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Guidelines for authors

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1. Introduction

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2. Methods

2.1. Particle Transport

2.1.1. Translation and Rotation

Three coordinate frames are generally used to describe the position and orientation of ellipsoidal particles, which can be referred to as the inertial frame $\mathbf{x}^{(in)} = [x^{(in)}, y^{(in)}, z^{(in)}]^T$, the co-moving frame $\mathbf{x}^{(cm)} = [x^{(cm)}, y^{(cm)}, z^{(cm)}]^T$ and the particle frame $\mathbf{x}^{(p)} = [x^{(p)}, y^{(p)}, z^{(p)}]^T$. The co-moving frame translates with the particle with its origin fixed at the particle centroid. The axes of the particle frame always coincide with the semi-axes of the ellipsoid. Thus, the particle frame also record particle rotation.

A point $\mathbf{x}^{(in)}$ in the inertial frame can be transformed to the co-moving frame by

$$\mathbf{x}^{(cm)} = \mathcal{T} \mathbf{x}^{(in)} \quad (2.1)$$

Here, the translation matrix \mathcal{T} is defined as

$$\mathcal{T} = \begin{bmatrix} 1 & 0 & 0 & -x_p \\ 0 & 1 & 0 & -y_p \\ 0 & 0 & 1 & -z_p \\ 0 & 0 & 0 & 1 \end{bmatrix}, \quad (2.2)$$

with (x_p, y_p, z_p) being the coordinates of the particle centroid in the inertial frame. Transformation between the co-moving frame and the particle can be given as

$$\mathbf{x}^{(cm)} = \mathcal{R} \mathbf{x}^{(in)}. \quad (2.3)$$

The rotation matrix \mathcal{R} can be expressed by Euler angles (ϕ, θ, ψ) or quaternions $(\varepsilon_1, \varepsilon_2, \varepsilon_3, \eta)$. In this study, we follow Chesnutt & Marshall (2009) and write \mathcal{R} in the form of quaternions

$$\mathcal{R} = \begin{bmatrix} 1 - 2(\varepsilon_2^2 + \varepsilon_3^2) & 2(\varepsilon_1\varepsilon_2 + \varepsilon_3\eta) & 2(\varepsilon_1\varepsilon_3 - \varepsilon_2\eta) \\ 2(\varepsilon_2\varepsilon_1 - \varepsilon_3\eta) & 1 - 2(\varepsilon_3^2 + \varepsilon_1^2) & 2(\varepsilon_2\varepsilon_3 + \varepsilon_1\eta) \\ 2(\varepsilon_3\varepsilon_1 + \varepsilon_2\eta) & 2(\varepsilon_3\varepsilon_2 - \varepsilon_1\eta) & 1 - 2(\varepsilon_1^2 + \varepsilon_2^2) \end{bmatrix}. \quad (2.4)$$

The initial values of quaternions are determined by

$$\varepsilon_1 = \cos \frac{\phi - \psi}{2} \sin \frac{\theta}{2}, \quad \varepsilon_2 = \sin \frac{\phi - \psi}{2} \sin \frac{\theta}{2}, \quad \varepsilon_3 = \sin \frac{\phi + \psi}{2} \cos \frac{\theta}{2}, \quad \eta = \cos \frac{\phi + \psi}{2} \cos \frac{\theta}{2}. \quad (2.5)$$

65 Then quaternions are evolved by the following equation

$$66 \quad \begin{bmatrix} d\varepsilon_1/dt \\ d\varepsilon_2/dt \\ d\varepsilon_3/dt \\ d\eta/dt \end{bmatrix} = \frac{1}{2} \begin{bmatrix} \eta\Omega_x^{(p)} - \varepsilon_3\Omega_y^{(p)} + \varepsilon_2\Omega_z^{(p)} \\ \varepsilon_3\Omega_x^{(p)} + \eta\Omega_y^{(p)} - \varepsilon_1\Omega_z^{(p)} \\ -\varepsilon_2\Omega_x^{(p)} + \varepsilon_1\Omega_y^{(p)} + \eta\Omega_z^{(p)} \\ -\varepsilon_1\Omega_x^{(p)} - \varepsilon_2\Omega_y^{(p)} - \varepsilon_3\Omega_z^{(p)} \end{bmatrix}, \quad (2.6)$$

67 where $\Omega_x^{(p)}$, $\Omega_y^{(p)}$ and $\Omega_z^{(p)}$ are the components of rotation rate in the particle frame.

68 The discrete element method (DEM) is employed to evolve particle movements. The
69 governing equations of the linear and angular momentum is given as

$$70 \quad m \frac{d\mathbf{v}_i^{(in)}}{dt} = \mathbf{F}_{E,i}^{(in)} + \sum_{j \neq i} \mathbf{F}_{C,j \rightarrow i}^{(in)}, \quad (2.7)$$

$$71 \quad I_x^{(p)} \frac{d\Omega_{x,i}^{(p)}}{dt} - \Omega_{y,i}^{(p)} \Omega_{z,i}^{(p)} (I_y^{(p)} - I_z^{(p)}) = M_{E,i,x}^{(p)} + \sum_{j \neq i} M_{C,j \rightarrow i,x}^{(p)}, \quad (2.8)$$

$$72 \quad I_y^{(p)} \frac{d\Omega_{y,i}^{(p)}}{dt} - \Omega_{z,i}^{(p)} \Omega_{x,i}^{(p)} (I_z^{(p)} - I_x^{(p)}) = M_{E,i,y}^{(p)} + \sum_{j \neq i} M_{C,j \rightarrow i,y}^{(p)}, \quad (2.9)$$

$$73 \quad I_z^{(p)} \frac{d\Omega_{z,i}^{(p)}}{dt} - \Omega_{x,i}^{(p)} \Omega_{y,i}^{(p)} (I_x^{(p)} - I_y^{(p)}) = M_{E,i,z}^{(p)} + \sum_{j \neq i} M_{C,j \rightarrow i,z}^{(p)}. \quad (2.10)$$

74 Here, $\mathbf{v}_i^{(in)}$ and $\boldsymbol{\Omega}_i^{(p)} = [\Omega_{x,i}^{(p)}, \Omega_{y,i}^{(p)}, \Omega_{z,i}^{(p)}]^T$ are the velocity and rotation rate of particle i . m is
75 the particle mass, $\mathbf{I}^p = [I_x^{(p)}, I_y^{(p)}, I_z^{(p)}]^T$ is the moment of inertia with $I_x^{(p)} = m(b^2 + c^2)/5$,
76 $I_y^{(p)} = m(c^2 + a^2)/5$ and $I_z^{(p)} = m(a^2 + b^2)/5$. $\mathbf{F}_{E,i}^{(in)}$ and $\mathbf{M}_{E,i}^{(in)}$ are the electrostatic force and
77 torque exerted on particle i . $\mathbf{F}_{C,j \rightarrow i}^{(in)}$ and $\mathbf{M}_{C,j \rightarrow i}^{(p)} = [M_{C,j \rightarrow i,x}^{(p)}, M_{C,j \rightarrow i,y}^{(p)}, M_{C,j \rightarrow i,z}^{(p)}]^T$ are
78 the contact force and torque acting on particle i by particle j .

