Journées de Probabilités 2021, Guidel, BZH

Euclidean Random Assignment
Problems: origin, state of the art
and some open problems in one
dimension

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Based on several papers in collaboration with

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Section 1 Background and Definition

Section 2 State of the art

Section 3 ERAPs at d = 1

Section 4 Two open problems



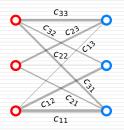
Section 1	Background and Definition
1	The Assignment Problem
2	Adding randomness to AP
3	Euclidean Random Assignment Prob- lems



The (linear sum) Assignment Problem (AP)

For a $n \times n$ cost matrix c, find a bijection π (a permutation) s.t. $E = \sum_i c_{i\pi(i)}$ is minimal. Let E_{\min} be the minimal value.

Example at n = 3:



$$c = \begin{pmatrix} 5 & 3.5 & 1 \\ 2 & 1.2 & 3 \\ 3 & 2 & 4 \end{pmatrix}$$

- Simple formulation → good model in applications
- P-complete with $\mathcal{O}(n^3)$ complexity [Munkres 1957]

The (linear sum) Assignment Problem (AP)

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Example at n = 3:

$$E_{\min} = 5$$



$$c = \begin{pmatrix} 5 & 3.5 & \boxed{1} \\ \boxed{2} & 1.2 & 3 \\ 3 & \boxed{2} & 4 \end{pmatrix}$$

Swap columns (rows) s.t. Tr(c) is minimal [Koopmans–Beckmann 1957]; Optimal mixed strategy in a "hide and seek" game [Von Neumann 1953, 1954]

- Simple formulation → good model in applications
- P-complete with $\mathcal{O}(n^3)$ complexity [Munkres 1957]

AP: an old problem



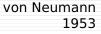
König 1916



Egérvary 1931

Canon simplicissimus.

	1	п	ш	ΙV	v	VI	VII
ī	25*	21	20	18	20	18	25
II	21	22*	21	21	13	21	22
ш	16	19	23*	22	17	14	16
īv	21	12	18	27*	13	14	24
v	25	22	22	27	31*	16	31
VI	10	18	23	21	19	23*	21
VII	5	14	10	27	31	20	40*





Kuhn 1955





"De investigando ordine systematis aequationum ..." [Jacobi 1860] See also [Ollivier 2009]

Section 1	Background and Definition
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The Random Assignment Problem I

c is a random matrix $(c_{ij} \text{ i.i.d. r.v. } \sim \rho(l) = l^r + o(l^r))$.

$$\mathbb{E}[E_{\min}]_n \underset{n \to \infty}{\sim} c_r n^{1 - \frac{1}{r+1}}.$$

- Pioneered in Physics in the 80s by Mézard–Parisi and Orland
- Entered Probability mostly through Aldous in the 90s

Result: only "short" edges are relevant for large n and r can be considered a "universal exponent".

Nice fact: at r = 0 (i.e. ρ is e.g. uniform or $Exp(\lambda)$ distribution)

$$c_0 = \zeta(2) = \frac{\pi^2}{6}$$
.



The Random Assignment Problem II

If $c_{ij} \sim \text{Exp}(1)$, Parisi conjectured (1998):

$$\mathbb{E}[E_{\min}]_n = \sum_{k=1}^n \frac{1}{k^2} = \frac{\pi^2}{6} - \frac{1}{n} + o\left(\frac{1}{n}\right)$$

1 Rectangular matrices

- [Coppersmith-Sorkin 1998]
- 2 Proof of $\zeta(2)$ limit (among other things)

[Aldous 2001]

3 Proof of Parisi conjecture

[Prabhakar-Sharma 2001]

4 Extension to the k-partite case (NP-hard for $k \ge 3$)

[Martin-Mézard-Rivoire 2004,2005]

5 3! solution to "cavity" equation

[Wästlund 2012, Larsson 2014, Salez 2015]

NOT discussed today...



