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Bioverse: The Habitable Zone Inner Edge Discontinuity as an Imprint of Runaway Greenhouse Climates on Exoplanet Demographics

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ABSTRACT

Long-term magma ocean phases on rocky exoplanets orbiting closer to their star than the runaway greenhouse threshold – the inner edge of the classical habitable zone – may offer insights into the physical and chemical processes that distinguish potentially habitable worlds from others. Thermal stratification of runaway planets is expected to significantly inflate their atmospheres, potentially providing observational access to the runaway greenhouse transition in the form of a "habitable zone inner edge discontinuity" in radius—density space. Here, we use Bioverse, a statistical framework combining contextual information from the overall planet population with a survey simulator, to assess the ability of ground- and space-based telescopes to test this hypothesis.

We find that the demographic imprint of the runaway greenhouse transition is likely detectable with high-precision transit photometry for sample sizes $\gtrsim 100$ planets if at least $\sim 10\,\%$ of those orbiting closer than the habitable zone inner edge harbor runaway climates. Our survey simulations suggest that in the near future, ESA's PLATO mission will be the most promising survey to probe the habitable zone inner edge discontinuity. We determine survey strategies that maximize the diagnostic power of the obtained data and identify as key mission design drivers: 1. A follow-up campaign of planetary mass measurements and 2. The fraction of low-mass stars in the target sample. Observational constraints on the runaway greenhouse transition will provide crucial insights into the distribution of atmospheric volatiles among rocky exoplanets, which may help to identify the nearest potentially habitable worlds.

1. INTRODUCTION

Despite recent advancements in observational tech-26 niques, our understanding of terrestrial-sized planets re-27 mains woefully limited, with fundamental aspects of 28 their nature, composition, and potential habitability 29 still largely unknown. Due to the inherent biases of 30 current exoplanet detection techniques, the best-studied 31 category of rocky exoplanets at present is that of hot or ³² warm, close-in planets (Jontof-Hutter 2019; Bean et al. 33 2021). These experience thermal states that are in some 34 aspects comparable to the ones of the inner solar sys-35 tem bodies at early stages of their evolution (Ikoma 36 et al. 2018; Chao et al. 2021), which likely profoundly 37 affected the distribution of volatiles between planetary 38 core, mantle, and atmosphere. Studying the geophysical 39 state of hot exoplanets can thus inform our understand-40 ing of the early evolutionary stages of Earth and other 41 habitable worlds (Lichtenberg et al. 2022; Krijt et al. 42 2022).

An example from the solar system for the potential significance of these early stages are the divergent atmo-

45 spheric evolutions of Venus and Earth (e.g., Kane et al. 46 2019, 2021; Salvador et al. 2023). While having accreted 47 from a similar mass reservoir (Raymond et al. 2020; 48 Kleine et al. 2020; Mezger et al. 2020; Zahnle & Carl-49 son 2020) and despite their similar bulk properties (Sm-50 rekar et al. 2018), they evolved into planets with very 51 different surface conditions (Donahue et al. 1982; Kast-52 ing 1988; Hamano et al. 2013; Kane et al. 2014; Way & 53 Del Genio 2020; Turbet et al. 2021). Both planets likely 54 underwent a giant impact phase (Raymond et al. 2020; 55 Gillmann et al. 2020; Liu et al. 2022) that melted their 56 mantles (Elkins-Tanton 2012; Schaefer & Elkins-Tanton 57 2018; Lichtenberg et al. 2022). Magma ocean states play 58 a substantial role in establishing the long-term geophys-59 ical and climatic regimes of rocky planets (Fegley et al. 60 2020), in particular owing to efficient heat and volatile 61 transfers between interior and atmosphere in the ab-62 sence of a stiff boundary separating them (Kite & Schae-63 fer 2021; Dorn & Lichtenberg 2021; Salvador et al. 2023). 64 Due to these similar formation sequences, it was com-65 monly assumed that the divergence of Venus and Earth ₆₆ – in particular Venus' water loss – occurred late in their evolution (e.g., Way & Del Genio 2020).

Yet, Hamano et al. (2013) suggested that the present-69 day dry conditions on Venus may have been directly in-70 herited from the early magma ocean stage. If a strongly 71 infrared-absorbing, condensable species such as water 72 was dominant in the atmosphere, the resulting strong 73 thermal blanketing effect would prevent the planet to 74 efficiently radiate to space and maintain the surface 75 molten (Ingersoll 1969; Kasting 1988; Pierrehumbert 76 2010; Goldblatt et al. 2013; Leconte et al. 2015; Salvador 77 et al. 2017). This runaway greenhouse state can extend 78 the magma ocean stage to hundreds of Myr (Schaefer 79 et al. 2016; Barth et al. 2021), enough to remove the 80 entire water reservoir from a rocky planet by H₂O pho-81 tolysis and subsequent hydrodynamic escape of hydro-82 gen (Wordsworth & Pierrehumbert 2013, 2014; Luger 83 & Barnes 2015). For Venus, the atmospheric composi-84 tion (Gillmann et al. 2020) and comprehensive analysis 85 of meteoritic samples across the Solar System (Alexan-86 der et al. 2018; Broadley et al. 2022) suggest that it 87 went through this phase. Although the past presence or 88 absence of a Venusian water ocean has not been defi-89 nitely established (Raymond et al. 2006, 2007; Hamano 90 et al. 2013; Way et al. 2016; Kane et al. 2019, 2021; 91 Turbet et al. 2021; Warren & Kite 2023), a transient 92 habitable phase cannot be conclusively ruled out (e.g., 93 Way et al. 2016; Salvador et al. 2017; Krissansen-Totton 94 et al. 2021).

The runaway greenhouse transition is a robust predic-⁹⁶ tion from climate models (Kasting 1988; Nakajima et al. 97 1992; Goldblatt & Watson 2012; Forget & Leconte 2014; 98 Boukrouche et al. 2021; Chaverot et al. 2022), and its 99 impact on planetary bulk properties has been shown to 100 be in the detectable range of current astronomical in-101 strumentation (Goldblatt 2015). In particular, planets in a runaway greenhouse state are expected to be thermally inflated (Turbet et al. 2019, 2020; Mousis et al. 104 2020), which directly increases their transit radii by an amount that is a function of the water content. However, dissolution of water (e.g., Elkins-Tanton & Seager 2008; 107 Hier-Majumder & Hirschmann 2017; Salvador & Samuel 108 2023) in the magma may decrease this effect (Dorn 109 & Lichtenberg 2021), and chemical exchange between 110 core and mantle material may influence the amount and 111 speciation of outgassed volatiles that are visible in the 112 atmosphere via transmission spectroscopy (Lichtenberg 113 2021; Schlichting & Young 2022). Astronomical ob-114 servations of planets that are currently in a runaway 115 greenhouse state may thus constrain properties of their 116 mantles and establish an observational connection between exoplanetary interiors and atmospheres (Lichtenberg et al. 2022; Wordsworth & Kreidberg 2022).

Of particular relevance is that the radiation-induced 120 transition between a runaway greenhouse state and 121 a temperate climate is thought to occur at a rela-122 tively sharp instellation threshold (Goldblatt et al. 2013; 123 Leconte et al. 2013a; Kopparapu et al. 2013). Conse-124 quently, the instellation at which the runaway green-125 house transition occurs is aptly considered to be the 126 inner boundary of the habitable zone (e.g., Ramirez 127 2018; Salvador et al. 2023). Its prevalent definition 128 refers to the possibility of a sustained liquid water 129 body on the surface of an Earth-like planet with an 130 oxidized CO₂/H₂O/N₂-rich atmosphere (Kasting et al. 131 1993; Kopparapu et al. 2013, 2014); its exact spatial 132 location and extent may be strongly influenced by the 133 interior and atmosphere oxidation state and resulting atmosphere composition (Pierrehumbert & Gaidos 2011; 135 Ramirez & Kaltenegger 2017, 2018; Katyal et al. 2019; 136 Graham & Pierrehumbert 2020; Graham et al. 2022; 137 Hakim et al. 2023). The fundamental concept of a hab-138 itable zone dates back centuries (Newton 1687; Whewell 139 1858; Shapley 1953; Huang 1959), and its modern form 140 has proven popular in the planetary literature¹. How-141 ever, it should be emphasized that the habitable zone, 142 as it stands today, is merely a concept based on the-143 oretical predictions and geochemical evidence from one 144 planet – Earth (Catling & Zahnle 2020) – and its general 145 validity remains controversial (e.g., Cockell et al. 2016; 146 Moore et al. 2017; Tuchow & Wright 2023). The ques-147 tion naturally emerges if a planetary habitable zone – in 148 its common form with boundaries defined by stellar irra-149 diation – is a predictive theoretical concept and how the 150 habitable zone hypothesis can be tested observationally. Observational tests being considered include searches 152 for direct evidence of liquid water conveyed by ocean 153 glint (Williams & Gaidos 2008; Robinson et al. 2010; 154 Lustig-Yaeger et al. 2018) or water vapor in planetary 155 atmospheres (Suissa et al. 2020). A different approach 156 relies on comparative planetology: aiming for a statisti-157 cal detection in a planet population provides robustness 158 against ambiguity that could otherwise arise from indi-159 vidual variations in a planet's composition or geophys-160 ical history (Checlair et al. 2019; Apai et al. 2019b). 161 Tests suggested in the literature include determining, 162 for a range of orbital distances, atmospheric H₂O and 163 CO₂ abundances, planetary albedos (Bean et al. 2017;

 $^{^1}$ At the time of writing, a search of the term "habitable zone" in the titles and abstracts of refereed articles in the National Aeronautics and Space Administration (NASA) Astrophysics Data System returned ~ 1700 results.

 164 Bixel & Apai 2021), or colors (Crow et al. 2011; Bixel 165 & Apai 2020), or testing the relationship between 166 partial pressure and incident flux (Lehmer et al. 2020). 167 All these tests require surveying a large enough sample 168 of terrestrial-sized planets with next-generation instruments, rendering them out of reach in the immediate 170 future.

Here, we explore the feasibility of a statistical test 171 172 of the habitable zone hypothesis by surveying plane-173 tary bulk properties close to its inner edge, the runaway 174 greenhouse transition. Our goal is to assess the ability of near-future transit surveys to test the hypothesis that 176 the runaway greenhouse effect causes a discontinuity of 177 planetary radii and bulk densities when ordered by re-178 ceiving instellation (Turbet et al. 2019). Our main tool 179 for this is Bioverse, a simulation framework for assess-180 ing the statistical power of exoplanet surveys (Bixel & 181 Apai 2021). It consists of a sample generator that populates stars from the Gaia catalog (Hardegree-Ullman et al. 2023; Prusti et al. 2016; Gaia Collaboration et al. ¹⁸⁴ 2022a) with planetary systems based on state-of-the-art occurrence rates (Bergsten et al. 2022), a flexible survey 186 simulator that allows for a broad range of trade studies, and a hypothesis testing module that quantifies the 188 survey's ability to detect a previously injected trend. 189 Trying different instrumentation and survey designs, we 190 use Bioverse to recover runaway greenhouse-induced 191 effects based on model predictions (Turbet et al. 2020; 192 Dorn & Lichtenberg 2021) that we inject into a baseline 193 planet population.

In particular, we test the capability of the *PLATO* (PLAnetary Transits and Oscillation of stars, Rauer et al. 2016) mission, which will measure the radii of a large number of terrestrial-sized planets, to detect the radius/density discontinuity and determine its sensitivity ity to model assumptions and fundamental processes. We then perform a parameter study to explore which trades in the survey design of a *PLATO*-like mission maximize its diagnostic power to test runaway greenhouse climate models through the detection of the habitable zone inner edge discontinuity.

We organize the paper as follows: Section 2 introduces the baseline model we use to produce synthetic star and planet samples. In Section 3, we describe the model component that produces runaway greenhouse-induced transit radius changes. Section 4 explains our survey simulations and hypothesis tests. We present our results in Section 5 before interpreting them in Section 6. Finally, we summarize our findings in Section 7.