79 2.2. Collision between Ellipsoidal Particles

80 2.2.1. Collision Detection

81 In the particle frame of the i th particle, the ellipsoid can be written as

$$82 \quad \mathbf{X}^{(p)T} \mathbf{Q}_i^{(p)} \mathbf{X}^{(p)} = 0. \quad (2.11)$$

83 Here, $\mathbf{X}^{(p)} = [x^{(p)}, y^{(p)}, z^{(p)}, 1]^T$ is the generalized position vector in the particle frame,
84 and the characteristic matrix of ellipsoid i is

$$85 \quad \mathbf{Q}_i^{(p)} = \begin{bmatrix} 1/a^2 & 0 & 0 & 0 \\ 0 & 1/b^2 & 0 & 0 \\ 0 & 0 & 1/c^2 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}. \quad (2.12)$$

86 For points in the inertial frame, the coordinates can be transformed to the particle frame
87 through $\mathbf{x}^{(p)} = \mathcal{RT} \mathbf{x}^{(in)}$. Thus, the ellipsoid can be given in the inertial frame as

$$(\mathcal{RT}\mathbf{X}^{(in)})^T \mathbf{Q}_i^{(p)} (\mathcal{RT}\mathbf{X}^{(in)}) = \mathbf{X}^{(in)T} \mathcal{T}^T \mathcal{R}^T \mathbf{Q}_i^{(p)} \mathcal{RT} \mathbf{X}^{(in)} = \mathbf{X}^{(in)T} \mathbf{Q}_i^{(in)} \mathbf{X}^{(in)} = 0, \quad (2.13)$$

where $\mathbf{Q}_i^{(in)} = \mathcal{T}^T \mathcal{R}^T \mathbf{Q}_i^{(p)} \mathcal{RT}$ is the 4×4 characteristic matrix of ellipsoid i in the inertial frame.

If a point $\mathbf{X}^{(in)}$ satisfies the equation of two different ellipsoids $\mathbf{Q}_1^{(in)}$ and $\mathbf{Q}_2^{(in)}$, then two ellipsoids intersect at $\mathbf{X}^{(in)}$. Multiplying 2.13 of ellipsoid 1 by λ and subtracting 2.13 of ellipsoid 2 thus yields

$$\mathbf{X}^{(in)T} (\lambda \mathbf{Q}_1^{(in)} - \mathbf{Q}_2^{(in)}) \mathbf{X}^{(in)} = 0. \quad (2.14)$$

When two ellipsoids overlap, a family of non-trivial solutions $\mathbf{X}^{(in)}$ exist to describe the intersection. Since $\mathbf{Q}_1^{(in)}$ is invertible, $\mathbf{Q}_1^{(in)-1} \mathbf{Q}_2^{(in)}$ should be singular. Thus, if two eigenvalues of $\mathbf{Q}_1^{(in)-1} \mathbf{Q}_2^{(in)}$ are complex conjugates, two ellipsoid intersect (Alfano & Greer 2003).

2.2.2. Contact Point

When two ellipsoid intersect, the contact point is identified to calculate the contact interactions. In this study, the method of level surfaces are applied for contact point identification (Schneider & Eberly 2002; Ting 1992). (Note: cite paper on geometric potential algorithms by Ning (1992)) 2.13 can be expressed in the quadratic form as

$$P_i(\mathbf{x}^{(in)}) = \mathbf{x}^{(in)T} \mathcal{S}_i^{(in)} \mathbf{x}^{(in)} + \mathbf{b}_i^{(in)T} \mathbf{x}^{(in)} + c_i^{(in)} = 0. \quad (2.15)$$

Here, $\mathcal{S}_i^{(in)}$, $\mathbf{b}_i^{(in)}$ and $c_i^{(in)}$ are defined by the components of $\mathbf{Q}_i^{(in)}$

$$\mathcal{S}_i^{(in)} = \begin{bmatrix} q_{11}^{(in)} & q_{12}^{(in)} & q_{13}^{(in)} \\ q_{12}^{(in)} & q_{22}^{(in)} & q_{23}^{(in)} \\ q_{13}^{(in)} & q_{23}^{(in)} & q_{33}^{(in)} \end{bmatrix} \quad (2.16)$$

$$\mathbf{b}_i^{(in)} = 2[q_{14}^{(in)}, q_{24}^{(in)}, q_{34}^{(in)}]^T, \quad (2.17)$$

$$c_i^{(in)} = q_{44}^{(in)}. \quad (2.18)$$

The contact point on ellipsoid 1 is defined as the tangent point of ellipsoid 1 on the innermost level surface of ellipsoid 2. The level surfaces of ellipsoid 2 is given by

$$P_2(\mathbf{x}^{(in)}) = \alpha, \quad (2.19)$$

where $\alpha < 0$ and $\alpha > 0$ corresponds to the interior and exterior of ellipsoid 2. Then finding the contact point is equivalent to finding the minimum value of α in 2.19 under the constraint of $P_1(\mathbf{x}^{(in)}) = 0$. By defining the Lagrangian function

$$\mathcal{L}(\mathbf{x}^{(in)}) = P_2(\mathbf{x}^{(in)}) + \tau P_1(\mathbf{x}^{(in)}) \quad (2.20)$$

for optimization, the contact point $\mathbf{x}_{C,1}^{(in)}$ is given by

$$\mathbf{x}_{C,1}^{(in)} = -\frac{1}{2}(\mathcal{S}_2^{(in)} + \tau \mathcal{S}_1^{(in)})^{-1}(\mathbf{b}_2^{(in)} + \tau \mathbf{b}_1^{(in)}) = \frac{1}{\Phi(\tau)} \mathbf{y}(\tau), \quad (2.21)$$

118 where $\Phi(\tau)$ is the determinant of $(S_2^{(in)} + \tau S_1^{(in)})$. τ is the Lagrangian multiplier that can be
 119 obtained from the following six-order polynomial (see Chesnutt & Marshall 2009)

$$120 \quad \mathbf{y}(\tau)^T S_1^{(in)} \mathbf{y}(\tau) + \Phi(\tau) \mathbf{b}_1^{(in)T} \mathbf{y}(\tau)^T + \Phi^2(\tau) c_1 = 0. \quad (2.22)$$

121 The above process can be repeated to identify the contact point of ellipsoid 2 on the level
 122 surfaces of ellipsoid 1.

