

Superconducting Qubits with Josephson junctions

Theoretical insights and real world applications

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Outline

Artificial atoms

Why
superconductors?

Making a non-
linear resonator:
the Josephson
junction

The SC qubits
family tree

The transmon
qubit

Qubits gates:
single and coupled

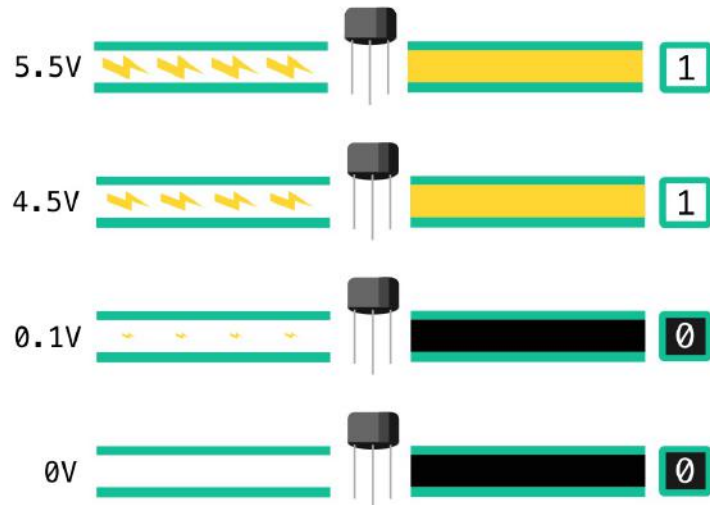
Qubit readout

Problems and
applications

Bit vs qubit

Classic bit (binary digit)

- Two levels system (0 or 1 logic state)
- Physically implementations:
 - electrical switches (*transistors*)
 - distinct levels of voltage/currents in a circuit
 - magnetization directions...
- A measurement does not affect the state

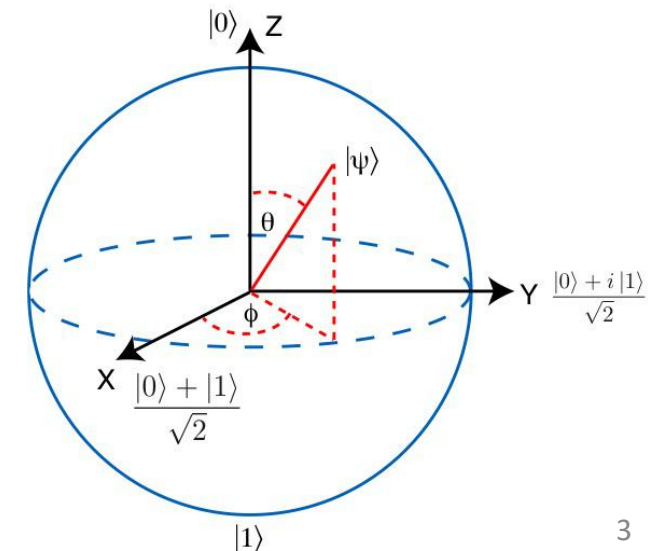


Qubit (quantum bit)

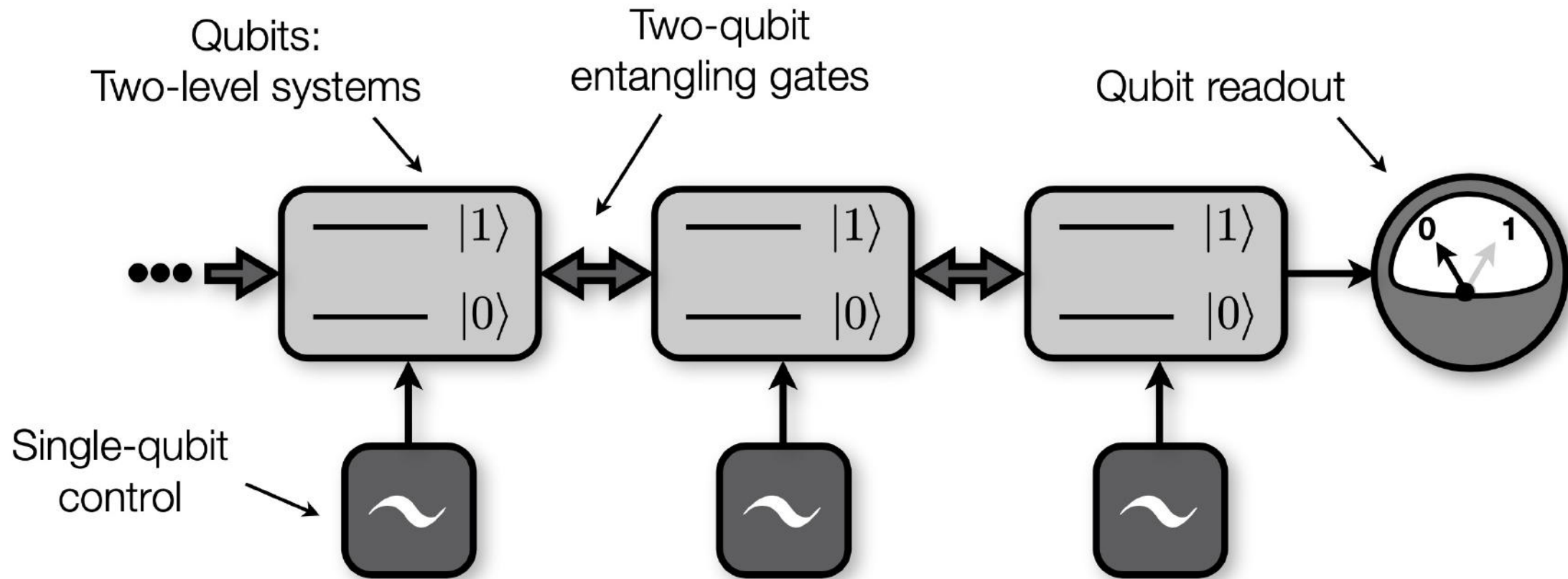
- Coherent superposition of two quantum states
- Wave vector pointing towards a position in the Bloch sphere:

$$\psi = \alpha|0\rangle + \beta|1\rangle$$

- A measurement destroys coherence

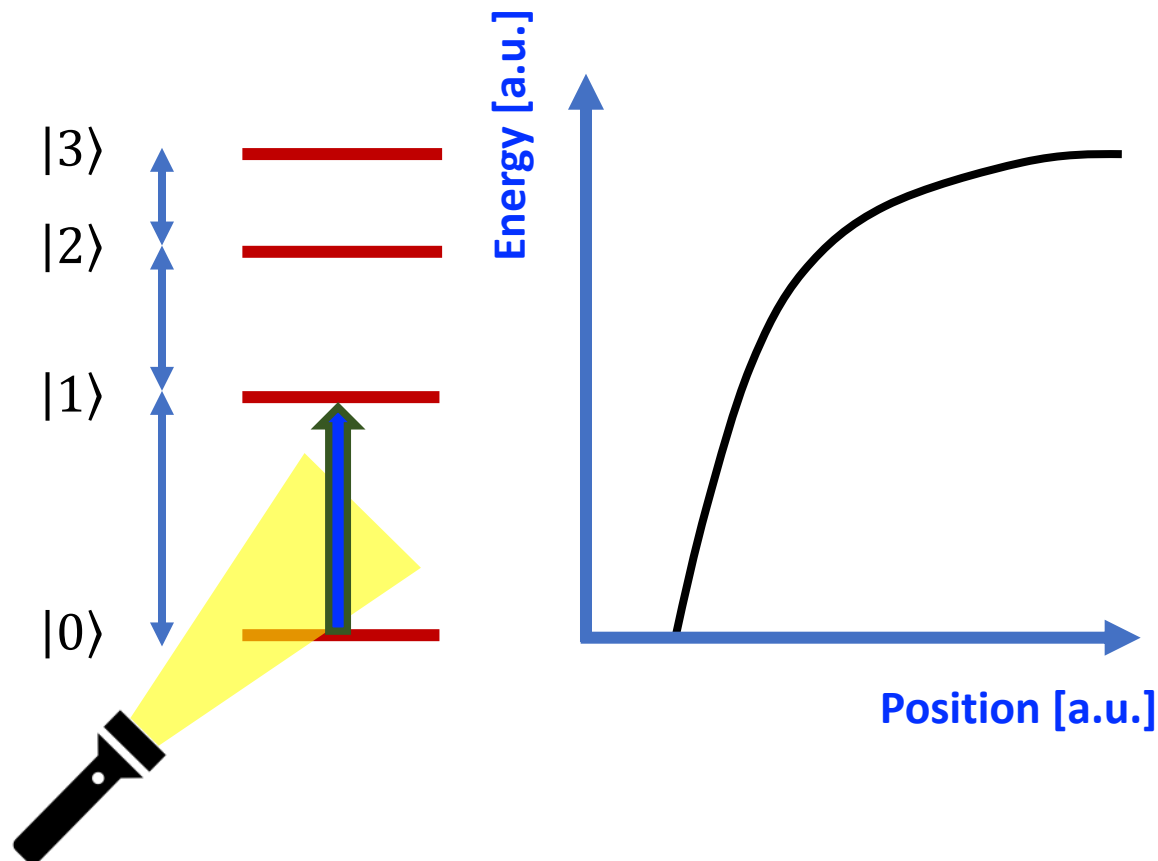


The challenge

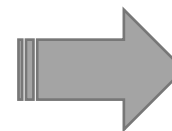


Nature “atomic” atoms

$$\hbar\omega_{01} = E_{01} \neq E_{12} = \hbar\omega_{12}$$



- Laser tuned at the right transition frequency ω_{01}



- $T_{decay} \sim \text{years}$
- $T_{decoh} \geq 10 \text{ seconds}$
- $T_{pulse} \sim 5 \mu\text{s}$

