Line current  $A = \frac{\mu_o}{4\pi} \int \frac{\vec{I}}{r} dl$ 

Surface current  $A = \frac{\mu_o}{4\pi} \int \frac{\vec{k}}{r} da$ 

 $1 q^2$ 

 $\frac{1}{4\pi\varepsilon_o} \frac{1}{2d}$ 

W = -

(a)  $\mathbf{k} = -\frac{\omega}{c}\hat{\mathbf{x}}; \ \hat{\mathbf{n}} = \hat{\mathbf{z}}. \ \mathbf{k} \cdot \mathbf{r} = \left(-\frac{\omega}{c}\hat{\mathbf{x}}\right) \cdot (x \hat{\mathbf{x}} + y \hat{\mathbf{y}} + z \hat{\mathbf{z}}) = -\frac{\omega}{c}x; \ \mathbf{k} \times \hat{\mathbf{n}} = -\hat{\mathbf{x}} \times \hat{\mathbf{z}} = \hat{\mathbf{y}}.$ 

 $\mathbf{E}(x,t) = E_0 \cos\left(\frac{\omega}{c}x + \omega t\right)\hat{\mathbf{z}}; \quad \mathbf{B}(x,t) = \frac{E_0}{c}\cos\left(\frac{\omega}{c}x + \omega t\right)\hat{\mathbf{y}}.$ 

1  $q^2$  $-\frac{1}{4\pi\varepsilon_o}\frac{q^2}{\left(2d\right)^2}\,\hat{z}$ 

Energy with 2 points charges & no conductor

 $\frac{1}{4\pi\varepsilon_o} \int x \frac{r}{r^2} dl \underline{\qquad} x = (\lambda or\sigma or\rho)$ 

E-field=0 inside conductor
Volume charge=0 inside conductor
Conductor net charge resides on surface
Hollow of conductor has e-field and charge just outside