123 2.2.3. Contact Forces and Torques

124 When two particles collide, the velocity at the contact point is

$$125 \quad \mathbf{v}_{C,i}^{(in)} = \mathbf{v}_i^{(in)} + \boldsymbol{\Omega}_i^{(in)} \times \mathbf{r}_{C,i}^{(in)}, \quad (2.23)$$

126 where $\mathbf{r}_{C,i}^{(in)} = \mathbf{x}_{C,i}^{(in)} - \mathbf{x}_{p,i}^{(in)}$ points from the ellipsoid centroid to the contact point. The normal
 127 velocity $\mathbf{v}_{rel,n}^{(in)}$ and tangential velocity $\mathbf{v}_{rel,t}^{(in)}$ are defined by

$$128 \quad \mathbf{v}_{rel,n}^{(in)} = (\mathbf{v}_{C,i}^{(in)} - \mathbf{v}_{C,j}^{(in)}) \cdot \mathbf{n}, \quad (2.24)$$

129 and

$$130 \quad \mathbf{v}_{rel,t}^{(in)} = (\mathbf{v}_{C,i}^{(in)} - \mathbf{v}_{C,j}^{(in)}) - (\mathbf{v}_{C,i}^{(in)} - \mathbf{v}_{C,j}^{(in)}) \cdot \mathbf{n}. \quad (2.25)$$

131 Here, the unit vector along the outward normal direction at the contact point are given by

$$132 \quad \mathbf{n}(\mathbf{x}_{C,i}^{(in)}) = \nabla P_i(\mathbf{x}_{C,i}^{(in)}) / |\nabla P_i(\mathbf{x}_{C,i}^{(in)})|, \quad (2.26)$$

133 while the tangent unit vector equals

$$134 \quad \mathbf{t} = \mathbf{v}_{rel,t}^{(in)} / |\mathbf{v}_{rel,t}^{(in)}|. \quad (2.27)$$

135 In each collision, particles are treated as soft spheres. The contact forces and torques are
 136 calculated according to the Hertz contact model (Marshall 2009).

$$137 \quad \mathbf{F}_{C,j \rightarrow i}^{(in)} = (F_{ne} + F_{nd})\mathbf{n} + F_t\mathbf{t} \quad (2.28)$$

138 Here, the normal force consists of two terms, i.e., the normal elastic force F_{ne} and the normal
 139 damping force F_{nd} . The normal elastic force can be expressed as

$$140 \quad F_{ne} = -k_N \delta_N. \quad (2.29)$$

141 $\delta_N = |\mathbf{x}_{C,i}^{(in)} - \mathbf{x}_{C,j}^{(in)}|$ is the normal overlap, and the stiffness k_N is written as

$$142 \quad k_N = \frac{4}{3} E \sqrt{R \delta_N}, \quad (2.30)$$

143 The effective radius R is defined by the mean curvature of two ellipsoids at their contact
 144 points as

$$145 \quad R = (K_{C,i} + K_{C,j})^{-1}, \quad (2.31)$$

146 with the local mean curvature K_i given by

$$147 \quad K_i = \frac{h^3}{2} \left[\frac{1}{a^2 b^2} \left(\frac{x_i^{(p)2}}{a^2} + \frac{y_i^{(p)2}}{b^2} \right) + \frac{1}{b^2 c^2} \left(\frac{y_i^{(p)2}}{b^2} + \frac{z_i^{(p)2}}{c^2} \right) + \frac{1}{c^2 a^2} \left(\frac{z_i^{(p)2}}{c^2} + \frac{x_i^{(p)2}}{a^2} \right) \right]; \quad (2.32a)$$

$$h = [(x_i^{(p)})^2/a^4 + (y_i^{(p)})^2/b^4 + (z_i^{(p)})^2/c^4]^{-1/2}. \quad (2.32b)$$

The effective elastic modulus E is defined as

$$\frac{1}{E} = \frac{1 - \nu_i^2}{E_i} + \frac{1 - \nu_j^2}{E_j}, \quad (2.33)$$

where E_i and ν_i are the elastic modulus and Poisson ratio of particle i . The normal damping force is proportional to the normal relative velocity

$$F_{nd} = -\eta_N \mathbf{v}_{rel} \cdot \mathbf{n}, \quad (2.34)$$

where the normal damping coefficient is defined as

$$\eta_N = \alpha_N (mk_N)^{1/2}. \quad (2.35)$$

Here, m is the particle mass, and α_N is related to the coefficient of restitution e (Marshall 2009). The tangential force is calculated based on the static friction model and written as

$$F_t = -\mu_F |F_n| \quad (2.36)$$

where $\mu_F = 0.3$ is the friction coefficient. Once the full contact force $\mathbf{F}_{C,j \rightarrow i}^{(in)}$ is obtained, the corresponding rotation torque is computed by

$$\mathbf{M}_{C,j \rightarrow i}^{(in)} = \mathbf{r}_{C,ij}^{(in)} \times \mathbf{F}_{C,j \rightarrow i}^{(in)}. \quad (2.37)$$

2.3. Induced Charge

2.3.1. Governing equation of surface charge

$$\mathcal{A}\sigma_b = b \quad (2.38)$$

$$\mathcal{A}\sigma_b = \bar{\kappa}\sigma_b + \varepsilon_0 \Delta \kappa \mathbf{E}_b \cdot \mathbf{n} \quad (2.39)$$

$$b = (1 - \bar{\kappa})\sigma_f - \varepsilon_0 \Delta \kappa \mathbf{E}_f \cdot \mathbf{n} \quad (2.40)$$

2.3.2. Surface discretization

$$\mathcal{A}_{ij} = \bar{\kappa}_i \delta_{ij} + \Delta \kappa_i \mathbf{n}_i \cdot \mathcal{I}_{ij} \mathbf{a}_j \quad (2.41)$$

$$\mathcal{I}_{ij} = (\mathbf{r}_i - \mathbf{r}_j)/4\pi |\mathbf{r}_i - \mathbf{r}_j|^3 \quad (2.42)$$

To the lowest order, $\mathcal{I}_{ii} = \mathbf{0}$.

