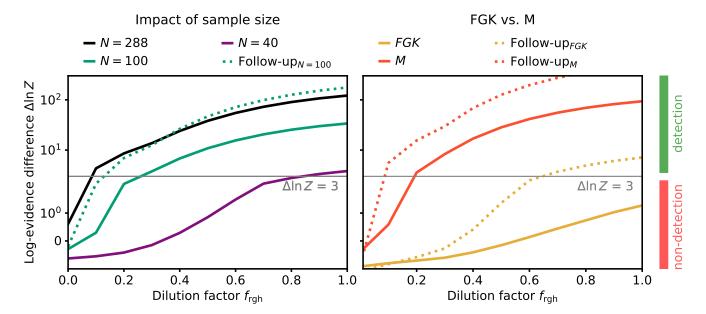
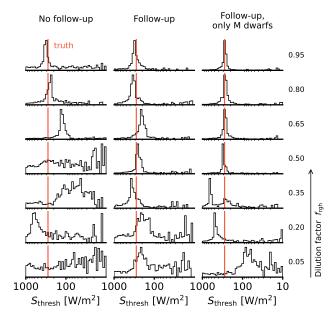


Figure 8. Expected delta-evidences as a function of the fraction of planets with runaway greenhouse climates for different versions of the PLATO survey. The median values of randomized survey simulations are shown;  $\Delta \ln Z > 3$  (gray horizontal line) is considered sufficient evidence to reject the null hypothesis. Left: For a large planet yield of  $N \approx 300$ , even small dilution factors  $\sim 0.1$  allow a detection. A sample of 100 planets is sufficient if their masses are constrained to within 10% (dotted green line). Without such follow-up measurements, sufficient diagnostic power can only be achieved with this sample if  $f_{\rm rgh} \gtrsim 0.2$ . Even smaller samples are unlikely to yield a significant detection. Right: Evidences when only FGK or only M dwarfs are considered. Only M dwarfs host enough planets on both sides of the threshold instellation to allow a reliable detection of the habitable zone inner edge discontinuity.

830 atmospheric layers suggest the existence of cooler pure-831 steam runaway greenhouse states that do not necessar-832 ily vield a molten surface (Selsis et al. 2023). Climate 833 state and atmospheric structure, as well as the frac-834 tion of planets fulfilling the requirements of a steam 835 atmosphere have an impact on the amplitude of the 836 demographic imprint of the runaway greenhouse transition. From a planet formation perspective, the incor-838 poration of water into planets in the terrestrial planet 839 zone is a standard expected outcome (e.g., Zeng et al. 840 2019; Venturini et al. 2020; Emsenhuber et al. 2021; 841 Schlecker et al. 2021a; Burn et al. 2021). But while 842 commonly considered volatile delivery channels suggest fraction of planets to be volatile-poor, the incorpora-844 tion of hydrogen into even the driest planetary mate-845 rials known in the solar system (McCubbin & Barnes 846 2019; Piani et al. 2020; Jin et al. 2021) suggests that 847 hydrogen is present in all rocky planets upon forma-848 tion. Accreted hydrogen in nominally dry planetary 849 materials react with mantle oxygen to form substan-850 tial amounts of water inside of the planet during the magma ocean phase (Ikoma et al. 2018; Kite & Schaefer 852 2021; Kimura & Ikoma 2020, 2022). Enhanced equilibration between the core, mantle, and atmosphere dur-854 ing magma ocean evolution of rocky exoplanets further 855 enhances this process (Lichtenberg 2021; Schlichting &

Young 2022). Therefore, even nominally dry planets generate substantial amounts of water during formation and early evolution. Reduced heating from short-lived radionuclides in extrasolar planetary systems increases the expected water abundance in exoplanet systems further, in particular for M dwarf systems (Lichtenberg & Clement 2022).

Planetary evolution and duration of the steam atmo-864 sphere phase: An inflated steam atmosphere can only be 865 sustained until the planet has lost its water. Depending 866 on host star spectral type, planetary mass, and compo-867 sition, planets can spend from a few Myr to several Gyr 868 in runaway greenhouse climates (Hamano et al. 2015; 869 Luger & Barnes 2015), and only planets observed dur-870 ing this phase will contribute to the habitable zone inner 871 edge discontinuity. However, the delivery uncertainty is 872 much greater than the predicted loss rates of water by 873 atmospheric escape, which typically is limited to in to-874 tal a few tens of terrestrial oceans per Gyr (Wordsworth et al. 2018) on even the most irradiated exoplanets (see 876 discussion in Lichtenberg & Clement 2022). The finite 877 duration of steam atmosphere phases is a main factor 878 for diluting the habitable zone inner edge discontinu-879 itv. and it leads to an increased likelihood of detecting 880 this imprint for younger planetary systems. Preferen-



**Figure 9.** Retrieved posterior distributions of the threshold instellation for different survey realizations. All cases assume  $x_{H_2O}=0.005$  and a planet sample size  $N=100\pm10$ ; The fraction of planets with runaway greenhouse climates varies across rows. Orange lines show the true value of the injected signal. Accuracy and precision of the constraint on  $S_{\rm thresh}$  generally improve with higher  $f_{\rm rgh}$ . Bulk density-based inferences improve the constraints, and M dwarf samples yield the highest accuracy and precision.

881 tial selection of younger systems could therefore be of 882 advantage in a targeted survey.

