1 Preambule

Theory box

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Formula box

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Lab/examples box

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2 Fluids as energy carriers

2.1 Fluid state variables and properties

Formulas

2.1.1 State variables

Density

$$\rho \triangleq \frac{m}{V} \left[\frac{kg}{m^3} \right] \tag{1}$$

Specific volume

$$v \triangleq \frac{V}{m} = \frac{1}{\rho} \left[\frac{m^3}{kg} \right] \tag{2}$$

2.1.2 Viscosity

Kinematic viscosity

$$\nu \triangleq \frac{\eta}{\rho} \left[\frac{m^2}{s} \right] \tag{3}$$

Dynamic viscosity

$$\eta \triangleq \nu \cdot \rho \left[Pa \cdot s = \frac{Ns}{m^2} = \frac{kg}{m \cdot s} \right]$$
(4)

2.1.3 Real and ideal fluid

Real fluid Ideal fluid

variable density $(\Delta \rho \neq 0)$ incompressible $(\Delta \rho = 0)$ friction $(\eta > 0, \nu > 0)$ frictionless $(\eta = 0, \nu = 0)$

2.1.4 Compressibility

Mach number

$$M \triangleq \frac{u}{c} \tag{5}$$

where:

- M is the Mach number [-] $M \lesssim 0.3$: incompressible flow
- u is the flow velocity [m/s]
- c is the speed of sound in the fluid [m/s]

and:

- $c_{\rm w}^{20^{\circ}} = 1484 \text{ m/s}$ $c_{\rm a}^{20^{\circ}} = 343 \text{ m/s}$

2.2 Laminar and turbulent flow

Reynolds number

$$Re = \frac{v \cdot L}{\nu} = \frac{\rho \cdot v \cdot L}{\eta} \left[- \right] \tag{6}$$

where:

- v is the mean flow velocity [m/s]
- L is the characteristic length [m]

Re values

- Re < 2000: laminar flow
- $Re \simeq 2300$: critical point
- 2000 < Re < 4000: transitional regime
- $Re \geq 4000$: turbulent flow

2.3 Pressure and velocity

Pressure

2.3.1 Total pressure

Added to the static pressure p_{stat} , there is also the dynamic pressure p_{dyn} and the total pressure p_{tot} :

$$p_{\text{tot}} = p_{\text{stat}} + p_{\text{dyn}} = \rho \left(gh + \frac{v^2}{2} \right)$$
 (7)

2.3.2 Absolute pressure

Absolute pressure p_{abs} refers to the pressure in a vacuum $p_{\text{vaacum}} = 0$ Pa while relative pressure p_{rel} can refer to any chosen reference pressure $p_{\rm ref}$.

$$p_{\rm abs} = p_{\rm rel} - p_{\rm ref} \tag{8}$$

2.3.3 Velocity

Velocity is a vector quantity:

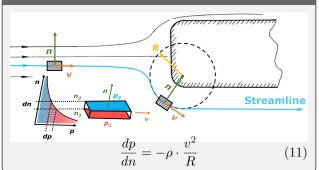
$$\vec{v} = (v_x v_y v_z) \tag{9}$$

The magnitude is given by:

$$v = \sqrt{v_x^2 + v_y^2 + v_z^2} \tag{10}$$

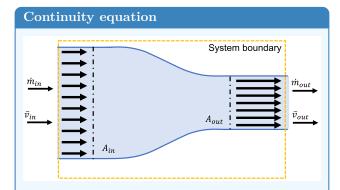
2.4 Curvature pressure formula

Deflection motion of a fluid element around a blunt body



3 Mass conservation

3.1 Continuity equation / Mass conservation



3.1.1 Steady mass-flow

$$\dot{m}_{\rm in} = \dot{m}_{\rm out} \tag{12}$$

3.1.2 Incompressible fluid

$$\dot{m} = \rho \, \dot{V} \implies \dot{V}_{\rm in} = \dot{V}_{\rm out}$$
 (13)