Section 1	Background and Definition
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The Euclidean Random Assignment Problem (ERAP)

Let $\mathcal{B}=(B_1,\ldots,B_n)$ be blue points and $\mathcal{R}=(R_1,\ldots,R_n)$ be red ones: n-samples of i.i.d. r.v. of pdf $\rho_{\mathcal{B}(\mathcal{R})}:\Omega\to\mathbb{R}$ ("disorder"), (Ω,\mathcal{D}) is a metric space (mostly an **Euclidean** space with \mathcal{D} **Euclidean** distance). For $p\in\mathbb{R}$ and an assignment (permutation) π , consider the *Hamiltonian*

$$\mathcal{H}(\pi) = \sum_{i=1}^{n} \mathcal{D}^{\rho}(B_i, R_{\pi(i)})$$

and the random variable (ground state energy)

$$\mathcal{H}_{\mathrm{opt},(n,d)}^{(p)} = \min_{\pi \in \mathcal{S}_n} \mathcal{H}(\pi) \quad (\pi_{\mathrm{opt}} = \arg\min_{\pi \in \mathcal{S}_n} \mathcal{H}(\pi)).$$

<u>Problem</u>: the rate of $E_{p,d}(n) := \mathbb{E}[\mathcal{H}_{\mathrm{opt},(n,d)}^{(p)}]$ as $n \to \infty$.

Three motivations: Physics, Mathematics and Computer Science

- **Spin Glasses** ERAP is a toy model of spin-glass (a **disordered** and **frustrated system**) in finite dimension, which is numerically simple in comparison to e.g. Edwards–Anderson spin glass [Mézard–Parisi 1988];
- **Optimal Transport** ERAP is a Monge-Kantorovitch problem associated to empirical measures $\rho_{\mathcal{B}}$, $\rho_{\mathcal{R}}$:

$$\mathcal{H}_{\text{opt}} = nW_p^p(\rho_{\mathcal{B}}, \rho_{\mathcal{R}})$$

where W_p is the *p*-Wasserstein distance [Villani 2009, Vershik 2013, Brezis 2018];

Computational Complexity Theory - ERAP is a small (but crucial) modification of random TSP, however finding π_{opt} is **easy** (recall that AP is P-complete).

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Section 2 State of the art

1 The phase diagram

2 State of the art

ERAP: the phase diagram

We shall put $\rho_{\mathcal{B}} = \rho_{\mathcal{R}} := \rho$. We wish to study

$$E_{p,d}(n) := \mathbb{E}[\mathcal{H}_{\text{opt},(n,d)}^{(p)}] \stackrel{?}{=} K_{p,d} n^{\gamma_{p,d}} (\ln n)^{\gamma'_{p,d}} (1 + o(1))$$

as $n \to \infty$, depending on (p, d) and the choice of ρ .

Phase diagram: $(\gamma_{p,d}, \gamma'_{p,d})$ are expected to be "universal", i.e. largely independent on the microscopic details (which may affect the constant $K_{p,d}$).

Remark: non-uniform disorder is more subtle!

Example: standard Gaussian disorder at (p, d) = (2, 1)

$$E_{2,1}(n) \sim 2 \ln \ln n$$
 (i.e. $\gamma_{2,1} = \gamma'_{2,1} = 0$).

[Caracciolo-D'A-Sicuro 2019, Bobkov-Ledoux 2019, Berthet-Fort 2020]

See [Benedetto–Caglioti 2020] for non-uniform case at d = 2.





Section 2 State of the art The phase diagram State of the art

$$d \ge 3$$
, $p \ge 1$

"Simple":

$$E_{p,d}(n)|_{d\geq 3} \sim c_{p,d} n^{\gamma_{LB}}$$

where

$$\gamma_{LB} := 1 - \frac{p}{d} = \gamma_{p,d}$$
 [Mézard-Parisi 1988]

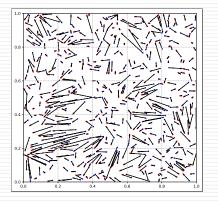
(if the disorder is uniform, otherwise <u>unknown</u>).