Mesh generation should be added here!

2.3.3. Electrostatic Force and Torque

$$\mathbf{F}_E = \int_S (\sigma_f + \sigma_b) \mathbf{E} dS, \quad \mathbf{M}_E = \int_S (\sigma_f + \sigma_b) \mathbf{r} \times \mathbf{E} dS \quad (2.43)$$

3. Authors responsibilities

Authors need to declare in their covering letter to the Editor and during the online submission process whether their manuscript had previously been considered for publication in the *Journal of Fluid Mechanics*. Questions and declarations to that effect must be answered truthfully. Editors, referees, readers and publishers have the right to assume that submitted (and published) manuscripts do not contain scientific dishonesty or fraud comprising, for example, fictitious or manipulated data, plagiarised material (either from previous work of the authors or that of other persons), reference omissions, false priority statements, 'hidden' multiple publication of the same data or incorrect authorship. Authors must not breach any copyright. The *Journal of Fluid Mechanics* uses the iThenticate software to detect instances of plagiarism in manuscripts submitted to it.

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Replication - All information required to replicate the study must be provided, within the body of the paper and/or publicly accessible repositories. Examples of what is required include but are not limited by:

- for analytical studies, the mathematically complete set of equations and boundary conditions, any theorems relied upon, appropriately referenced;
- for numerical studies, the mathematically complete set of equations and boundary conditions, sufficient descriptions of the algorithms or packages used to solve them, appropriately referenced, and the resolution used with respect to the independent variables;
- for laboratory experiments, the dimensions and construction of any apparatus, the materials used including their relevant physical properties, the protocol adopted for the running of the experiments, the measurement tools used including their resolution and accuracy, including appropriate calibration standards;
- for field studies, the raw data collected or used, any protocols or tools used to access the data (e.g. data-mining tools) or to process it.

Verification - Most studies can be verified or falsified provided that sufficient detail is given for them to be replicated (see above). Where data is manipulated (for example, bringing together multiple data sets by scaling) the raw (dimensional) data relating to the primary measurements (laboratory) or outputs (numerical) should be provided together with the protocols or tools used to process them.

Comparison - All graphical information should be supplemented with numerical data or precise algorithms to reproduce it. For example, data points should be provided in a spreadsheet and curves should be defined either explicitly with an equation or as resulting from a precisely defined algorithm.

4. Types of paper

4.1. Standard papers

Regular submissions to JFM are termed 'standard papers'. Note that such papers must contain original research. Papers should be written in a concise manner; though JFM has no page limit, each paper will be judged on its own merits, and those deemed excessive in length will be rejected or will require significant revision.

4.2. *JFM Rapids*

JFM Rapids is devoted to the rapid publication of short, high-impact papers across the full range of fluid mechanics. Manuscripts submitted as *JFM Rapids* must be strictly 10 or fewer printed pages, and must be submitted in L^AT_EX using the `jfm.cls` class file, so as to ensure that they meet the page limit and to expedite their production. As with standard papers, the principal and over-riding objective is to publish papers of the highest scientific quality.

Once a paper is submitted, reviewers are asked to provide reports with a short turnaround. In order to be accepted for publication in *JFM Rapids*, such papers must be strongly endorsed by the referees and should require only minor revisions to improve clarity, usually without recourse to a second round of reviewing. In this case, and at the discretion of the editor, some additional pages may be allowed to address specific points raised by the reviewers, such as the addition of an extra figure or some explanatory text.

Papers that are rejected having been submitted to Rapids are rejected on behalf of the whole Journal and may not be submitted for consideration by another associate editor of JFM, whether for Rapids or as a Standard paper.

In cases where the editor, guided by the reviewers, judges that a paper has merit but requires substantial revision that will require significant reviewing, a decision of ‘revise and resubmit’ will be given. On re-submission, such papers will be handled as standard JFM papers and if accepted will not subsequently appear as *JFM Rapids*.

JFM Rapids will be published online within one month of final acceptance. They will appear within a designated section on the *Journal of Fluid Mechanics* website. Each *Rapid* will be cited and indexed as a JFM article but with a distinctive *Rapids* identifier, and will be assigned to a JFM volume.

4.3. *JFM Perspectives*

Review papers are published under *JFM Perspectives* and are by invitation only.

5. File types

Authors are strongly encouraged to compose their papers in L^AT_EX, using the `jfm.cls` style file and supporting files provided, with the `jfm-instructions.tex` file serving as a template (please note that this is mandatory for *JFM Rapids*). A PDF of the L^AT_EX file should then be generated and submitted via the submission site. For the review process the pdf file should be no more than 10MB. There is no need to submit the L^AT_EX source files alongside the PDF, but upon provisional acceptance of the paper, the L^AT_EX source files, along with individual figure files and a PDF of the final version, will need to be submitted for typesetting purposes. Authors may also compose standard papers in Word, though this will lead to the paper spending a longer period in production. If using Word, please note that equations must NOT be converted to picture format and the file must be saved with the option ‘make equation editable’. All submitted video abstract files should be formatted as MP4 (H.264). MP4 has full compatibility across commonly used browsers, whereas other video formats will only work on selected browsers. This will ensure the greatest possible dissemination of this work.