Prevalent water inventory: The magnitude of radius 883 change at the runaway greenhouse threshold is sensitive to the water mass fraction. As a result, the statistical 886 abundance of water in terrestrial planets impacts the 887 strength of the demographic pattern: The higher the water content and the higher the fraction of planets in 889 runaway greenhouse climates, the greater the likelihood 890 that a discontinuity is detectable. However, dissolution into the magma ocean and atmospheric inflation show 892 non-linear coupling: the more water there is in the atmosphere, the more will be dissolved (Dorn & Lichtenberg 2021), leading to a convergence in radii until saturation. Most important for the present work, however, 896 is the qualitative dichotomy between sub-runaway and unaway planets, which outcompetes radius variations within each of these climate states due to a changing water mass fraction (Turbet et al. 2020). If a planet hosts significant water inventory that is not easily stripped by atmospheric escape (Johnstone 2020), then the statistical power of our model is high (see Figure 7).

Size and composition of the planetary sample: The significance of a statistical trend increases with a larger

905 sample size. In addition, the sample must include plan-906 ets on both sides of the instellation threshold. This is 907 only likely for low-mass host stars due to the strongly 908 distance-dependent detection bias associated with the 909 transit method and the temporal coverage of upcoming 910 transit missions.

Radius measurement precision: The more precise individual planet radii can be determined, the more pronounced the discontinuity will be. Good accuracy is less important, as long as it does not have a systematic error scaling with stellar irradiance.

Availability and precision of mass measurements: For simple geometric reasons ( $\rho \propto R^{-3}$ ), the expected dissorbination on the habitable zone inner edge is stronger when measured in bulk density than it is in planet radius space. If transiting planets can be followed up to obtain mass measurements, the statistical significance increases.

Besides these main factors, uncertainties in the measured instellations can influence the result, although they are typically small due to the very precise orbital period measurements available for transiting planets. This can be different for young host stars when their ages cannot be well constrained; in particular, the long pre-main sequence phase of M dwarfs shows a large variation in bolometric luminosity (see Figure 2).

#### 6.2.2. False positive scenarios

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Runaway greenhouse climates are not the only physi-933 cal mechanism that may cause a change in transit radius 934 for a subset of planets. Alternatives include atmospheric 935 loss due to either photoevaporation through high-energy 936 radiation by the host star (e.g., Ikoma & Hori 2012; 937 Owen & Wu 2013; Jin et al. 2014; Mordasini 2020) or 938 due to residual heat from the planet's interior shortly 939 after formation (Ginzburg et al. 2016, 2018; Gupta & 940 Schlichting 2019). Both processes are being discussed as 941 potentially sculpting the observed bimodality in the ra-942 dius distribution of small exoplanets (Fulton et al. 2017; 943 Van Eylen et al. 2018), and both lead to a decrease of 944 planet radius for planets close to their host star (Pas-945 cucci et al. 2019; Bergsten et al. 2022). This is distinct 946 from the radius inflation introduced by runaway green-947 house climates. For example, the innermost planet in 948 the K2-3 system has an increased radius compared to 949 its outer siblings, contrary to what would be expected 950 from atmospheric escape (Diamond-Lowe et al. 2022).

Other false positive contributions may stem from potential unknown occurrence rate gradients in radiusinstellation space, especially if these variations are simlar to the expected habitable zone inner edge discontinuity. Although an abrupt pattern at the expected location of the transition seems unlikely, examples of steep occurrence rate density changes exist. An example is the "Neptune desert", a triangular region in period-radius space of low planet occurrence (Szabó & Kiss 2011; Mazeh et al. 2016; Dreizler et al. 2020). The shape of this region is such that smaller planets become less frequent the closer to the star they are, which to some degree resembles the pattern introduced by the instellation dependency of the runaway greenhouse transition. However, the Neptune desert occurs at smaller orbital periods and is sensitive to the planet radius (Szabó & Kiss 2011), which is not expected for the runaway greenhouse transition.

Luque & Pallé (2022) found that small planets or-970 biting red dwarfs can be classified into three density 971 regimes with a particularly strong separation between 972 planets consistent with a pure rocky and those consis-973 tent with a water-rich composition. This trend does 974 not represent a false positive scenario for the habitable 975 zone inner edge discontinuity, since no strong depen-976 dency on instellation has been found or is expected. 977 A population of "water worlds" with low bulk densi-978 ties on a wide range of orbits would merely attenuate the statistical runaway greenhouse imprint. If the 980 dichotomy forms primordially through migrated plan-981 ets that accreted from different regions of their protoplanetary disk (Venturini et al. 2020; Burn et al. 2021; 983 Schlecker et al. 2021a,b) or from inter-system variations 984 in the desiccation of volatile-rich planetesimals (Lichtenberg et al. 2019; Lichtenberg et al. 2021b; Lichtenberg 986 & Krijt 2021; Lichtenberg & Clement 2022; Bonsor et al. 987 2023), systems of all ages can be affected by this atten-988 uation.