Remark: the constant $c_{p,d}$ is **unknown** (upper and lower bounds in [Talagrand 1992]).

- Almost-sure limits of Euclidean functionals of finite random point sets [Barthe–Bordenave 2013 and refs. therein];
- Recurrent interest in Optimal Transport [Goldman–Trevisan 2020].

d = 2: a challenge for both mathematicians and physicists

Example configuration for $\Omega = [0, 1]^2$ and \mathcal{D} Euclidean distance:



Optimal assignment typically involves $O(\ln n)$ -nearest-neighbors: $(\gamma_{p,d}, \gamma'_{p,d}) = (\gamma_{LB}, \frac{p}{2})$ if $p \ge 1$ [Ajtai–Komlós–Tusnády 1984]

Recent developments in Mathematics and Physics

2014 Caracciolo–Lucibello–Parisi–Sicuro (Phys. Rev. E): using a (classical) field-theoretical approach, predicted

$$K_{2,2}=\frac{1}{2\pi};$$

- 2019 Ambrosio–Stra–Trevisan (PTRF): proof of $K_{2,2} = \frac{1}{2\pi}$ (among other things) via PDE methods;
- 2020 Ambrosio–Glaudo (JEP): refinement on the remainder term (among other things);
- 2021 Benedetto–Caglioti-Caracciolo–**D'A**–Sicuro–Sportiello (JStatPhys): exact energy differences for ERAPs on two manifolds Ω , Ω '.

See https://www.youtube.com/watch?v=4RcOiW20C_E for a discussion of the latter results in the light of Weyl's law in spectral theory (and extension to ERAPs at d=3).

Section 1 Background and Definition

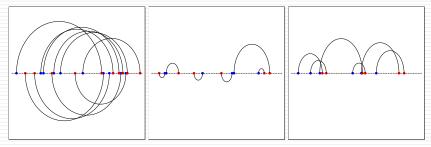
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d = 1: properties of the solution

p = 0 and p = 1 separate three qualitative regimes:



If realizations are sorted in natural order, π_{opt} is:

- p < 0 **cyclical**: $\pi_{opt} = i + k \pmod{n}$ [Caracciolo-**D'A**-Sicuro 2017];
- $p \in (0,1)$ **non-crossing**: intervals "covered by edges" are either disjoint or one is included into the other [McCann 1999];
 - $p \ge 1$ **ordered**: π_{opt} is the identity permutation.

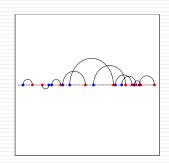
Remark: combinatorial properties hold for any disorder ρ .

Section 3	ERAPs at $d=1$
1	Convex regime $p \ge 1$
2	Concave regime $p \in (0, 1)$
3	Anomalous scalings

$$p \ge 1$$
, $\Omega = \mathbb{R}$, $[0, 1]$, \mathbb{R}^+ and $\mathcal{D} = | |$

Convexity and strict monotonicity of $\mathcal{D}^p \implies$ if realizations (b_1, \ldots, b_n) and (r_1, \ldots, r_n) are sorted in **natural order**, then

$$H_{\text{opt},(n,1)}^{(p)}|_{p\geq 1} = \sum_{i=1}^{n} |b_i - r_i|^p.$$

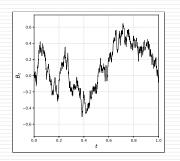


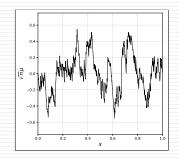
 \implies a path to the Brownian world.

$$p \ge 1$$
, $\Omega = \mathbb{R}$, $[0, 1]$, \mathbb{R}^+ and $\mathcal{D} = |\cdot|$: Brownian Bridge

Let the **transport field** be $\mu_i := b_i - r_i$. Then

$$\sqrt{n}\mu_i \xrightarrow[weakly]{} B_t$$
, the Brownian Bridge





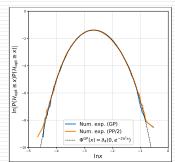
Thus $\mathcal{H}_{\text{opt,}(n,p)} = \sum |\mu_i|^p$, energy \sim moments of B_t .