The goal of this study is to determine – for dif-215 ferent configurations of near-future exoplanet surveys ₂₁₆ – the confidence level with which the runaway green-217 house threshold can be detected statistically. Our basic 218 methodology was as follows: We expanded the Bioverse 219 framework (Bixel & Apai 2020, 2021)² to generate syn-220 thetic samples of stars that host planets according to 221 the observed exoplanet demographics. We then adapted 222 planetary bulk properties as predicted from models of 223 runaway greenhouse atmospheres, simulated observa-224 tions of the planets, and computed Bayesian evidences 225 in favor of a habitable zone inner edge discontinuity 226 (see diagram in Figure 1). In this section, we review 228 the source of the stellar sample, the modeled luminosity 229 evolution, the generation of a synthetic planet sample, 230 and the orbital parameters of the planets. An overview 231 of our key assumptions and model parameters can be 232 found in Table 1.

2.1. Stellar sample from Gaia DR3

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The original Bioverse stellar catalog was generated randomly from the Chabrier (2003) stellar mass function. Improved parallax and photometric data from the Gaia mission made it possible to generate a homogeneous and complete stellar catalog out to about log for Bioverse (Hardegree-Ullman et al. 2023). Here, we briefly describe how we derived the stellar effective temperature $T_{\rm eff}$, luminosity L_{\star} , stellar radius R_{\star} , and stellar mass M_{\star} .

Hardegree-Ullman et al. (2023) used the Gaia Catlateral alogue of Nearby Stars (hereafter GCNS, Smart et al.
lateral 2021) as the basis for deriving stellar parameters for
the Bioverse catalog. The GCNS identified stars out
to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and includes Gaia DR3 parallaxes and pholateral to 120 pc and Gaia Collaboration et al. 2022b) and Gaia DR3 parallaxes and pholateral 2014 parameters from 2015 parallaxes and photometry from 2015 parallaxes and absolute magnitudes and applied initial colorlateral magnitude cuts to remove non-main sequence stars.

We derived effective temperatures primarily from colorlateral parameters table from Pecaut & Mamajek (2013).

Luminosities were computed from absolute G-band mag-

² Bioverse is actively maintained and documented open source software written in Python. Its latest version and documentation can be found at https://github.com/danielapai/bioverse.

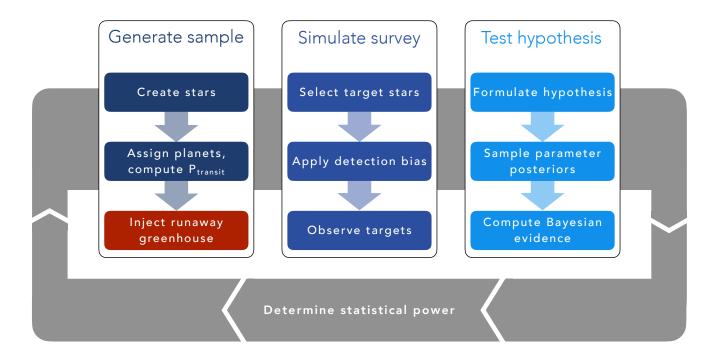


Figure 1. Workflow of our hypothesis testing with **Bioverse**. First, we generate a sample of stars and populate them with planets based on *Kepler* demographics. A fraction of them are then assigned a runaway greenhouse climate based on the model described in Sect. 3. We then simulate an exoplanet survey, whereby selection effects and detection biases are introduced. Finally, we test the runaway greenhouse hypothesis based on data from the survey simulation. By iterating through these steps, we compute the statistical power of testing the hypothesis for different survey designs.

259 nitudes and a derived bolometric correction. We com-260 puted stellar radii with the effective temperatures and 261 luminosities using the Stefan-Boltzmann law or using 262 absolute K_S -band magnitudes and an empirical radiusluminosity relation from Mann et al. (2015) for targets within the absolute magnitude range of M dwarfs. Fi-265 nally, we derived masses from the mass-luminosity rela-266 tion of Torres et al. (2010) for stars with $M_{\star} \gtrsim 0.7 \, M_{\odot}$, 267 and from that of Mann et al. (2019) for targets within 268 the absolute magnitude range of M dwarfs. The derived 269 stellar parameters were compared to measured param-270 eters for all known exoplanet hosts from the literature 271 and were found to be consistent within 1%, 3%, and 272 5.5% for $T_{\rm eff}$, R_{\star} , and M_{\star} , respectively, which are all 273 below the typical measurement uncertainties of 3.3%, 6.8%, and 7.9%, respectively (Hardegree-Ullman et al. 275 2023). From this catalog, Bioverse samples stars within 276 an isotropic distance from the solar system as required 277 by the planetary sample size.

2.2. Stellar luminosity evolution

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Planetary systems are hosted by stars of a wide range of ages, and stellar luminosities evolve with time. Since the occurrence of a runaway greenhouse state is highly

²⁸² dependent on the amount of radiation received by the planet, and thus on the luminosity of the host star, we assigned age-dependent luminosities to our synthetic stars.

While stellar ages are notoriously poorly constrained (e.g., Adams et al. 2005), the age distribution of planet host stars in the Solar neighborhood was shown to be broadly consistent with uniform (Reid et al. 2007; Gaidos et al. 2023). For our synthetic stars, we thus drew random ages from a uniform distribution from OGyr to 10 Gyr. We then assigned each star a luminosity from the mass-dependent evolutionary models of Baraffe et al. (1998). Figure 2 shows the corresponding luminosity evolution as a function of stellar mass and

2.3. Synthetic planet sample

Next, we assigned to the stellar sample planetary systems with frequencies, orbital parameters, and bulk properties derived from the *Kepler* mission. We adopted the model from Bergsten et al. (2022), which defines the occurrence rate of small planets in radius and orbital period. Following Youdin (2011), their inferred occurrence

Table 1. Key assumptions and model parameters used in our simulation setup

Parameter	Value	Unit	Description
Stellar sample			
$G_{ m max}$	16		Maximum Gaia magnitude
$M_{\star,\mathrm{max}}$	1.5	M_{\odot}	Maximum stellar mass
Luminosity evolution			Baraffe et al. (1998)
Planetary parameters			
$M_{ m P}$	0.1 - 2.0	${ m M}_{\oplus}$	Planetary mass range
$R_{ m P,min}$	0.75	R_{\oplus}	Minimum planet radius
Baseline mass-radius relation			Zeng et al. (2016) $100\% \text{ MgSiO}_3^a$
δ_{\min}	80	ppm	Minimum transit depth
P_{\max}	500	d	Maximum orbital period [d]
S	10 - 2000	${ m Wm^{-2}}$	Net instellation range
$S_{ m thresh}$	280	${ m Wm^{-2}}$	Threshold instellation for runaway greenhouse
Runaway greenhouse model			
Runaway greenhouse atmospheric models			Turbet et al. (2020); Dorn & Lichtenberg (2021)
x_{H_2O}	$10^{-5} - 0.1$		Bulk water mass fraction (fiducial case: 0.005)
$f_{ m rgh}$	0-1		Dilution factor (fiducial case: 0.8)
Priors			
$\Pi(S_{ ext{thresh}})$	[10, 1000]	${ m Wm^{-2}}$	uniform
$\Pi(x_{H_2O})$	$[10^{-5}, 0.1]$		log-uniform
$\Pi(f_{\mathrm{rgh}})$	[0, 1]		uniform
$\Pi(\langle R_{ m P} angle_{ m out})$	[0, 15]	R_{\oplus}	Mean radius of non-runaway planets, uniform

 $_{305}$ rate density can be expressed in the form

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$$\frac{\mathrm{d}^2 n}{\mathrm{d}R \,\partial P} = F_0 C_n g(R, P, M_\star),\tag{1}$$

where F_0 represents the average number of planets per star, C_n is a normalization constant, and the shape function $g(R,P,M_\star)$ describes the distribution of planets in radius, orbital period, and stellar host mass. Bioverse generates planets based on the above occurrence rate density and assigns them to the previously generated stars.

2.4. Orbit parameters and planet masses

Eccentric orbits alter the probability of a planet to transit (e.g., Barnes 2007). The distribution of eccentricities e of exoplanets has been found to resemble a Beta function (Kipping 2013), which we chose to draw synthetic eccentricities from. Following Kipping (2013), we used a Beta distribution with parameters a=0.867 and b=3.03, and truncated the distribution at e=0.8. ³²² Assuming isotropic alignments of orbits, we assigned ³²³ each planet an inclination drawn from a distribution uni-³²⁴ form in $\cos(i)$.

To assign masses to our planets, we use the semigraph assuming a pure MgSiO $_3$ composition from Zeng et al. (2016) (see green line in Figure 3). This represents the baseline bulk density before any climate-related effects are applied.

2.5. Transit probability

We model the occurrence of transits by assuming isotropic orientations of planetary orbits and calculating the impact parameters $b=a\cos(i)/R_{\star}$. Following the approach in Bixel & Apai (2021), we further consider only planets with |b|<1. For these cases we calculate the transit depth

$$\delta = \left(\frac{R_{\rm P}}{R_{\star}}\right)^2,\tag{2}$$

^aFor a comparison with alternative interior compositions, see Appendix A.3.

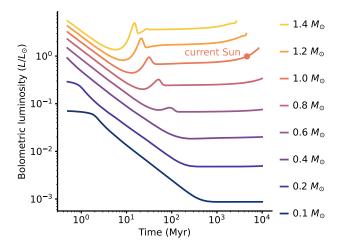


Figure 2. Bolometric luminosity tracks of stars with different masses, computed from stellar evolution models of Baraffe et al. (1998). Low-mass stars, which make up the majority of stars in the solar neighborhood, undergo an extended early phase of several magnitudes higher luminosity before entering a lifetime of relative faintness.

338 which is relevant for the detection probability of the 339 respective planet (see Sect. 4.1). Excluding all non-340 transiting planets diminishes the sample to 1.5% of its 341 original size.

3. RUNAWAY GREENHOUSE MODEL

The climate state of a planet has a direct influ-343 ence on its apparent size measured by transit photome-344 345 try (Turbet et al. 2019, 2020; Mousis et al. 2020; Aguichine et al. 2021). With even a fraction of the Earth's water inventory, a planet absorbing more flux than 348 the radiation limit of steam atmospheres will enter a runaway greenhouse state resulting in a global magma ocean (Lichtenberg et al. 2021; Boukrouche et al. 2021). We use predictions on transit atmospheric thickness from geophysical models to derive the change in transit 353 radius and bulk density that planets with instellation-354 induced runaway greenhouse climates experience, de-355 pending on the distribution of water between planetary 356 interior and atmosphere, and on the resulting thermal 357 atmospheric structure (Dorn & Lichtenberg 2021; Salvador & Samuel 2023). 358

The net absorbed stellar fluxes of planets are a funcmonotonic fintrinsic atmospheric properties such as their monotonic albedo, which are generally poorly constrained for planmonotonic ets outside the solar system (e.g., Angerhausen et al. monotonic planet are considered as a system (e.g., Angerhausen et al. monotonic planet are considered as a planet with a high albedo may monotonic sustain temperate conditions closer to the star than the monotonic same planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further monotonic planet with a lower albedo and located further 368 tion of incoming flux and a fixed Bond albedo of 0.3, 369 comparable to Earth's (Haar & Suomi 1971). We do 370 not take into account additional heating sources such as 371 tidal effects (e.g., Barnes et al. 2013).

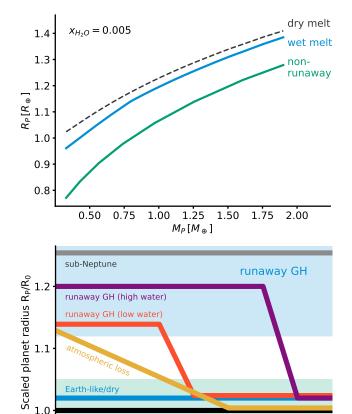
While we search for the signature of runaway greenhouse climates in demographic quantities such as avrage planet radii, the injected changes happen on the
planetary level: We changed each planet's transit radius based on its individual set of properties and the
syr associated predictions from steam atmosphere and wamass M, its net instellation S, and its bulk water inventory expressed as the total planetary water mass fraction x_{H_2O} . We consider the following cases (see Figure 3):

Non-runaway planets retain the radius assigned based on exoplanet occurrence rates (see Sect. 2.3). This case serves as our null hypothesis.