## 6.2.3. Atmospheric spectral signatures

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We currently see four potential lines of discriminat-990 991 ing atmospheric signatures associated with runaway cli-992 mates in the exoplanet population. (i) Detecting the 993 atmospheric windows of water vapor in the near- to 994 mid-infrared. In a runaway greenhouse atmosphere, ab-995 sorption is dominated by the opacity of water vapor, 996 which has two prominent spectral features: one at 3.5-997 4.5  $\mu$ m, and one between 8–20  $\mu$ m (e.g., Boukrouche 998 et al. 2021). Probing these features requires an in-999 strument covering these wavelength ranges, for instance 1000 ELT METIS (Brandl et al. 2021), JWST MIRI (Rieke 1001 et al. 2015), or future missions such as LIFE (Bonati 1002 et al. 2019; Quanz et al. 2022; Dannert et al. 2022). Po-1003 tentially the two-band filter capabilities of PLATO may 1004 offer insight into particularly pronounced spectral fea-1005 tures in this runaway greenhouse regime (Grenfell et al. 1006 2020). However, detecting water features in steam at1007 mospheres may be obscured by high-altitude clouds that 1008 could potentially mute a transit signal as measured from peak to trough of an observed spectral line (Suissa et al. 1010 2020; Fauchez et al. 2019). A better understanding of 1011 the role of clouds in this problem may be critical for 1012 observational constraints on the runaway climate of in-1013 dividual exoplanets. (ii) Post-runaway planets may be 1014 detectable through O<sub>3</sub> absorption due to build-up of abi-1015 otic oxygen, leftover from photochemical dissociation of 1016 water, and hydrogen loss (Wordsworth & Pierrehum-1017 bert 2014; Luger & Barnes 2015). (iii) Water loss in 1018 runaway greenhouse episodes would increase the atmo-1019 sphere's D/H isotopic ratio, akin to the enhancement 1020 in Venus' present-day atmosphere (Kane et al. 2019, 1021 2021). This may be detectable in high-resolution ob-1022 servations focusing on isotope-sensitive transitions (Lin-1023 cowski et al. 2019; Mollière & Snellen 2019). (iv) Dise-1024 quilibrium chemistry in tidally-locked runaway planets. 1025 It has been suggested that the atmospheric depth and 1026 the presence or absence of oceans on "sub-Neptunes" 1027 could be probed via the abundance of species that are photochemically destroyed in the upper atmosphere, and 1029 replenished from either thermochemical layers (Yu et al. 1030 2021; Tsai et al. 2021) or at an ocean-atmosphere inter-1031 face (Loftus et al. 2019; Hu et al. 2021). Atmospheric 1032 nitrogen and carbon compounds can be partitioned into magma (Grewal et al. 2022; Bower et al. 2022). There-1034 fore, it may be possible to discern the presence of an un-1035 derlying magma ocean if the presence of an atmosphere 1036 on a rocky exoplanet can be confirmed. This has been 1037 suggested to be done via (v) eclipse photometry (Mans-1038 field et al. 2019; Koll et al. 2019) through the presence 1039 of a high albedo, which is expected to differ from the 1040 crystallized rock of a solidified magma ocean (Essack 1041 et al. 2020; Fortin et al. 2022). Spectral information 1042 from atmospheres of highly irradiated planets may also 1043 help to distinguish classical runaway greenhouse states 1044 from other climate regimes, such as a moist bistabil-1045 ity (Leconte et al. 2013b).

## 6.2.4. Detection in multi-planet systems

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As suggested by Turbet et al. (2019), an alterna1048 tive approach for detecting the runaway greenhouse1049 induced radius inflation is to search for its "local"
1050 imprints in multi-planet systems. Systems harbor1051 ing planets on both sides of the transition, such as
1052 TRAPPIST-1 (Gillon et al. 2016, 2017; Luger et al.
1053 2017; Agol et al. 2021), K2-3 (Diamond-Lowe et al.
1054 2022), or Kepler-138 (Piaulet et al. 2022), may show
1055 the predicted abrupt radius and density change, pro1056 vided the initial volatile content was sufficient and com1057 plete desiccation has not yet occurred. While degen-

eracies remain in interpreting bulk density fluctuations within individual systems (e.g., Turbet et al. 2020; Dorn & Lichtenberg 2021), the detection of a consistent pation tern in several such systems could be a convincing statistical evidence of the runaway greenhouse transition. The current sample of suitable systems is sparse: The California-Kepler Survey catalog (Fulton & Petigura 2018) contains only six planets with instellations  $<2\,S_{\oplus}$  and smaller than  $2\,R_{\oplus}$  in five multi-planet systems. Future additions to the multi-planet sample through missions such as PLATO are needed.

## 6.3. Diagnostic power of near-future exoplanet missions

Confirming or disproving the predicted habitable zone 1071 1072 inner edge discontinuity will depend on the significance with which the null hypothesis can be excluded, which is a function of instrumentation and survey strategy. As discussed in Section 6.2.1, key drivers from a mission 1076 design perspective are sample size, photometric precision, and the availability of planets around low-mass 1078 host stars. We found that along these axes, PLATO will be the most favorable among the upcoming transit missions. The PLATO team has released an estimate on the number of exoplanets that will be characterized 1082 in the course of the main survey mission. With an expected transit radius precision of 3\% (ESA 2017) for hundreds of planets (Rauer 2021), the PLATO mission is comparable to the optimistic survey (see Section 5.2) in terms of sample size and precision. If successful, it should readily detect the predicted statistical imprint or, in case of a non-detection, provide strong upper limits on the occurrence rate of runaway greenhouse planets. The latter depends on the lifetimes of runaway greenhouse phases, which are a function of the initial water 1092 inventory of the planets (Hamano et al. 2015). Overall, 1093 it seems feasible to derive the typical water content of low-mass exoplanets from these occurrence estimates. 1094

What other planned missions are suited to probe the habitable zone inner edge discontinuity? Kepler and K2 have contributed a large number of discovered terrestrial-sized planets, but few of them are in the habitable zone and their host stars are typically too faint for RV follow-up with current instrumentation (Dressing & 1101 Charbonneau 2015).

Similarly, the Transiting Exoplanet Survey Satellite (*TESS*, Ricker et al. 2014) planet sample lacks temperate, small planets around bright host stars (Ment & Charbonneau 2023), as was expected from planet yield calculations (Barclay et al. 2018). As of March 22, 2023, <sup>1107</sup> the NASA Exoplanet Archive<sup>4</sup> lists 40 *TESS* candidate <sup>1108</sup> or confirmed planets smaller than  $4\,R_\oplus$  with lower esti- <sup>1109</sup> mated instellation than Earth's.

