Exercise 2.2.7 (not in book, but needed). Let Δ = det G be the **determinant of** G, where $G = (g_{ab})$ is the metric tensor matrix. The **minor**, m^{ab} , for element g_{ab} is the determinant of the submatrix that excludes row a and column b. The **cofactor** for element g_{ab} is $c^{ab} = (-1)^{a+b} m^{ab}$. Show that

$$\partial_c \Delta = c^{ab} \partial_c g_{ab} = \Delta g^{ab} \partial_c g_{ab} .$$

Solution Note: This derivation is valid for any matrix g_{ab} , not just the metric tensor.

By definition, Δ is the sum of all possible products of the form $(-1)^{a_1+b_1} g_{a_1b_1} \cdot ... \cdot g_{a_Nb_N}$ where the factors $g_{a_kb_k}$ are in distinct rows and columns. This definition is somewhat complicated to state and even more complicated to express mathematically, which we do now.

There are two formulas for Δ that we use here. The first one is to sum the products $g_{ab}c^{ab}$ across row a. However, we have to be careful on how to express this in Einstein notation. In non-Einstein notation, we can write

$$\Delta = \sum_{b=1}^{N} g_{ab} c^{ab}.$$

However, in Einstein notation, $g_{ab} c^{ab}$ sums over not just row a, but over all rows, which results in $\Delta + ... + \Delta$, N times. Thus, in Einstein notation we express this sum as

$$g_{ab} c^{ab} = N \Delta. (a)$$

A related formula is:

$$g_{ab} g^{ab} = N$$
 (b) $g_{ab} g^{ab} = g_{ab} g^{ad} \delta_d^b \stackrel{(1.78)}{=} \delta_b^d \delta_d^b = N.$

Putting these two formulas together we get:

$$c^{ab} = \Delta g^{ab} \qquad : \qquad (c)$$

$$g_{ab} c^{ab} \stackrel{(a)}{=} N \Delta \stackrel{(b)}{=} g_{ab} g^{ab} \Delta \Rightarrow c^{ab} = g^{ab} \Delta \checkmark$$

The second definition of Δ is more convoluted. We express Δ mathematically exactly as stated in words, above:

$$\Delta \equiv \epsilon^{b_1 \dots b_N} g_{1 b_1} \dots g_{N b_N}, \tag{d}$$

by taking advantage of the **Levi-Civita symbol**, $\epsilon^{b_1 \dots b_N}$, which we now explain.

This formula for Δ uses Einstein notation to sum over all of the different combinations of indices b_k . The indices b_k range over the integers 1 - N for all k. The symbol $\epsilon^{b_1 \dots b_N}$ is defined to be zero unless each integer appears exactly once; that is, unless $b_1, ..., b_N$ is a permutation of 1, 2, ..., N. For permutations, its value is defined as +1 for even permutations and -1 for odd permutations. A permutation is even if the number of pairwise swaps required to restore the natural order is even, and odd, otherwise. For example, 1, 3, 2, 4 is an odd permutation because it requires just one swap. 1, 3, 4, 2 is an even permutation. So, in the determinant for N = 3, the product g_{11} g_{22} g_{33} is preceded by a plus sign, represented by $\epsilon^{b_1b_2b_3}=\epsilon^{123}$, and the product g_{11} g_{23} g_{32} is preceded by a negative sign. ϵ^{132} .

We call $e^{b_1 \dots b_N}$ the Levi-Civita symbol because it is a pseudo-tensor, almost a tensor except that it reverses sign under Jacobian transformations having negative determinant. It is not quite proper to call it the Levi-Civita tensor.

The cofactor, also, can be defined in Levi-Civita terminology since it is a determinant. Consider the following attempt at a cofactor expression:

$$c^{a\,b} \equiv \epsilon^{k_1\,\ldots\,k_{a-1}\,b\;\,k_{a+1}\,\ldots\,k_N} \, g_{1\,k_1}\,\ldots\,g_{a-1\,k_{a-1}} \, g_{a+1\,k_{a+1}}\,\ldots\,g_{N\,k_N}.$$

Every row appears in $g_{1k_1} \dots g_{a-1k_{a-1}} g_{a+1k_{a+1}} \dots g_{Nk_N}$ except row a. The exponent of ϵ ranges over all strings $k_1 \dots k_{a-1} b k_{a+1} \dots k_N$, and ϵ is non-zero only for strings that are permutations of 1 - N. So, none of the non-zero terms have any k_a equal to b. Thus, when a string is a permutation, every column except b appears in $g_{1k_1} ... g_{a-1k_{a-1}} g_{a+1k_{a+1}} ... g_{Nk_N}$. So far, so good.