We consider as runaway planets those planets that ab-386 sorb a stellar flux higher than a dayside-averaged thresh- $_{387}$ old instellation $S_{
m thresh}.$ For all planets absorbing an instellation exceeding $S_{\text{thresh}} = 280 \,\text{W}\,\text{m}^{-2}$, we assume 389 an inflated transit radius due to a steam atmosphere. While the actual instellation threshold for a runaway 391 climate depends on planetary albedo, surface gravity, 392 and clouds (Pluriel et al. 2019; Turbet et al. 2021; Pier-393 rehumbert 2022), this value was found to be a typical 394 limit for the flux a planet can emit in a runaway green-395 house situation (Goldblatt et al. 2013; Kopparapu et al. 396 2013; Leconte et al. 2013a; Hamano et al. 2015; Salyador et al. 2017; Katyal et al. 2019; Boukrouche et al. 398 2021; Lichtenberg et al. 2021a). This choice translates 399 to about 1.18 times the instellation of present-day Earth 400 with a fixed albedo of 0.3, which compares favourably to 401 previous climate simulations (Leconte et al. 2013a; Wolf 402 & Toon 2015). We adopt the same threshold instellation 403 for all host star spectral types.

To quantify the radius change, we applied the massradius relationships derived by Turbet et al. (2020) using and a 1D inverse radiative-convective model (Turbet et al. 2019). Their calculations rely on the same mass-radius relations for rocky interiors that we apply for our nonrunaway planets (Zeng et al. 2016). For each planet above the instellation threshold, we assigned the predicted radius for the given water mass fraction and planet mass.

Nominally, the above models assume a *dry melt* without dissolved volatiles. Here, however, we consider a 415 wet melt magma ocean and take into account a radius 416 decrease from retention of water in the melt (Dorn & 417 Lichtenberg 2021). The impact of the water distribu-418 tion between melt and atmosphere on the change of the 419 transit radius depends on the planet's mass and water



Top: Mass-radius relationships for a bulk water mass fraction $x_{H_2O} = 0.005$ and different planet states. Green: planets with a solid mantle and no steam atmosphere. Dashed: planets with steam atmospheres. Blue: planets with steam atmospheres and including the effect of water retention in the melt. Steam atmospheres cause a significant radius increase, which is slightly reduced when water retention in the melt is considered. *Bottom*: Radius evolution of different planet types, illustrating degeneracies and potential for confusion among planet classes. Shown is a schematic time evolution of the transit radius normalized to the atmosphere-free radius R_0 for different scenarios. Planets can move between planet classes through processes such as atmospheric loss and desiccation, which ultimately ends a runaway greenhouse phase on a timescale dependent on a plaret's water content.

no atmosphere

10

content and is generally small compared to the radius inflation from the steam atmosphere. We computed radius deviations between a wet magma ocean and a solid mantle following Dorn & Lichtenberg (2021) and assuming a tropopause pressure $P_{\rm iso} = 0.1$ bar. We then added the (in almost all cases negative) radius deviations to the inflated planet radii computed for the dry melt case.

We illustrate the mass-radius relations of the three asses in Figure 3 where a fiducial bulk water mass frac-

 $x_{H_2O} = 0.005$ is assumed. In the following, we only distinguish between the non-runaway greenhouse and wet melt scenarios.

To account for planets unable to sustain a steam atmo-433 sphere over extended time spans, as well as evolutionary 434 effects such as desiccation through water photodissocia-435 tion and H escape (see bottom panel of Figure 3), we in-436 troduce a dilution parameter $f_{\rm rgh}$. It represents the frac-437 tion of planets above the instellation threshold whose 438 atmospheres are currently inflated due to a runaway 439 greenhouse climate. Our simulation setup is such that 440 all planets receiving a net instellation $S < S_{\text{thresh}}$ follow the non-runaway greenhouse relation, and a fraction f_{rgh} 442 of the planets with $S > S_{\rm thresh}$ follow the wet melt rela-443 tion. The choice $f_{\rm rgh} < 1$ reflects the fact that planets – 444 despite a high irradiation – can evade a runaway green-445 house climate. This situation may, for instance, arise in 446 the absence of an atmosphere or of volatiles that could 447 form a steam atmosphere. In the following, we test if 448 and under what conditions this parametrization causes 449 a demographic trend that is large enough to be detected 450 with high significance.

4. EXOPLANET SURVEY SIMULATIONS AND HYPOTHESIS TESTING

451

452

non-runaway GH

1000

100

Time [Myr]

The survey module of Bioverse converts the synthetic planet sample into a set of uncertainty-laden measurements on a subset of that sample. This task includes
selection of the targets, application of detection biases,
and conducting simulated measurements, all of which
are specific to the particular survey. For each planetlevel measurement such as transit radius or instellation,
we draw the measured value from a normal distribution
centered on the true value with a standard deviation set
by the survey's precision. We then follow a Bayesian hypothesis testing approach to assess various realizations
of simulated surveys in terms of their ability to detect
and characterize the runaway greenhouse transition.

4.1. Detection bias, target selection, and sensitivity

Not all transiting planets are detectable with the same likelihood and detection biases have an impact on the demographic measurements we are interested in. A detailed characterization of the detection biases of individual missions would not be justifiable given the uncertainties of the theoretical predictions. Instead, we derived generic observing limits that reflect the limitations of state-of-the-art transit surveys.

A successful transit detection requires a sufficient signal-to-noise ratio, which is sensitive to the achieved photometric precision. PLATO (PLAnetary Transits and Oscillation of stars) is an ESA mission designed

to characterize terrestrial planets in the habitable zones of Sun-like stars via long-term high-precision photometric monitoring of a sample of bright stars (Rauer et al. 2016). In line with this requirement, PLATO is designed to enable the detection of a 80 ppm transit signal (Matuszewski et al. 2023; ESA 2017). To reflect its sensitivity, we chose a minimum transit depth of 80 ppm as a detection limit and consider only measurements of planets ets exceeding this threshold. We further exclude target stars with Gaia magnitudes $M_{\rm G} > 16$.

The runaway greenhouse effect becomes obsolete both for very small instellations and where no atmosphere can be maintained due to proximity to the host star and resulting atmospheric erosion. Ensuring to stay well clear of such regions, we clear our sample from all planets with a net instellation $S < 10\,\mathrm{W\,m^{-2}}$ or $S > 2000\,\mathrm{W\,m^{-2}}$. We further consider only rocky planets with masses below $2\,\mathrm{M_{\oplus}}$.

4.2. Measurements and their uncertainties

Under real-world conditions, the planetary properties in question can only be probed with a finite precision that is specific to each exoplanet mission. PLATO's definition study report (ESA 2017) states precision requirements for planet radii (3%), planet masses through radial velocity (RV) follow-up (10%), and stellar masses, radii, and ages (10%). We adopted these estimates and assumed a 10% error on instellation measurements.

Since planetary bulk density $\rho \propto R_{\rm P}^{-3}$, we expect a stronger runaway greenhouse signal when measured through bulk density instead of transit radius. We thus simulated measurements of planetary densities assuming the mass-radius relation defined above. For uncertainties in bulk density measurements, we propagated the errors of the mass measurements assuming $\sigma_{\rm M_P} = 10\,\%$.

4.3. Hypothesis tests

We now turn to quantifying the ability of the simulated surveys to detect the habitable zone inner edge discontinuity and to constrain parameters associated to the runaway greenhouse transition. To do this, we rely on a Bayesian hypothesis testing approach where we quantify the evidence of a hypothesis over another based on the (simulated) data. For our specific problem, this implies comparing evidences for a demographic imprint of the runaway greenhouse effect to its absence. As a null hypothesis, we consider the case where the planetary radius distribution is independent of the instellation,

$$H_0(\theta, S) = \theta, \tag{3}$$

where θ is the set of parameters defining the radius distribution. We further define an alternative hypothesis

that describes radius changes due to runaway greenhouse climates and inflated steam atmospheres. As motivated above, this hypothesis takes the form of a step function in net instellation S, where the step occurs at the outer edge of the runaway greenhouse region. Our main observable shall be the average transit radius in the planet population on either side of this threshold. The runaway greenhouse hypothesis is then defined as

$$H_{\text{rgh}}(\theta, S) = \begin{cases} H_0, & S \leq S_{\text{thresh}} \\ \langle R_{\text{P}} \rangle (f_{\text{rgh}}, \Delta R_{\text{stm}}, \Delta R_{\text{wtr}}), & S > S_{\text{thresh}}. \end{cases}$$
(4)

Here, $f_{\rm rgh}$ is the fraction of planets above the instellation threshold experiencing a runaway greenhouse effect. $\Delta R_{\rm stm}$ and $\Delta R_{\rm wtr}$ are predicted radius changes from the steam atmosphere and water retention models, respectively. They are assumed to act additively on the planet radii and thus on their average $\langle R_{\rm P} \rangle$.

The only free parameter of the null hypothesis, which assumes the average transit radius to be independent of instellation, is the predicted mean radius $\langle R_{\rm P} \rangle$. The functional form of the runaway greenhouse hypothesis is more complex: Besides the mean radius of planets outside the threshold $\langle R_{\rm P} \rangle_{\rm out}$, which is a nuisance parameter necessary to define the hypothesis, it relies on the threshold instellation for the "step" $S_{\rm thresh}$, the planetary bulk water mass fraction x_{H_2O} , and the dilution factor $f_{\rm rgh}$. For hypothesis tests based on bulk density instead of radius, we proceeded in the same way and substituted $R_{\rm P}$ by the bulk density ρ .

A sensible choice of priors is central for evidence estimation via nested sampling. As the parameters of interest are poorly constrained by previous data, we used relatively uninformative priors to sample the entire physically plausible parameter space. For $S_{\rm thresh}$, we chose a uniform prior in $[10,1000]~{\rm W\,m^{-2}}$. We sampled x_{H_2O} from a log-uniform distribution to imply scale-invariant ignorance. Its boundaries $[10^{-5},0.1]$ are motivated by the water mass fractions covered by the geophysical models (Sect. 3). For $f_{\rm rgh}$, we chose a uniform prior in [0,1]. Finally, we adopted a broad, uniform prior for $\langle R_{\rm P} \rangle_{\rm out}$ bound by $[0.1,15]~{\rm R}_{\oplus}$. In the case of measuring bulk densities instead of transit radii, we drew uniformly from $[1,6]~{\rm g\,cm^{-3}}$.

The measured radii $R_{\rm P,i}$ or bulk densities $\rho_{\rm i}$ cannot be directly used for the hypothesis tests as they include intrinsic scatter that is not caused by measurement errors. $H_{\rm rgh}$ and H_0 should thus be tested against a statistical estimator that represents the population mean. To avoid binning and the artificial patterns it may introduce, we those to test our hypotheses against a simple moving average SMA along the instellation axis with a window

of size 25 centered around each measurement.³ We fur-578 ther computed the uncertainty of this moving average 579 by propagating the individual measurement errors and 580 applying a rolling standard error of the mean.

As our procedure involves random sampling of the model parameters θ , we need to define the probability 583 of obtaining a data set given the model parameters, i.e., ⁵⁸⁴ a likelihood function \mathcal{L} . We assumed here that the in-585 dividual moving averages SMA_i are measured with a 586 normally distributed uncertainty σ_{SMA_i} and adopted a 587 normal distribution

$$\mathcal{L}(SMA \mid \boldsymbol{\theta}) = \prod_{i}^{N} \frac{1}{\sqrt{2\pi\sigma_{SMA_{i}}^{2}}}$$
 (5)

$$\mathcal{L}(SMA \mid \boldsymbol{\theta}) = \prod_{i}^{N} \frac{1}{\sqrt{2\pi\sigma_{SMA_{i}}^{2}}}$$

$$\times \exp\left(-\frac{(SMA_{i} - H(\boldsymbol{\theta}, S_{i}))^{2}}{2\sigma_{SMA_{i}}^{2}}\right).$$
 (6)

Here, $H(\theta, S_i)$ corresponds to the functional form of the ⁵⁹¹ runaway greenhouse or null hypothesis.