[Boniolo–Caracciolo–Sportiello 2014, Caracciolo–Sicuro 2014, Caracciolo–**D'A**–Sicuro 2017)]

$$\Omega = \mathbb{S}_1$$
, $\mathcal{D} = \text{arc-distance}$, $p = 2$: distribution of \mathcal{H}_{opt}

The asymptotic CDF of $\mathcal{H}_{\text{opt},(n,2)}$ on the unit circle can be obtained via **Fourier analysis** [**D'A**, PhD Thesis 2020]:

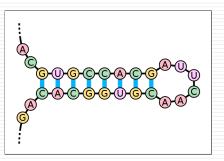
$$\Pr[\mathcal{H}_{\text{opt},(n,2)} \le x] \stackrel{n \to \infty}{=} 9_4(0, e^{-2\pi^2 x})$$
$$:= \sum_{s \in \mathbb{Z}} e^{i\pi s} e^{-2\pi^2 s^2 x}$$

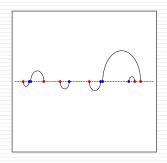


Section 3	ERAPs at $d=1$
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ERAPs at $p \in (0, 1)$: further motivation in biology

Besides economics, interest due to the non-crossing property of the solution: toy-model for the **secondary structure of RNA** (discarding pseudo-knots).

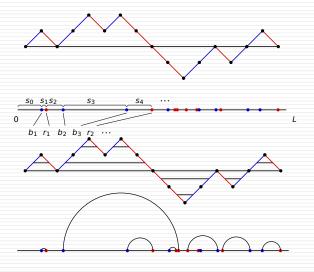




Despite this, poorly understood. [McCann 1999]

IDEA: try a canonical construction, an approximate solution sharing non-crossing property with $\pi_{opt} \rightarrow Dyck$ matchings!

The Dyck matching [Caracciolo-**D'A**-Erba-Sportiello 2020]



The Dyck Upper Bound at $p \in (0, 1)$

Expected energy of Dyck matchings grows asymptotically as

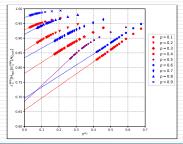
$$\mathbb{E}_n(\mathcal{H}_{\mathsf{Dyck}}) \simeq \left\{ \begin{array}{ll} n^{1-p} & \text{if} & 0 \leq p < \frac{1}{2} \\ \sqrt{n} \ln n & \text{if} & p = \frac{1}{2} \\ n^{\frac{1}{2}} & \text{if} & \frac{1}{2} < p \leq 1 \end{array} \right.$$

Conjecture [Caracciolo-**D'A**-Erba-Sportiello 2020]:

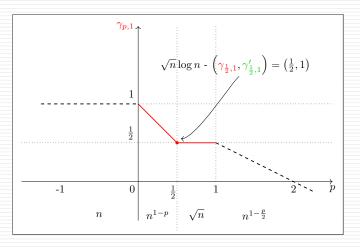
$$\lim_{n\to\infty}\frac{\mathbb{E}_n(\mathcal{H}_{\text{opt}})}{\mathbb{E}_n(\mathcal{H}_{\text{Dyck}})}=k_p$$

for some constant $0 < k_p < 1$.

Numerical evidence \implies



Section of the Phase Diagram at d = 1



Product formula for number of solutions at d = p = 1 [Caracciolo–Erba–Sportiello 2021].

Section 3	ERAPs at $d=1$
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Beyond uniform disorder: bulk and anomalous scaling

Recall the Selberg result for points $\sim U[0,1]$ and $p \geq 1$ (see [Caracciolo *et al.* 2019]):

$$\begin{split} E_{p,1}^{\mathsf{U}}(n) &= n \frac{\Gamma\left(1 + \frac{p}{2}\right)}{p+1} \frac{\Gamma(n+1)}{\Gamma(n+1 + \frac{p}{2})} \\ &= c_p n^{1 - \frac{p}{2}} \left(1 + o(1)\right) = c_p n^{\gamma_{p,1}} \left(1 + o(1)\right) \; (\textit{Bulk scaling}). \end{split}$$

What if ρ is non-uniform?