Index b is in the ath superscript position. Will that produce the correct sign for the term $g_{1k_1} \dots g_{a-1k_{a-1}} g_{a+1k_{a+1}} \dots g_{Nk_N}$? It requires |b-a| index swaps to move bfrom position b to position a. So, $\epsilon^{1 \dots a-1 b} = (-1)^{b-a}$. The term $g_{11} \dots$ $g_{a-1} = g_{a+1} = g_{a$ of $(-1)^{b-a}$, we must precede this term by another $(-1)^{b-a}$. We therefore modify our first attempt and try

$$c^{a\,b} \equiv (-1)^{b-a} \, \epsilon^{k_1 \, \dots \, k_{a-1} \, b \, k_{a+1} \, \dots \, k_N} \, g_{1\,k_1} \, \dots \, g_{a-1\,k_{a-1}} \, g_{a+1\,k_{a+1}} \, \dots \, g_{N\,k_N}.$$

This expression has the correct sign as well as the correct rows and columns for the cofactor c^{ab} . However, in proper tensor notation, when a and b are superscripts on LHS then they must be superscripts on RHS, but index a appears in the RHS subscripts as well as the RHS superscripts. We can try to remedy this with a Kronecker delta, like

$$(-1)^{b-a} \delta_a^s \ \epsilon^{k_1 \, \ldots \, k_{a-1} \, b \, \, k_{a+1} \, \ldots \, k_N} \, g_{1 \, k_1} \, \ldots \, g_{s-1 \, k_{s-1}} \, g_{s+1 \, k_{s+1}} \, \ldots \, g_{N \, k_N}.$$

In this attempt, index a no longer appears among the subscripts of g, but it now appears as a subscript in the Kronecker delta. To solve that problem, we toss in a summation sign:

$$c^{ab} \equiv (-1)^{b-a} \sum_{s=1}^{N} \delta_{s}^{a} \epsilon^{k_{1} \dots k_{s-1} b k_{s+1} \dots k_{N}} g_{1 k_{1}} \dots g_{s-1 k_{s-1}} g_{s+1 k_{s+1}} \dots g_{N k_{N}}.$$
 (e)

In this summation we throw away all terms except when s = a, making this identical to the prior expression while adhering to the Einstein summation convention.

Since this is a little confusing, we give one example.

Example To illustrate cofactor definition (e), let a = 3, b = 5, and N = 7. We can replace the sum by letting s = 3:

$$c^{3\times5} \equiv (-1)^{5-3} \epsilon^{k_1 k_2 5 k_4 \dots k_7} g_{1 k_1} g_{2 k_2} g_{4 k_4} \dots g_{7 k_7}$$

Einstein summation occurs over every k_a except k_3 and so includes sequences with repeated numbers, but the only non-zero values of ϵ are when the k_a 's include one each of the values 1, 2, 3, 4, 6, 7 in any order. Thus, as desired, $g_{1\,k_1}\,g_{2\,k_2}\,g_{4\,k_4}\,...\,g_{7\,k_7}$ pairs every possible column with every possible row, excluding row 3 and column 5.

This expression is preceded by $(-1)^2 = +1$ because the number of pairwise swaps for b=5 to be the fifth superscript differs by two swaps from it being the 3rd superscript.

Finally, we can compute $\partial_c \Delta$ from equation (d), which we re-express here:

$$\Delta \equiv \epsilon^{k_1 \dots k_N} g_{1 k_1} \dots g_{a-1 k_{a-1}} g_{a k_a} g_{a+1 k_{a+1}} \dots g_{N k_N}.$$
 (f)

Then

$$\begin{array}{lll} \partial_{c}\Delta & = & (-1)^{b-1} \; \partial_{c}g_{1\,b} \; \epsilon^{b\,k_{2}...\,k_{N}}g_{2\,k_{2}}...\; g_{N\,k_{N}} \\ & + \; ... \\ & + & (-1)^{b-a} \; \partial_{c}g_{a\,b} \; \epsilon^{k_{1}...\,k_{a-1}\,b\;k_{a+1}...\,k_{N}} \; g_{1\,k_{1}}...g_{a-1\,k_{a-1}} \; g_{a+1\,k_{a+1}}...\; g_{N\,k_{N}} \\ & + \; ... \\ & + & (-1)^{b-N} \; \partial_{c}g_{N\,b} \; \epsilon^{k_{1}...\,k_{N-1}\,b}g_{1\,k_{1}}...\; g_{N-1\,k_{N-1}} \\ & = \sum\limits_{s=1}^{N} \delta_{s}^{a} \; (-1)^{b-s} \; \partial_{c}g_{s\,b} \; \epsilon^{k_{1}...\,k_{s-1}\,b\;k_{s+1}...\,k_{N}} \; g_{1\,k_{1}}...g_{s-1\,k_{s-1}} \; g_{s+1\,k_{s+1}}...\; g_{N\,k_{N}}. \end{array} \tag{g}$$

Observe that we needed to toss in $(-1)^{b-a}$ terms again because when index b is in position a, we have to again compensate for the fact that ϵ contains a $(-1)^{b-a}$ factor.

We can simplify equation (g) by moving $(-1)^{b-s} \partial_c g_{sb}$ to the left of the summation sign. It becomes $\partial_c g_{ab}$ (-1)^{b-a} since the summations are all zero except when s = a, and this completes the solution:

$$\partial_{c}\Delta = \partial_{c}g_{ab} (-1)^{b-a} \sum_{s=1}^{N} \delta_{s}^{a} \epsilon^{k_{1} \dots k_{b-1} b k_{b+1} \dots k_{N}} g_{1 k_{1} \dots g_{s-1} k_{s-1}} g_{s+1 k_{s+1} \dots} g_{N k_{N}}$$

$$\stackrel{(e)}{=} \partial_{c}g_{ab} c^{ab} \qquad \checkmark$$

$$\stackrel{(c)}{=} \Delta g^{ab} \partial_{c}g_{ab} \qquad \checkmark$$