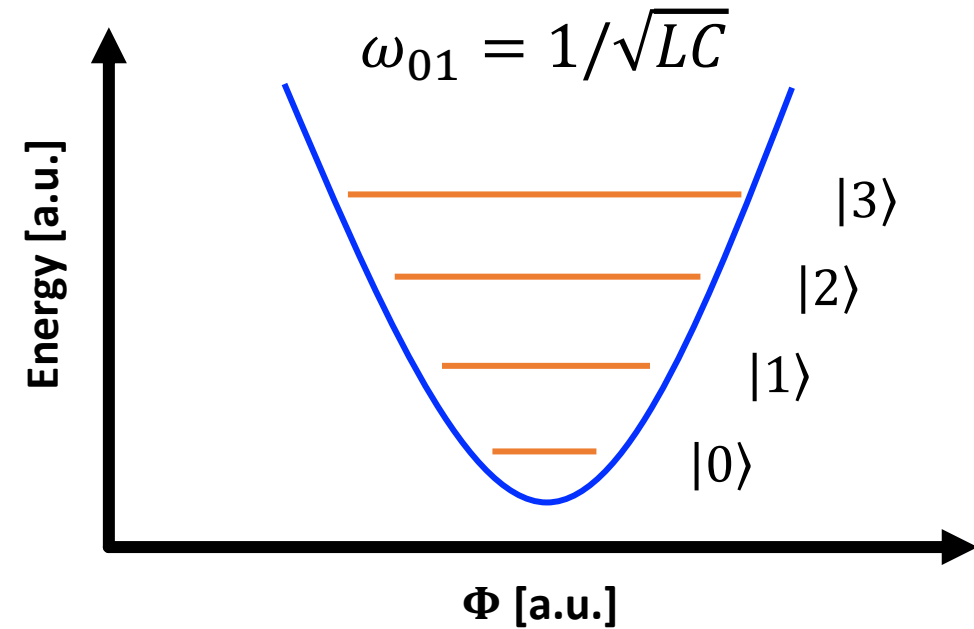
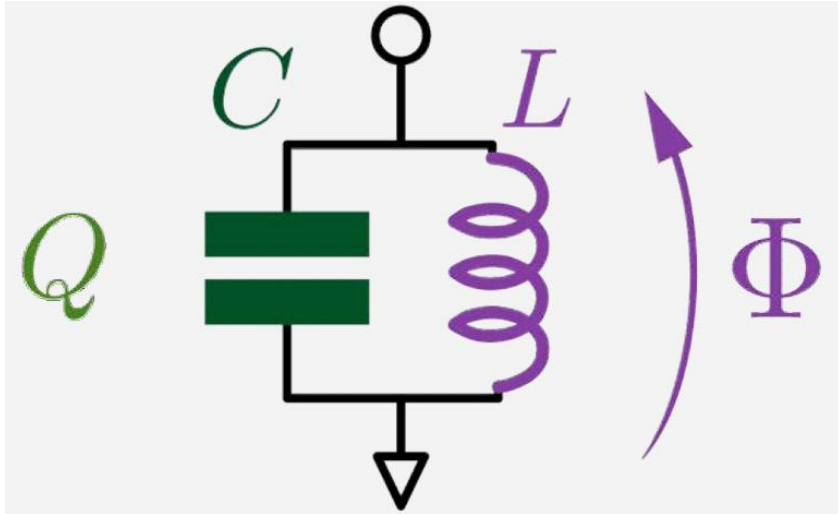


*Low error per
gate $\sim 0.48 \%$*

“Artificial” atoms – LC harmonic oscillator



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$$H = \frac{Q^2}{2C} + \frac{\Phi^2}{2L}$$