If $\gamma_{p,1} \neq 1 - \frac{p}{2}$ or in presence of a logarithmic correction, we talk about an **anomalous scaling**.

"Reduction to quadratures" in the bulk scaling

Let R be the cdf of ρ . Set $\Psi^{(\rho)} := \rho \circ R^{-1}$. Then

$$E_{p,1}^{(\rho)}(n) = \frac{2^p}{\sqrt{\pi}} \Gamma\left(\frac{p+1}{2}\right) \int_0^1 \left[\frac{\sqrt{s(1-s)}}{\psi(\rho)(s)}\right]^p ds \, n^{1-\rho/2} + o(n^{1-\rho/2}).$$

Caracciolo—**D'A**—Sicuro (JStatPhys 2018): non-rigorous **regularization methods** inspired by cutoff regularization in QFT. Analogous problem "in the continuum" by Bobkov–Ledoux (AMS 2019).

Example: $\rho(x) = e^{-x}$, $\Psi^{\text{exp}}(s) = 1 - s$.

• Cutoff method:
$$E_{p,1}^{\text{exp}}(n) = \frac{2^p}{\sqrt{\pi}} \Gamma\left(\frac{p+1}{2}\right) \int_0^{1-\epsilon/n} \left(\frac{s}{1-s}\right)^{\frac{p}{2}} ds$$
.
At $p = 2$, we get $E_2(n) = 2 \ln n - 2 \log c - 2 + o(1)$;

Exact calculation: at p = 2 (via Beta integrals): $E_{2,1}^{\text{exp}}(n) = 2 \sum_{k=1}^{n} \frac{1}{k} = 2 \ln n + 2 \gamma_{\text{E}} + o(1).$

(by comparison
$$c = e^{-\gamma_E - 1} = 0.20655...$$
)

On **D'A**-Sportiello 2020 (in preparation)

Emergence of logarithmic corrections to scaling \implies detailed study of contribution of k-th edge at fixed n, then $n \rightarrow \infty$.

General strategy: 1) Binomial expand $\mathcal{H}_{\mathrm{opt},(n,1)}^{(p)}|_{p\geq 1}$ at fixed (n,p) (for p>1 even); 2) exploit linearity to compute expectation w.r.t. a certain Beta distrib.; 3) asymptotic analysis (case by case depending on ρ).

More precisely

$$E_{p,1}^{(p)}(n) = \sum_{k=1}^{n} \mathbb{E}[(b_k - r_k)^p] = \sum_{k=1}^{n} \sum_{q=0}^{p} \binom{p}{q} (-1)^{p-q} M_{n,k,q}^{(p)} M_{n,k,p-q}^{(p)}$$

for

$$M_{n,k,\ell}^{(\rho)} = \langle R_{\rho}^{-1}(u)^{\ell} \rangle_{P_{n,k}}, \quad R_{\rho}^{-1}(u)$$
 the quantile function,

where $\langle ... \rangle_{P_{n,k}}$ is expectation w.r.t. Beta(k, n-k+1):

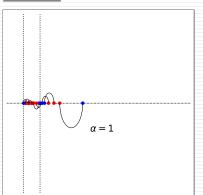
$$P_{n,k}(u)du := \frac{n!}{(k-1)!(n-k)!} u^{k-1} (1-u)^{n-k} du.$$

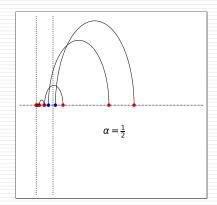
Example: stretched exponential distribution

Let us consider the following pdf (θ is Heaviside's function)

$$\rho_{ie,\alpha}(x) = \alpha x^{\alpha-1} \exp(-x^{\alpha})\theta(x).$$

Example: solutions at n = 10





Thus
$$R_{\alpha}(x) = \exp(-x^{\alpha})\theta(x) \implies R_{\alpha}^{-1}(u) = (-\ln u)^{\frac{1}{\alpha}}$$
.

