Unifying Gompertzian growth with the communicable disease spreading paradigm

A PREPRINT

May 18, 2022

ABSTRACT

Recently, a number of studies have shown that cumulative mortality followed a Gompertz curve in the initial Coronavirus epidemic period, March-April 2020. We show that the Gompertz curve is 2 incompatible with the traditional communicable disease spreading hypothesis, and propose a new 3 theory which better explains the nature of the mortality characteristics based on an environmental stressor. Second, we show that for the Gompertz curve to emerge, the stressor has to act on everyone simultaneously, rejecting the possibility of a disease propagation stage. Third, we show that the 6 population acts like a coherent organism under growth/depletion. Finally, we connect the Susceptible-7 Infected-Recovered (SIR) model with our new theory and show that the SIR model is compatible 8 with Gompertzian growth only when all nodes in the transmission network communicate with infinite 9 speed and interaction. 10

Keywords Gompertz · Coherence · Covid · Coronavirus · Network Analysis · Stochastic

2 1 Introduction

Traditional communicable disease spreading theory assumes a pathogen which infects the population through a network of transmission. Following this line of reasoning it can be theoretically shown that in the early stages of an epidemic, growth follows a logistic-like curve, where the very beginning exhibits exponential growth, and for which analytic and semi-analytic solutions have recently emerged [1–4]. A large body of research has emerged that explores how these models fit the recent mortality seen due to the Coronavirus epidemic [5–10].

However, instead of showing logistic-like growth, observed cumulative mortality in the initial period March-April 2020 exhibits almost perfect resemblance to Gompertzian growth [11, 12] where the log-transformed cumulative mortality, or log-mortality for short, is exponentially *decreasing* in time,

$$\frac{d}{dt}\ln Y(t) = -\beta \ln \frac{Y(t)}{\tilde{V}} + \nu,\tag{1}$$

with constants \tilde{Y} , β , and ν , and whose solution is given as

$$Y(t) = Y_{\infty} \left(\frac{Y_0}{Y_{\infty}}\right)^{e^{-\beta t}},\tag{2}$$

where $Y_0=Y(t=0)$ and $Y_\infty=Y(t o\infty)=\tilde{Y}e^{
u/eta}.$

This phenomenon is recorded by others [13–17], showing Gompertz curves at national levels instead of the traditionally predicted logistic curves. What is causing such a discrepancy between reality and the current theory of communicable diseases [18]?

One could approach this conundrum from parametrization of the Susceptible-Infected-Recovered (SIR) model under a time-dependent infection/recovery rate ratio, $\phi(t)$ [19]. How would $\phi(t)$ have to behave? Start by recalling the SIR

model for a pool of susceptible people of size N evolving between the three states: susceptible, S(t), recovered, R(t)and infected, I(t),

$$\frac{dS}{dt} = -\beta \frac{IS}{N} \tag{3}$$

$$\frac{dS}{dt} = -\beta \frac{IS}{N}$$

$$\frac{dI}{dt} = \beta \frac{IS}{N} - \alpha I.$$
(3)

$$\frac{dR}{dt} = \alpha I,\tag{5}$$

omitting the argument t in each variable for brevity and where α and β in this context signify recovery and infection 26 27

A line of reasoning employed by Rypdal and Rypdal [14] to obtain the Gompertz curve is to linearize the SIR model by assuming both the number of infected, I(t), and cumulative infected, Y(t), is much less than the initial pool of susceptible people, $N \gg Y > I$, which seems well founded at the national level, viz.

$$\frac{dY}{dt} = \beta I \tag{6}$$

$$\frac{dI}{dt} = (\beta - \alpha)I = \alpha(\phi(t) - 1)I,\tag{7}$$

and where the number of recovered, R(t), is under this linearization decoupled from the other variables.

They further assume the number of diseased is proportional to the number of cumulative infected, offset by a time lag, which allows us to use the same set of linearized equations to model the number of cumulative people diseased without 30 loss of generality.