The ongoing CHEOPS mission was designed as a follow-up mission to search for transits of planets discovered with other techniques, in particular with radial velocity measurements (Benz et al. 2021). As such, it will provide precise radius constraints on a sample of small planets; however, only a small number of plantagely limits CHEOPS' coverage to planets within the largely limits CHEOPS' coverage to planets within the runaway greenhouse regime, preventing a detection of the transition.

As CHEOPS, the Atmospheric Remote sensing In-1121 frared Exoplanet Large survey (Ariel, Puig et al. 2016) 1122 will be a follow-up mission that is not designed to pro-1123 vide a large number of new radius measurements. Ariel's 1124 primary targets are larger planets in the range of sub-1125 Neptune to Jupiter-like planets. We thus do not expect 1126 a significant contribution to statistically exploring the inner edge of the habitable zone for Earth-sized planets. While not primarily designed to detect transiting 1129 planets, the Galactic Bulge Time Domain Survey of the 1130 Nancy Grace Roman Space Telescope (Spergel et al.  $_{1131}$  2015) is expected to yield  $\sim 10^5$  transiting planets on 1132 short orbits and constrain their radii in the course of its mission (Montet et al. 2017).  $\mathcal{O}(1000)$  planets smaller 1134 than Neptune could be found around early to mid-1135 M dwarfs, however, only a small fraction of them will 1136 reach into the habitable zone (Tamburo et al. 2023). We 1137 thus conclude that the Nancy Grace Roman Space Tele-1138 scope could provide a useful sample to explore the run-1139 away greenhouse transition, albeit with a predominant 1140 focus on water-rich (sub-)Neptunes (e.g., Pierrehumbert 1141 2022).

Looking further ahead, the Nautilus Space Observatory concept (Apai et al. 2019a) represents a statistical
mission able to provide precise radius measurements of
tals a large sample ( $\sim 1000$ ) of small exoplanets. It employs
tals a constellation of  $\sim 35$  large-diameter ( $D \sim 8.5 \,\mathrm{m}$ ) teletals ements (Milster et al. 2020) with the primary goal to
tals study the atmospheres of transiting exoplanets. Optals erating in an array mode, Nautilus would achieve the
tals equivalent light-collecting area of a 50 m telescope. Its
tals expected 1 ppm photometric precision (Apai et al. 2022)
tals would enable precise radius measurements of a large
tals sample, also through a low number of required visits
tals per object. If realized, Nautilus will be a valuable instru-

<sup>&</sup>lt;sup>4</sup> https://exoplanetarchive.ipac.caltech.edu

1156 ment for characterizing the runaway greenhouse transi-

Other missions have been proposed that focus on char-1159 acterizing exoplanet habitability, most notably the Hab-1160 itable Worlds Observatory concept, which will build on the two precursor direct imaging concepts LU-VOIR (The LUVOIR Team 2019) and HabEx (Gaudi et al. 2020). Similar science objectives are pursued by 1164 the Large Interferometer For Exoplanets (LIFE, Quanz et al. 2022) initiative, a mission concept utilizing a 1166 space-based mid-infrared nulling interferometer. Direct 1167 imaging surveys do not directly measure planetary radii and are primarily useful for providing context through 1169 atmospheric measurements of individual planets. How-1170 ever, because mid-infrared retrievals feature reduced degeneracy between cloud albedo and changes in sur-1172 face area, the planet radius can be constrained in midinfrared wavelengths (Defrère et al. 2018; Quanz et al. 1174 2021). Mid-infrared direct imaging techniques, in par-1175 ticular, enable to study much deeper atmospheric layers 1176 than possible in reflected light (Wordsworth & Kreid-1177 berg 2022). Hence, the atmospheric structure can be 1178 retrieved for a wider variety of thermal and atmospheric scenarios (Alei et al. 2022; Konrad et al. 2022). Since 1180 the peak thermal emission in runaway greenhouse atmospheres will substantially decrease the star-to-planet 1182 flux ratio, mid-infrared wavelengths offer the possibility to probe the diversity of runaway climates in systems across different ages (Lupu et al. 2014; Bonati et al. 1185 2019). Finally, mid-infrared surveys such as LIFE show preference for M star planets (Quanz et al. 2022), 1187 which is beneficial for detecting the predicted habitable 1188 zone inner edge discontinuity (see Sect. 5.4.2). With a sample size of a few tens of planets crossing the runaway greenhouse transition, direct imaging missions will thus enable key insights into the compositional inventory of 1192 atmospheric volatiles and climate states (Hinkley et al. 1193 2021; Currie et al. 2022), adding important details to a 1194 potential runaway greenhouse detection purely via tran-1195 sit radii.

As for exoplanets missions in their implementation phase, however, *PLATO* overall remains to be the most promising mission for an empirical confirmation or falsification of the runaway greenhouse transition at this time.

#### 6.4. Mission design trade studies

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To explore the impact of mission trades on the de-1203 tectability of the habitable zone inner edge discontinu-1204 ity, we simulated different survey designs and strategies 1205 and measured their capability to recover the trend and 1206 constrain its parameters. We assessed this capability 1207 based on two determinants: the likelihood that the mis-1208 sion is able to detect the injected trend, and the preci-1209 sion with which it can constrain the parameters of that 1210 trend.