E=

here here $\frac{1}{r}$ , and $\frac{1}{r}$ ,	TOT SAVART LAW $ \frac{\mu_0}{4\pi} \int \frac{\mathbf{I} \mathbf{x} \hat{\mathbf{h}}}{\mathbf{r}^2} \hat{\mathbf{d}} $ $ \frac{\mu_0}{4\pi} \int \frac{\mathbf{I} \mathbf{x} \hat{\mathbf{h}}}{\mathbf{r}^2} \hat{\mathbf{d}} $ $ \frac{\mathbf{I} \mathbf{x} \hat{\mathbf{h}}}{4\pi} \int \frac{\mathbf{I} \mathbf{x} \hat{\mathbf{h}}}{\mathbf{r}^2} \hat{\mathbf{d}} $ MAXWELLS FIQ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{B} = 0 $ MAXWELLS FIQ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ $ \nabla \cdot \mathbf{E} = \rho / \epsilon_0 $ To the theorem for the cost of the sphere of th
surface current density $K = \frac{1}{2\pi a} J = \frac{1}{2\pi a}$ volume s.6  Problem 5.6  K at a distance r from center $K = \sigma \omega r$ . Current density J point $(r, \theta, \phi) \omega r \sin \theta \phi$ , Problem 5.8 Bfield center of a square loop $\frac{\sqrt{2\mu_0 I} / \pi R}{(r + \frac{1}{2})^3}  ^{-2\pi L} $ Problem 5.10 Force square loop $\frac{\mu_0 I}{(r + \frac{1}{2})^3}  ^{-2\pi L} $ Problem 5.10 Force square loop $\frac{\mu_0 I}{(r + \frac{1}{2})^3}  ^{-2\pi L} $ Problem 5.10 Force square loop $\frac{\mu_0 I}{2\pi R} \sin(\pi/n).$ Problem 2.14 field inside a sphere $\rho = kr$ Problem 2.15 A long coaxial cable $\rho$ on inner urface charge on the outer $\frac{\mu_0 I}{2\pi R} \sin(\pi/n)$ out $E = \frac{\pi}{2} \pi kr^2 f$ Problem 2.30 field inside and outside a long ollow cylindrical tube, out $\frac{\pi}{6}$ is $m = 0$ .  Problem 3.34 point charge released from rest ow long will it take $t = (\pi d/q) \sqrt{2\pi \epsilon_0 m d}$	MACCETIC VECTOR POTINIAL     MACCETIC VECTOR POTINIAL     A = $\frac{\mu_0}{4\pi} \int_{\Gamma} \frac{K}{I} da = \frac{\mu_0}{4\pi} \int_{\Gamma} \frac{K}{I} da = \frac{\mu_0}{I} \int_{\Gamma} \frac{K}{I} da $
Problem 7.17 long solenoid looped by a wire resistance $\sum_{Z\pi} \sin(\omega t) \ln \binom{\alpha}{\alpha} 2t$ . longitudinal Problem 7.17 long solenoid looped by a wire resistance $\sum_{Z\pi} \frac{2\pi\alpha^2 J_0 n I}{R}$ . If square loop distance $s$ from infinite straight wire someone cuts the wire $I(t) = \begin{cases} (1-\alpha t)I. & \text{for } 0 \leq t \leq 1/\alpha. \\ 0. & \text{for } t > 1/\alpha. \end{cases}$ $Q = \frac{I_{LOG} \ln 2}{2\pi R}  \text{induced current flows}$ $Q = \frac{2\pi R}{2\pi R}  \text{counterclockwise}$ Problem 7.45 conducting spherical shell rotates $E = 2$ $B = B_0 2.  \text{cmf} = \frac{1}{2}B_{00}\alpha^2.$ Problem 2.46 The electric potential $e^{-\lambda t}$ $E = Ae^{-\lambda t}(1+\lambda t) - \frac{k}{r} = -\lambda t $ $E = Ae^{-\lambda t}(1+\lambda t) - \frac{k}{r} = -\lambda t$	ANDEREYS LAW V-B=0 $\oint B$ -da=0   $\oint B$ -da=0   $\oint B$ -da=0   $\oint B$ -da=  $\bigvee A$   $\bigvee A$
IN: $\frac{E}{2\pi R^3} \frac{\partial}{\partial} OUI: \frac{E}{2\pi R^3} \frac{\partial}{\partial} \frac{\partial}{\partial R} \frac{(1.05 \times 10^{-3})}{2\pi R} = \frac{\mu_0 N^2 N\omega}{2\pi R} \ln(b) \frac{\partial}{\partial R}$ Problem 7.20 small loop of wire distance z above the center of a large loop by $\frac{\partial}{\partial R} = \frac{\mu_0 \pi I \alpha^2 b^2}{2(b^2 + z^2)^{3/2}}$ .  Current I flows in the little loop $\frac{\partial}{\partial R} = \frac{\mu_0 \pi I \alpha^2 b^2}{2(b^2 + z^2)^{3/2}}$ mutual inductance $\frac{(\mu_0 \pi \alpha^2 b^2)}{2(b^2 + z^2)^{3/2}}$ .  Problem 7.22 self-inductance per unit length of a long solenoid $\hat{L} = \frac{\mu_0 \pi I \alpha^2 b^2}{4(b^2 + z^2)^{3/2}}$ .  Problem 7.24 $I_0 \cos(\omega t)$ (amplitude 0.5 A, frequency solenoid $\hat{L} = \frac{\mu_0 \pi I \alpha^2 b^2}{4(b^2 + z^2)^{3/2}}$ .  Problem 2.15 A hollow spherical shell carries charge density $\rho = \frac{k}{r^2}$ in between $\frac{k}{G}$ ( $\frac{E}{r^2}$ ) $\hat{r}$ . E $\frac{k}{G}$ ( $\frac{b - a}{r^2}$ ) $\hat{r}$ .  Problem 3.28 spherical shell of radius $R$ $\sigma = k \cos \theta$ .  Part $\frac{k}{R^3}$ $\frac{E}{R^3}$ as $\frac{k}{R^3}$ $\frac{E}{R^3}$ Dipole.  Problem 4.15 A thick spherical shell $\frac{k}{r}$ $\frac{E}{r}$ $\frac{k}{G}$ $\frac{E}{r}$	EMWAYES S= WC S=
$A \cdot (A \times B) = 0 \qquad A \cdot (B \times C) = (A \times B) \cdot C \qquad A \times B = -(B \times A)$ $\nabla (fg) = f \nabla g + g \nabla f \qquad \nabla \cdot fA = f(\nabla \cdot A) + A \cdot (\nabla f)$ $\nabla (A \cdot B) = A \times (\nabla xB) + B \times (\nabla xA) + (A \cdot \nabla)B + (B \cdot \nabla)A$ $\nabla \cdot (AxB) = B \cdot (\nabla xA) - A \cdot (\nabla xB) \qquad \nabla x (fA) = f(\nabla xA) - Ax(\nabla f)$ $\nabla x (AxB) = (B \cdot \nabla)A - (A \cdot \nabla)B + A(\nabla \cdot B) - B(\nabla \cdot A)$ $\nabla x (\nabla s) = 0 \qquad \nabla (\nabla \cdot v) = v \text{ field} \qquad \nabla \cdot (\nabla s) = \text{Laplacian}$ $\nabla \cdot (\nabla x v) = 0 \qquad \nabla x (\nabla x v) = \nabla(\nabla \cdot v) - \nabla^{4}v$ $\int \nabla T \cdot dI = T(b) - T(a) \qquad \text{flux } v = \int (v \cdot \hat{\mathbf{n}}) d\mathbf{a} = \int v \cdot d\mathbf{a}$	$\begin{aligned} \mathbf{F} \mathbf{E} \cdot \mathbf{da} &= \mathbf{Q}_{\text{enc}} / \varepsilon_{o}  \nabla \cdot \mathbf{E} = \rho / \varepsilon_{o} \\ \nabla \mathbf{x} \mathbf{E} &= 0  \frac{\sigma}{\varepsilon_{o}} \text{ between plates} \\ \mathbf{P} \text{ ont charge } \mathbf{E} &= \mathbf{k} \mathbf{Q} \\ \mathbf{D} \text{ infinite line } \mathbf{E} &= \frac{\lambda}{2\pi \varepsilon_{o}z} \end{aligned} $ $\begin{aligned} \mathbf{D} \mathbf{H} \mathbf{M} \mathbf{S} \mathbf{L} \mathbf{A} &= \mathbf{A} \mathbf{F} \mathbf{a} \mathbf{A} \end{aligned} $ $\begin{aligned} \mathbf{D} \mathbf{H} \mathbf{D} \mathbf{M} \mathbf{S} \mathbf{L} \mathbf{A} \mathbf{B} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} \mathbf{C} C$