6. Preparing your manuscript

Authors should write their papers clearly and concisely in English, adhering to JFM’s established style for mathematical notation, as provided in Section 12. We encourage the submission of online supplementary material alongside the manuscript where appropriate (see Section 6.3). Metric units should be used throughout and all abbreviations must be

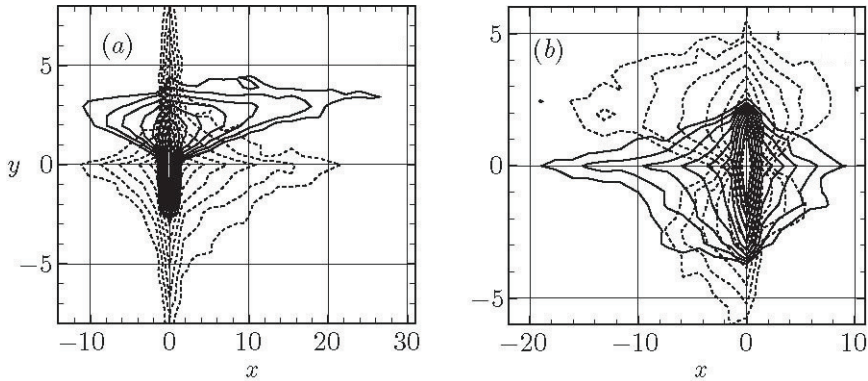


Figure 1: Trapped-mode wavenumbers, kd , plotted against a/d for three ellipses: —, $b/a = 1$; \cdots , $b/a = 1.5$.

defined at first use, even those deemed to be well known to the readership. British spelling must be used, and should follow the *Shorter Oxford English Dictionary*.

6.1. Figures

All authors need to acquire the correct permissions and licences to reproduce figures, which should be submitted with the production files. Further information on applying for permission to reuse figures can be found [here](#). Images should be submitted in EPS or high-resolution TIFF format (1200 dpi for lines, 300 dpi for halftone and colour in RGB format, and 600 dpi for a mixture of lines and halftone) and all labels should be editable. Unless very large, vector graphics are preferred to ensure image sharpness regardless of sizing. The minimum acceptable width of any line is 0.5pt. Each figure should be accompanied by a single caption, to appear beneath, and must be cited in the text. Figures should appear in the order in which they are first mentioned in the text and figure files must be named accordingly ('Abstract.eps', 'Fig1.eps', 'Fig2a.tiff', etc) to assist the production process (and numbering of figures should continue through any appendices). Words *figure 1*, *table 1* and *movie 1* should be lower case. For example see figures 1 and 2. Failure to follow figure guidelines may result in a request for resupply and a subsequent delay in the production process. Note that *all* figures will be relabelled by the typesetter, so please ensure all figure labels are carefully checked against your originals when you receive your proofs.

6.2. Tables

Tables, however small, must be numbered sequentially in the order in which they are mentioned in the text. Words *table 1*, *table 2* should be lower case throughout. See table 1 for an example.

6.3. Online supplementary material

Relevant material which is not suitable for inclusion in the main article, such as movies or numerical simulations/animations, can be uploaded as part of the initial submission. Movies must be submitted in .mp4 format and have the file designation of 'Movie'. Each movie must be numbered in the order they are mentioned and titled movie 1, movie 2 etc and accompanied by a separate caption. To ensure maths terms display correctly they should be bounded by $\$$ and written in TeX, e.g. movie 1. Side view of numerical Schlieren contours from case E1N at $z = L_z/2\$$. Each movie should be no more than 50MB. Upon publication these materials will then be hosted online alongside the final published article. Likewise, should

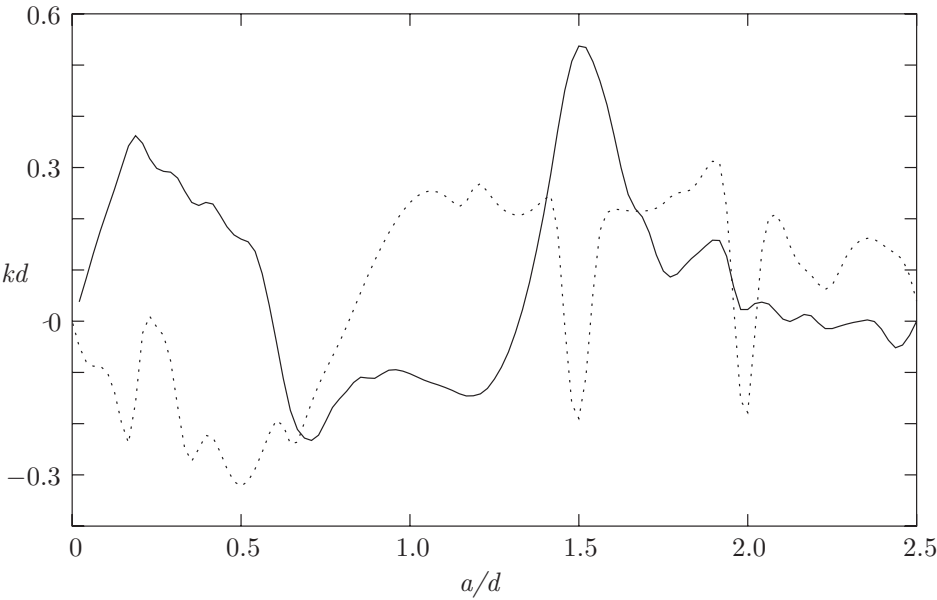


Figure 2: The features of the four possible modes corresponding to (a) periodic and (b) half-periodic solutions.

a/d	$M = 4$	$M = 8$	Callan <i>et al.</i>
0.1	1.56905	1.56	1.56904
0.3	1.50484	1.504	1.50484
0.55	1.39128	1.391	1.39131
0.7	1.32281	10.322	1.32288
0.913	1.34479	100.351	1.35185

Table 1: Values of kd at which trapped modes occur when $\rho(\theta) = a$.

295 there be detailed mathematical relations, tables or figures which are likely to be useful only
296 to a few specialists or take up excessive space in the article, these should also be published
297 online as supplementary material [designated as ‘Other supplementary material’]. Note that
298 supplementary material is published ‘as is’, with no further intervention made during the
299 Production process, all ‘draft’ information should be removed.

300 **7. Editorial decisions**

301 *7.1. Revision*

302 If a revision is requested, you should upload revised files following the same procedure as
303 for submitting a new paper. You begin by clicking on ‘Manuscripts with decision’ in your
304 Corresponding Author Centre, and then on ‘Create a revision’. (Note that if you abandon the
305 process before completing the submission, to continue the submission, you must click on
306 ‘Revised manuscripts in draft’.) There is a new first page showing the decision letter and a
307 space for your reply to the reviewer’s/editor’s comments. You also have the opportunity at
308 this stage to upload your reply to the comments as separate files. All the values filled in on
309 original submission are displayed again. The ID number of the paper will be appended ‘.R1’.