#### 6.4.1. The value of follow-up campaigns

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The constraining power changes when additional in-1213 formation beyond planet radii is available for the char-1214 acterized planet population. As runaway greenhouse 1215 phases leave a stronger imprint on bulk density than on planet radius (see Sect. 5.4.1), it would be benefi-1217 cial to obtain constraints on planetary masses and test 1218 the runaway greenhouse hypothesis in density space in-1219 stead of radius space. This way, useful results can be 1220 obtained under more pessimistic conditions, e.g., a low 1221 predominant water content of planetary surfaces and at-1222 mospheres or a smaller available planet sample. For a 1223 mission design similar to PLATO a density-based hy-1224 pothesis test on about a third of the overall sample is 1225 equivalent to a pure radius-based analysis. At a fixed 1226 sample size, key parameters of the runaway greenhouse 1227 models can be more narrowly constrained when addi-1228 tional mass measurements are available.

Precise ground-based radial velocity measurements will be needed to provide these data, and a number of instruments are already successfully employed in characterizing terrestrial-sized exoplanets (e.g., Queloz et al. 2021; Pepe et al. 2010; Johnson et al. 2010; Ribas et al. 2023) and confronting these results with planet formation theory (e.g., Miguel et al. 2020; Burn et al. 2021; Za-unid wadzki et al. 2021; Schlecker et al. 2022). A new generation of instruments on extremely large telescopes such as G-CLEF on the Giant Magellan Telescope (Szentgy-unid as G-CLEF) on the Giant Magellan Telescope (Szentgy-unid as G-CLEF) (Marcantonio et al. 2022), or MODHIS Large Telescope (Marcantonio et al. 2022), or MODHIS unid the Thirty Meter Telescope (Mawet et al. 2019) will upone up the discovery space even further.

Recently, NASA and the National Science Foundation (NSF) commissioned an "Extreme Precision Radial Velocity Initiative" (Crass et al. 2021) to develop methods and facilities for precise mass measurements of temperizer ate terrestrial planets. Their findings highlight that such measurements are costly, and therefore follow-up efforts may only be available for a subsample of the targets of a mission of *PLATO*'s scale. The diagnostic power of the hypothesis tests we demonstrated here may be improved by simultaneously fitting for the habitable zone inner edge discontinuity in the subsample without RV follow-up. An optimized mission in search for the inner edge of the habitable zone will further enhance its information content via an informed selection of follow-up

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targets, i.e., balancing objects located on either side of the expected instellation threshold  $S_{\rm thresh}$ .

## 6.4.2. The importance of M dwarfs in the target list

To date, the majority of planets with radius measure-1260 ments orbit FGK dwarfs, and, based on the instellation they receive, most of them lie in the runaway greenhouse regime (Thompson et al. 2018). Obviously, a radius/density discontinuity in the exoplanet demographics like the habitable zone inner edge discontinuity can-1266 not be constrained well if only one side of the disconti-1267 nuity is being sampled. This, however, is the situation for planetary systems around Sun-like stars – their habitable zones are so distant that transiting planets within them are very rare due to pure geometrical reasons. 1271 It was, among other reasons, the sharp drop in tran-1272 sit probability with orbital distance that has prompted number of recent transit surveys to specifically target M dwarfs (e.g., Irwin et al. 2009; Obermeier et al. 2016; Delrez et al. 2018; Sebastian et al. 2021; Dietrich et al. 2023), but the sample of terrestrial planets orbiting them is still small (e.g., Berger et al. 2020; Hardegree-Ullman et al. 2020). 1278

M dwarf systems are also key for detecting the run-1279 1280 away greenhouse transition: Our calculations with different spectral types (Sect. 5.4.2) show that the information content of M dwarfs in a sample dominates 1283 the hypothesis tests. Besides their large number in a volume/magnitude-limited sample, transiting M dwarf planets are more likely to be located near the threshold instellation and in particular on orbits further out, i.e., in the optimistic habitable zone. In fact, we showed that the FGK part of the planet sample barely con-1289 tributes to the statistical power. Furthermore, the transit depth difference at the transition is expected to be larger for M dwarfs ( $\sim 100 \, \mathrm{ppm}$  for early,  $\sim 1000 \, \mathrm{ppm}$ 1292 for late M stars, Turbet et al. 2019), enhancing the demographic signal it leaves. An additional advantage of 1294 targeting M dwarfs are the extended runaway green-1295 house phases of their planets that can last on the order 1296 of gigayears (Luger & Barnes 2015). This increases the probability of observing any given planet in the sample 1298 during the runaway greenhouse phase, essentially driv-1299 ing  $f_{\rm rgh}$  to higher values. Therefore, in addition to the 1300 high scientific value of boosted detections of potentially habitable planets, M dwarfs are also indispensable for 1302 the discovery and characterization of the runaway green-1303 house transition. As with a pure volume-limited sample, 1304 a targeted M dwarf survey, too, profits from follow-up 1305 measurements of planetary masses with an order of mag-1306 nitude increase in evidence.

## 6.5. Constraining planetary habitability

A potential for liquid water on the surface of a planet 1309 is commonly used as an environmental marker to assess 1310 its surface habitability (Huang 1959; Hart 1978; Kasting et al. 1993; Kaltenegger & Sasselov 2011; Koppa-1312 rapu et al. 2013). The runaway greenhouse transition 1313 represents an upper bound on received irradiation for 1314 this condition. Its detection would thus not only em-1315 pirically confirm the habitable zone concept but also 1316 help to locate it in the observationally available plan-1317 etary parameter space. In Sect. 5.4.3, we show that the 1318 threshold instellation at which the runaway greenhouse 1319 transition occurs can be reasonably constrained without 1320 imposing overly optimistic conditions on the underlying 1321 planet population, instrumentation, or survey strategy. 1322 A mission like PLATO is well equipped to perform this measurement; the constraining power is directly propor-1324 tional to the proportion of characterized planets around 1325 M dwarfs and to the number of planets for which masses can be determined.