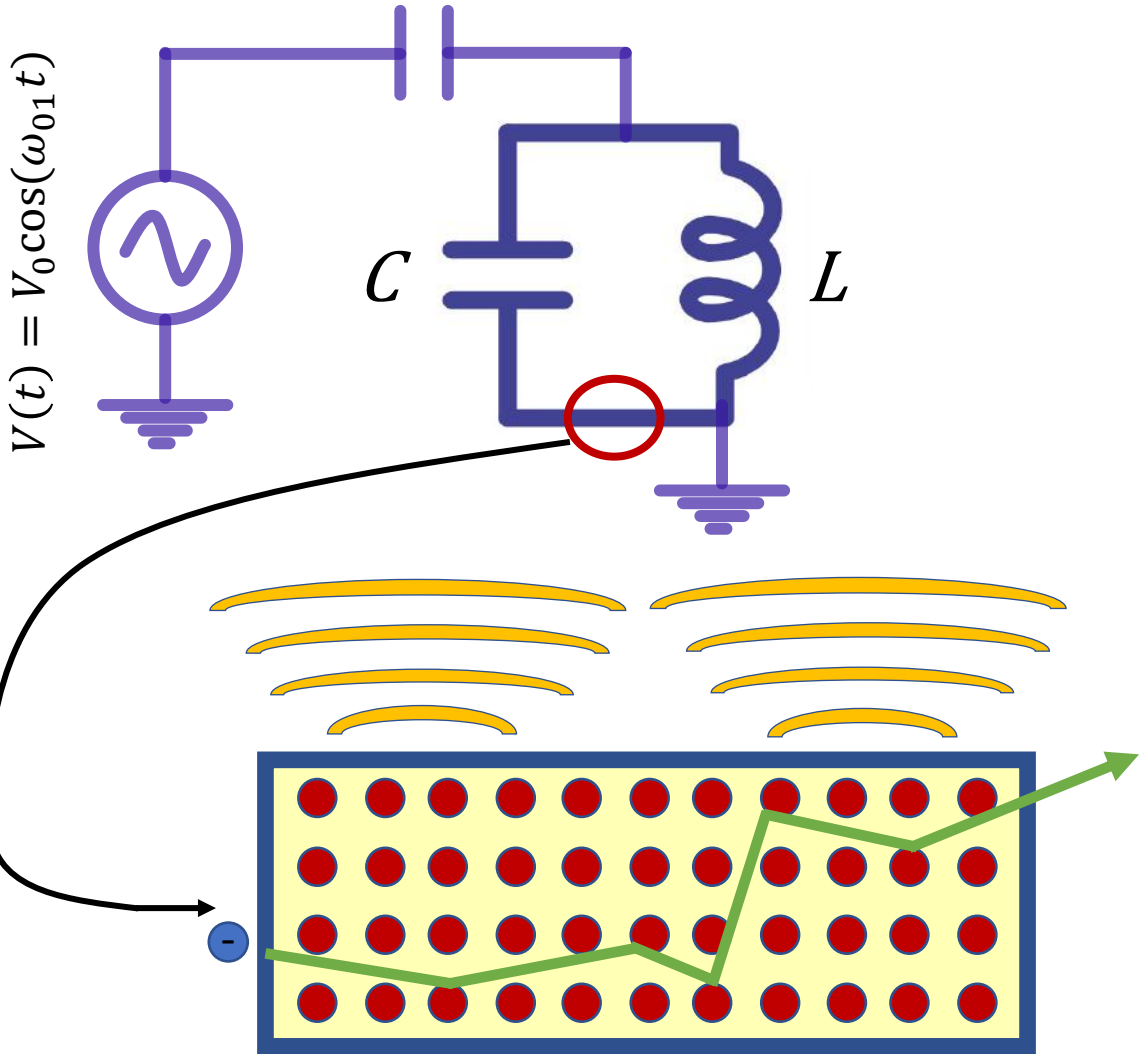
Capacitive term

Inductive term

How to make a quantum circuit



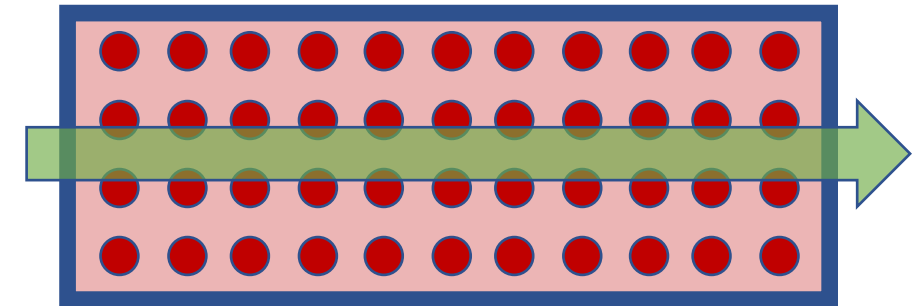
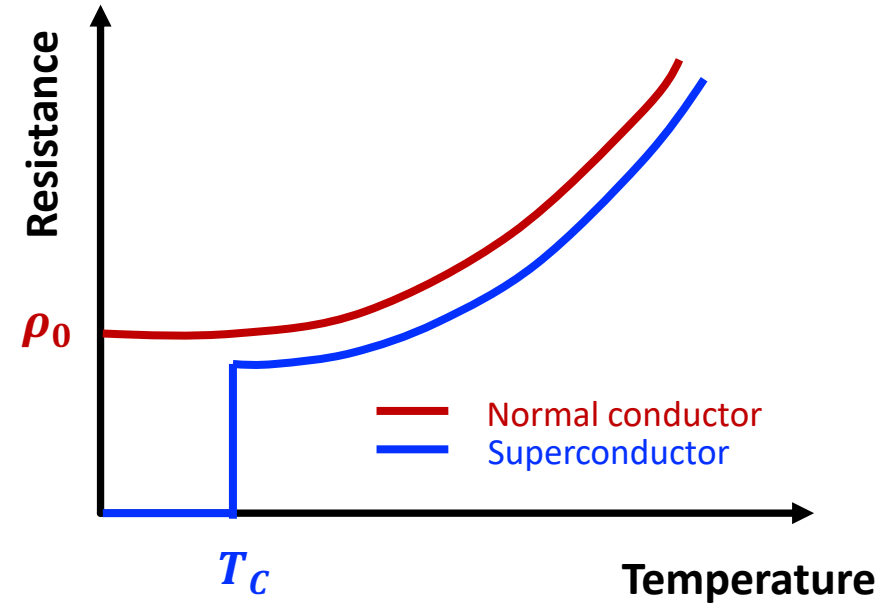
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Normal conductor



Joule heating, leakage
of information



Superconducting wire



quantum mechanical
behavior is preserved!

The Helium dilution refrigerator

$$\omega_{01} = \frac{1}{\sqrt{LC}} \sim 10 \text{ GHz}$$

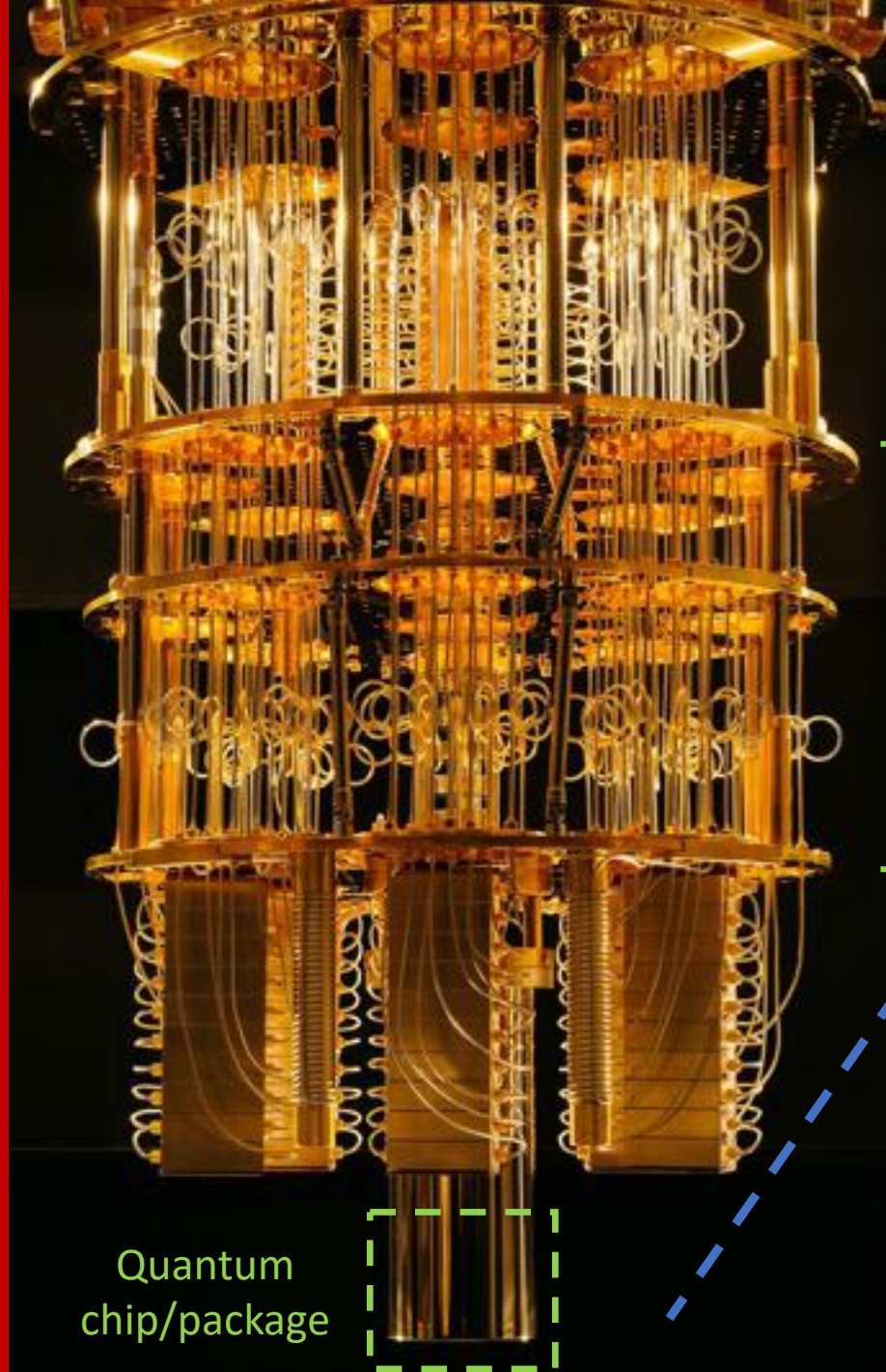
Quantum condition:

$$T \ll \frac{\hbar\omega_{01}}{K_B} \sim 0.3 \text{ K}$$

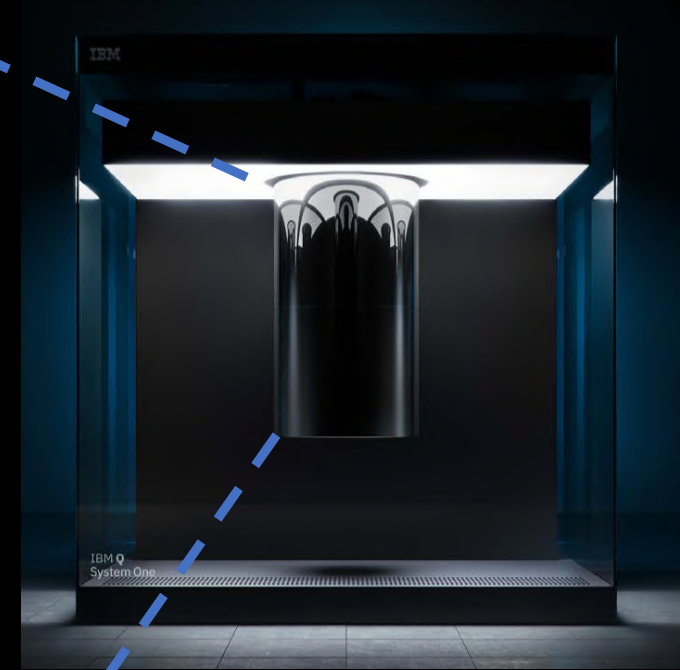


Commercial He dilution refrigerators:

$$T = 20 \text{ mK}$$



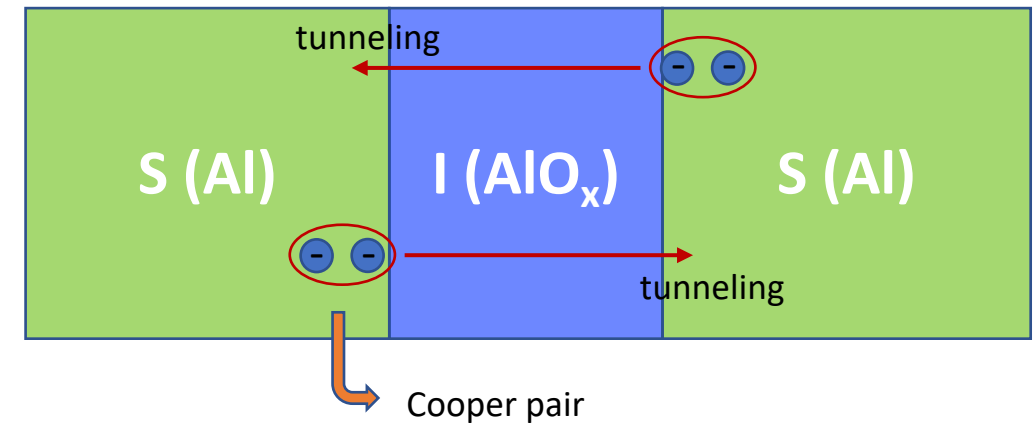
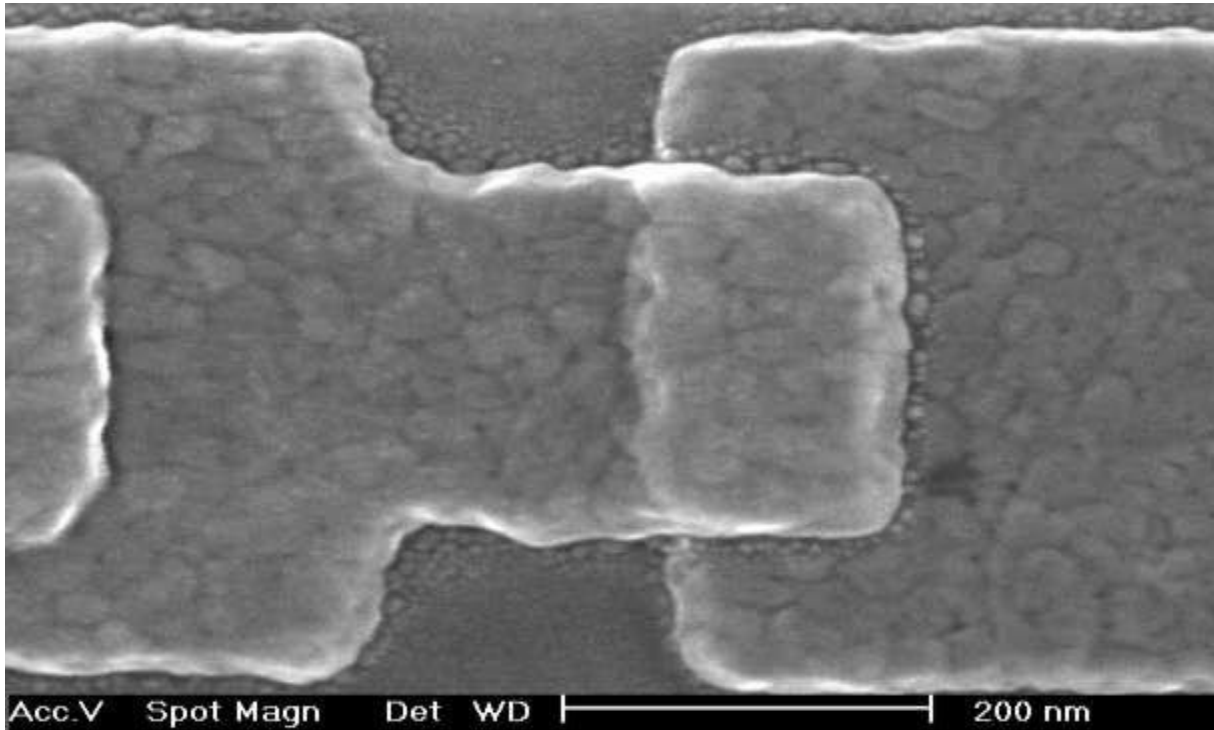
Quantum chip/package



Staggered temperature stages

Illustration: IBM, Qiskit Textbook

The Josephson junction – *anharmonic* oscillator (I)

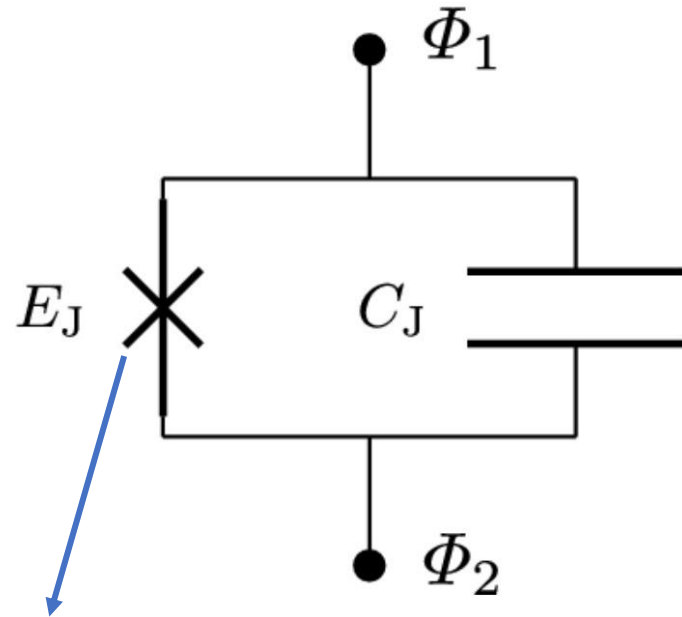


$$I = I_0 \sin(2\pi\Phi/\Phi_0)$$

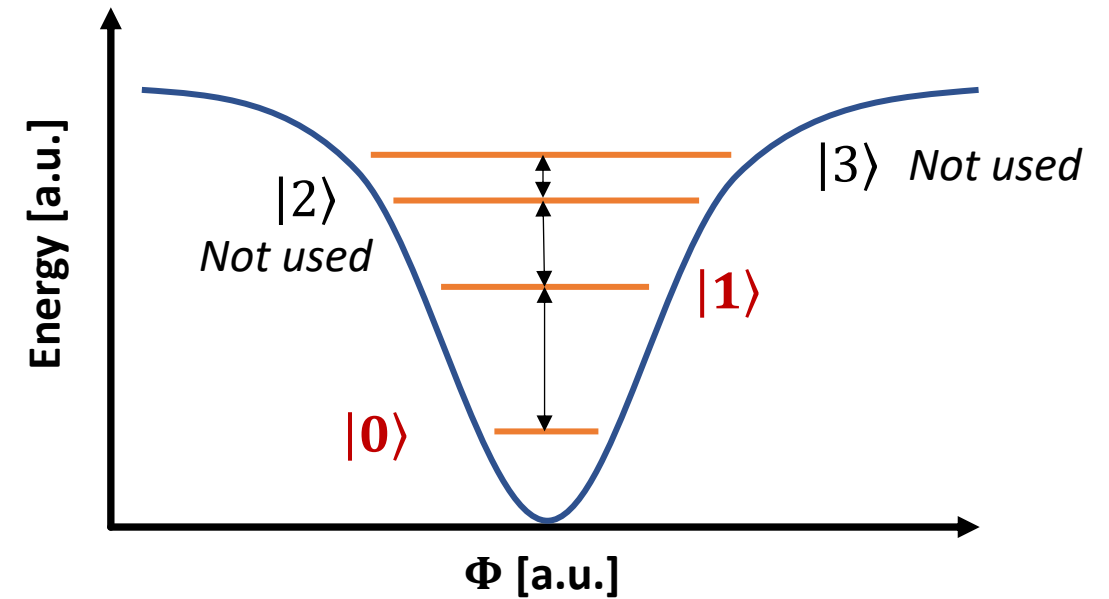


$$L_J = \left(\frac{\partial I}{\partial \Phi} \right)^{-1} = \frac{\Phi_0}{2\pi I_0} \frac{1}{\cos(2\pi\Phi/\Phi_0)}$$

The Josephson junction – *anharmonic* oscillator (II)



Josephson junction
symbol



$$H = \frac{Q^2}{2C} - E_J \cos(2\pi\Phi/\Phi_0)$$

Josephson-junction based qubits

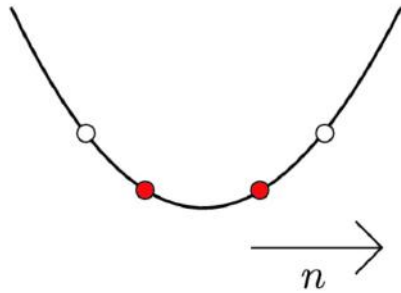
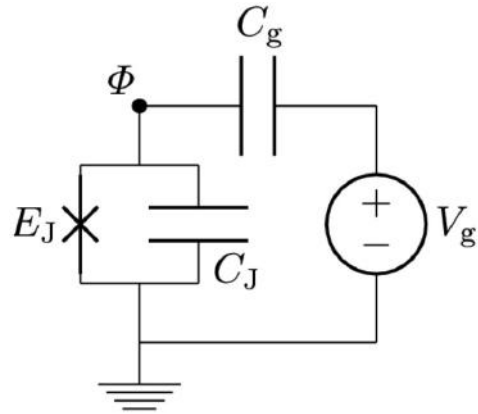


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$$[\phi, n] = i \quad \Rightarrow \quad \Delta\phi\Delta n \geq 1 \quad \text{Heisenberg uncertainty}$$

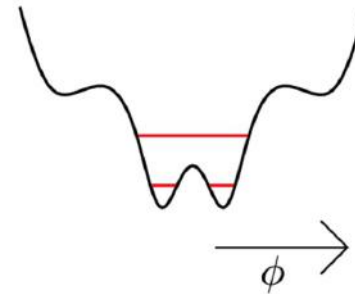
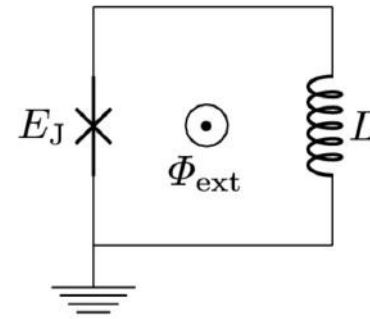
Charge qubit



$$\frac{E_J}{E_C} \ll 1$$

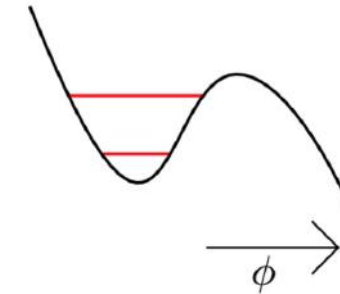
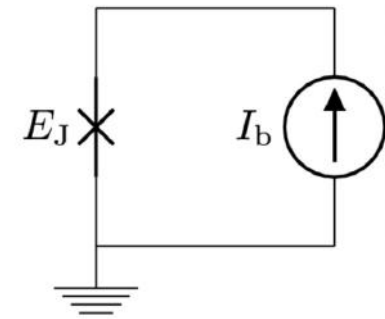
- large ϕ quantum fluctuations
- well defined charge number n

