

ENERGY CALCULATIONS

MATTHEW KOWALSKI AND JAMES HOGAN

Theorem 0.1 (Energy of a Soliton). *Consider an N -soliton of the form*

$$u(t, x) = \frac{P(t, x)}{Q(t, x)}$$

where $Q(t)$ is polynomial in x of degree N and $P(t)$ is polynomial in x of degree at most $N - 1$. Then

$$E(u) = \frac{1}{2} \left\| \frac{P_x}{Q} \right\|_{L_x^2}^2$$

Proof. We recall that by the definition of a N -soliton,

$$|P|^2 = i(Q_x \bar{Q} - \bar{Q}_x Q).$$

Then by definition of the energy,

$$\begin{aligned} E(u) &= \frac{1}{2} \int |u_x - i\Pi_+(|u|^2)u|^2 \\ &= \frac{1}{2} \int \left| \frac{QP_x - PQ_x}{Q^2} + \Pi_+ \left(\frac{Q_x}{Q} - \frac{\bar{Q}_x}{\bar{Q}} \right) \frac{P}{Q} \right|^2 \end{aligned}$$

Since Q (resp. \bar{Q}) has N zeros in the lower (resp. upper) half plane, the definition of Π_+ implies that

$$\begin{aligned} E(u) &= \frac{1}{2} \int \left| \frac{QP_x - PQ_x}{Q^2} + \frac{Q_x P}{Q^2} \right|^2 \\ &= \frac{1}{2} \left\| \frac{P_x}{Q} \right\|_{L_x^2}^2 \end{aligned}$$

as desired. □

Theorem 0.2 (Energy of Traveling Wave). *Consider a traveling solitary wave*

$$u(t, x) = e^{i\omega t} \mathcal{R}_{v,w}(x - vt) = e^{i\omega t} e^{\frac{i}{2}v(x-vt)} e^{i\theta} \frac{\sqrt{2\lambda}}{\lambda x - \lambda vt + y + i}$$

for $\omega, \theta, v, y \in \mathbb{R}$ and $\lambda > 0$. Then

$$E(u) = \frac{\pi|v|^2}{4}.$$

Proof. Define

$$\tilde{u}(t, x) = e^{i\omega t} e^{\frac{-i}{2}v^2 t} e^{i\theta} \frac{\sqrt{2\lambda}}{\lambda x - \lambda vt + y + i}$$

so that $u(t, x) = e^{\frac{i}{2}vx} \tilde{u}(t, x)$. Note that for fixed t , $\tilde{u}(t), u(t)$ have the profiles of 1-solitons. In particular, $\tilde{u}_x - i\Pi_+(|\tilde{u}|^2)\tilde{u} = 0$ and $\|u\|_2^2 = 2\pi$. This implies

$$\begin{aligned} E(u) &= \frac{1}{2} \int \left| \frac{i}{2}vu + e^{\frac{i}{2}vx} \tilde{u}_x - \Pi_+(|\tilde{u}|^2) e^{\frac{i}{2}vx} \tilde{u} \right|^2 \\ &= \frac{|v|^2}{8} \int |u|^2 \\ &= \frac{\pi|v|^2}{4} \end{aligned}$$

as desired. \square

Theorem 0.3 (2-Soliton Energy). *Consider a 2-soliton $u(t, x)$ of the general form*

$$u(t, x) = \frac{e^{i\varphi} \sqrt{2\rho}(\gamma_1 + 2\lambda t + i\lambda^{-1} - x)}{x^2 - (\gamma_0 - i\rho + \gamma_1 + 2\lambda t)x + (\gamma_0 - i\rho)(\gamma_1 + 2\lambda t) - \lambda^{-2}}$$

for $\varphi, \gamma_0, \gamma_1 \in \mathbb{R}$, $\rho > 0$, and $\lambda \in \mathbb{R} \setminus \{0\}$. Then

$$E(u) = \pi\lambda^2.$$

Proof. Let $z_{\pm}(t)$ denote the zeros of the denominator of $u(t, x)$. Recall that asymptotically,

$$z_+(t) \rightarrow \gamma_0 - i\rho, \text{ and } z_-(t) \rightarrow 2\lambda t + i\frac{-\rho}{4\lambda^4 t^2} + O(1).$$

Then by direct calculation and theorem ??,

$$E(u) = \frac{1}{2} \left\| \frac{-e^{i\varphi} \sqrt{2\rho}}{(x - z_+)(x - z_-)} \right\|_{L_x^2}^2 = \rho \left\| \frac{1}{(x - z_+)(x - z_-)} \right\|_{L_x^2}^2$$

A standard contour shift then yields

$$\begin{aligned} E(u) &= \rho \int \frac{1}{(x - z_+)(x - z_-)(x - \bar{z}_+)(x - \bar{z}_-)} \\ &= 2\pi i \rho \left(\frac{1}{(\bar{z}_+ - z_+)(\bar{z}_+ - z_-)(\bar{z}_+ - \bar{z}_-)} + \frac{1}{(\bar{z}_- - z_+)(\bar{z}_- - z_-)(\bar{z}_- - \bar{z}_+)} \right) \end{aligned}$$

Since $E(u)$ is conserved, we may work asymptotically to find

$$\begin{aligned} E(u) &= \lim_{t \rightarrow \infty} 2\pi i \rho \left(\frac{1}{(2i\rho)(2\lambda t + O(1))(2\lambda t + O(1))} + \frac{1}{(2\lambda t + O(1))(\frac{2i\rho}{4\lambda^4 t^2} + O(1))(2\lambda t + O(1))} \right) \\ &= \pi\lambda^2 \end{aligned}$$

as desired. \square

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF CALIFORNIA, LOS ANGELES, CA 90095, USA
Email address: mattkowalski@math.ucla.edu

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF CALIFORNIA, LOS ANGELES, CA 90095, USA
Email address: jameshogan@math.ucla.edu