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4.4. Bayesian model comparison

We can now assess the relative plausibility of $H_{\rm rgh}$ 593 ₅₉₄ and H_0 given the synthetic data we have generated, 595 assigning equal a priori probabilities to these models. This is done by comparing the Bayesian evidence $\mathcal Z$ of the models, which we estimated with the nested sampling (Skilling 2004) algorithm dynesty (Speagle 2020). We initialized the sampler with the priors defined above 600 to let it estimate the evidence and sample the posterior distributions. Our criterion to reject the null hypothesis 602 is

$$\Delta \ln \mathcal{Z} = \ln \mathcal{Z}_{rgh} - \ln \mathcal{Z}_0 > 3. \tag{7}$$

5. RESULTS

5.1. Statistical signature of the runaway greenhouse threshold

To characterize the population-level imprint of indi-608 vidual radius changes, we generated a generic planet 609 population with an injected runaway greenhouse effect assuming a water fraction $x_{\rm H_2O} = 0.005$. Figure 4 shows 612 the resulting planetary radii. When ordered in orbital 613 period space, the different planet types overlap, dilut-614 ing the demographic imprint. With net instellation as 615 an independent variable, planets above and below the 616 runaway greenhouse threshold separate: The runaway 617 greenhouse-induced radius inflation introduces a discon-618 tinuity of average planet radii and bulk densities as a 619 function of stellar irradiation. We also show the pre-620 dictions of observable average planet radii from the sta-621 tistical hypotheses defined above. Within the runaway

 $_{622}$ greenhouse regime, an average radius change of 15 % oc-623 curs. This pattern is consistent with the injected radius 624 inflation as predicted from the atmospheric models (see 625 Appendix A.2 for an investigation of the interplay be-626 tween model predictions and our synthetic planet pop-627 ulation).

5.2. Testability of the runaway greenhouse hypothesis

Figure 5 shows a prototypical statistical detection of 629 630 the habitable zone inner edge discontinuity. Before we 631 study limiting cases of such a detection below, we first 632 demonstrate the interpretation process based on an op-633 timistic scenario where the sample of characterized plan-634 ets is large (N = 500) and the measurement uncertainties are small ($\sigma_R = 2\%, \sigma_S = 5\%$). Here, we 636 assumed that the fraction of those planets irradiated 637 stronger than S_{thresh} that have runaway greenhouse climates is $f_{\rm rgh}=0.8$, and we chose a water mass fraction of $x_{H_2O} = 0.005$ for each planet. In this case, the hab-640 itable zone inner edge discontinuity was detected with high significance ($\Delta \ln \mathcal{Z} \approx 100$).

With such a strong signal, we can attempt an inference 643 of the parameters defining the injected effect. Figure 6 shows the posterior distributions of S_{thresh} , x_{H_2O} , and $f_{\rm rgh}$ as determined by the nested sampler. The thresh-646 old instellation can be accurately constrained. Both a 647 higher water mass fraction and a higher dilution factor 648 lead to larger average radii, thus these parameters are 649 strongly correlated.

Figure 7 explores the statistical power of the hypothe-651 sis test achieved in the above scenario for different com-₆₅₂ binations of the poorly constrained parameters x_{H_2O} $_{653}$ and $f_{\rm rgh}$. It is highest for large water inventories and 654 large dilution factors. For all but very low water frac-655 tions, $f_{\rm rgh}$ dominates this trend: It enters linearly into 656 the average planet radius, whereas the contribution of x_{H_2O} - as predicted by the geophysical models - is sub-658 linear with a power-law exponent of ~ 0.3 . Within the framework of our model and as long as $f_{
m rgh}$ is larger 660 than ~ 0.2 , a sample size of 500 is sufficient for a 50 % detection rate even for water ratios as low as 10^{-3} .

5.3. Detecting the runaway greenhouse transition with PLATO

To simulate the transit survey of ESA's PLATO mis-665 sion, we considered a volume-limited sample with a size 666 according to projections and including all stellar spec-667 tral types. Given PLATO's expected radius precision, we find that a yield of ~ 300 is needed for a significant 669 detection if the fraction of runaway greenhouse planets $f_{\rm rgh} = 0.1$ (see Figure 8). The minimum needed fraction ₆₇₁ rises to 0.2 for N = 100. For much smaller samples, only 672 an optimistically strong signal is likely to be detected.

³ See Appendix A.1 for a robustness test using a different estima-

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Synthetic planets

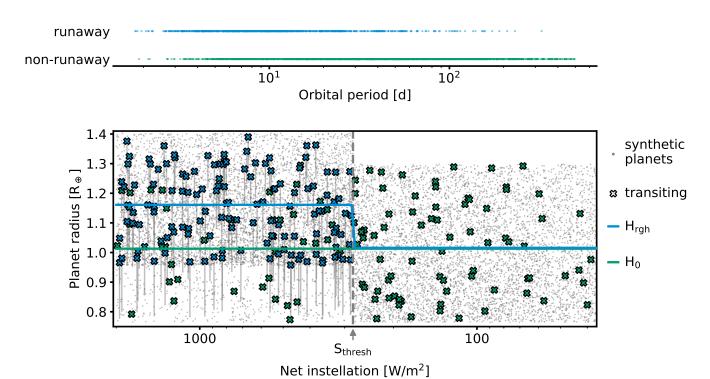


Figure 4. Synthetic planets above and below the runaway greenhouse threshold. Top: Planet state as a function of orbital period. Planets with and without a runaway greenhouse climate mix and are not distinguishable in orbital period space. Bottom: Transit radii of synthetic planets with injected radius deviation as a function of net instellation. Only the planets marked as transiting are observable. Above the runaway greenhouse threshold $S_{\text{thresh}} = 280 \,\text{W m}^{-2}$, some planets maintain their original radii (green crosses) whereas some have their transit radius inflated (blue crosses) by the amount indicated with gray lines. The sharp boundary at S_{thresh} causes a discontinuity in the average planet radius (blue line). This runaway greenhouse hypothesis call be tested against the null hypothesis H_0 (green line), where average radii are independent of instellation.

5.4. Statistical power of different mission designs 5.4.1. Additional planet mass measurements

Comparing a measurement of the habitable zone infrom her edge discontinuity in radius space with a measurefrom ment in density space (which requires planetary mass from measurements), a stronger detection occurs in the latter from case: With an optimistic choice of geophysical paramefrom ters (see Sect. 5.2), the average measured radius change from is 15% whereas the average density change is -33%.

Figure 8 shows that investigations of the discontinuity are more constraining when radius measurements can be augmented with mass measurements: At unchanged sample size, the difference in Bayesian evidence can be up to an order of magnitude larger. Consequently, we achieve a statistically significant detection with smaller samples or lower dilution factors $f_{\rm rgh}$. A density-based survey of 100 targets is roughly equivalent to a radius-based survey of 300 targets. At N=100, pure radius measurements require $f_{\rm rgh}\gtrsim 0.2$ whereas bulk density measurements enable a detection from $f_{\rm rgh}\gtrsim 0.1$.

5.4.2. Dependence on host star spectral type

Since the incident radiation at a given orbital distance 695 depends on the spectral type of the host star, the relative 696 number of planets on either side of $S_{
m thresh}$ is different for 697 FGK and M dwarfs. We tested the detectability of the 698 habitable zone inner edge discontinuity when only FGK 699 or only M dwarfs are considered (see Figure 8). All other 700 parameters of the climate models are kept the same. 701 The samples are volume and magnitude-limited to re-702 flect the target counts of PLATO's provisional Long-703 duration Observation Phase fields (15996 FGK stars in 704 the P1 and P2 samples, 33948 M stars in the P4 sam-705 ple, Nascimbeni et al. 2022). The resulting M dwarf 706 planet sample is significantly larger with 228 ± 14 plan-707 ets compared to 40 ± 6 planets in the FGK sample. No 708 significant detection is possible in the pure FGK sample, 709 independent of the assumed geophysical parameters. In 710 the M dwarf sample, the evidence threshold is reached 711 around $f_{\rm rgh} \sim 0.2$, similar to the case above where all 712 spectral types are considered.

Optimistic survey (500 planets)

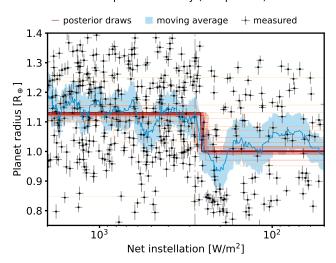


Figure 5. Detection of the runaway greenhouse threshold. From simulated radius and instellation measurements of a large (N=500) survey, we compute the moving average (blue confidence intervals) and fit the runaway greenhouse hypothesis to it (Eqn. 4, random draws from the posterior in \mathbb{R}^{2} d). The pattern is detected with high significance.

5.4.3. Constraining the threshold instellation

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Although the exact location of the runaway green-715 house transition depends on the stellar energy actually $_{716}$ absorbed by the planet, a measurement of $S_{
m thresh}$ would 717 be a key constraint resulting from a detection of the 718 habitable zone inner edge discontinuity. Here, we as-719 sess the ability of different mission concepts explored 720 above to constrain this parameter. Figure 9 shows posterior distributions from a grid of inferences together 722 with the true values of the injected signal for different fractions $f_{\rm rgh}$. We consider three cases: only radius 724 measurements, radius and mass measurements, and ra-725 dius and mass measurements of only planets orbiting 726 M dwarfs. The simulations otherwise represent the sim-727 ulated PLATO example described above. We chose a planetary sample size of $N=100\pm10$, as this was found 729 to be a threshold case in Sect. 5.3.

We find that retrievals from radius measurements alone require high fractions of greenhouse climate-bearing planets to achieve an accurate constraint on $S_{\rm thresh}$: dilution factors $f_{\rm rgh}\gtrsim 0.5$ yield posterior probability distributions that are condensed at the order of magnitude of the true value; accurate constraints to within $\pm 0.25\,{\rm dex}$ of the truth are reached only from $f_{\rm rgh}\approx 0.8$.

In contrast, if planet masses are available and the hypothesis test is conducted in bulk density space, useful

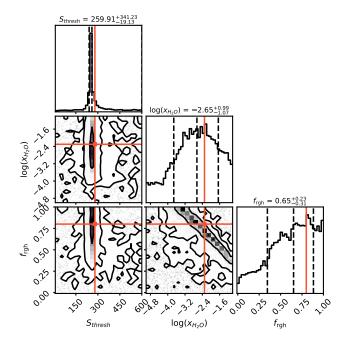


Figure 6. Retrieved posterior distribution of key parameters in the optimistic scenario. The density maps in each panel show relationships between and marginalized distributions of the threshold instellation $S_{\rm thresh}$, the bulk water mass fraction x_{H_2O} , and the dilution factor $f_{\rm rgh}$ as they could be retrieved with a high-precision transit survey and a sample of 500 planets. True values of the parameters for the injected effect are shown in orange. The threshold instellation can be reasonably constrained; the predominant water fraction and the fraction of planets with runaway greenhouse climates are degenerate.

740 constraints emerge already from about $f_{\rm rgh} \approx 0.35$. Ac-741 curacy and precision of the retrievals are improved.

A sample containing only planets around M dwarf yields still better performance and result in the overall best accuracy and precision. The reason is that these planets contribute most to the statistical power of the retrieval.

6. DISCUSSION

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6.1. Statistical imprint of exoplanet climates

We showed in Sect. 5.1 that injecting the theoretically predicted radius inflation effect into a synthetic planet population following the currently known demographics leaves a distinct pattern in the radius and density distribution as a function of instellation. The transity at radius inflation was suggested before as an observational diagnostic to probe the runaway greenhouse transition (Turbet et al. 2019), and our simulations show quantitatively how the contributions of the individual radius changes combine to create a significant demographic signal. Its strength depends on largely unknown

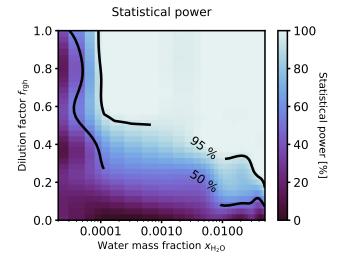


Figure 7. Statistical power of the runaway greenhouse hypothesis test as a function of model parameters. For a sample size N=500, the color code shows the fraction of simulations resulting in a sound detection $(\Delta \ln \mathcal{Z} > 3)$ for different combinations of bulk water mass fraction and dilution factor. Higher values in either parameter result in a more reliable detection. For water mass fractions $\gtrsim 10^{-4}$, the statistical power largely depends on the fraction of greenhouse climate placets in the sample.