Due to this linearization, the infection/recovery ratio, $\phi(t)$, will have to change as a function of time to accommodate for the boundary conditions. And since I is a function of Y, we can combine (6) via an instantaneous relative growth rate, $\gamma(t) = dY(t)/(Ydt) = \beta I/Y$, in turn parameterized by a scaling factor, θ , representing the shape of the growth,

$$\frac{dY}{dt} = \gamma(t)Y(t) \tag{8a}$$

$$\gamma(t) = \frac{\gamma_{\infty}}{\theta} \left[1 - \left(\frac{Y}{Y_{\infty}} \right)^{\theta} \right], \tag{8b}$$

where $\gamma_{\infty} = \gamma(t \to \infty)$. This parametrization is the commonly used Richard's growth curve [20], also called θ -logistic growth, and has been used by others [21]. Although not immediately justified in the communicable disease theory, one 33 could imagine that θ represents non-linear network behavior [22]. Note that at $\theta = 1$, the traditional logistic growth 34 curve is obtained, while at $\theta \to \infty$ we recover the exponential (Malthusian) explosion. 35

The observed Gompertzian mortality curves are realized in the limit $\theta \to 0$, with the relative growth rate,

$$\gamma(t) = \lim_{\theta \to 0} \frac{\gamma_{\infty}}{\theta} \left[1 - \left(\frac{Y}{Y_{\infty}} \right)^{\theta} \right] = \gamma_{\infty} \ln \frac{Y_{\infty}}{Y}$$
 (9)

At this limit the relative growth rate approaches infinity as $Y \to 0$, which seems odd under the hypothesis that the 37 pathogen has just started spreading. The Gompertzian limit also implies a decreasing relative growth rate from the very 38 first time point, which under the SIR model seems unlikely given the large pool of susceptible people in the beginning. 39 One would rather expect a near-constant relative growth rate in the beginning due to a disease propagation stage. Rypdal and Rypdal [14] suggest that the decreasing relative growth rate is caused by social and political mitigating efforts, but 41 these hardly justify such coherent and consistent mortality characteristics across countries. 42

Perhaps a more likely scenario from which a Gompertz curve would emerge is the selective infection of central nodes 43 in the transmission network resulting in an immediate decrease in relative growth. On the other hand, an infection of peripheral nodes should cause immediate exponential growth. Herrmann and Schwartz [23] studied a networked SIR 45 model on a variety of networks, but did not elaborate on a possible fit to a Gompertz curve. Although it may be possible 46 to realize Gompertzian growth from a special network, firm theoretical work has yet to be done to show this connection. 47

We will touch upon how the Gompertz curve emerges from one such network below, under some caveats.

49 1.1 The SIR model family is almost Gompertzian

Without the linearization and the somewhat arbitrary θ -parametrization, one can still obtain Gompertzian growth from the communicable disease models, i.e. a straight line under a double-log transform of (2), viz.

$$\ln\left(\ln\left(Y_{\infty}/Y\right)\right) = -\beta t + k,\tag{10}$$

with $k=\ln\ln\frac{Y_{\infty}}{Y_0}$. To show this, we follow Carletti et al. [5] and consider the extended version of the SIR the model by including a group of diseased, D(t), such that N=S(t)+I(t)+R(t)+D(t), also called the SIRD model. This requires an addition to our original set of equations (3), with a extra equation for the diseased group's growth at some rate δ relative to the current infected group,

$$\frac{dD}{dt} = \delta I. \tag{11}$$

56 Some algebraic manipulations reveal that the diseased group is described by a single equation, viz.

$$\frac{dD}{dt} = \eta (1 - e^{-\xi D}) - \kappa D,\tag{12}$$

where $\eta = \delta N$, $\xi = \beta/(\delta N)$, $\kappa = \delta + \alpha$, and where N is again synonymous to the initial susceptible pool of people. This pool of people cannot be obtained from deaths alone, but can be inferred by assuming a known ratio between mortality and recovery, δ and α . When only modeling the diseased however, knowledge of N is irrelevant and the three parameters in (12) are sufficient.

Note here that small values of D follow logistic growth, viz.

62

69

of log-mortality,

$$\frac{dD}{dt} = cD(1 - D/K),\tag{13}$$

and SEIR model exhibits similar properties and their discrepancy with initial observations related to the Coronavirus epidemic has been noted by others [24].