Sketch of the computation

For integer $s = \frac{1}{\alpha}$, we wish to evaluate

$$M_{n,k;s} := \int_0^1 du P_{n,k}(u) (-\ln u)^s.$$

Lemma I (**D'A**-Sportiello 2020)

$$M_{n,k;s} = s! h_s(A_{k,n}),$$

for the alphabet $A_{k,n} := \left\{ \frac{1}{k}, \frac{1}{k+1}, \dots, \frac{1}{n} \right\}$, where h_s is the complete homogeneous symmetric function of degree s^* .

Two hints: 1) use the representation $-\ln u = \lim_{x\to 0} \frac{u^{-x}-1}{x}$ and 2) recall that for a polynomial $A(q) = a_p q^p + a_{p-1} q^{p-1} + \dots + a_0$,

$$\sum_{p=0}^{p} \binom{p}{q} (-1)^{p-q} A(q) = p! \, \alpha_p.$$

^{*}I.e. for alphabet $(x_1, ..., x_m)$, $h_s(x_1, ..., x_m) := \sum_{1 \le j_1 \le ... \le j_s \le m} x_{j_1} x_{j_2} \cdots x_{j_s}$.

Sketch of the computation

The contribution $E_{(s,p),n}(k)$ of the k-th edge in the solution to the total energy $E_{(s,p)}(n) = \sum_{k=1}^{n} E_{(s,p),n}(k)$ is thus

$$E_{(s,p),n}(k) = \sum_{q=0}^{p} {p \choose q} (-1)^{p-q} M_{n,k;sq} M_{n,k;s(p-q)}$$

$$= \sum_{q=0}^{p} (-1)^{q} {p \choose q} (sq)! (s(p-q))! h_{sq}(A_{k,n}) h_{s(p-q)}(A_{k,n}).$$

In the next step, we make use of generating functions.





Exact result in the case s = 1 (exponential)

For exp. distrib. points, we have a remarkable simplification. Set

$$H(t,X) = \sum_{\ell \geq 0} h_{\ell}(X) t^{\ell} = \prod_{x \in X} (1-tx)^{-1}.$$

Then

$$E_{(1,p),n}(k) = p! \sum_{q=0}^{p} (-1)^q h_q(A_{k,n}) h_{p-q}(A_{k,n})$$
$$= p! [t^p] H(t, A_{k,n}) H(-t, A_{k,n}).$$

Obviously for any alphabet X, $H(t,X)H(-t,X) = H(t^2,X^2)$, so that

$$E_{(1,p),n}(k) = p![t^p]H(t^2, A_{k,n}^2) = p!h_{p/2}(A_{k,n}^2).$$

Asymptotic result at integer s > 1 (stretched-exponential)

For the power-sum functions $p_k(X) = \sum_{x \in X} x^k$, define

$$P(t,X) = \sum_{\ell \geq 1} \ell^{-1} p_\ell(X) t^\ell = \log H(t,X).$$

Fact: $(p_{\ell}(A_{k,n})^{1/\ell})_{\ell}$ is monotonic decreasing (due to the generalised-mean ineq.). In a regime in which $p_1 \gg \sqrt{p_2}$ we consider a **perturbative expansion**. By setting $P_+(t,X) = P(t,X) - tp_1(X)$ we then have

$$h_k(X) = [t^k]H(t, X) = [t^k] \exp(P(t, X))$$
$$= \sum_{\ell=0}^k \frac{1}{(k-\ell)!} p_1(X)^{k-\ell} [t^\ell] \exp(P_+(t, X)).$$

After some computation, we get

$$E_{(s,p),n}(k) \simeq p! \, s^p p_1(A_{k,n})^{(s-1)p} [t^p] \left(\exp(P_+(t,A_{k,n})) \exp(P_+(-t,A_{k,n})) \right).$$

Asymptotic result at s > 1 integer (stretched-exponential)