⁷⁶⁰ factors such as the planetary volatile content, but turned out to be well in the detectable range under reasonable assumptions of these factors.

Our finding that the expected habitable zone inner 764 edge discontinuity is strong except for extremely low 765 water mass fractions gives reason for optimism regard-766 ing its detection. It also presents a potential for con-767 straints on the water inventory of terrestrial planets in 768 the case of a non-detection. Such constraints will provide insight into whether the initial water content of 770 rocky planets is varied by systematic inter-system ef-771 fects (e.g., Raymond et al. 2004; Mulders et al. 2015; 772 Sato et al. 2016; Lichtenberg et al. 2019; Lichtenberg 773 & Krijt 2021; Lichtenberg & Clement 2022) or if there 774 is a predominant pattern of volatile-enrichment across 775 planetary systems that is only modified by intra-system 776 effects such as planet migration (Schlecker et al. 2021a) 777 or atmospheric escape (Owen & Wu 2016). Ultimately, 778 such measurements will shed light on how delivery and 779 loss effects during rocky planet formation and evolution shape the diversity of exoplanetary climates. Overall, 781 probing the runaway greenhouse discontinuity appears 782 to be the most promising approach toward a first em-783 pirical test of the habitable zone concept.

6.2. Detectability of the habitable zone inner edge discontinuity

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We showed in Sect. 5.2 that under favorable conditions, a sufficiently large (500 planets) photometric surtions, a sufficiently large (500 planets) photometric surtions runaway greenhouse transition and accurately constrain
the associated threshold instellation. Of course, even astions under the survey design, our current models may not capture fully the
complexity of trends in the geophysics and demographtics of small planets. The detectability of the runaway
tics of small planets. The detectability of the runaway
tion and the diverse outcomes of rocky planet formation and evolution. In the following, we will explore
the key factors that may influence the emerging demo-

6.2.1. Key factors influencing tests of the runaway greenhouse hypothesis

What influences the probability of correctly rejecting a false null hypothesis? We revidentified six drivers of the diagnostic power for detecting the runaway greenhouse transition with a transit survey. Many of these factors directly influence the fraction of planets currently in runaway greenhouse states, which in our model is represented by the dilution factor $f_{\rm rgh}$.

- Occurrence rate of planets forming steam atmospheres
- Planetary evolution and duration of the steam atmosphere phase
 - Prevalent water inventory

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- Size and composition of the planetary sample
 - Radius measurement precision
 - Availability and precision of mass measurements

 $_{\mbox{\scriptsize 817}}$ We will now briefly explore the above drivers.

Occurrence rate of planets forming steam atmospheres:
The runaway greenhouse climate relies on sufficient amounts of atmospheric water vapor that can act as a greenhouse gas. However, already about ~10–20 bar of water vapor – corresponding to a minor fraction of one Earth ocean and thus the lower limit of water on Earth – is enough to sustain sufficiently high surface temperatures to keep the planet in a magma ocean stage (Boukrouche et al. 2021; Lichtenberg et al. 2021). The fraction of planets fulfilling these requirements has an impact on the amplitude of the demographic imprint of the runaway greenhouse transition. From a planet formation perspective, the incorporation of water into planets in the terrestrial planet zone is a standard expected outcome (e.g., Zeng et al. 2019; Venturini et al. 2020;

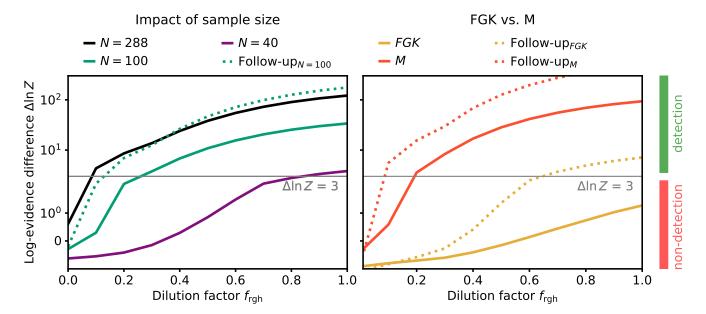


Figure 8. Expected delta-evidences as a function of the fraction of planets with runaway greenhouse climates for different versions of the PLATO survey. The median values of randomized survey simulations are shown; $\Delta \ln Z > 3$ (gray horizontal line) is considered sufficient evidence to reject the null hypothesis. Left: For a large planet yield of $N \approx 300$, even small dilution factors ~ 0.1 allow a detection. A sample of 100 planets is sufficient if their masses are constrained to within 10% (dotted green line). Without such follow-up measurements, sufficient diagnostic power can only be achieved with this sample if $f_{\rm rgh} \gtrsim 0.2$. Even smaller samples are unlikely to yield a significant detection. Right: Evidences when only FGK or only M dwarfs are considered. Only M dwarfs host enough planets on both sides of the threshold instellation to allow a reliable detection of the habitable zone inner edge discontinuity.

Emsenhuber et al. 2021; Schlecker et al. 2021a; Burn 834 et al. 2021). But while commonly considered volatile 835 delivery channels suggest a fraction of planets to be volatile-poor, the incorporation of hydrogen into even 837 the driest planetary materials known in the solar sys-838 tem (McCubbin & Barnes 2019; Piani et al. 2020; Jin et al. 2021) suggests that hydrogen is present in all rocky 840 planets upon formation. Accreted hydrogen in nomi-841 nally dry planetary materials react with mantle oxygen 842 to form substantial amounts of water inside of the planet 843 during the magma ocean phase (Ikoma et al. 2018; Kite 844 & Schaefer 2021; Kimura & Ikoma 2020, 2022). En-845 hanced equilibration between the core, mantle, and at-846 mosphere during magma ocean evolution of rocky exoplanets further enhances this process (Lichtenberg 2021; 848 Schlichting & Young 2022). Therefore, even nominally 849 dry planets generate substantial amounts of water dur-850 ing formation and early evolution. Reduced heating 851 from short-lived radionuclides in extrasolar planetary 852 systems increases the expected water abundance in ex-853 oplanet systems further, in particular for M dwarf systems (Lichtenberg & Clement 2022).

Planetary evolution and duration of the steam atmosphere can only be sphere phase: An inflated steam atmosphere can only be sustained until the planet has lost its water. Depending on host star spectral type, planetary mass, and composite

859 sition, planets can spend from a few Myr to several Gyr 860 in runaway greenhouse climates (Hamano et al. 2015; 861 Luger & Barnes 2015), and only planets observed dur-862 ing this phase will contribute to the habitable zone inner 863 edge discontinuity. However, the delivery uncertainty is 864 much greater than the predicted loss rates of water by 865 atmospheric escape, which typically is limited to in to-866 tal a few tens of terrestrial oceans per Gyr (Wordsworth 867 et al. 2018) on even the most irradiated exoplanets (see 868 discussion in Lichtenberg & Clement 2022). The finite 869 duration of steam atmosphere phases is a main factor 870 for diluting the habitable zone inner edge discontinu-871 ity, and it leads to an increased likelihood of detecting 872 this imprint for younger planetary systems. Preferen-873 tial selection of younger systems could therefore be of 874 advantage in a targeted survey.

Prevalent water inventory: The magnitude of radius change at the runaway greenhouse threshold is sensitive to the water mass fraction. As a result, the statistical abundance of water in terrestrial planets impacts the strength of the demographic pattern: The higher the water content and the higher the fraction of planets in runaway greenhouse climates, the greater the likelihood that a discontinuity is detectable. However, dissolution into the magma ocean and atmospheric inflation show non-linear coupling: the more water there is in the atmo-

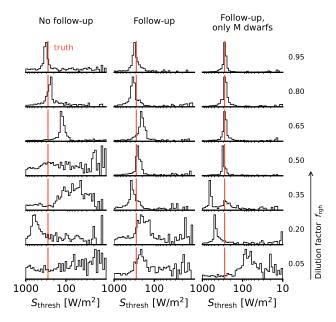


Figure 9. Retrieved posterior distributions of the threshold instellation for different survey realizations. All cases assume $x_{H_2O}=0.005$ and a planet sample size $N=100\pm10$; The fraction of planets with runaway greenhouse climates varies across rows. Orange lines show the true value of the injected signal. Accuracy and precision of the constraint on $S_{\rm thresh}$ generally improve with higher $f_{\rm rgh}$. Bulk density-based inferences improve the constraints, and M dwarf samples yield the highest accuracy and precision.

sphere, the more will be dissolved (Dorn & Lichtenberg 2021), leading to a convergence in radii until saturation. Most important for the present work, however, is the qualitative dichotomy between sub-runaway and runaway planets, which outcompetes radius variations within each of these climate states due to a changing water mass fraction (Turbet et al. 2020). If a planet hosts a significant water inventory that is not easily stripped by atmospheric escape (Johnstone 2020), then the statistical power of our model is high (see Figure 7).

Size and composition of the planetary sample: The significance of a statistical trend increases with a larger sample size. In addition, the sample must include planesses on both sides of the instellation threshold. This is only likely for low-mass host stars due to the strongly distance-dependent detection bias associated with the transit method and the temporal coverage of upcoming transit missions.

Radius measurement precision: The more precise individual planet radii can be determined, the more pronounced the discontinuity will be. Good accuracy is less important, as long as it does not have a systematic error scaling with stellar irradiance. Availability and precision of mass measurements: For simple geometric reasons ($\rho \propto R^{-3}$), the expected dissociation continuity at the habitable zone inner edge is stronger when measured in bulk density than it is in planet radius space. If transiting planets can be followed up to obtain mass measurements, the statistical significance increases.

Besides these main factors, uncertainties in the measured instellations can influence the result, although they are typically small due to the very precise orbital period measurements available for transiting planets. This can be different for young host stars when their ages cannot be well constrained; in particular, the long pre-main sequence phase of M dwarfs shows a large variation in bolometric luminosity (see Figure 2).

6.2.2. False positive scenarios

Runaway greenhouse climates are not the only physi-925 cal mechanism that may cause a change in transit radius 926 for a subset of planets. Alternatives include atmospheric 927 loss due to either photoevaporation through high-energy 928 radiation by the host star (e.g., Ikoma & Hori 2012; 929 Owen & Wu 2013; Jin et al. 2014; Mordasini 2020) or 930 due to residual heat from the planet's interior shortly 931 after formation (Ginzburg et al. 2016, 2018; Gupta & 932 Schlichting 2019). Both processes are being discussed as 933 potentially sculpting the observed bimodality in the ra-934 dius distribution of small exoplanets (Fulton et al. 2017; 935 Van Eylen et al. 2018), and both lead to a decrease of 936 planet radius for planets close to their host star (Pas-937 cucci et al. 2019; Bergsten et al. 2022). This is distinct 938 from the radius inflation introduced by runaway green-939 house climates. For example, the innermost planet in 940 the K2-3 system has an increased radius compared to 941 its outer siblings, contrary to what would be expected 942 from atmospheric escape (Diamond-Lowe et al. 2022).