Flux qubit



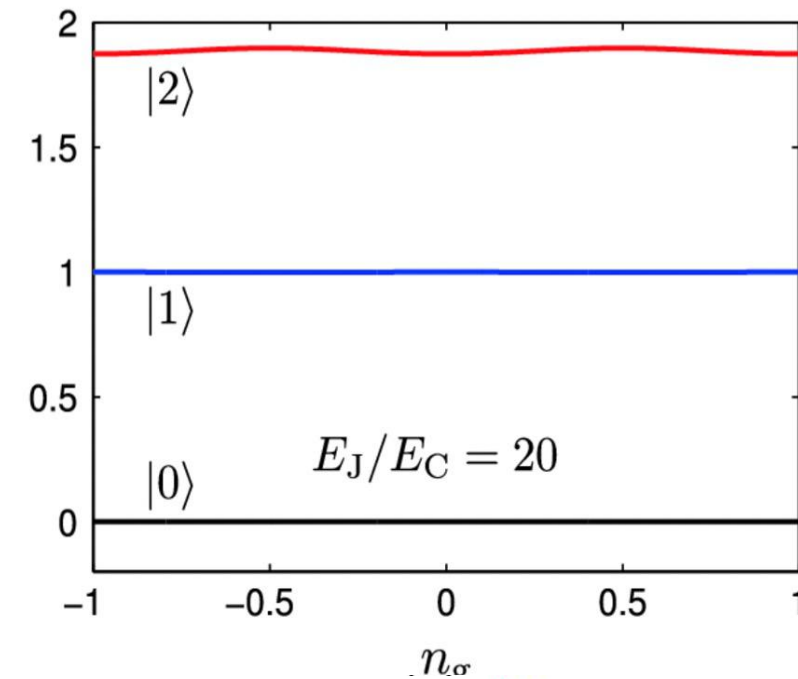
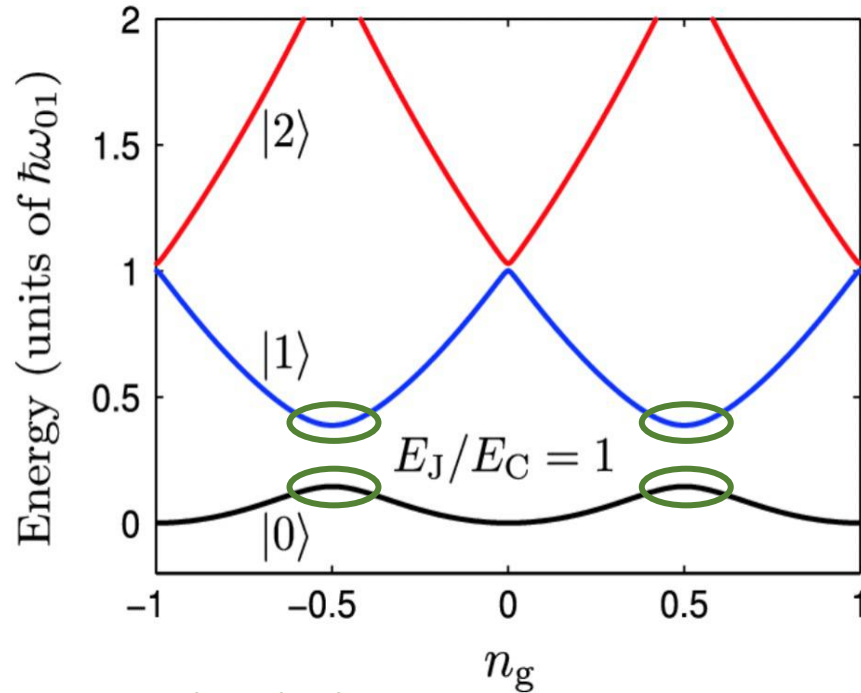
$$\frac{E_J}{E_C} \gg 1$$

Phase qubit



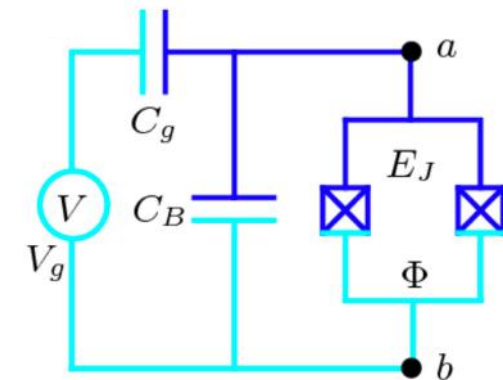
Opposite behavior
w.r.t. Charge qubit

The transmon qubit



sweet spots at
 $n_g = m + \frac{1}{2}$,
 $m \in \mathbb{Z}$

$$H = 4 E_C (n - n_g)^2 + E_J \cos(\Phi)$$



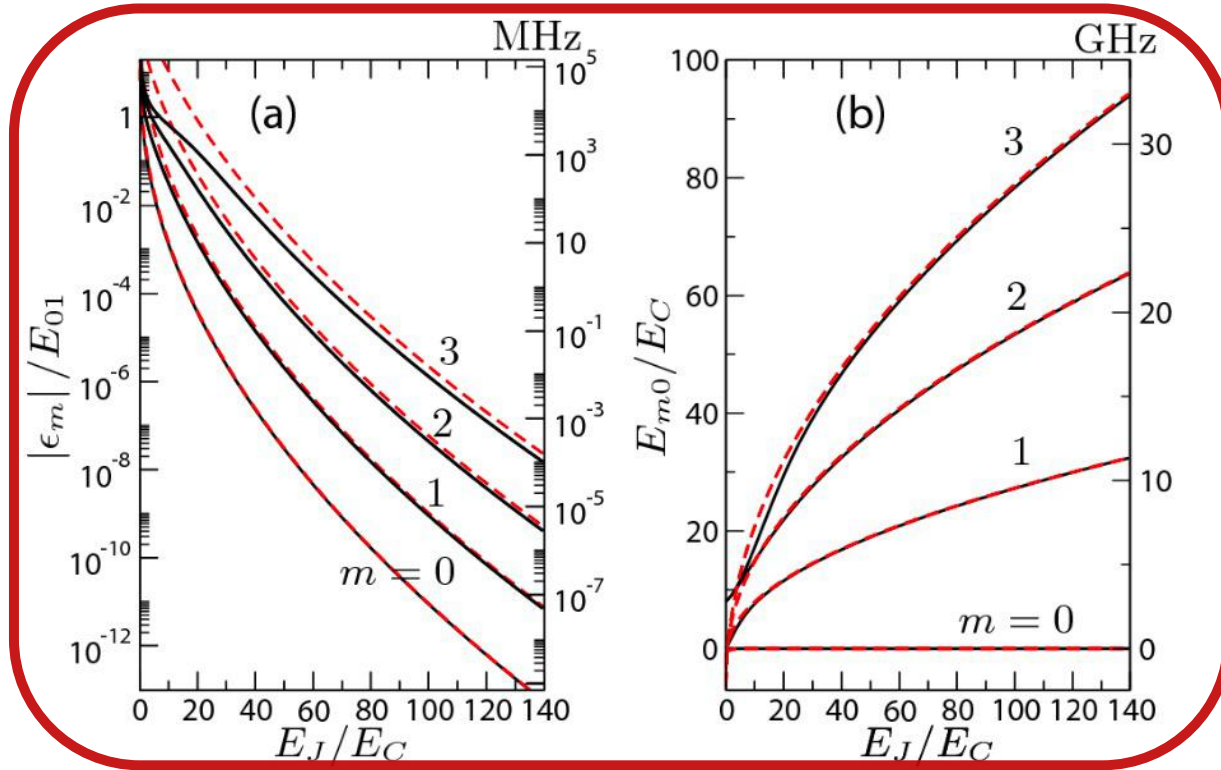
Charge dispersion vs anharmonicity at degeneracy points



