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5.60 Thermodynamics & Kinetics Spring 2008

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Reversible Adiabatic Expansion (or compression) of an Ideal Gas

1 mole gas $(V_1, T_1) = 1$ mole gas (V_2, T_2)

adiabatic
$$\Rightarrow$$
 $dq = 0$ Reversible \Rightarrow $dw = -pdV$
Ideal gas \Rightarrow $dU = C_vdT$

 \therefore From 1st Law $dU = -pdV \Rightarrow C_vdT = -pdV$ along path

$$\mathcal{C}_{V}dT = -pdV \quad \underset{p = \frac{RT}{V}}{\Longrightarrow} \quad \mathcal{C}_{V}\frac{dT}{T} = -R\frac{dV}{V}$$

$$\mathcal{C}_{V}\int_{T_{1}}^{T_{2}}\frac{dT}{T} = -R\int_{V_{1}}^{V_{2}}\frac{dV}{V} \quad \Rightarrow \quad \left(\frac{T_{2}}{T_{1}}\right) = \left(\frac{V_{1}}{V_{2}}\right)^{R/C_{V}} \qquad \xrightarrow{C_{p}-C_{V}=R \text{ for i.g.}} \quad \left(\frac{T_{2}}{T_{1}}\right) = \left(\frac{V_{1}}{V_{2}}\right)^{C_{p}-1}$$

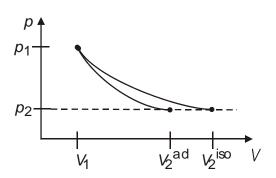
Define $\gamma \equiv \frac{C_p}{C_V} \implies \left| \left(\frac{T_2}{T_1} \right) = \left(\frac{V_1}{V_2} \right)^{\gamma - 1} \right|$

In an adiabatic <u>expansion</u> $(V_2 > V_1)$, the gas cools $(T_2 > T_1)$. And in an adiabatic <u>compression</u> $(V_2 < V_1)$, the gas heats up.

For an ideal gas (one mole) $T = \frac{pV}{R} \Rightarrow \left(\frac{p_2}{p_1}\right) = \left(\frac{V_1}{V_2}\right)^{\gamma} \Rightarrow \boxed{p_1V_1^{\gamma} = p_2V_2^{\gamma}}$

 pV^{γ} is constant along a reversible adiabat

For an isothermal process $T = constant \Rightarrow pV = constant$



 $V_2^{
m adiabat}$ < $V_2^{
m isotherm}$ because the gas cools during reversible adiabatic expansion

• <u>Irreversible</u> Adiabatic Expansion of an ideal gas against a constant external pressure

1 mol gas
$$(p_1, T_1) = 1$$
 mol gas (p_2, T_2) $(p_{ext}=p_2)$

adiabatic
$$\Rightarrow$$
 $dq = 0$

Constant $p_{\text{ext}} = p_2$ \Rightarrow $dw = -p_2 dV$

Ideal gas \Rightarrow $dU = C_v dT$
 $du = -p_2 dV$

$$\therefore \quad C_{v} d T = - p_{2} d V$$

Integrating: $C_v(T_2 - T_1) = -p_2(V_2 - V_1)$

Using
$$pV = RT$$

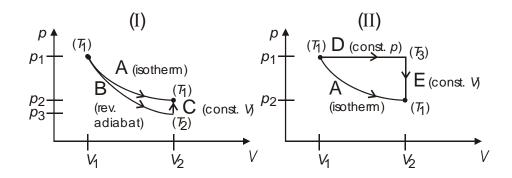
$$T_2(C_V + R) = T_1\left(C_V + \frac{p_2}{p_1}R\right)$$

Note $p_2 < p_1 \Rightarrow T_2 < T_1$ Again, expansion cools

Note also $(-w_{rev}) > (-w_{irrev})$ Less work is recovered through an irreversible process

Some Thermodynamic Cycles

• Reversible Ideal Gas processes: Find ΔU , ΔH , q, w, $\int \frac{dq}{T}$



[A] 1 mol gas $(p_1, V_1, T_1) \stackrel{\text{const. } T}{=} 1 \text{ mol gas } (p_2, V_2, T_1)$

Ideal gas isotherm:

$$\Delta U_{A} = 0$$

$$\Delta H_{A} = 0$$

$$W_A = -RT_1 \ln \frac{V_2}{V_1}$$
 $q_A = RT_1 \ln \frac{V_2}{V_1}$ $\int \frac{dq}{T} = R \ln \frac{V_2}{V_1}$

[B] 1 mol gas $(p_1, V_1, T_1) \stackrel{\text{rev.adiabat}}{=} 1 \text{ mol gas } (p_3, V_3, T_2)$