Other false positive contributions may stem from po-944 tential unknown occurrence rate gradients in radius-945 instellation space, especially if these variations are sim-946 ilar to the expected habitable zone inner edge discon-947 tinuity. Although an abrupt pattern at the expected 948 location of the transition seems unlikely, examples of 949 steep occurrence rate density changes exist. An example 950 is the "Neptune desert", a triangular region in period-951 radius space of low planet occurrence (Szabó & Kiss 952 2011; Mazeh et al. 2016; Dreizler et al. 2020). The shape 953 of this region is such that smaller planets become less 954 frequent the closer to the star they are, which to some 955 degree resembles the pattern introduced by the instella-956 tion dependency of the runaway greenhouse transition. 957 However, the Neptune desert occurs at smaller orbital 958 periods and is sensitive to the planet radius (Szabó &

959 Kiss 2011), which is not expected for the runaway green-960 house transition.

Luque & Pallé (2022) found that small planets or-962 biting red dwarfs can be classified into three density regimes with a particularly strong separation between 964 planets consistent with a pure rocky and those consis-965 tent with a water-rich composition. This trend does 966 not represent a false positive scenario for the habitable 967 zone inner edge discontinuity, since no strong dependency on instellation has been found or is expected. population of "water worlds" with low bulk densi-970 ties on a wide range of orbits would merely attenu-971 ate the statistical runaway greenhouse imprint. If the 972 dichotomy forms primordially through migrated plan-973 ets that accreted from different regions of their proto-974 planetary disk (Venturini et al. 2020; Burn et al. 2021; 975 Schlecker et al. 2021a,b) or from inter-system variations 976 in the desiccation of volatile-rich planetesimals (Lichtenberg et al. 2019; Lichtenberg et al. 2021b; Lichtenberg 978 & Krijt 2021; Lichtenberg & Clement 2022; Bonsor et al. 979 2023), systems of all ages can be affected by this atten-980 uation.

6.2.3. Atmospheric spectral signatures

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We currently see four potential lines of discriminat-982 983 ing atmospheric signatures associated with runaway cli-984 mates in the exoplanet population. (i) Detecting the 985 atmospheric windows of water vapor in the near- to 986 mid-infrared. In a runaway greenhouse atmosphere, ab-987 sorption is dominated by the opacity of water vapor, 988 which has two prominent spectral features: one at 3.5-989 4.5 μ m, and one between 8–20 μ m (e.g., Boukrouche 990 et al. 2021). Probing these features requires an in-991 strument covering these wavelength ranges, for instance 992 ELT METIS (Brandl et al. 2021), JWST MIRI (Rieke 993 et al. 2015), or future missions such as LIFE (Bonati et al. 2019; Quanz et al. 2022; Dannert et al. 2022). Po-995 tentially the two-band filter capabilities of PLATO may 996 offer insight into particularly pronounced spectral features in this runaway greenhouse regime (Grenfell et al. 998 2020). However, detecting water features in steam at-999 mospheres may be obscured by high-altitude clouds that could potentially mute a transit signal as measured from peak to trough of an observed spectral line (Suissa et al. 1002 2020; Fauchez et al. 2019). A better understanding of the role of clouds in this problem may be critical for observational constraints on the runaway climate of individual exoplanets. (ii) Post-runaway planets may be 1006 detectable through O₃ absorption due to build-up of abi-1007 otic oxygen, leftover from photochemical dissociation of water, and hydrogen loss (Wordsworth & Pierrehum-1009 bert 2014; Luger & Barnes 2015). (iii) Water loss in

1010 runaway greenhouse episodes would increase the atmo-1011 sphere's D/H isotopic ratio, akin to the enhancement 1012 in Venus' present-day atmosphere (Kane et al. 2019, 1013 2021). This may be detectable in high-resolution ob-1014 servations focusing on isotope-sensitive transitions (Lin-1015 cowski et al. 2019; Mollière & Snellen 2019). (iv) Dise-1016 quilibrium chemistry in tidally-locked runaway planets. 1017 It has been suggested that the atmospheric depth and 1018 the presence or absence of oceans on "sub-Neptunes" 1019 could be probed via the abundance of species that are 1020 photochemically destroyed in the upper atmosphere, and 1021 replenished from either thermochemical layers (Yu et al. 1022 2021; Tsai et al. 2021) or at an ocean-atmosphere inter-1023 face (Loftus et al. 2019; Hu et al. 2021). Atmospheric 1024 nitrogen and carbon compounds can be partitioned into 1025 magma (Grewal et al. 2022; Bower et al. 2022). There-1026 fore, it may be possible to discern the presence of an un-1027 derlying magma ocean if the presence of an atmosphere 1028 on a rocky exoplanet can be confirmed. This has been 1029 suggested to be done via (v) eclipse photometry (Mans-1030 field et al. 2019; Koll et al. 2019) through the presence 1031 of a high albedo, which is expected to differ from the 1032 crystallized rock of a solidified magma ocean (Essack 1033 et al. 2020; Fortin et al. 2022). Spectral information 1034 from atmospheres of highly irradiated planets may also 1035 help to distinguish classical runaway greenhouse states 1036 from other climate regimes, such as a moist bistabil-1037 ity (Leconte et al. 2013b).

6.2.4. Detection in multi-planet systems

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As suggested by Turbet et al. (2019), an alterna-1040 tive approach for detecting the runaway greenhouse-1041 induced radius inflation is to search for its "local" 1042 imprints in multi-planet systems. Systems harbor-1043 ing planets on both sides of the transition, such as 1044 TRAPPIST-1 (Gillon et al. 2016, 2017; Luger et al. 1045 2017; Agol et al. 2021), K2-3 (Diamond-Lowe et al. 1046 2022), or Kepler-138 (Piaulet et al. 2022), may show 1047 the predicted abrupt radius and density change, provided the initial volatile content was sufficient and com-1049 plete desiccation has not yet occurred. While degen-1050 eracies remain in interpreting bulk density fluctuations within individual systems (e.g., Turbet et al. 2020; Dorn 1052 & Lichtenberg 2021), the detection of a consistent pat-1053 tern in several such systems could be a convincing sta-1054 tistical evidence of the runaway greenhouse transition. 1055 The current sample of suitable systems is sparse: The California-Kepler Survey catalog (Fulton & Petigura 1057 2018) contains only six planets with instellations $< 2 S_{\oplus}$ 1058 and smaller than 2 R⊕ in five multi-planet systems. Fu-1059 ture additions to the multi-planet sample through mis-1060 sions such as PLATO are needed.

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6.3. Diagnostic power of near-future exoplanet missions

Confirming or disproving the predicted habitable zone 1063 inner edge discontinuity will depend on the significance 1064 with which the null hypothesis can be excluded, which is a function of instrumentation and survey strategy. As discussed in Section 6.2.1, key drivers from a mission design perspective are sample size, photometric preci-1069 sion, and the availability of planets around low-mass 1070 host stars. We found that along these axes, PLATO will be the most favorable among the upcoming transit 1072 missions. The PLATO team has released an estimate 1073 on the number of exoplanets that will be characterized 1074 in the course of the main survey mission. With an ex-1075 pected transit radius precision of 3 % (ESA 2017) for 1076 hundreds of planets (Rauer 2021), the PLATO mission 1077 is comparable to the optimistic survey (see Section 5.2) 1078 in terms of sample size and precision. If successful, it 1079 should readily detect the predicted statistical imprint or, 1080 in case of a non-detection, provide strong upper limits on the occurrence rate of runaway greenhouse planets. The latter depends on the lifetimes of runaway green-1083 house phases, which are a function of the initial water inventory of the planets (Hamano et al. 2015). Overall, 1085 it seems feasible to derive the typical water content of 1086 low-mass exoplanets from these occurrence estimates.

What other planned missions are suited to probe 1087 1088 the habitable zone inner edge discontinuity? and K2 have contributed a large number of discovered terrestrial-sized planets, but few of them are in the habitable zone and their host stars are typically too faint for 1092 RV follow-up with current instrumentation (Dressing & Charbonneau 2015).

Similarly, the Transiting Exoplanet Survey Satel-1095 lite (TESS, Ricker et al. 2014) planet sample lacks temperate, small planets around bright host stars (Ment & Charbonneau 2023), as was expected from planet yield calculations (Barclay et al. 2018). As of March 22, 2023, the NASA Exoplanet Archive⁴ lists 40 TESS candidate or confirmed planets smaller than $4R_{\oplus}$ with lower esti-1101 mated instellation than Earth's.

The ongoing CHEOPS mission was designed as a 1102 1103 follow-up mission to search for transits of planets discovered with other techniques, in particular with radial velocity measurements (Benz et al. 2021). As such, it 1106 will provide precise radius constraints on a sample of small planets; however, only a small number of plan-1108 ets with orbital periods >50 d are being observed. This 1109 largely limits CHEOPS' coverage to planets within the 1110 runaway greenhouse regime, preventing a detection of 1111 the transition.

As CHEOPS, the Atmospheric Remote sensing In-1113 frared Exoplanet Large survey (Ariel, Puig et al. 2016) 1114 will be a follow-up mission that is not designed to pro-1115 vide a large number of new radius measurements. Ariel's 1116 primary targets are larger planets in the range of sub-1117 Neptune to Jupiter-like planets. We thus do not expect 1118 a significant contribution to statistically exploring the inner edge of the habitable zone for Earth-sized planets. While not primarily designed to detect transiting planets, the Galactic Bulge Time Domain Survey of the 1122 Nancy Grace Roman Space Telescope (Spergel et al. $_{1123}$ 2015) is expected to yield $\sim 10^5$ transiting planets on 1124 short orbits and constrain their radii in the course of its mission (Montet et al. 2017). $\mathcal{O}(1000)$ planets smaller 1126 than Neptune could be found around early to mid-1127 M dwarfs, however, only a small fraction of them will 1128 reach into the habitable zone (Tamburo et al. 2023). We 1129 thus conclude that the Nancy Grace Roman Space Tele-1130 scope could provide a useful sample to explore the run-1131 away greenhouse transition, albeit with a predominant 1132 focus on water-rich (sub-)Neptunes (e.g., Pierrehumbert 1133 2022).

Looking further ahead, the Nautilus Space Observa-1135 tory concept (Apai et al. 2019a) represents a statistical 1136 mission able to provide precise radius measurements of 1137 a large sample (~ 1000) of small exoplanets. It employs a constellation of ~ 35 large-diameter ($D \sim 8.5\,\mathrm{m}$) tele-1139 scopes using ultralight diffractive-refractive optical el-1140 ements (Milster et al. 2020) with the primary goal to 1141 study the atmospheres of transiting exoplanets. Op-1142 erating in an array mode, Nautilus would achieve the 1143 equivalent light-collecting area of a 50 m telescope. Its 1144 expected 1 ppm photometric precision (Apai et al. 2022) 1145 would enable precise radius measurements of a large 1146 sample, also through a low number of required visits per object. If realized, Nautilus will be a valuable instru-1148 ment for characterizing the runaway greenhouse transi-1149 tion.

Other missions have been proposed that focus on characterizing exoplanet habitability, most notably the Hab-1152 itable Worlds Observatory concept, which will build on the two precursor direct imaging concepts LU-VOIR (The LUVOIR Team 2019) and HabEx (Gaudi 1155 et al. 2020). Similar science objectives are pursued by 1156 the Large Interferometer For Exoplanets (LIFE, Quanz 1157 et al. 2022) initiative, a mission concept utilizing a 1158 space-based mid-infrared nulling interferometer. Direct 1159 imaging surveys do not directly measure planetary radii and are primarily useful for providing context through 1161 atmospheric measurements of individual planets. How-

⁴ https://exoplanetarchive.ipac.caltech.edu

1162 ever, because mid-infrared retrievals feature reduced 1163 degeneracy between cloud albedo and changes in sur-1164 face area, the planet radius can be constrained in midinfrared wavelengths (Defrère et al. 2018; Quanz et al. 1166 2021). Mid-infrared direct imaging techniques, in particular, enable to study much deeper atmospheric layers 1168 than possible in reflected light (Wordsworth & Kreid-1169 berg 2022). Hence, the atmospheric structure can be 1170 retrieved for a wider variety of thermal and atmospheric scenarios (Alei et al. 2022; Konrad et al. 2022). Since 1172 the peak thermal emission in runaway greenhouse at-1173 mospheres will substantially decrease the star-to-planet flux ratio, mid-infrared wavelengths offer the possibility 1175 to probe the diversity of runaway climates in systems 1176 across different ages (Lupu et al. 2014; Bonati et al. 1177 2019). Finally, mid-infrared surveys such as LIFE show preference for M star planets (Quanz et al. 2022), which is beneficial for detecting the predicted habitable zone inner edge discontinuity (see Sect. 5.4.2). With a sample size of a few tens of planets crossing the runaway greenhouse transition, direct imaging missions will thus enable key insights into the compositional inventory of atmospheric volatiles and climate states (Hinkley et al. 2021; Currie et al. 2022), adding important details to a 1186 potential runaway greenhouse detection purely via tran-1187 sit radii.