Also note that if a manuscript is submitted as a *JFM Rapid*, but requires substantial revision, it will be re-designated as a standard paper, and the ID and paper type will be amended to reflect this.

7.2. Provisional acceptance

If the paper is accepted as suitable for publication you will be sent a provisional acceptance decision. This enables you to upload the final files required for production: (1) the final PDF or word version of the paper, designated as a 'Main Document'; (2) any source files (see section 5) which must be designated as 'Production Files' and uploaded as a single .zip or .tar file; (3) a completed author publishing agreement form which is available to download at [Cambridge Core](#). For Open Access there is a one-off fee, further information can be found at [JFM open access - FAQs](#). If your research is publicly funded and your organisation comes under one of Cambridge University Press's Read and Publish agreements you may be entitled to free Open Access. Please check your eligibility [here](#).

7.3. Acceptance

On receipt of the production files you will be sent an email indicating completion of the acceptance process.

8. Publication process

Once a paper has been accepted for publication and the source files have been uploaded, the manuscript will be sent to Cambridge University Press for copyediting and typesetting, and will be assigned a digital object identifier (doi). When the proof is ready, authors will receive an email alert containing a link to the PDF of the proof, and instructions for its correction and return. It is imperative that authors check their proofs closely, particularly the equations and figures, which should be checked against the accepted file, as the production schedule does not allow for corrections at a later stage. Once ready, papers will be published online on [Cambridge Core](#) in the current 'open' volume. Each volume will be kept open for approximately two weeks. Note that the PDF published online is the Version of Record and no further alterations/corrections to this document will be allowed. The corresponding author is emailed a link to the published article when it is published online.

9. Corrigenda

The Journal will publish corrigenda that alter significant conclusions made in a paper. Such corrigenda should be submitted to an associate editor, who will consider the submission similarly to a new paper and may consult referees if appropriate. When published, corrigenda are clearly linked with the original articles to which they refer, and the articles to them.

The Journal does not normally publish corrigenda to amend typographical errors, so it is extremely important that authors make very careful checks of their manuscript at every stage, including the reading of proofs, prior to publication.

10. Obtaining help

Technical support for the online submission system is available by clicking on the 'Get Help Now' link at the top-right corner of each page of the submission site. Any other questions relating to the submission or publication process should be directed to the JFM Editorial Assistant, Mrs Amanda Johns, at JFMEditorial@cambridge.org.

11. Cambridge Author Services - in partnership with American Journal Experts

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12. Notation and style

Generally any queries concerning notation and journal style can be answered by viewing recent pages in the Journal. However, the following guide provides the key points to note. It is expected that Journal style and mathematical notation will be followed, and authors should take care to define all variables or entities upon first use. Also note that footnotes are not normally accepted. Abbreviations must be defined at first use, glossaries or lists/tables of abbreviations are not permitted.

12.1. Mathematical notation

12.1.1. Setting variables, functions, vectors, matrices etc

- **Italic font** should be used for denoting variables, with multiple-letter symbols avoided except in the case of dimensionless numbers such as *Re*, *Pr* and *Pe* (Reynolds, Prandtl, and Péclet numbers respectively, which are defined as `\Rey`, `\Pran` and `\Pen` in the template).

- **Upright Roman font** (or upright Greek where appropriate) should be used for:

- (i) (vI) label, e.g. *T*, *t* (transpose)

- (ii) Fixed operators: *sin*, *log*, *d*, Δ , *exp* etc.

- (iii) Constants: *i* ($\sqrt{-1}$), π (defined as `\upi`), *e* etc.

- (iv) Special Functions: *Ai*, *Bi* (Airy functions, defined as `\Ai` and `\Bi`), *Re* (real part, defined as `\Real`), *Im* (imaginary part, defined as `\Imag`), etc.

- (v) Physical units: *cm*, *s*, etc.

- (vi) Abbreviations: *c.c.* (complex conjugate), *h.o.t.* (higher-order terms), *DNS*, etc.

- **Bold italic font** (or bold sloping Greek) should be used for vectors (with the centred dot for a scalar product also in bold): $\mathbf{i} \cdot \mathbf{j}$

- **Bold sloping sans serif font**, defined by the `\mathsfbsi` macro, should be used for tensors and matrices: ***D***

- **Calligraphic font** (for example \mathcal{G} , \mathcal{R}) can be used as an alternative to italic when the same letter denotes a different quantity use `\mathcal` in \LaTeX

12.1.2. *Other symbols*

Large numbers that are not scientific powers should not include commas, but should use a non-breaking space, and use the form 1600 or 16 000 or 160 000. Use O to denote ‘of the order of’, not the L^AT_EX O .

The product symbol (\times) should only be used to denote multiplication where an equation is broken over more than one line, to denote a cross product, or between numbers. The \cdot symbol should not be used, except to denote a scalar product of vectors specifically.

13. Citations and references

All papers included in the References section must be cited in the article, and vice versa. Citations should be included as, for example “It has been shown (Rogallo 1981) that...” (using the `\citep` command, part of the natbib package) “recent work by Dennis (1985)...” (using `\citet`). The natbib package can be used to generate citation variations, as shown below.

```
\citet[pp. 2-4]{Hwang70}:
Hwang & Tuck (1970, pp. 2-4)
\citep[p. 6]{Worster92}:
(Worster 1992, p. 6)
\citep[see][]{Koch83, Lee71, Linton92}:
(see Koch 1983; Lee 1971; Linton & Evans 1992)
\citep[see][p. 18]{Martin80}:
(see Martin 1980, p. 18)
\citep{Brownell04,Brownell07,Ursell50,Wijngaarden68,Miller91}:
(Brownell & Su 2004, 2007; Ursell 1950; van Wijngaarden 1968; Miller 1991)
(Briukhanov et al. 1967)
Bouguet (2008)
(Joseph & Saut 1990)
```

The References section can either be built from individual `\bibitem` commands, or can be built using BibTeX. The BibTeX files used to generate the references in this document can be found in the zip file [jfm-ifcs](#).

