Projection and Variational Monte Carlo simulations

A presentation for the course in Computer Simulation

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SISSA - DOCTORATE SCHOOL IN CONDENSED MATTER



Code at:
https://github.com/matteosecli/PMC-VMC

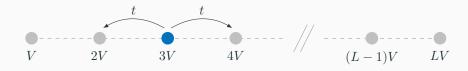
Projection Monte Carlo

The system



Single particle in a 1D lattice with:

- L sites
- A linear (site-dependent) potential $V(x) = V \cdot x$
- \cdot A hopping amplitude t
- · Open boundary conditions



$$H = -t \sum_{\langle xx' \rangle} c_{x'}^{\dagger} c_x + \sum_x V(x) c_x^{\dagger} c_x \tag{1}$$



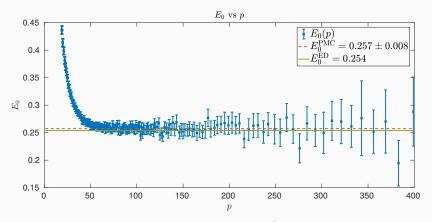


Figure 1: E_0 for L=20, t=1, V=1 and $N=10^7$. The mean position is $x_0=1.47\pm0.04$.



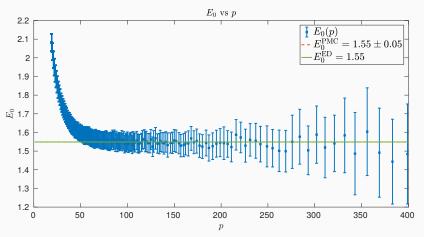


Figure 2: E_0 for L=20, t=1, V=2 and $N=10^7$. The mean position is $x_0=1.20\pm0.04$.



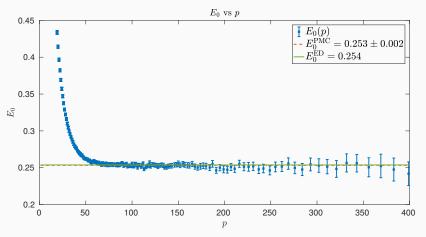


Figure 3: E_0 for L=20, t=1, V=1 and $N=10^8$. The mean position is $x_0=1.48\pm0.01$.



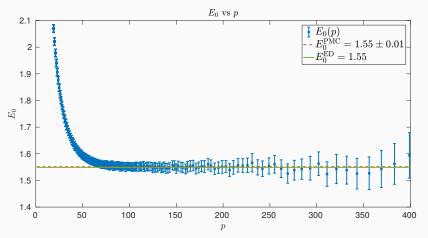


Figure 4: E_0 for L=20, t=1, V=2 and $N=10^8$. The mean position is $x_0=1.21\pm0.01$.

Summary



	$E_0[t]$			x_0		
	ED	PMC	VMC	ED	PMC	VMC
V/t = 1	0.25381	0.253(2)	0.3306(1)	1.461 27	1.48(1)	1.3747(2)
V/t = 2	1.549 13	1.55(1)	1.5773(1)	1.186 22	1.21(1)	1.1566(1)

Table 1: Comparison of results for the ground state energy and mean position obtained with three different methods.

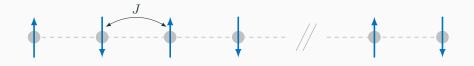
Variational Monte Carlo

The system



Antiferromagnetic spin-1/2 Heisenberg model on a 1D lattice with:

- L sites
- · Periodic boundary conditions
- · Zero magnetization



$$H = J \sum_{i=1}^{L} \vec{S}_i \cdot \vec{S}_{i+1} = J \sum_{i=1}^{L} \left\{ S_i^z S_{i+1}^z + \frac{1}{2} \left(S_i^+ S_{i+1}^- + S_i^- S_{i+1}^+ \right) \right\}$$
 (2)

with J > 0 in order to be antiferromagnetic.