As for exoplanets missions in their implementation phase, however, *PLATO* overall remains to be the most promising mission for an empirical confirmation or fal-sification of the runaway greenhouse transition at this time.

6.4. Mission design trade studies

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To explore the impact of mission trades on the de1195 tectability of the habitable zone inner edge discontinu1196 ity, we simulated different survey designs and strategies
1197 and measured their capability to recover the trend and
1198 constrain its parameters. We assessed this capability
1199 based on two determinants: the likelihood that the mis1200 sion is able to detect the injected trend, and the preci1201 sion with which it can constrain the parameters of that
1202 trend.

6.4.1. The value of follow-up campaigns

The constraining power changes when additional information beyond planet radii is available for the characterized planet population. As runaway greenhouse
phases leave a stronger imprint on bulk density than
no planet radius (see Sect. 5.4.1), it would be benefino cial to obtain constraints on planetary masses and test
the runaway greenhouse hypothesis in density space inno the runaway greenhouse hypothesis in

predominant water content of planetary surfaces and atmospheres or a smaller available planet sample. For a
mission design similar to *PLATO* a density-based hynote pothesis test on about a third of the overall sample is
requivalent to a pure radius-based analysis. At a fixed
sample size, key parameters of the runaway greenhouse
models can be more narrowly constrained when addimodels can be more narrowly constrained when addimodels can be more narrowly constrained when addi-

Precise ground-based radial velocity measurements will be needed to provide these data, and a number of instruments are already successfully employed in characterizing terrestrial-sized exoplanets (e.g., Queloz et al. 2021; Pepe et al. 2010; Johnson et al. 2010; Ribas et al. 2023) and confronting these results with planet formation theory (e.g., Miguel et al. 2020; Burn et al. 2021; Zaton theory (e.g., Miguel et al. 2020; Burn et al. 2021; Zaton of instruments on extremely large telescopes such as G-CLEF on the Giant Magellan Telescope (Szentgy-1231 orgyi et al. 2016), ANDES on the European Extremely Large Telescope (Marcantonio et al. 2022), or MODHIS on the Thirty Meter Telescope (Mawet et al. 2019) will open up the discovery space even further.

Recently, NASA and the National Science Foundation (NSF) commissioned an "Extreme Precision Radial Ve-1237 locity Initiative" (Crass et al. 2021) to develop methods and facilities for precise mass measurements of temper-1239 ate terrestrial planets. Their findings highlight that such 1240 measurements are costly, and therefore follow-up efforts 1241 may only be available for a subsample of the targets of ₁₂₄₂ a mission of *PLATO*'s scale. The diagnostic power of 1243 the hypothesis tests we demonstrated here may be im-1244 proved by simultaneously fitting for the habitable zone 1245 inner edge discontinuity in the subsample without RV 1246 follow-up. An optimized mission in search for the inner 1247 edge of the habitable zone will further enhance its in-1248 formation content via an informed selection of follow-up 1249 targets, i.e., balancing objects located on either side of 1250 the expected instellation threshold S_{thresh} .

6.4.2. The importance of M dwarfs in the target list

To date, the majority of planets with radius measurements orbit FGK dwarfs, and, based on the instellation
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1264 sit probability with orbital distance that has prompted 1265 a number of recent transit surveys to specifically tar1266 get M dwarfs (e.g., Irwin et al. 2009; Obermeier et al. 1267 2016; Delrez et al. 2018; Sebastian et al. 2021; Dietrich 1268 et al. 2023), but the sample of terrestrial planets orbiting 1269 them is still small (e.g., Berger et al. 2020; Hardegree1270 Ullman et al. 2020).

M dwarf systems are also key for detecting the run-1271 1272 away greenhouse transition: Our calculations with dif-1273 ferent spectral types (Sect. 5.4.2) show that the information content of M dwarfs in a sample dominates the hypothesis tests. Besides their large number in a volume/magnitude-limited sample, transiting M dwarf planets are more likely to be located near the threshold 1278 instellation and in particular on orbits further out, i.e., 1279 in the optimistic habitable zone. In fact, we showed that the FGK part of the planet sample barely contributes to the statistical power. Furthermore, the transit depth difference at the transition is expected to be $_{1283}$ larger for M dwarfs ($\sim 100\,\mathrm{ppm}$ for early, $\sim 1000\,\mathrm{ppm}$ 1284 for late M stars, Turbet et al. 2019), enhancing the demographic signal it leaves. An additional advantage of targeting M dwarfs are the extended runaway greenhouse phases of their planets that can last on the order 1288 of gigayears (Luger & Barnes 2015). This increases the 1289 probability of observing any given planet in the sample 1290 during the runaway greenhouse phase, essentially driv f_{rgh} ing f_{rgh} to higher values. Therefore, in addition to the 1292 high scientific value of boosted detections of potentially 1293 habitable planets, M dwarfs are also indispensable for 1294 the discovery and characterization of the runaway green-1295 house transition. As with a pure volume-limited sample, targeted M dwarf survey, too, profits from follow-up measurements of planetary masses with an order of mag-1298 nitude increase in evidence.

6.5. Constraining planetary habitability

A potential for liquid water on the surface of a planet is commonly used as an environmental marker to assess 1302 its surface habitability (Huang 1959; Hart 1978; Kast-1303 ing et al. 1993; Kaltenegger & Sasselov 2011; Koppa-1304 rapu et al. 2013). The runaway greenhouse transition 1305 represents an upper bound on received irradiation for this condition. Its detection would thus not only em-1307 pirically confirm the habitable zone concept but also 1308 help to locate it in the observationally available plan-1309 etary parameter space. In Sect. 5.4.3, we show that the 1310 threshold instellation at which the runaway greenhouse 1311 transition occurs can be reasonably constrained without 1312 imposing overly optimistic conditions on the underlying 1313 planet population, instrumentation, or survey strategy.

measurement; the constraining power is directly proportional to the proportion of characterized planets around M dwarfs and to the number of planets for which masses can be determined.

The situation is different for the planetary water inventory and the fraction of planets with runaway greenhouse climates: Since these parameters are degenerate,
they cannot be well constrained without independent
measurements. This degeneracy could be lifted if incompositions
can be made. For example, detections of water vapor in
planets above the threshold instellation, combined with
precise radius measurements, would constrain the premeasurements of terrestrial planets.

Once a runaway greenhouse region is identified in the parameter space, the community will have a tool at hand to discern potentially habitable planets from Venusian worlds on an empirical basis. Together with atmospheric measurements (see Sect. 6.2.3), we will be able to put a number on the probability of an individual planet to harbor sufficient surface water to sustain life.

6.6. Impact of assumptions on our findings

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The prospects for probing the runaway greenhouse transition depends on astro- and geophysical factors, as well as on the specific instrumentation and survey stratage of a particular mission. Our state-of-the-art models approximate the situation and offer testable predictions. In the following, we review a few considerations that future models may include to refine these predictions.

6.6.1. Structures in the planet occurrence rate density

The baseline occurrence rate density in radius-period space that governs the generation of synthetic planets might influence our findings, especially if it contains any features that coincide with the injected demographic feature. This is not the case in the model from Berg-1550 sten et al. (2022) that we adopted: Its occurrence rate density varies smoothly in the domain relevant for the runaway greenhouse hypothesis; transitions only occur at smaller instellations ($< 50\,\mathrm{W\,m^{-2}}$) and larger radii habitable zone inner edge discontinuity. If any currently lass habitable zone inner edge discontinuity. If any currently unknown sharp features in the distribution of terres-1558 trial planets emerge, they should be considered in future studies.

6.6.2. Baseline mass-radius relationship

Our baseline mass-radius relationship assuming pure 1362 MgSiO $_3$ interiors (Zeng et al. 2016) might not be representative of the rocky planet population. However, while

interior composition may introduce an offset to the raisis dius habitable zone inner edge discontinuity, we do not
expect a change of its structure. Since the magnitude of
the radius inflation effect is expected to be larger for an
Earth-like interior composition with an iron core-silicate
mantle structure (Zeng et al. 2016; Noack & Lasbleis
ision 2020; Bonati et al. 2021), we consider our mass-radius
relationship a conservative case. We performed a sanity check to assess the impact of varying our baseline
model (see Appendix A.3) and found general agreement
between different interior compositions.

Future self-consistent modeling of interior-atmosphere interactions may include constraints on additional radius increases due to a molten interior (Bower et al. 2019) and any potential effects stemming from a deviating gas exchange between atmosphere and interior in runaway greenhouse planets due to different redox conditions (Ikoma et al. 2018; Lichtenberg et al. 2021; Bower et al. 2022; Gaillard et al. 2021) or water outgassing efficiency (e.g., Hier-Majumder & Hirschmann 2017; Ikoma et al. 2018; Salvador & Samuel 2023).

6.6.3. Bulk water mass fraction

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The predominant mass fractions of water, which sen-1386 sitively controls the atmospheric state of a rocky exo-1388 planet, is poorly constrained. Inferred water contents 1389 in the literature range from upper limits on the order 10^{-5} to "water worlds" with tens of percent mass fraction (e.g., Rogers & Seager 2010; Unterborn et al. 2018; 1392 Mousis et al. 2020; Agol et al. 2021; Luque & Pallé 1393 2022), all of which are within the realm of theoretical predictions (Selsis et al. 2007; Mulders et al. 2015; Sato 1395 et al. 2016; Jin & Mordasini 2018; Lichtenberg et al. 1396 2019; Bitsch et al. 2019; Venturini et al. 2020; Emsenhuber et al. 2021; Schlecker et al. 2021a; Lichtenberg & Clement 2022; Izidoro et al. 2022). Our nominal case as- $_{\rm 1399}$ sumes a bulk water mass fraction of $x_{\rm H_2O}=0.005.$ This 1400 can be considered a conservative choice that is unlikely to introduce a systematic overestimation of the habit-1402 able zone inner edge discontinuity. Cases of pure rocky 1403 composition and very low volatile contents can be con-1404 sidered absorbed by the dilution factor $f_{\rm rgh}$. Assuming 1405 a distribution of water mass fractions instead of a fixed 1406 value would thus not significantly change our results.

6.6.4. "Sharpness" of the habitable zone inner edge discontinuity

The habitable zone inner edge discontinuity may be affected by several processes that are challenging to quantify: Planets that lack an atmosphere, sufficient volatiles, or have non-water-dominated outgassed compositions cannot bear steam atmospheres, and those that that do eventually move to the non-runaway greenhouse

1415 category due to desiccation (Watson et al. 1981; Kasting 1416 & Pollack 1983; Hamano et al. 2013) or evolution of their 1417 host star (Luger & Barnes 2015). Hydrogen/Helium-1418 dominated planets may disguise as inflated rocky planets 1419 and not contribute to the demographic signal, although 1420 a runaway greenhouse radius inflation effect was sug-1421 gested for water-dominated sub-Neptunes (Pierrehum-1422 bert 2022; Innes et al. 2023). A subset of such gas-1423 rich planets will experience atmospheric loss via photo-1424 evaporation (Owen & Wu 2013) or core-powered mass loss (Ginzburg et al. 2018), reducing their transit radius. 1426 Intrinsic variation in the threshold instellation is caused by differences in planetary features influencing the on-1428 set of a runaway climate such as albedo, atmospheric 1429 composition, clouds, or surface gravity (Salvador et al. 1430 2017; Turbet et al. 2021; Lichtenberg et al. 2021; Pier-1431 rehumbert 2022; Innes et al. 2023). The choice of a sta-1432 tistical estimator for the hypothesis tests may further 1433 influence the recovered discontinuity; we compare our 1434 nominal running mean approach with a binned statistic $_{1435}$ in Appendix A.1.