With the full 3-parameter single ODE of the diseased group in the SIRD model, one can obtain a curve quite close to a straight line in the double-log domain even in the initial observations, although there will always be a non-zero concavity (Fig. 1 and Supplemental Information). Meanwhile, even though the Gompertz model can be fit with a 3-parameter model as shown in (1), it can also be simplified to a 2-parameter model estimated through linear regression

obtained with a second order expansion of the exponential term in (12) and with $c = \eta \xi - \kappa$ and $K = 2b/\eta \xi^2$. The SIR

$$\frac{d}{dt}\ln Y(t) = -\beta \left[\ln Y(t) - \ln Y_{\infty}\right]. \tag{14}$$

Thus, through linear regression estimates, the Gompertz model mitigates the possibility of non-identifiability issues of the parameters [25].

Note also that the SIRD model above considers average macroscopic behavior of an ensemble of microscopic units justified through mean-field theory [26] which does not consider network effects explicitly. Rather, all entities are connected and communicating instantaneously as shown by Mombach et al. [27]. However, even with such strong assumptions, it is odd that observed mortality never exhibits a stage of near-exponential growth as this macroscopic SIRD model predicts, but rather a constant negative slope in the double-log domain. We are thus prompted to look for another model which can explain the observed mortality patterns.

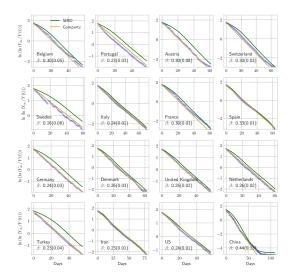


Figure 1: Cumulative number of diseased from the Coronavirus, transformed by $g(Y) = \ln \ln (Y_{\infty}/Y(t))$ plotted against number of days elapsed after $Y(t)/Y_{max} > 0.005$, comparing a SIRD model (green) with a Gompertz model (yellow) for a variety of countries in the period Jan-May 2020. Both models are fit using non-linear least squares according to equations (1) and (12) in the text. Although the Gompertz curve can also be obtained with a simple linear fit using (14), it is here obtained using non-linear least squares to put both models on equal footing. The temporal evolution of the SIRD model is obtained from the Runge-Kutta algorithm, while the Gompertz has a closed form for its temporal evolution. Each plot is annotated with an inferred transmission rate in the SIRD model with 1 standard deviation in brackets, assuming bi-normally distributed parameter estimates. Fitting is done with Python-Scipy and a 6-day moving average of deaths. Observations are taken from the Github repository compiled by the Center for Systems Science and Engineering (CSSE) at Johns Hopkins University, Baltimore, USA [28].

78 2 An alternative theory for observed mortality

An alternative line of reasoning that does not rely on the framework of communicable diseases is that the biosphere was perturbed by an external stressor, initiating a stress response to eventually bring mortality rates back to stability.

This could explain the immediate dampening of mortality growth observed. It could also explain the lack of correlation

between population density and mortality (or infection) rates observed by some [29–34].

Inspired by De Lauro et al. [35], the stressor can be modeled as a multiplicative stochastic dampening term along with a countering force of the immune system. As with all multiplicative processes, it is convenient to log-transform mortality and work in a dimensionless space such that $F(t) = Y(t)/Y_{\infty}$ and $Z(t) = \ln F(t)$, from which a natural perturbation model emerges,

$$dZ(t) = -\beta Z(t)dt + \sqrt{\sigma}dW(t), \tag{15}$$

where dW(t) is a delta-correlated Wiener process with zero mean, making Z(t) a stochastic process. The first term on the right hand side represents the growth due to the stressor, while the second represents the stress response.

This perturbation model is called an Ornstein-Uhlenbeck process [36]. It is also interpreted as Newton's equation of motion with friction and a random force (Langevin's equation), and a continuous version of an Auto-Regressive(1) model [37]. The diffusion coefficient, σ , represents the strength of the perturbation, which we directly see if we recast this equation in terms of mortality while introducing a new parameter $K = \exp\left(-\frac{\sigma}{2R}\right)$, viz.

$$dF(t) = \left\{ \frac{\sigma}{2} F(t) - \beta F(t) \ln \left[\frac{F(t)}{K} \right] \right\} dt + \sqrt{\sigma} F(t) dW(t). \tag{16}$$