The situation is different for the planetary water in1328 ventory and the fraction of planets with runaway green1329 house climates: Since these parameters are degenerate,
1330 they cannot be well constrained without independent
1331 measurements. This degeneracy could be lifted if in1332 dependent measurements of atmospheric compositions
1333 can be made. For example, detections of water vapor in
1334 planets above the threshold instellation, combined with
1335 precise radius measurements, would constrain the pre1336 dominant water content of terrestrial planets.

Once a runaway greenhouse region is identified in the parameter space, the community will have a tool at hand to discern potentially habitable planets from Venusian worlds on an empirical basis. Together with atmospheric measurements (see Sect. 6.2.3), we will be able to put a number on the probability of an individual planet to harbor sufficient surface water to sustain life.

## 6.6. Impact of assumptions on our findings

The prospects for probing the runaway greenhouse transition depends on astro- and geophysical factors, as well as on the specific instrumentation and survey stratages egy of a particular mission. Our state-of-the-art models approximate the situation and offer testable predictions. In the following, we review a few considerations that future models may include to refine these predictions.

#### 6.6.1. Structures in the planet occurrence rate density

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The baseline occurrence rate density in radius-period space that governs the generation of synthetic planets might influence our findings, especially if it contains any features that coincide with the injected demographic feature. This is not the case in the model from Berg-

 $_{1358}$  sten et al. (2022) that we adopted: Its occurrence rate  $_{1359}$  density varies smoothly in the domain relevant for the  $_{1360}$  runaway greenhouse hypothesis; transitions only occur  $_{1361}$  at smaller instellations (<  $50\,\mathrm{W\,m^{-2}}$ ) and larger radii  $_{1362}$  (>  $1.6\,R_\oplus$ ). We thus do not expect the model underly- $_{1363}$  ing our planet sample to affect the investigation of the  $_{1364}$  habitable zone inner edge discontinuity. If any currently  $_{1365}$  unknown sharp features in the distribution of terres- $_{1366}$  trial planets emerge, they should be considered in future  $_{1367}$  studies.

#### 6.6.2. Baseline mass-radius relationship

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Our baseline mass-radius relationship assuming pure MgSiO<sub>3</sub> interiors (Zeng et al. 2016) might not be representative of the rocky planet population. However, while interior composition may introduce an offset to the radius habitable zone inner edge discontinuity, we do not expect a change of its structure. Since the magnitude of the radius inflation effect is expected to be larger for an Earth-like interior composition with an iron core-silicate mantle structure (Zeng et al. 2016; Noack & Lasbleis Rock Bonati et al. 2021), we consider our mass-radius relationship a conservative case. We performed a sanity check to assess the impact of varying our baseline model (see Appendix A.3) and found general agreement between different interior compositions.

Future self-consistent modeling of interior-atmosphere interactions may include constraints on additional radius increases due to a molten interior (Bower et al. 2019) and any potential effects stemming from a deviating gas exchange between atmosphere and interior in runaway greenhouse planets due to different redox conditions (Ikoma et al. 2018; Lichtenberg et al. 2021; Bower et al. 2022; Gaillard et al. 2021) or water outgassing efficiency (e.g., Hier-Majumder & Hirschmann 2017; Ikoma et al. 2018; Salvador & Samuel 2023).

#### 6.6.3. Bulk water mass fraction

The predominant mass fractions of water, which sensitively controls the atmospheric state of a rocky exoplanet, is poorly constrained. Inferred water contents
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 $_{1409}$  to introduce a systematic overestimation of the habitable zone inner edge discontinuity. Cases of pure rocky  $_{1411}$  composition and very low volatile contents can be considered absorbed by the dilution factor  $f_{\rm rgh}$ . Assuming  $_{1413}$  a distribution of water mass fractions instead of a fixed  $_{1414}$  value would thus not significantly change our results.

## 6.6.4. "Sharpness" of the habitable zone inner edge discontinuity

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The habitable zone inner edge discontinuity may be 1418 affected by several processes that are challenging to 1419 quantify: Planets that lack an atmosphere, sufficient 1420 volatiles, or have non-water-dominated outgassed com-1421 positions cannot bear steam atmospheres, and those that do eventually move to the non-runaway greenhouse 1423 category due to desiccation (Watson et al. 1981; Kasting <sup>1424</sup> & Pollack 1983; Hamano et al. 2013) or evolution of their 1425 host star (Luger & Barnes 2015). Hydrogen/Helium-1426 dominated planets may disguise as inflated rocky planets 1427 and not contribute to the demographic signal, although 1428 a runaway greenhouse radius inflation effect was sug-1429 gested for water-dominated sub-Neptunes (Pierrehum-1430 bert 2022; Innes et al. 2023). A subset of such gas-1431 rich planets will experience atmospheric loss via photo-1432 evaporation (Owen & Wu 2013) or core-powered mass loss (Ginzburg et al. 2018), reducing their transit radius. 1434 Intrinsic variation in the threshold instellation is caused by differences in planetary features influencing the on-1436 set of a runaway climate such as albedo, atmospheric 1437 composition, clouds, or surface gravity (Salvador et al. 1438 2017; Turbet et al. 2021; Lichtenberg et al. 2021; Pier-1439 rehumbert 2022; Innes et al. 2023). The choice of a sta-1440 tistical estimator for the hypothesis tests may further 1441 influence the recovered discontinuity; we compare our 1442 nominal running mean approach with a binned statistic 1443 in Appendix A.1.