Since still obviously $\exp(P_+(t,X)) \exp(P_+(-t,X)) = H(t^2,X^2)$, we get

$$E_{(s,p),n}(k) \simeq p! \, s^p p_1(A_{k,n})^{(s-1)p} \, h_{p/2}(A_{k,n}^2)$$

which recovers the exact result at s=1. As $p_1(A_{k,n}) \sim \ln(n) - \ln(k) + \mathcal{O}(1)$ whenever $n \gg k$, while $p_h(A_{k,n}) \sim \frac{1}{h-1}(k^{1-h} - n^{1-h})$ if $n \gg k \gg 1$, in any case $p_1/\sqrt{p_2}$ is at least of order $\ln(n)$ whenever $n \gg k$. We can thus replace $p_1(A_{k,n})$ with $\ln(n)$ to get

$$E_{(s,p),n}(k) = p! \, s^p \ln(n)^{(s-1)p} \, h_{p/2}(A_{k,n}^2) \big(1 + \mathcal{O}((1+\ln k)/\ln n) \big).$$

In order to estimate $E_{(s,p)}(n) = \sum_k E_{(s,p),n}(k)$, we are left with the task of estimating $(q = \frac{p}{2})$

$$F_{n,q} := \sum_{k=1}^{n} h_q(A_{k,n}^2).$$

Emergence of **multiple** ζ^* **values** as $n \to \infty$

Observe that each monomial contributing to $F_{n,q}$ is of the form $\frac{1}{l_1^2 l_2^2 \dots l^2}$, and enters exactly i_1 times in the sum. Hence

$$F_{n,q} = \sum_{1 \le i_1 \le i_2 \le \dots \le i_q \le n} \frac{1}{i_1 \, i_2^2 \cdots i_q^2} = \sum_{k=1}^n \frac{1}{k} h_{q-1}(A_{k,n}^2).$$

In this series[†] we recognize a special kind of "**multiple** ζ^* **values**"(MZSVs), defined as

$$\zeta^*(a_1,\ldots,a_r) := \sum_{1 \le i_1 \le i_2 \le \cdots \le i_r} \frac{1}{i_1^{a_1} \cdots i_r^{a_r}}.$$

MZSVs are related to **multiple** ζ **values** arising e.g. in **Vassiliev–Kontsevich theory** [Zagier 1992]. Not new in **Physics**: scattering amplitudes in perturbative quantum field theory [Broadhurst–Kreimer 1997].

[†]Convergence as $n \to \infty$ holds since $h_{q-1}(A_{\nu,n}^2) \sim k^{-(q-1)}$ for $n \gg k$.

Result for $E_{(p,1)}^{(s)}(n)$

For our special MZSVs a **further remarkable simplification holds** [Ohno–Wakabayashi 2006,Ohno–Zudilin 2008, Borwein–Zudilin 2011]

$$F_{+\infty,q} = \zeta^*(1, \underbrace{2, \dots, 2}_{q-1}) = 2\zeta(2q-1)$$

where ζ is Riemann's function. In conclusion, we have

Lemma II (D'A-Sportiello 2020)

In the ERAP with stretched-exponentially distributed points with pdf $\rho_{\mathrm{ie},\alpha}(x)=\alpha x^{\frac{1}{s}-1}\exp(-x^{\frac{1}{s}})$, for p>1 even and s integer

$$E_{p,1}^{(s)}(n) \simeq \begin{cases} 2s^2 \ln(n)^{2s-1} & p = 2\\ 2\zeta(p-1)s^p p! \ln(n)^{(s-1)p} & p \ge 4 \end{cases}$$

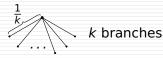
Remark: the <u>same asymptotics</u> holds for $\rho_{ie,\alpha}^+(x) = \alpha x^{\alpha-1} \exp(1-x^{\alpha})$ (the argument goes with minor modifications).