3.1.3 Streamline theory

$$\dot{V} = \bar{v} A \implies \bar{v}_{\rm in} A_{\rm in} = \bar{v}_{\rm out} A_{\rm out}$$
 (14)

4 Energy conservation

4.1 Fluid mechanical energy conservation

Derivation of the Bernoulli equation

$$\dot{m}_1 \left(\frac{p_1}{\rho} + \frac{v_1^2}{2} + gz_1 \right) = \dot{m}_2 \left(\frac{p_2}{\rho} + \frac{v_2^2}{2} + gz_2 \right) \tag{15}$$

This derivation is based on the assumption that the system has:

- steady flow ideal fluid
- no work in or out of the system
- adiabatic process
- 1D streamline flow

4.1.1 Energy flow

$$\frac{dE}{dt} = \underbrace{\sum_{\text{Energy flow across system boundary}}}_{\text{Energy transfer mass in}} + \underbrace{\sum_{in} \left[\dot{m}^{\checkmark} \cdot \left(h^{\checkmark} + \frac{v^{2\checkmark}}{2} + gz^{\checkmark} \right) \right]}_{\text{Energy transfer mass in}} - \underbrace{\sum_{out} \left[\dot{m}^{\nearrow} \cdot \left(h^{\nearrow} + \frac{v^{2\nearrow}}{2} + gz^{\nearrow} \right) \right]}_{\text{Energy transfer}} \tag{16}$$

4.1.2 Outflow formula according to Torricelli

$$gz_1 = \frac{v_2^2}{2} \Longrightarrow v_2 = \sqrt{2g\Delta z} \tag{17}$$

4.2 Bernoulli equation

$$\frac{p_1}{\rho} + \frac{v_1^2}{2} + gz_1 = \frac{p_2}{\rho} + \frac{v_2^2}{2} + gz_2 = \text{const.} \left[\frac{J}{kg} \right]$$
 (18)

4.2.1 Alternative forms

Pressure equation

$$p_1 + \frac{\rho v_1^2}{2} + \rho g z_1 = p_2 + \frac{\rho v_2^2}{2} + \rho g z_2 = \text{const.} [Pa]$$
(19)

Height equation

$$\frac{p_1}{\rho g} + \frac{v_1^2}{2g} + z_1 = \frac{p_2}{\rho g} + \frac{v_2^2}{2g} + z_2 = \text{const.}[m]$$
 (20)

True energy equation

The Bernoulli equation states that the sum of these energies is constant along a streamline.

4.2.2 Pressure energy

$$E_p = m \cdot \frac{p}{\rho} [J] \tag{21}$$

4.2.3 Kinetic energy

$$E_{\rm kin} = m \cdot \frac{v^2}{2} \left[J \right] \tag{22}$$

4.2.4 Potential energy

$$E_{\text{pot}} = m \cdot g \cdot z [J] \tag{23}$$

4.2.5 Energy conservation

$$E_{p,1} + E_{\text{kin},1} + E_{\text{pot},1} = E_{p,2} + E_{\text{kin},2} + E_{\text{pot},2}$$

$$m\left(\frac{p_1}{\rho} + \frac{v_1^2}{2} + gz_1\right) = m\left(\frac{p_2}{\rho} + \frac{v_2^2}{2} + gz_2\right)$$
 (24)

4.3 Hydrostatics

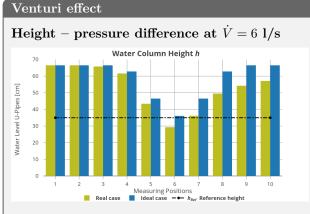
Fundamental law of hydrostatics

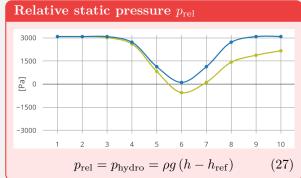
$$p = p_0 + \rho g h = \text{const.} [Pa] \tag{25}$$

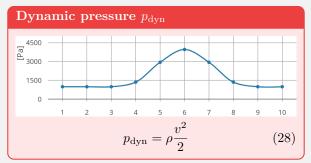
derived from:

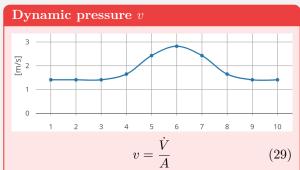
$$p = p_0 + \frac{F_g}{A} = p_0 + \frac{mg}{A} = p_0 + \frac{\rho h Ag}{A}$$
 (26)

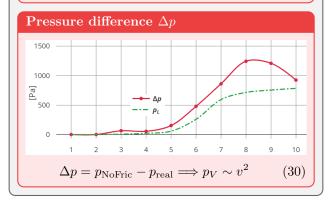
4.4 Venturi effect experiment

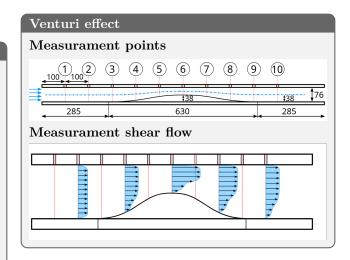




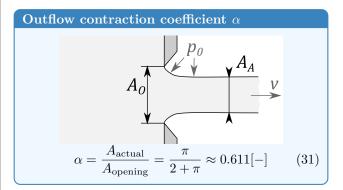




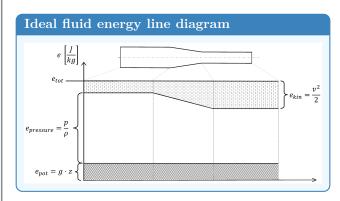


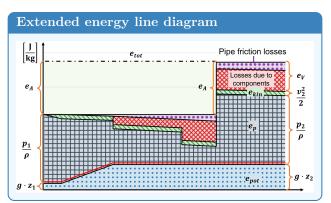


4.5 Contraction coefficient



4.6 Energy line diagram





4.7 Extended Bernoulli equation

Extension of the Bernoulli equation

$$\frac{p_1}{\rho} + \frac{v_1^2}{2} + gz_1 + e_A = \frac{p_2}{\rho} + \frac{v_2^2}{2} + gz_2 + e_V \left[\frac{J}{kg} \right]$$

$$E_{p,1} + K_1 + U_1 + E_A = E_{p,2} + K_2 + U_2 + E_V [J]$$
(32)

4.7.1 Additional terms

Work term e_A

$$e_A = \frac{p_A}{\rho} = gz_A = \frac{E_A}{m} = \frac{P_A}{\dot{m}} \left[\frac{J}{kg} \right]$$
 (33)

where:

 e_A : work term [J/kg] E_A : energy difference [J] p_A : pressure diff [Pa] P_A : power difference [W]

 z_A : height difference [m]

If energy is added to the fluid along a streamline from point 1 to point 2 (eg. a pump), the total energy at point 2 becomes higher than at point 1.

Sign convention

 $e_A > 0$: work is done on the fluid

 \rightarrow energy is added to the fluid (eg. pump);

 $e_A < 0$: work is done by the fluid \rightarrow energy is extracted from the fluid (eg. turbine).

Pump and turbine work Y

In the pressure equation, the pressure p_A increase (or decrease with a turbine) can be read directly at the working term, hence:

$$e_w = Y = \frac{W_A}{\dot{m}} = \frac{E_A}{m} = H \cdot g = \frac{p_A}{\rho} \left[\frac{J}{kq} \right]$$
 (34)

The hydraulic power P_{hyd} is then given by:

$$P_{\text{hyd}} = \dot{m} \cdot Y = \dot{V} \cdot \rho \cdot Y = \rho \cdot \dot{V} \cdot g \cdot H[W]$$
 (35)

Specific loss term e_V

$$e_V = \frac{p_V}{\rho} = gz_V = \frac{E_V}{m} = \frac{P_V}{\dot{m}} \left[\frac{J}{kg} \right]$$
 (36)

where:

 e_V : loss term [J/kg] E_V : energy loss [J] p_V : pressure diff [Pa] P_V : power loss [W]

 z_V : height loss [m]

The effects of a viscous fluid along a stramline from point 1 to point 2 are taken into account by e_V .