Where there are up to ten authors, all authors’ names should be given in the reference list. Where there are more than ten authors, only the first name should appear, followed by *et al.*

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Supplementary data. Supplementary material and movies are available at <https://doi.org/10.1017/jfm.2019...>

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Funding. Please provide details of the sources of financial support for all authors, including grant numbers. For example, “This work was supported by the National Science Foundation (grant number XXXXXXXX)”. Multiple grant numbers should be separated by a comma and space, and where research was funded by more than one agency the different agencies should be separated by a semi-colon, with ‘and’ before the final funder. Grants held by different authors should be identified as belonging to individual authors by the authors’ initials. For example, “This work was supported by the Deutsche Forschungsgemeinschaft (A.B., grant numbers XXXX, YYYY), (C.D., grant number ZZZZ); the Natural Environment Research Council (E.F., grant number FFFF); and the Australian Research Council (A.B., grant number GGGG), (E.F., grant number HHHH)”.

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Data availability statement. The data that support the findings of this study are openly available in [repository name] at [http://doi.org/\[doi\], reference number \[reference number\]](http://doi.org/[doi], reference number [reference number]).

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Author contributions. Authors may include details of the contributions made by each author to the manuscript, for example, "A.G. and T.F. derived the theory and T.F. and T.D. performed the simulations. All authors contributed equally to analysing data and reaching conclusions, and in writing the paper."

14. Appeals process

The *Journal of Fluid Mechanics* has an appeal procedure which provides authors with the opportunity to respond to the editorial decision on their manuscript, should they think that their manuscript was treated in an unfair manner during the peer-review process. Authors have the right to appeal to the Editor or Editor-in-Chief against any decision taken on their manuscript at any stage. An appeal will be considered at the direction of the Editorial Board of the Journal.

14.1. How do I appeal?

Step 1. Requests to have the decision on a submission re-considered should be made in the first instance to the Associate Editor who handled the submission and made the decision. Send a rebuttal letter to the Associate Editor, explaining clearly why you disagree with the decision on your manuscript and including a detailed response to any points of contention in the referees' reports. The Associate Editor will consider your appeal and either invite you to submit a revised paper or confirm the original decision.

Step 2. In case you remain unsatisfied with the Associate Editor's response after Step 1 or at any stage should you consider that your submission was treated unfairly, you should send a letter of appeal to the Editor-in-Chief via the Journal email (JFMEditorial@cambridge.org). Your letter should explain clearly the grounds for your appeal.

Step 3. The Editor-in-Chief will consider the grounds of your appeal and if he considers there to be a *prima facie* case to consider may assign one of the Deputy Editors to consider the appeal in detail. All appeal requests are handled on a case by case basis and the Deputy Editor's or Editor-in-Chief's decision is final. Appeals are normally considered on the basis of whether or not the process of review was conducted appropriately. Papers will not routinely be sent for further review.

Appendix A.

This appendix contains sample equations in the JFM style. Please refer to the \LaTeX source file for examples of how to display such equations in your manuscript.

$$(\nabla^2 + k^2)G_s = (\nabla^2 + k^2)G_a = 0 \quad (\text{A } 1)$$

$$\nabla \cdot \mathbf{v} = 0, \quad \nabla^2 P = \nabla \cdot (\mathbf{v} \times \mathbf{w}). \quad (\text{A } 2)$$

$$G_s, G_a \sim 1/(2\pi) \ln r \quad \text{as} \quad r \equiv |P - Q| \rightarrow 0, \quad (\text{A } 3)$$

$$\left. \begin{aligned} \frac{\partial G_s}{\partial y} &= 0 \quad \text{on} \quad y = 0, \\ G_a &= 0 \quad \text{on} \quad y = 0, \end{aligned} \right\} \quad (\text{A } 4)$$

$$-\frac{1}{2\pi} \int_0^\infty \gamma^{-1} [\exp(-k\gamma|y-\eta|) + \exp(-k\gamma(2d-y-\eta))] \cos k(x-\xi)t \, dt, \quad 0 < y, \quad \eta < d, \quad (\text{A } 5)$$

$$\gamma(t) = \begin{cases} -i(1-t^2)^{1/2}, & t \leq 1 \\ (t^2-1)^{1/2}, & t > 1. \end{cases} \quad (\text{A } 6)$$

$$-\frac{1}{2\pi} \int_0^\infty B(t) \frac{\cosh k\gamma(d-y)}{\gamma \sinh k\gamma d} \cos k(x-\xi)t \, dt$$

$$G = -\frac{1}{4}i(H_0(kr) + H_0(kr_1)) - \frac{1}{\pi} \int_0^\infty \frac{e^{-k\gamma d}}{\gamma \sinh k\gamma d} \cosh k\gamma(d-y) \cosh k\gamma(d-\eta) \quad (\text{A } 7)$$

Note that when equations are included in definitions, it may be suitable to render them in line, rather than in the equation environment: $\mathbf{n}_q = (-y'(\theta), x'(\theta))/w(\theta)$. Now $G_a = \frac{1}{4}Y_0(kr) + \widetilde{G}_a$ where $r = \{[x(\theta) - x(\psi)]^2 + [y(\theta) - y(\psi)]^2\}^{1/2}$ and \widetilde{G}_a is regular as $kr \rightarrow 0$. However, any fractions displayed like this, other than $\frac{1}{2}$ or $\frac{1}{4}$, must be written on the line, and not stacked (ie 1/3).

$$\begin{aligned} \frac{\partial}{\partial n_q} \left(\frac{1}{4}Y_0(kr) \right) &\sim \frac{1}{4\pi w^3(\theta)} [x''(\theta)y'(\theta) - y''(\theta)x'(\theta)] \\ &= \frac{1}{4\pi w^3(\theta)} [\rho'(\theta)\rho''(\theta) - \rho^2(\theta) - 2\rho'^2(\theta)] \quad \text{as} \quad kr \rightarrow 0. \end{aligned} \quad (\text{A } 8)$$

$$\frac{1}{2}\phi_i = \frac{\pi}{M} \sum_{j=1}^M \phi_j K_{ij}^a w_j, \quad i = 1, \dots, M, \quad (\text{A } 9)$$

where

$$K_{ij}^a = \begin{cases} \partial G_a(\theta_i, \theta_j)/\partial n_q, & i \neq j \\ \partial \widetilde{G}_a(\theta_i, \theta_i)/\partial n_q + [\rho'_i \rho''_i - \rho_i^2 - 2\rho_i'^2]/4\pi w_i^3, & i = j. \end{cases} \quad (\text{A } 10)$$