Adiabat: $q_B = 0$

Ideal gas: $\Delta U_{\text{B}} = C_{\text{V}} (T_2 - T_1)$ $\Delta \mathcal{H}_{\text{B}} = C_{\text{p}} (T_2 - T_1)$

1st Law: $w_B = C_V (T_2 - T_1)$

$$\int \frac{\mathrm{d}q_{\mathrm{B}}}{T} = 0$$

[C] 1 mol gas
$$(p_3, V_2, T_2) \stackrel{\text{reversible}}{=} 1$$
 mol gas (p_2, V_2, T_1)

Constant V:

$$w_c = 0$$

1st Law:

$$q_{c} = C_{V} \left(T_{1} - T_{2} \right)$$

$$\int \frac{\mathrm{d}q_{\mathcal{C}}}{\mathcal{T}} = \mathcal{C}_{\mathcal{V}} \ln \left(\frac{\mathcal{T}_1}{\mathcal{T}_2} \right)$$

[A] vs.
$$[B] + [C]$$

$$\Delta U_{A} = 0$$

$$\Delta H_{B} + \Delta U_{C} = 0 = \Delta U_{A}$$

$$\Delta H_{B} + \Delta H_{C} = 0 = \Delta H_{A}$$

$$\begin{aligned} q_{\mathrm{A}} &= R \, T_{1} \ln \frac{V_{2}}{V_{1}} & q_{\mathrm{B}} + q_{\mathrm{C}} &= \mathcal{C}_{V} \left(T_{1} - T_{2} \right) \neq q_{\mathrm{A}} \\ w_{\mathrm{A}} &= -R \, T_{1} \ln \frac{V_{2}}{V_{1}} & w_{\mathrm{B}} + w_{\mathrm{C}} &= \mathcal{C}_{V} \left(T_{2} - T_{1} \right) \neq w_{\mathrm{A}} \\ \int \frac{\mathrm{d}q_{\mathrm{A}}}{T} &= R \ln \frac{V_{2}}{V_{1}} & \int \frac{\mathrm{d}q_{\mathrm{B}}}{T} + \int \frac{\mathrm{d}q_{\mathrm{C}}}{T} &= R \ln \left(\frac{V_{2}}{V_{1}} \right) = \int \frac{\mathrm{d}q_{\mathrm{A}}}{T} \end{aligned}$$

This result suggests that $\left(\int \frac{\mathrm{d}q_{\mathrm{rev}}}{\mathcal{T}}\right)$ is a state function!

$$\begin{bmatrix} \Delta U_{\mathsf{D}} = \mathcal{C}_{\mathsf{V}} \left(\mathcal{T}_{\mathsf{3}} - \mathcal{T}_{\mathsf{1}} \right) \\ \Delta \mathcal{H}_{\mathsf{D}} = \mathcal{C}_{\mathsf{p}} \left(\mathcal{T}_{\mathsf{3}} - \mathcal{T}_{\mathsf{1}} \right) \end{bmatrix} \qquad \qquad \boxed{q_{\mathsf{D}} = \mathcal{C}_{\mathsf{p}} \left(\mathcal{T}_{\mathsf{3}} - \mathcal{T}_{\mathsf{1}} \right)}$$

$$\int \frac{\mathrm{d}q_{\mathsf{E}}}{\mathcal{T}} = \mathcal{C}_{\mathsf{V}} \ln \left(\frac{\mathcal{T}_{\mathsf{1}}}{\mathcal{T}_{\mathsf{3}}} \right)$$

$$[A] \qquad \text{vs.} \qquad [D] + [E]$$

$$\Delta U_{A} = 0 \qquad \Delta U_{D} + \Delta U_{E} = \Delta U_{A}$$

$$\Delta H_{A} = 0 \qquad \Delta H_{D} + \Delta H_{E} = \Delta H_{A}$$

$$q_{A} = R T_{1} \ln \frac{V_{2}}{V_{1}} \qquad q_{D} + q_{E} = R (T_{3} - T_{1}) \neq q_{A}$$

$$w_{A} = -R T_{1} \ln \frac{V_{2}}{V_{1}} \qquad w_{D} + w_{E} = -R (T_{3} - T_{1}) \neq w_{A}$$

$$\int \frac{dq_{A}}{T} = R \ln \frac{V_{2}}{V_{1}} \qquad \int \frac{dq_{D}}{T} + \int \frac{dq_{E}}{T} = R \ln \left(\frac{V_{2}}{V_{1}}\right) = \int \frac{dq_{A}}{T}$$

Here again $\left(\int \frac{\mathrm{d}q_{\mathrm{rev}}}{T}\right)$ looks like a state function!