The VMC procedure I



• Variational wavefunction on the basis of the configurations $\{|x\rangle\}$ with definite S_i^z :

$$\psi(x) = \operatorname{Sign}_{M}(x)e^{\frac{\alpha}{2}\sum_{i \neq j} v_{i,j}^{z}(2S_{i}^{z})(2S_{ij}^{z})}$$
(3)

with $\mathrm{Sign}_M(x) = (-1)^{\sum_{i=1}^{L/2} (S^z_{2i} + 1/2)}$ and $v^z_{i,j} = 2\log(|2\sin(\pi(i-j)/L)|).$

· Local energy:

$$E_L(x) = \frac{\langle x \mid H \mid \psi \rangle}{\langle x \mid \psi \rangle} = \langle x \mid H \mid x \rangle + \sum_{x' \neq x} \langle x \mid H \mid x' \rangle \frac{\langle x' \mid \psi \rangle}{\langle x \mid \psi \rangle}$$
(4)

· Staggered magnetization:

$$m^2 = \left\langle \left(\frac{1}{L} \sum_{i=1}^{L} (-1)^i S_i^z\right)^2 \right\rangle \tag{5}$$

The VMC procedure II



· First neighbor correlation function:

$$C_1 = \left\langle \frac{1}{L} \sum_{i=1}^{L} S_i^z S_{i+1}^z \right\rangle. \tag{6}$$

Exact results from Bethe ansatz in $1D^1$ ($L \to \infty$):

$$E_0 = -|J| \left(\log 2 - \frac{1}{4}\right) \simeq -0.4431471|J|$$

and

$$C_1 = \frac{1}{12} (1 - 4 \log 2) \simeq -0.14771573.$$

¹See Minoru Takahashi, *Thermodynamics of One-Dimensional Solvable Models*. Cambridge University Press, 2005. ISBN 9780521019798.



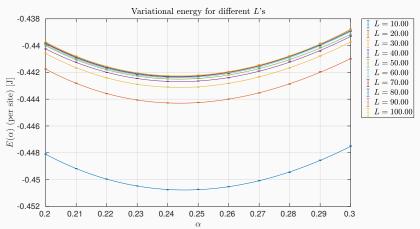


Figure 5: Variational energy. The best value for the variational parameter is estimated to be $\alpha = 0.244 \pm 0.001$.



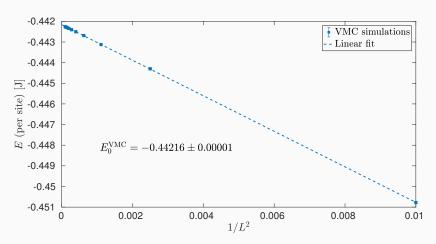


Figure 6: Extrapolation for $L \to \infty$.



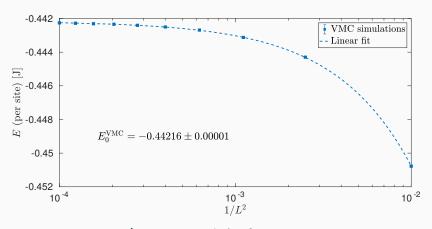


Figure 7: Extrapolation for $L \to \infty$.



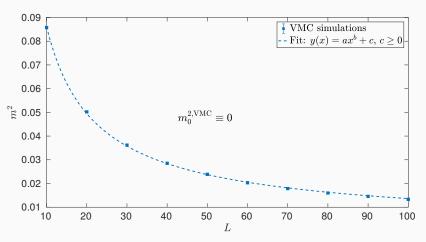


Figure 8: Extrapolation for $L \to \infty$.



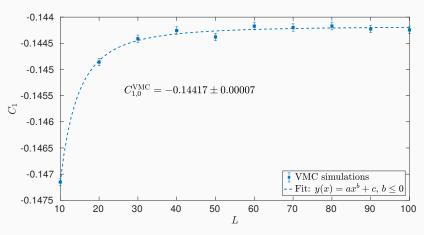


Figure 9: Extrapolation for $L \to \infty$.



Method	$E_0[J]$	m^2	C_1
Bethe ansatz ²	-0.443147	0	-0.147716
VMC^3	-0.44216(1)	0	-0.14417(7)
ED/Lanczos ⁴	-0.44304(5)	-	_

 $^{^2}L \to \infty$, exact.

 $^{^3}L \to \infty$, upper bound.

 $^{^4\}text{Extrapolated to}\ L\to\infty$ from a set of calculations for $L=\{8,10,12,14,16\}$, with max 32 steps for the Lanczos. Code: http://www.fisica.uniud.it/~giannozz/Corsi/MQ/Software/C/heisenberg_exact.c

Conclusions



Hope this was at least a little bit interesting...

Thanks for your attention!







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