$$\rho_l = \lim_{\zeta \rightarrow Z_l^-(x)} \rho(x, \zeta), \quad \rho_u = \lim_{\zeta \rightarrow Z_u^+(x)} \rho(x, \zeta) \quad (\text{A } 11a, b)$$

$$(\rho(x, \zeta), \phi_\zeta \zeta(x, \zeta)) = (\rho_0, N_0) \quad \text{for} \quad Z_l(x) < \zeta < Z_u(x). \quad (\text{A } 12)$$

$$\tau_{ij} = (\overline{u_i u_j} - \bar{u}_i \bar{u}_j) + (\overline{u_i u_j^{SGS}} + u_i^{SGS} \bar{u}_j) + \overline{u_i^{SGS} u_j^{SGS}}, \quad (\text{A } 13a)$$

$$\tau_j^\theta = (\overline{u_j \bar{\theta}} - \bar{u}_j \bar{\theta}) + (\overline{u_j \theta^{SGS}} + u_j^{SGS} \bar{\theta}) + \overline{u_j^{SGS} \theta^{SGS}}. \quad (\text{A } 13b)$$

$$\mathbf{Q}_C = \begin{bmatrix} -\omega^{-2} V'_w & -(\alpha^t \omega)^{-1} & 0 & 0 & 0 \\ \frac{\beta}{\alpha \omega^2} V'_w & 0 & 0 & 0 & i\omega^{-1} \\ i\omega^{-1} & 0 & 0 & 0 & 0 \\ iR_\delta^{-1}(\alpha^t + \omega^{-1} V''_w) & 0 & -(i\alpha^t R_\delta)^{-1} & 0 & 0 \\ \frac{i\beta}{\alpha \omega} R_\delta^{-1} V''_w & 0 & 0 & 0 & 0 \\ (i\alpha^t)^{-1} V'_w & (3R_\delta^{-1} + c^t(i\alpha^t)^{-1}) & 0 & -(\alpha^t)^{-2} R_\delta^{-1} & 0 \end{bmatrix}. \quad (\text{A } 14)$$

$$\boldsymbol{\eta}^t = \hat{\boldsymbol{\eta}}^t \exp[i(\alpha^t x_1^t - \omega t)], \quad (\text{A } 15)$$

$$\text{where } \hat{\boldsymbol{\eta}}^t = \mathbf{b} \exp(i\gamma x_3^t).$$

$$\text{Det}[\rho \omega^2 \delta_{ps} - C_{pqrs}^t k_q^t k_r^t] = 0, \quad (\text{A } 16)$$

$$\langle k_1^t, k_2^t, k_3^t \rangle = \langle \alpha^t, 0, \gamma \rangle \quad (\text{A } 17)$$

$$\mathbf{f}(\theta, \psi) = (g(\psi) \cos \theta, g(\psi) \sin \theta, f(\psi)). \quad (\text{A } 18)$$

$$f(\psi_1) = \frac{3b}{\pi[2(a+b\cos\psi_1)]^{3/2}} \int_0^{2\pi} \frac{(\sin\psi_1 - \sin\psi)(a+b\cos\psi)^{1/2}}{[1-\cos(\psi_1-\psi)](2+\alpha)^{1/2}} dx, \quad (\text{A } 19)$$

$$\begin{aligned} g(\psi_1) = & \frac{3}{\pi[2(a+b\cos\psi_1)]^{3/2}} \int_0^{2\pi} \left(\frac{a+b\cos\psi}{2+\alpha} \right)^{1/2} \left\{ f(\psi) [(\cos\psi_1 - b\beta_1)S + \beta_1 P] \right. \\ & \times \frac{\sin\psi_1 - \sin\psi}{1-\cos(\psi_1-\psi)} + g(\psi) \left[\left(2+\alpha - \frac{(\sin\psi_1 - \sin\psi)^2}{1-\cos(\psi-\psi_1)} - b^2\gamma \right) S \right. \\ & \left. \left. + \left(b^2\cos\psi_1\gamma - \frac{a}{b}\alpha \right) F\left(\frac{1}{2}\pi, \delta\right) - (2+\alpha)\cos\psi_1 E\left(\frac{1}{2}\pi, \delta\right) \right] \right\} d\psi, \end{aligned} \quad (\text{A } 20)$$

$$\alpha = \alpha(\psi, \psi_1) = \frac{b^2[1-\cos(\psi-\psi_1)]}{(a+b\cos\psi)(a+b\cos\psi_1)}, \quad \beta - \beta(\psi, \psi_1) = \frac{1-\cos(\psi-\psi_1)}{a+b\cos\psi}. \quad (\text{A } 21)$$

$$\left. \begin{aligned} H(0) &= \frac{\epsilon \bar{C}_v}{\tilde{v}_T^{1/2}(1-\beta)}, \quad H'(0) = -1 + \epsilon^{2/3} \bar{C}_u + \epsilon \hat{C}'_u; \\ H''(0) &= \frac{\epsilon u_*^2}{\tilde{v}_T^{1/2} u_P^2}, \quad H'(\infty) = 0. \end{aligned} \right\} \quad (\text{A } 22)$$

LEMMA 1. Let $f(z)$ be a trial Batchelor (1971, pp. 231–232) function defined on $[0, 1]$.
Let Λ_1 denote the ground-state eigenvalue for $-d^2g/dz^2 = \Lambda g$, where g must satisfy $\pm dg/dz +$

520 $\alpha g = 0$ at $z = 0, 1$ for some non-negative constant α . Then for any f that is not identically
 521 zero we have

$$522 \quad \frac{\alpha(f^2(0) + f^2(1)) + \int_0^1 \left(\frac{df}{dz}\right)^2 dz}{\int_0^1 f^2 dz} \geq \Lambda_1 \geq \left(\frac{-\alpha + (\alpha^2 + 8\pi^2\alpha)^{1/2}}{4\pi}\right)^2. \quad (\text{A } 23)$$

523 COROLLARY 1. Any non-zero trial function f which satisfies the boundary condition
 524 $f(0) = f(1) = 0$ always satisfies

$$525 \quad \int_0^1 \left(\frac{df}{dz}\right)^2 dz. \quad (\text{A } 24)$$

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