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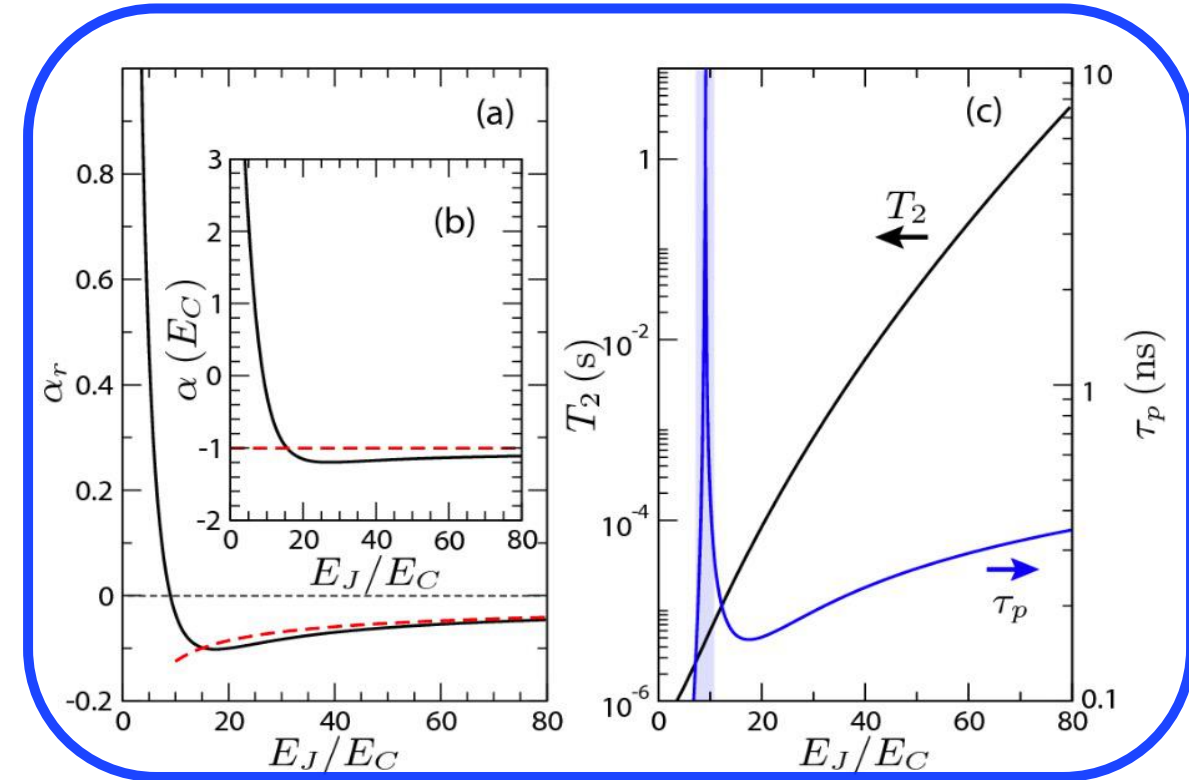


Charge dispersion



$$\epsilon_m \propto e^{-\sqrt{8E_J/E_C}} \text{ for } E_J/E_C \gg 1$$

Anharmonicity



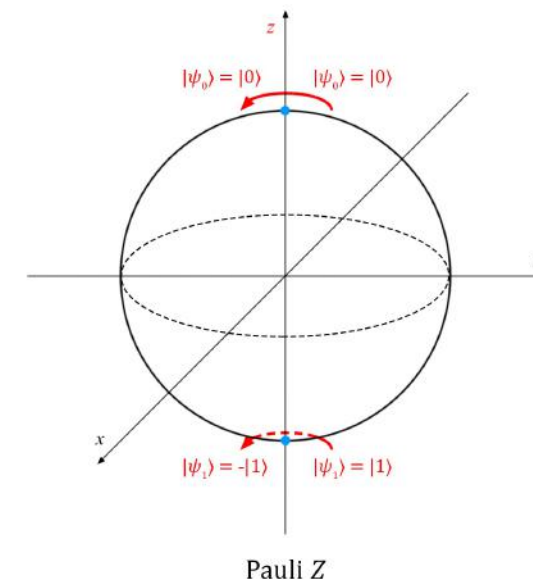
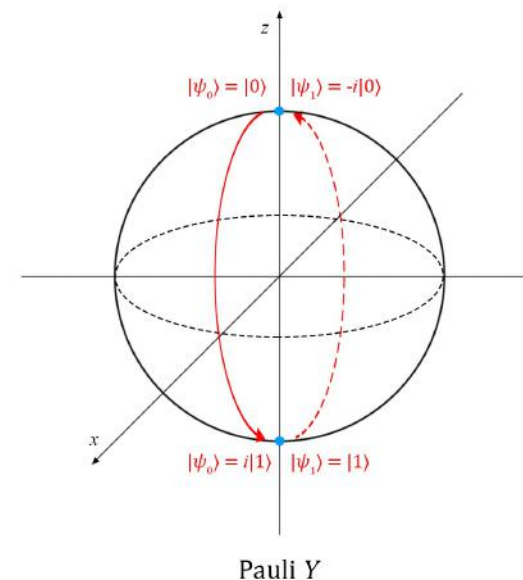
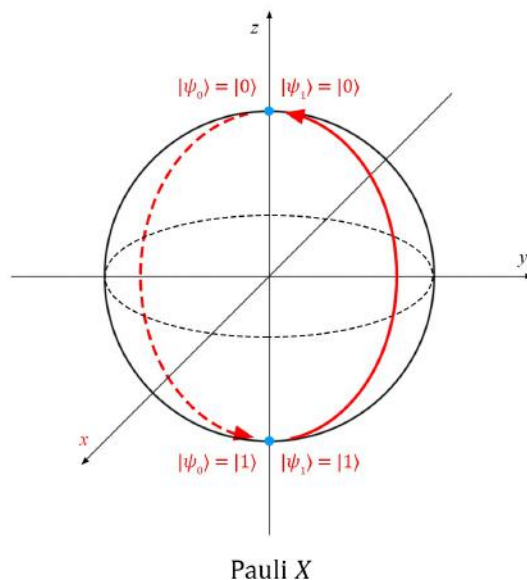
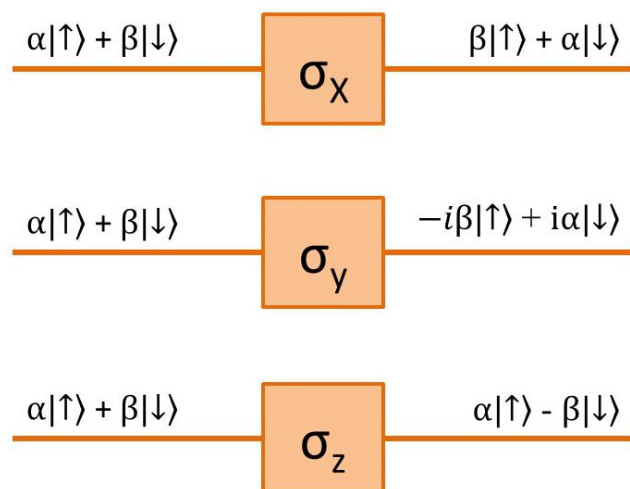
$$\alpha_r \propto -\sqrt{8E_J/E_C} \text{ for } E_J/E_C \gg 1$$

The DiVincenzo criteria



1. **Qubits**: The state must be described as a normalized superposition of two states and the internal Hamiltonian of the system must be known
2. **Initialization**: it must be possible to initialize these qubits to a simple and known state
3. **Gates**: it must be possible to perform both single and two qubit gates on the qubits with high fidelity
4. **Readout**: it must be possible to measure the states of the qubits.
5. **Coherence**: the coherence times of the qubits must be long enough to allow a large number of gates to be performed in sequence before a significant loss of quantum coherence occurs

Qubit single gates



Pauli matrices

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Coupling J-J qubits



Josephson junction coupling



Direct connection:

- capacitively
- inductively

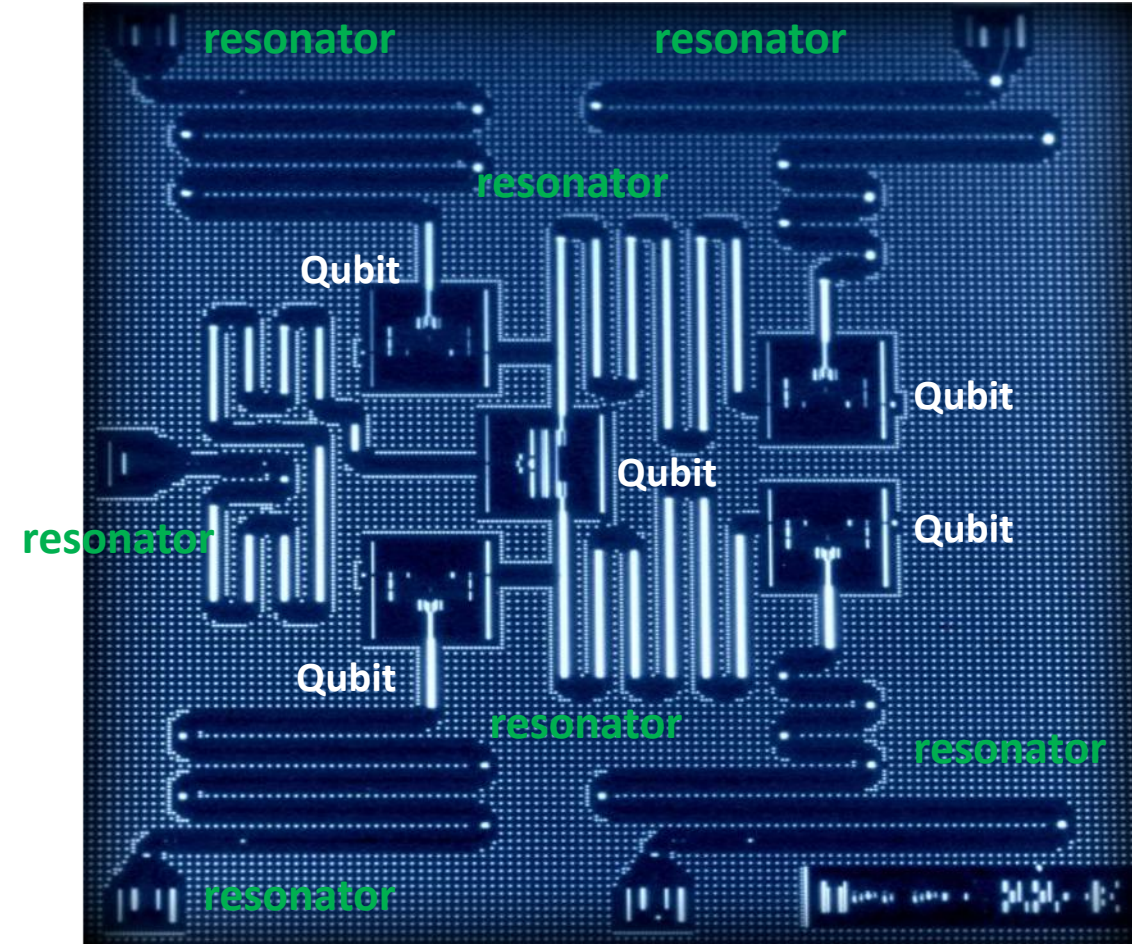


Connection by
intermediate
electrical coupling
circuit:

- LC resonator
- Another J-J qubit
- DC-Squid



Intermediate
quantum bus



Turn off qubit coupling: tuning transition frequencies far from resonance with each other

Qubit readout

Qubit

+

Resonator

+

Interaction

Frequency ω_q

Frequency ω_r

Interaction strength g

$$H = H_{Qub} + H_{Res} + H_{Int}$$

Dispersive interaction
condition:

$$\omega_r - \omega_q \gg g$$

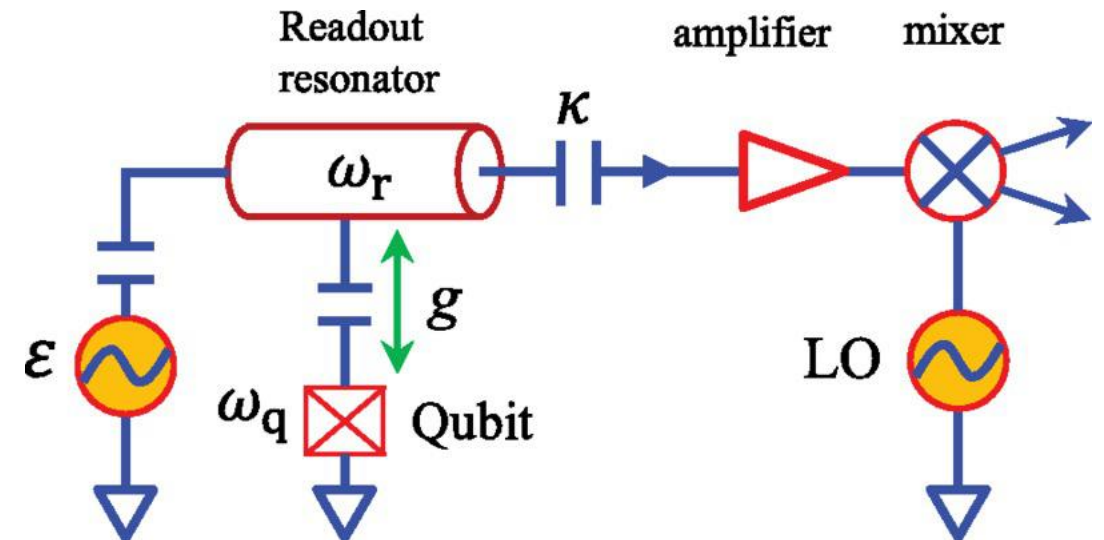
$$H = \hbar \frac{g^2}{\omega_r - \omega_q} (|1\rangle\langle 1| + a^\dagger a \sigma_z)$$

Time evolution:

$$e^{-iHt/\hbar} |0\rangle |\alpha\rangle = |0\rangle |\alpha\rangle e^{-ig^2 t / (\omega_r - \omega_q)}$$

$$e^{-iHt/\hbar} |1\rangle |\alpha\rangle = |1\rangle |\alpha\rangle e^{ig^2 t / (\omega_r - \omega_q)}$$

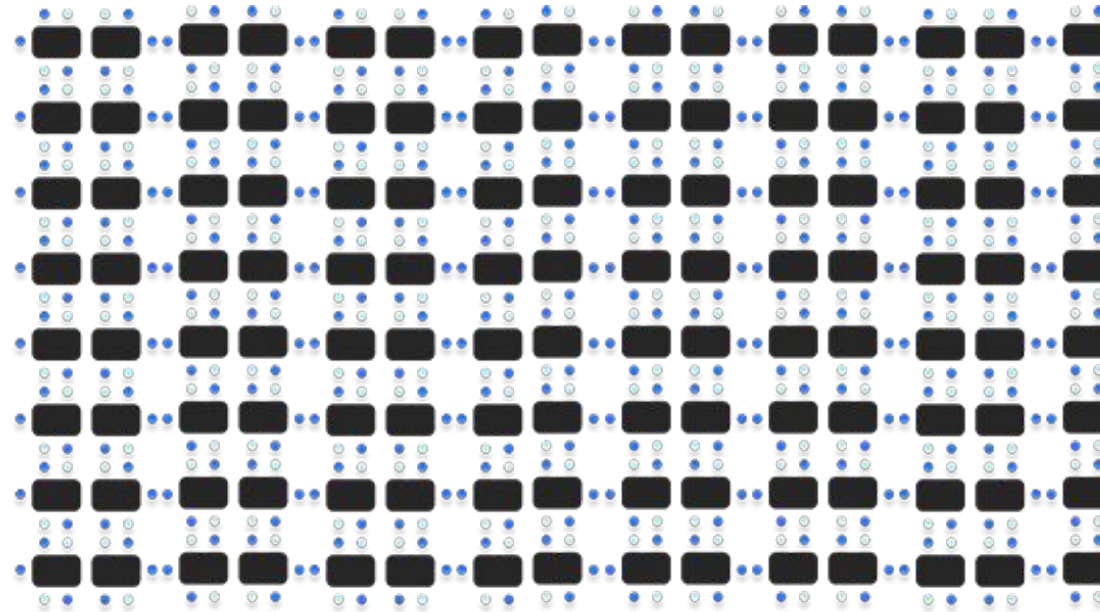
Jaynes-Cummings Hamiltonian



Supercomputer vs quantum computer

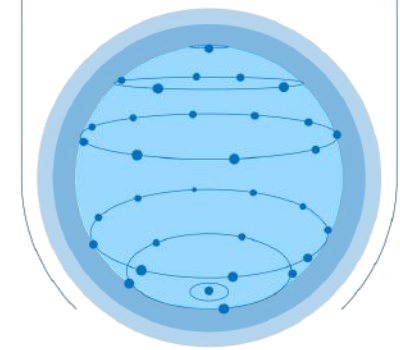


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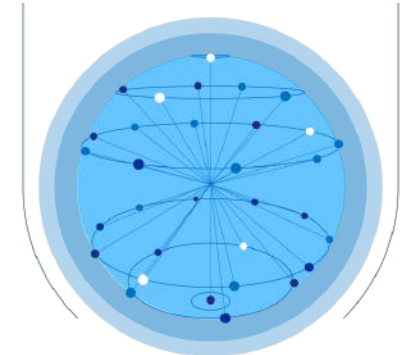


- 5 people – 120 combinations
- 10 people – 3.628.800 combinations!