While these factors may offset the signal's amplitude, they preserve its general shape. The "dilution factor" $f_{\rm rgh}$ in our model embodies our ignorance of the magnitude of this offset. In a real survey, additional contextual information about planets in the sample may be available.

7. CONCLUSIONS

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Significant inflation of rocky planet radii is a ro1444 bust prediction of runaway greenhouse models. Using
1445 Bioverse, a quantitative hypothesis testing framework,
1446 we have explored the potential of contemporary exo1447 planet missions to statistically detect a radius/density
1448 discontinuity resulting from this inflation in the exo1449 planet population. Our key findings are as follows:

- The predicted runaway greenhouse transition causes a discontinuity in the radius and density distribution of small exoplanets with respect to their irradiation.
- 2. This habitable zone inner edge discontinuity should be detectable with high-precision transit measurements. For a planet sample $\gtrsim 100$, a detection is likely if radius inflation occurs on at least 10% of the observed planets and if typical bulk water mass fractions are above $\sim 10^{-3}$.
- 3. We find that the planned *PLATO* transit survey will provide a sufficient sample and the required precision to confirm or reject the predicted trend. Assuming the projected photometric pre-

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cision, *PLATO* will be able to test the runaway greenhouse hypothesis for planet yields $\gtrsim 100$.

4. The diagnostic power of transit missions in testing this hypothesis can be increased through a follow-up campaign providing planet mass measurements. This can reduce the required planet yield by about a factor of three. Only an adequate sample of planets orbiting M dwarfs will ensure sufficient targets on both sides of the expected threshold instellation.

5. Testing the runaway greenhouse hypothesis on a population level can provide constraints on the water inventory of rocky exoplanets and thus make an important contribution to assessing their habitability. A detection will provide an empirical confirmation of the habitable zone concept and localize its inner edge.

The habitable zone concept is widely employed in tar1482 get prioritization for exoplanet missions, and it will pro1483 vide context for interpreting potential signatures of life.
1484 As we have demonstrated, it appears realistic that an
1485 empirical test of the habitable zone hypothesis is im1486 minent. The confirmation or rejection of the habitable
1487 zone inner edge discontinuity will be a key contribution
1488 to understanding the diversity of exoplanet climates and
1489 the search for extraterrestrial life in the Universe.

1490 The authors thank Gabriele Cugno, Yann Fabel, Brad 1491 Foley, Lisa Kaltenegger, Ravi Kopparapu, Eric Ma-1492 majek, Megan Mansfield, Paul Molliere, Gijs Mulders, 1493 Matthew Murphy, Sukrit Ranjan, and Terry-Ann Suer 1494 for insightful discussions. We are grateful to Mar-1495 tin Turbet for providing the mass-radius relationship 1496 data for planets harboring steam atmospheres. thank the anonymous referees for providing insight-1498 ful reports that helped to improve this manuscript. 1499 This material is based upon work supported by the 1500 National Aeronautics and Space Administration un-1501 der Agreement No. 80NSSC21K0593 for the program 1502 "Alien Earths". The results reported herein benefited 1503 from collaborations and/or information exchange within 1504 NASA's Nexus for Exoplanet System Science (NExSS) 1505 research coordination network sponsored by NASA's 1506 Science Mission Directorate. This work has made 1507 use of data from the European Space Agency (ESA) 1508 mission Gaia (https://www.cosmos.esa.int/gaia), pro-1509 cessed by the Gaia Data Processing and Analysis 1510 Consortium (DPAC, https://www.cosmos.esa.int/web/ ₁₅₁₁ gaia/dpac/consortium). Funding for the DPAC has 1512 been provided by national institutions, in particular 1513 the institutions participating in the Gaia Multilateral 1514 Agreement. T.L. was supported by a grant from the 1515 Branco Weiss Foundation. G.B. acknowledges sup-1516 port from the NASA Astrophysics Data Analysis Pro-1517 gram under Grant No. 80NSSC20K0446. A.S. acknowl-1518 edges support from NASA's Habitable Worlds Program 1519 (No. 80NSSC20K0226).

AUTHOR CONTRIBUTIONS

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M.S., D.A., and T.L. conceived the project, planned its implementation, and interpreted the results. D.A. leads the "Alien Earths" program through which this project is funded and helped to guide the strategy of the project. T.L. and A.S. provided expertise on runder away greenhouse climates and exoplanet interiors. M.S. carried out the hypothesis tests and statistical analyses. M.S. wrote the manuscript; T.L., G.B., K.H.-1529 U., and A.S. provided text contributions. G.B. implemented the planet generator in the Bioverse frame-1531 work. All authors provided comments and suggestions on the manuscript.

REPRODUCIBILITY

This study uses the reproducibility framework "showyourwork" (Luger et al. 2021). All code responding quired to reproduce our results, figures, and this artiself is available at https://github.com/matiscke/spanner-edge-discontinuity. The code to reproduce a

1539 figure can be accessed via the icon link next to the re-1540 spective figure caption. Data sets associated with this 1541 work are available at doi:10.5281/zenodo.7080391 (stel-1542 lar luminosity tracks from Baraffe et al. (1998)) and doi:10.5281/zenodo.7946446 (results from model grid runs of Bioverse).

Software: Bioverse (Bixel & Apai 2021), As-1546 tropy (Astropy Collaboration et al. 2018), NumPy (Har-1547 ris et al. 2020), SciPy (Virtanen et al. 2020), 1548 corner.py (Foreman-Mackey 2016), dynesty (Speagle 1549 2020).

1550 APPENDIX

A. ROBUSTNESS TESTS

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A.1. Alternative statistics for the average radius or bulk density

The hypothesis tests introduced in Sect. 4 rely on a statistical estimator for the variation of planetary radii robust or bulk densities as a function of net instellation, and we star chose a moving average for this estimator in our nominal setup. Here, we explore how robust our results are against this choice by demonstrating the recovery of the runaway greenhouse signal in the case of the optimistic survey (Sect. 5.2) with an alternative estimator: instead of computing moving averages, we used a binned statistic.

We first binned the data of the simulated survey in instellation space, choosing the number of bins via the rule of Freedman & Diaconis (1981) and using logarithmic binning. In each bin, we computed the arithmetic mean of the planet radius and its standard deviation. Then, we assigned each planet the mean radius according to the instellation bin it occupies and used this as the measure for testing the runaway greenhouse hypothesis.

Figure 10 shows the simulated data together with binned, average planet radii and draws from the posterior of the hypothesis test. A clear detection resulted, although with somewhat lower significance ($\Delta \ln Z \approx 30$) compared to the nominal setup. The accuracy of the restriction covered instellation threshold is comparable. This test demonstrates that our results are not sensitive to the choice of statistical estimator to test the hypotheses against.

It is conceivable that with very large sample sizes and very sharp runaway greenhouse transition a binned solution would perform better. For a search with real data, both approaches, and possibly other alternatives, should be considered.

A.2. Statistical imprint of runaway greenhouse atmospheres

The predicted runaway greenhouse-induced planet ra-1589 dius changes are a function of instellation, planet mass, 1590 and bulk water mass fraction (compare Sect. 3). In or-

Optimistic survey (binned)

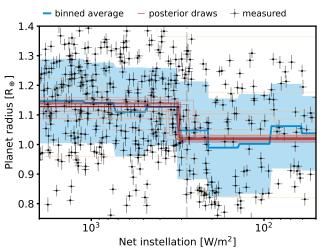


Figure 10. Detection of the runaway greenhouse threshold with a binned statistic. Using simulated data from the optimistic survey case (N=500, compare Figure 5), we tested the runaway greenhouse hypothesis. Instead of using a running mean, we computed a binned statistic (arithmetic mean and standard deviation in blue) to assign planets the average radius or bulk density in their neighborhood in instellation space. Random draws from the posterior of the runaway greenhouse hypothesis (Eqn. 4) are shown in red. As in the nominal case, the pattern is detected with high significance.

1591 der to better understand the interaction between the 1592 planetary populations underlying our simulations and 1593 these predictions, we compared the latter to the aver-1594 age radius and bulk density changes we measured in the 1595 synthetic population. We used the "optimistic" scenario 1596 with a sample size of 500.

Figure 11 shows this comparison for a range of plan1598 etary masses and water mass fractions. The complex
1599 dependence of the radius inflation on these parameters
1600 is evident, but no significant abrupt changes capable
1601 of causing spurious signals occur. Differences between
1602 model prediction and population can be explained by
1603 the wide dispersion in planet mass in the population.

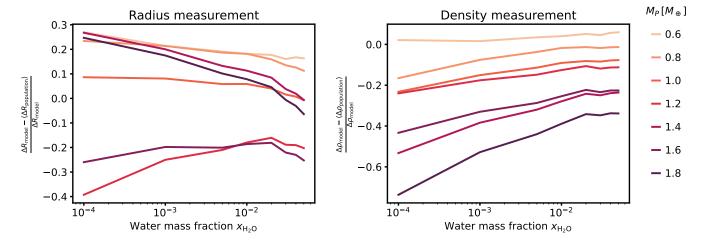


Figure 11. Comparison of radius and bulk density changes predicted by the atmospheric models to the average changes measured in the synthetic planet population. Model predictions ΔR_{model} and $\Delta \rho_{\text{model}}$ depend on individual planet masses; $\langle \Delta R_{\text{population}} \rangle$ and $\langle \Delta \rho_{\text{population}} \rangle$ are averaged measurements of the overall population. Significant differences are thus expected.

A.3. Influence of different mass-radius relationships on our results

We assessed how a different choice of baseline massradius relation influences our results. Focussing on the 1607 detectability of the statistical runaway greenhouse signal 1608 and the ability to constrain the threshold instellation, we 1610 repeated the hypothesis test in Sect. 5.2 with alternative mass-radius relationships. Instead of assuming a pure 1611 MgSiO₃ composition, we assigned planet masses using either the probabilistic relationship in Wolfgang et al. (2016) or a semi-empirical, two-layer relation assuming 1614 an Earth-like (32.5 % Fe + 67.5 % MgSiO₃) composi-1615 tion (Zeng et al. 2016).

Figure 12 shows how the two alternative baseline 1617 mass-radius relations influence the significance of a detection and the ability to constrain the threshold instellation. The relation of Wolfgang et al. (2016) includes intrinsic scatter, which impedes a detection at low dilution factors. In this regime, a layered, Earthlike composition leads to more significant detections and a narrower, although biased, constraint on S_{thresh} . 1625 Both mass-radius relations agree and recover the in-1626 jected value where the null hypothesis can be rejected with high significance. This is consistent with our nominal mass-radius relation (compare Sect. 2.4). We conclude that the underlying core and mantle composition of planets may affect the detectability of the transition if the fraction of planets with runaway greenhouse climates 1632 is low; however, the overall trends that our experiments 1633 revealed appear robust.

We caution that this analysis may serve only as a san-1635 ity check and should not be taken as a result in itself: 1636 The atmospheric model from Turbet et al. (2020) we 1637 adopted relies on a silicate interior composition for its

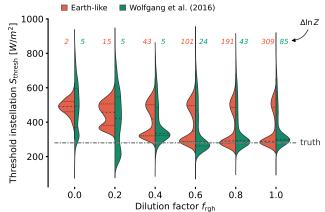


Figure 12. Posterior threshold instellations as a function of dilution factor for two alternative baseline mass-radius relations. For each grid step in $f_{\rm rgh}$, we show kernel density estimates of retrieved posteriors (averaged over 50 iterations) assuming an Earth-like (32.5 % Fe + 67.5 % MgSiO₃) composition or the probabilistic relation from Wolfgang et al. (2016) and otherwise following the optimistic scenario in Sect. 5.2. Lines within the violins show quartiles of the distributions, and the gray line indicates the injected threshold instellation of 280 W m⁻². Higher log-evidence differences $\Delta \ln Z$ correspond to more significant rejections of the null hypothesis. In the regime of strong detections, both mass-radius relations lead to similar, accurate constraints on $S_{\rm thresh}$. Differences occur at low dilution factors, where the Earth-like relation leads to narrower estimates.

1638 transit radius prediction. Therefore, only our nominal 1639 procedure throughout the main body of the paper rep-1640 resents a self-consistent treatment.

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