While these factors may offset the signal's amplitude, they preserve its general shape. The "dilution factor"  $f_{\rm rgh}$  in our model embodies our ignorance of the magnitude of this offset. In a real survey, additional contextual information about planets in the sample may be available.

## 7. CONCLUSIONS

Significant inflation of rocky planet radii is a ro1452 bust prediction of runaway greenhouse models. Using
1453 Bioverse, a quantitative hypothesis testing framework,
1454 we have explored the potential of contemporary exo1455 planet missions to statistically detect a radius/density
1456 discontinuity resulting from this inflation in the exo1457 planet population. Our key findings are as follows:

1. The predicted runaway greenhouse transition causes a discontinuity in the radius and density

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distribution of small exoplanets with respect to their irradiation.

- 2. This habitable zone inner edge discontinuity should be detectable with high-precision transit measurements. For a planet sample  $\gtrsim 100$ , a detection is likely if radius inflation occurs on at least 10% of the observed planets and if typical bulk water mass fractions are above  $\sim 10^{-3}$ .
- 3. We find that the planned PLATO transit survey will provide a sufficient sample and the required precision to confirm or reject the predicted trend. Assuming the projected photometric precision, PLATO will be able to test the runaway greenhouse hypothesis for planet yields  $\gtrsim 100$ .
  - 4. The diagnostic power of transit missions in testing this hypothesis can be increased through a follow-up campaign providing planet mass measurements. This can reduce the required planet yield by about a factor of three. Only an adequate sample of planets orbiting M dwarfs will ensure sufficient targets on both sides of the expected threshold instellation.
  - 5. Testing the runaway greenhouse hypothesis on a population level can provide constraints on the water inventory of rocky exoplanets and thus make an important contribution to assessing their habitability. A detection will provide an empirical confirmation of the habitable zone concept and localize its inner edge.

The habitable zone concept is widely employed in tar1490 get prioritization for exoplanet missions, and it will pro1491 vide context for interpreting potential signatures of life.
1492 As we have demonstrated, it appears realistic that an
1493 empirical test of the habitable zone hypothesis is im1494 minent. The confirmation or rejection of the habitable
1495 zone inner edge discontinuity will be a key contribution
1496 to understanding the diversity of exoplanet climates and
1497 the search for extraterrestrial life in the Universe.

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### AUTHOR CONTRIBUTIONS

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M.S., D.A., and T.L. conceived the project, planned its implementation, and interpreted the results. D.A. leads the "Alien Earths" program through which this project is funded and helped to guide the strategy of the project. T.L. and A.S. provided expertise on run-1534 away greenhouse climates and exoplanet interiors. M.S. carried out the hypothesis tests and statistical anal-1536 yses. M.S. wrote the manuscript; T.L., G.B., K.H.-1537 U., and A.S. provided text contributions. G.B. im-1538 plemented the planet generator in the Bioverse frame-1539 work. All authors provided comments and suggestions on the manuscript.

#### REPRODUCIBILITY

This study uses the reproducibility framework "showyourwork" (Luger et al. 2021). All code re1544 quired to reproduce our results, figures, and this arti1545 cle itself is available at https://github.com/matiscke/
1546 hz-inner-edge-discontinuity, and the repository state
1547 at the time of paper acceptance can be found at

1548 doi:10.5281/zenodo.8251077. The code to reproduce a 1549 figure can be accessed via the icon link next to the re-1550 spective figure caption. Data sets associated with this 1551 work are available at doi:10.5281/zenodo.7080391 (stel-1552 lar luminosity tracks from Baraffe et al. (1998)) and

doi:10.5281/zenodo.7946446 (results from model grid runs of Bioverse).

Software: Bioverse (Bixel & Apai 2021), As-1556 tropy (Astropy Collaboration et al. 2018), NumPy (Har-1557 ris et al. 2020), SciPy (Virtanen et al. 2020), 1558 corner.py (Foreman-Mackey 2016), dynesty (Speagle 1559 2020).

#### 1560 APPENDIX

#### A. ROBUSTNESS TESTS

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## A.1. Alternative statistics for the average radius or bulk density

The hypothesis tests introduced in Sect. 4 rely on a statistical estimator for the variation of planetary radii or bulk densities as a function of net instellation, and we star chose a moving average for this estimator in our nominal setup. Here, we explore how robust our results are against this choice by demonstrating the recovery of the runaway greenhouse signal in the case of the optimistic survey (Sect. 5.2) with an alternative estimator: instead of computing moving averages, we used a binned statistic.

We first binned the data of the simulated survey in instellation space, choosing the number of bins via the rule of Freedman & Diaconis (1981) and using logarithmic binning. In each bin, we computed the arithmetic mean of the planet radius and its standard deviation. Then, we assigned each planet the mean radius according to the instellation bin it occupies and used this as the measure for testing the runaway greenhouse hypothesis.

Figure 10 shows the simulated data together with binned, average planet radii and draws from the postetion of the hypothesis test. A clear detection resulted, although with somewhat lower significance ( $\Delta \ln \mathcal{Z} \approx 30$ ) compared to the nominal setup. The accuracy of the respective covered instellation threshold is comparable. This test demonstrates that our results are not sensitive to the choice of statistical estimator to test the hypotheses against.

1591 It is conceivable that with very large sample sizes and 1592 a very sharp runaway greenhouse transition a binned 1593 solution would perform better. For a search with real 1594 data, both approaches, and possibly other alternatives, 1595 should be considered.