ERAPs on the k-star graph (W.i.p. with Caracciolo, Liu and Sportiello 2021-)

Choice of the triple $((\Omega, \mathcal{D}), (\rho_{\mathcal{B}}, \rho_{\mathcal{R}}), p)$:

1 $\Omega = \mathbb{K}_{1,k}$ (i.e. "star" with k branches)



2 Distance \mathcal{D} :

$$l_2 \bigvee_{l_1} \mathcal{D} = l_2 - l_1$$

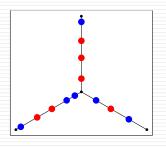
$$l_2 = l_2 + l_1$$

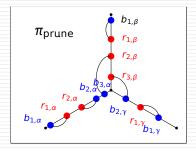
- \mathfrak{B} and \mathcal{R} are $\sim U[0,\frac{1}{k}] \otimes \mathcal{U}\{1,k\}$ random variables;
- 4 *p* ≥ 1.

On the configuration π_{prune}

IDEA: try a canonical construction, an approximate solution sharing local ordering property with $\pi_{\text{opt}} \to \pi_{\text{prune}}!$

Example: k = 3 and n = 6





On the configuration π_{prune}

Conjecture :

$$\lim_{n \to +\infty} \mathbb{E} \Big[\mathcal{H}_{\mathrm{opt}}^{(p)} \Big] = \lim_{n \to +\infty} \mathbb{E} \Big[\mathcal{H}_{\mathrm{prune}}^{(p)} \Big].$$

Consequence: $\gamma_{p,1}^{(k)} = 1 - \frac{p}{2}$, $\gamma_{p,1}^{',(k)} = 0$, explicit (and simple!) formulas for limit constants.

Example: at p = 2

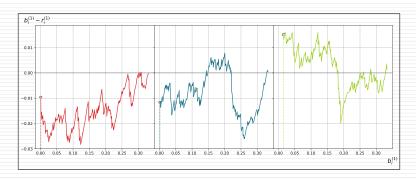
$$\lim_{n \to +\infty} \mathbb{E} \Big[\mathcal{H}_{\text{opt}}^{(2)} \Big] = \frac{3k-2}{3k^2}.$$

On the configuration π_{prune}

Claim:

$$\lim_{n \to +\infty} \mathbb{E} \Big[\mathcal{H}_{\text{opt}}^{(2)} \Big] = \frac{3k-2}{3k^2}.$$

Sketch of proof: moments of k independent Brownian Bridges stopped at time $\frac{1}{k}$.



Section 1 Background and Definition

Section 2 State of the art

Section 3 ERAPs at d = 1

Section 4 Two open problems

1 2 3 4

Open problem 1: the Dyck conjecture

Proof of the Dyck conjecture for $p \in (0, 1)$:

$$\lim_{n\to\infty} \frac{\mathbb{E}_n(\mathcal{H}_{\text{opt}})}{\mathbb{E}_n(\mathcal{H}_{\text{Dyck}})} = k_p.$$

Open problem 2: implications of Generalised Selberg Integrals

Take only $\Omega = [0, 1]$, $\mathcal{D} = |.|$ and $\rho_{\mathcal{B}} = \rho_{\mathcal{R}} := \rho = \mathbb{I}_{[0,1]}(x)$. From Generalised Selberg Integrals, we know [Caracciolo *et al.* 2019], for integer ℓ ,

$$\mathbb{E}[|b_k-r_k|^\ell] = \frac{\Gamma^2(n+1)\Gamma(k+\frac{\ell}{2})\Gamma(n-k+1+\frac{\ell}{2})\Gamma(1+\ell)}{\Gamma(k)\Gamma(n-k+1)\Gamma(n+1+\frac{\ell}{2})\Gamma(n+1+\ell)\Gamma(1+\frac{\ell}{2})}.$$

For the usual cost function $f = |.|^p$, we have a nice formula for E(n).

Problem: are there choices of a more general cost function f = f(|.|) so that a "<u>nice expression</u>" for E(n) (i.e. not necessarily involving hypergeometric functions) can be obtained upon resummation?[‡]

[‡]This question was raised by N.Enriquez during a talk given by the author at CIRM Marseilles - Luminy in March 2021.

Thank you for your attention!