Pressure loss Δp_V

$$\Delta p_V = e_V \cdot \rho = \frac{E_V \cdot \rho}{m} = g \cdot z_V \cdot \rho = \zeta \cdot \rho \cdot \frac{v^2}{2} [Pa]$$
(37)

4.8 Loss behavior in turbolent flows

Zeta value

$$\zeta = \frac{2 \cdot \Delta p_V}{\rho \cdot v^2} \tag{38}$$

Total pressure loss

If multiple losses occur in a system due to sequentially connected hydraulic components, the ttal loss $\Delta p_{V,\mathrm{tot}}$ is given by the sum of the individual losses:

$$\Delta p_{V,\text{tot}} = \sum_{i} \Delta p_{V,i} = \sum_{i} \zeta_i \cdot \rho \cdot \frac{v_i^2}{2} \left[Pa \right]$$
 (39)

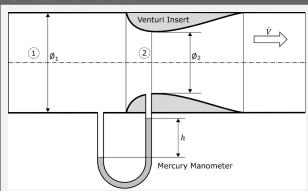
$$\Delta p_{V,\text{tot}} = \rho \cdot \frac{v^2}{2} \cdot \sum_{i} \zeta_i = \rho \cdot \frac{v^2}{2} \cdot \zeta_{\text{tot}} \left[Pa \right] \quad (40)$$

Pressure head (prevalenza)

The pressure head H is the (energy) height corresponding to its specific potential energy e_A :

$$H = \frac{e_A}{g} = \frac{\Delta p_A}{\rho \cdot g} [m] \tag{41}$$

U-Tube manometer



$$h = \frac{\rho \left(v_2^2 - v_1^2\right)}{2g\left(\rho_{\rm Hg} - \rho_w\right)} \tag{42}$$

4.9 Efficiency

Efficiency factor η

$$\eta = \frac{P_{\text{out}}}{P_{\text{in}}} = \frac{\text{Benefit}}{\text{Effort}}$$
(43)

$$\eta_{\text{hyd}} = \frac{P_{\text{real}}}{P_{\text{ideal}}} = \frac{\dot{m} \cdot e_{\text{real}}}{\dot{m} \cdot e_{\text{ideal}}} = \frac{e_A - e_V}{e_A}$$

$$\eta_{\text{hyd}} = \left(= \frac{\Delta e_k + \Delta e_{\text{pot}} + \Delta e_p}{e_A} \right) \tag{44}$$

4.9.1 Volumetric efficiency $\eta_{\rm vol}$

$$\eta_{\text{vol}} = \frac{\dot{m}_{\text{real}}}{\dot{m}_{\text{ideal}}} = \frac{\dot{V}_{\text{real}}}{\dot{V}_{\text{ideal}}}$$
(45)

Efficiency factor η

4.9.2 Efficiency of a pump-driven system

$$\eta_{\text{pump}} = \frac{P_{\text{hyd}}}{P_{\text{mech}}} = \frac{\dot{m} \cdot Y}{M \cdot \omega}$$
(46)

$$\eta_{\text{tot}} = \underbrace{\eta_{\text{el}} \cdot \eta_{\text{mech}} \cdot \eta_{\text{vol}}}_{\text{Pump}} \cdot \eta_{\text{hyd}}^{\text{system}}$$
(47)

In the case of an eletrically driven pump, the effective power transferred to the fluid is thus:

$$P_{\text{eff}} = P_{\text{el}} \cdot \eta_{\text{tot}} \tag{48}$$

4.9.3 Efficiency of a turbine-driven system

$$\eta_{\rm turbine} = \frac{P_{\rm mech}}{P_{\rm hyd}} = \eta_{\rm mech} \cdot \eta_{\rm hyd}$$
(49)

$$\eta_{\text{tot}} = \eta_{\text{turbine}} \cdot \eta_{\text{el}} = \eta_{\text{mech}} \cdot \eta_{\text{hyd}} \cdot \eta_{\text{el}}$$
(50)