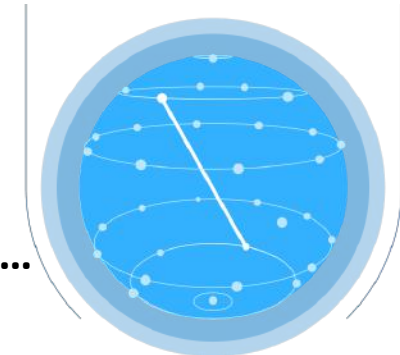
1. **Activation:**
 2^n states



2. **Encoding:**
applying gates



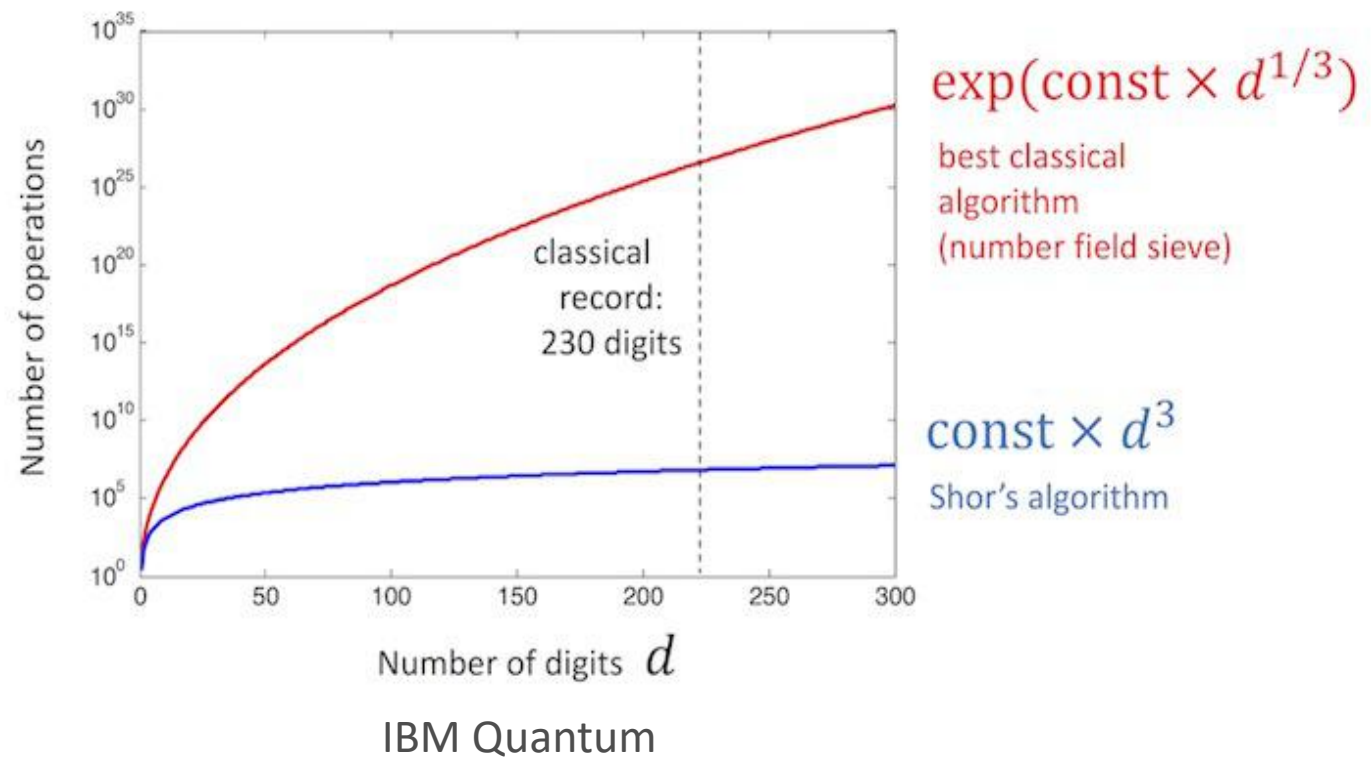
3. **Solution:** using
interference and
translating to 010011...



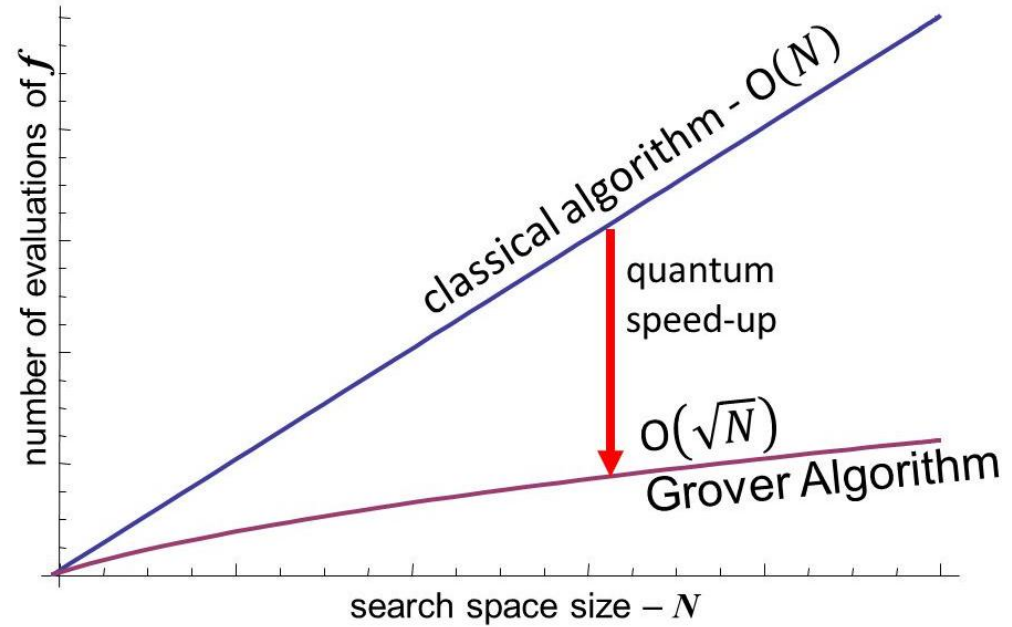
Primes factorization and searching algorithms



Primes factorization



Searching from a list



$M = p \cdot q$
 p, q prime numbers

\longrightarrow

Classical	$t \sim \exp(O(d^3))$	$2.8 \cdot 10^{22}$ years
Quantum	$t \sim O(d^3)$	100 s

Applications



Combinatorial optimization problems:

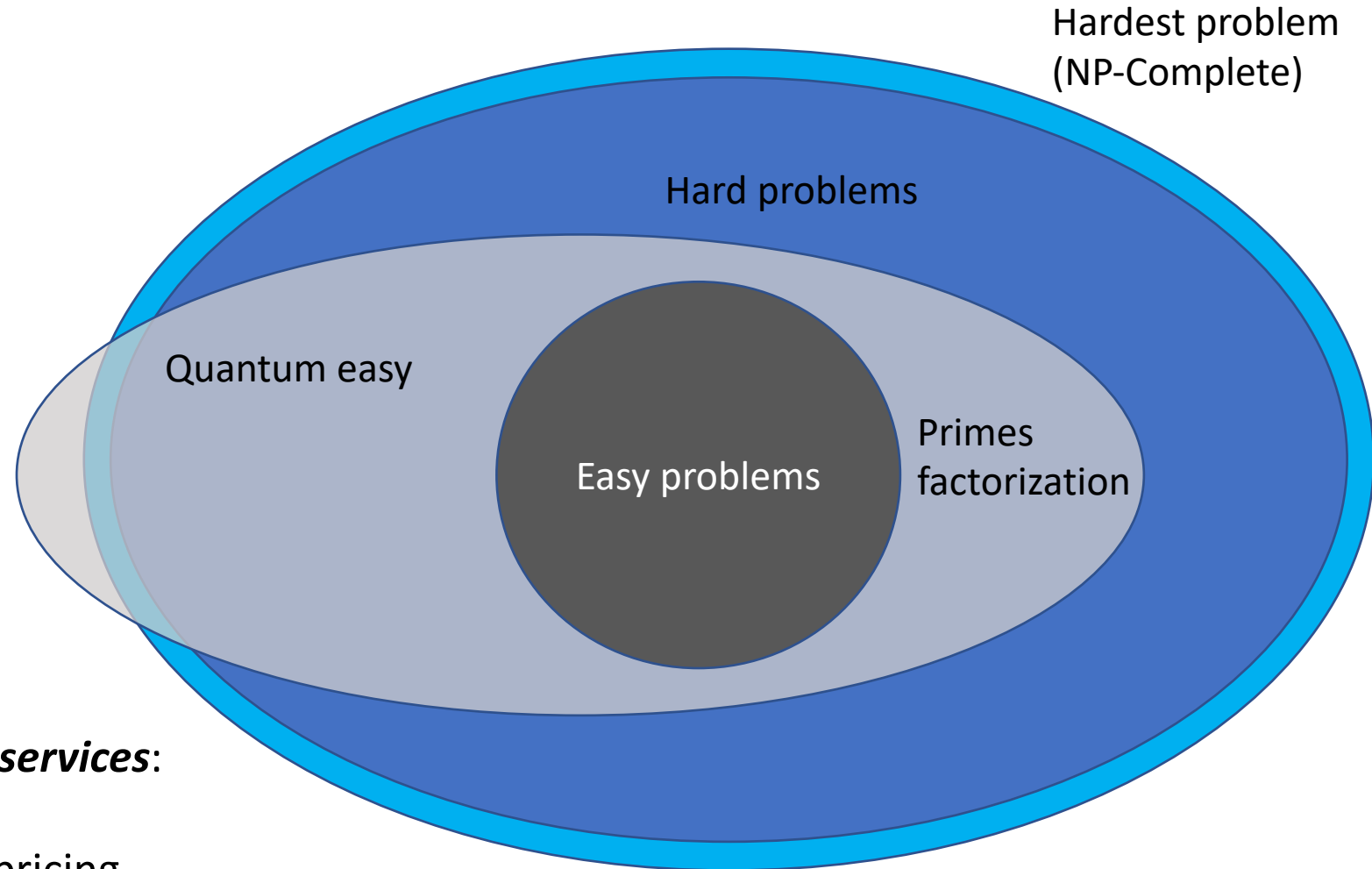
- Assignment problem
- Closure problem
- Constraint satisfaction problem
- Cutting stock problem
- Integer programming
- Knapsack problem
- Minimum spanning tree
- Scheduling problem
- Traveling salesman problem
- Vehicle routing and rescheduling

Chemistry models:

- Material design
- Drug discovery

Financial services:

- Asset pricing
- Risk analysis
- Rare events simulation





Nature isn't classical, dammit, and if
you want to make a simulation of
nature, you'd better make it
quantum mechanical, and by golly
it's a wonderful problem, because it
doesn't look so easy.

— *Richard P. Feynman* —