## A.2. Statistical imprint of runaway greenhouse atmospheres

The predicted runaway greenhouse-induced planet ra-1599 dius changes are a function of instellation, planet mass, 1600 and bulk water mass fraction (compare Sect. 3). In or-

## Optimistic survey (binned)

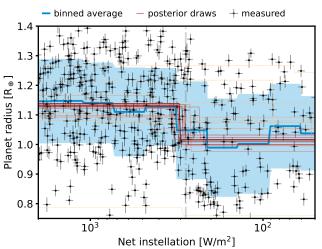


Figure 10. Detection of the runaway greenhouse threshold with a binned statistic. Using simulated data from the optimistic survey case (N=500, compare Figure 5), we tested the runaway greenhouse hypothesis. Instead of using a running mean, we computed a binned statistic (arithmetic mean and standard deviation in blue) to assign planets the average radius or bulk density in their neighborhood in instellation space. Random draws from the posterior of the runaway greenhouse hypothesis (Eqn. 4) are shown in red. As in the nominal case, the pattern is detected with high significance.

1601 der to better understand the interaction between the 1602 planetary populations underlying our simulations and 1603 these predictions, we compared the latter to the aver-1604 age radius and bulk density changes we measured in the 1605 synthetic population. We used the "optimistic" scenario 1606 with a sample size of 500.

Figure 11 shows this comparison for a range of plan-1608 etary masses and water mass fractions. The complex 1609 dependence of the radius inflation on these parameters 1610 is evident, but no significant abrupt changes capable 1611 of causing spurious signals occur. Differences between 1612 model prediction and population can be explained by 1613 the wide dispersion in planet mass in the population.

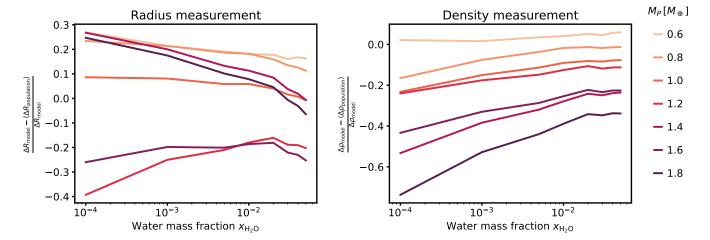


Figure 11. Comparison of radius and bulk density changes predicted by the atmospheric models to the average changes measured in the synthetic planet population. Model predictions  $\Delta R_{\rm model}$  and  $\Delta \rho_{\rm model}$  depend on individual planet masses;  $\langle \Delta R_{\rm population} \rangle$  and  $\langle \Delta \rho_{\rm population} \rangle$  are averaged measurements of the overall population. Significant differences are thus expected.

## A.3. Influence of different mass-radius relationships on our results

We assessed how a different choice of baseline massradius relation influences our results. Focussing on the 1617 1618 detectability of the statistical runaway greenhouse signal and the ability to constrain the threshold instellation, we 1620 repeated the hypothesis test in Sect. 5.2 with alternative mass-radius relationships. Instead of assuming a pure 1621 MgSiO<sub>3</sub> composition, we assigned planet masses using either the probabilistic relationship in Wolfgang et al. (2016) or a semi-empirical, two-layer relation assuming 1624 an Earth-like (32.5 % Fe + 67.5 % MgSiO<sub>3</sub>) composi-1625 tion (Zeng et al. 2016). 1626

Figure 12 shows how the two alternative baseline 1627 mass-radius relations influence the significance of a de-1628 tection and the ability to constrain the threshold instellation. The relation of Wolfgang et al. (2016) includes intrinsic scatter, which impedes a detection at low dilution factors. In this regime, a layered, Earthlike composition leads to more significant detections and a narrower, although biased, constraint on  $S_{\text{thresh}}$ . 1635 Both mass-radius relations agree and recover the in-1636 jected value where the null hypothesis can be rejected with high significance. This is consistent with our nominal mass-radius relation (compare Sect. 2.4). We conclude that the underlying core and mantle composition of planets may affect the detectability of the transition if the fraction of planets with runaway greenhouse climates 1642 is low; however, the overall trends that our experiments 1643 revealed appear robust.

We caution that this analysis may serve only as a san-1645 ity check and should not be taken as a result in itself: 1646 The atmospheric model from Turbet et al. (2020) we 1647 adopted relies on a silicate interior composition for its

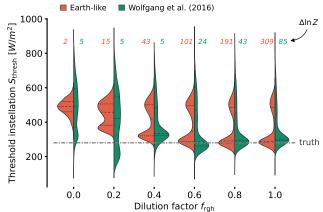


Figure 12. Posterior threshold instellations as a function of dilution factor for two alternative baseline mass-radius relations. For each grid step in  $f_{\rm rgh}$ , we show kernel density estimates of retrieved posteriors (averaged over 50 iterations) assuming an Earth-like (32.5 % Fe + 67.5 % MgSiO<sub>3</sub>) composition or the probabilistic relation from Wolfgang et al. (2016) and otherwise following the optimistic scenario in Sect. 5.2. Lines within the violins show quartiles of the distributions, and the gray line indicates the injected threshold instellation of 280 W m<sup>-2</sup>. Higher log-evidence differences  $\Delta \ln Z$  correspond to more significant rejections of the null hypothesis. In the regime of strong detections, both mass-radius relations lead to similar, accurate constraints on  $S_{\rm thresh}$ . Differences occur at low dilution factors, where the Earth-like relation leads to narrower estimates.

1648 transit radius prediction. Therefore, only our nominal 1649 procedure throughout the main body of the paper rep-1650 resents a self-consistent treatment.

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