5 Pipe flows

5.1 Flow characteristics

Reynolds number in pipes

$$Re = \frac{v_m \cdot d}{\nu} \tag{51}$$

Pipe flows

5.1.1 Laminar pipe flow

The pressure loss of a laminar pipe flow is described by the Hagen-Poiseuille:

$$v(r) = \frac{p_1 - p_2}{4\eta \cdot l} \left(R^2 - r^2 \right)$$
 (52)

$$v_m = \frac{v_{\text{max}}}{2} = \frac{p_1 - p_2}{8\eta \cdot l} \cdot R^2$$

$$v_m = \frac{p_1 - p_2}{32\eta \cdot l} \cdot d^2$$

$$\Delta p = 32\eta \cdot v_m \cdot \frac{l}{d^2} \tag{53}$$

5.1.2 Turbolent flow / Pressure lost in pipelines

Flow losses in pipeline systems consist of pressure losses in straight or curved pipes as well as in fittings.

$$\Delta p = \lambda \cdot \frac{l}{d} \cdot \rho \cdot \frac{v_m^2}{2} \tag{54}$$

where:

 λ : resistance coeff. [-] d: pipe diameter [m] l: pipe length [m] v_m : mean flow velocity [m/s]

Resistance coefficient λ

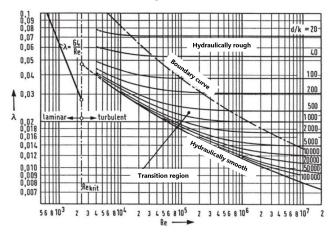
$$\lambda \cdot \frac{l}{d} \cdot \rho \cdot \frac{v_m^2}{2} = 32\eta \cdot v_m \cdot \frac{l}{d^2}$$

$$\lambda = \frac{64\eta}{v_m \cdot d \cdot \rho} = \frac{64}{Re}$$

5.2 Straight pipes

5.2.1 Moody diagram

The resistance coefficient λ depends on the flow characteristics (quantified by the Reynolds number Re) and the relative wall roughness.

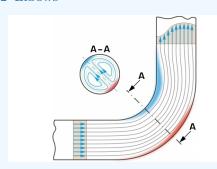


Pipe fittings

In pipeline systems, a portion of the pressure losses is caused by fittings:

$$\Delta p = \zeta \cdot \rho \cdot \frac{v_m^2}{2} \tag{55}$$

5.2.2 Elbows



$$\Delta p = \zeta \cdot \rho \cdot \frac{v^2}{2} \tag{56}$$

$$\zeta = f_{Re} \cdot \zeta_u \tag{57}$$

where (given from individual diagrams):

- ζ_u is the geometric resistance coefficient;
- f_{Re} is the Reynolds correction factor.

5.2.3 Diffuser

A diffuser is a section in a pipeline with a continuous increase in cross-sectional area.

The frictional losses Δp_v in a diffuser are given by:

$$\Delta p_v = \frac{\zeta \rho v_1^2}{2} \tag{58}$$

$$p_2 - p_1 = \Delta p_B - \Delta p_v \tag{59}$$

where Δp_B is the Bernoulli pressure (frictionless).

Pipe fittings

The diffuser efficiency η_D according to Bernoulli:

$$\eta_D = \frac{p_2 - p_1}{\Delta p_B} = 1 - \zeta \frac{1}{1 - \left(\frac{A_1}{A_2}\right)^2}$$
(60)

The various coefficients are stated as:

$$c_p = \frac{2(p_2 - p_1)}{\rho v_1^2} = \eta_D \cdot c_{p,id}$$
 (61)

$$c_{p,id} = 1 - \left(\frac{A_1}{A_2}\right)^2 \tag{62}$$

$$\zeta_1 = c_{p,id} - c_p \tag{63}$$

The opening angle of the diffuser can be calculated

$$\tan(\theta) = \frac{d_2 - d_1}{2L}$$

$$\varphi = 2\theta$$
(64)
$$(65)$$

$$\varphi = 2\theta \tag{65}$$

The optimal angle φ_{opt} is between 6-20 degrees.

5.2.4 Inlets and outlets