To obtain the deterministic observable, we first take the average in the log-domain (15) and transform back to the original domain,

$$d\langle \ln F(t) \rangle = -\beta \langle \ln F(t) \rangle dt, \tag{17}$$

where the bra-ket notation signifies the averaging operation. Then we use the property that the average of log-quantities is the logarithm of the median quantity, where we denote the median of F(t) as M(t),

$$d\ln M(t) = -\beta \ln M(t)dt,\tag{18}$$

which corresponds with the familiar deterministic Gompertz differential equation in (1) with $Y(t) = Y_{\infty}M(t)$. 97 Comparing the stochastic stressor σ in (16) with the deterministic growth equation in (1) gives $\nu = \sigma/2$, suggesting 98 that the stressor is indeed the source of growth, while β is the growth-limiting factor. We also see that by comparing 99 (8a) and (9) with the stochastic counterpart in (16) that the final growth level is governed by the stressor magnitude, 100

$$\sigma = 2\gamma_{\infty} \ln Y_{\infty}. \tag{19}$$

Thus, a more parsimonious interpretation of the observations not reliant on a transmission network is that mortality was 101 caused by a planetary perturbation, modeled as a random process, to which organisms gradually develop resistance at 102 a geometric rate in the log-transformed domain [38, 39], which is the natural transformation for many processes in 103 nature [40]. Under this model, the distribution of the abundance of F(t) is log-normal, a result that can be obtained by directly solving the stochastic equation in (16) [41, 42], or from thermodynamic principles [43–45]. Intuitively, this is seen by noting that the solution to the perturbation model in the log-domain (15) is Gaussian in the variable Z(t), thus 106 suggesting a log-normal distribution of F(t). This implies that the central limit theorem applies to the log-domain. 107

Unifying the SIR model and the Gompertz model

The remarkable observation that the log-transformed domain is the natural one merits closer study. First, juxtapose the logistic model with the Gompertz model, 110

$$\dot{M} = \frac{d}{dt}M(t) = \beta M(t)(1 - M(t)) \quad \text{Logistic}$$

$$\frac{d}{dt}\ln(M(t)) = -\beta\ln(M(t)) \quad \text{Gompertz},$$
(20a)

$$\frac{d}{dt}\ln\left(M(t)\right) = -\beta\ln\left(M(t)\right) \qquad \text{Gompertz}, \tag{20b}$$

where M(t) is deterministic. 111

108

In the logistic model, we recognize the rightmost side of the logistic equation as the transmission term in a SIR model, 112 but also as a linear interaction term between the two macroscopic states. As mentioned earlier, this procedure is a mean field approximation with an implied average interaction between the variables. Thus, all dynamics are governed by macroscopic deterministic variables parametrized by a transmission rate. 115

A microscopic solution could be modeled by splitting the system into N microscopic deterministic units, $M(t) \rightarrow$ 116 $x_1(t), x_2(t), ..., x_N(t)$, where the lower case x_i emphasizes the microscopic quality of the variables, and is here 117 interpreted as probability of infection. From these microscopic units, macroscopic growth could be obtained by taking 118 their arithmetic average, 119

$$M(t) = \frac{1}{N} \sum_{i}^{N} x_i(t).$$
 (21)

Naturally, this simple arithmetic average treats each microscopic unit as an independent variable contributing to the 120 macroscopic observable. 121

One could add network interaction necessitating a corresponding matrix version of (20a) 122

$$\frac{dx_i}{dt} = \beta(1 - x_i) \sum_{j} a_{ij} x_j \quad \forall i,$$
(22)

using the shorthanded $x_i = x_i(t)$ and with a fixed correlation governed by the network's growth or infection rate, β , and adjacency matrix, $\{a_{i,j}\} = \mathbf{A} \in \mathbb{R}^{N \times N}$, a binary matrix with ones where the i^{th} and j^{th} nodes are connected, and zeroes otherwise. Notice that there is an implied causality from the infected to the susceptible, which will become 123 124 125 relevant below. Furthermore, a linear correlation between variables is seen as the partial derivatives of the instantaneous growth rate with respect to pairs of microscopic variables,

$$\frac{\partial^2 \dot{M}}{\partial x_i \partial x_j} = -\frac{\beta}{N} a_{i,j}. \tag{23}$$

Still, no Gompertz curve will emerge at the onset of the growth process. Estrada and Bartesaghi [46] provide illuminating analysis on this topic.

In contrast, as discussed in Section 2, the Gompertz model is implied by a multiplicative stochastic process with a log-normal distribution in its abundance at any given point in time. A log-normal distribution of abundance implies that the log-domain is the natural domain in which the central limit theorem applies, thus implying correlated mortality growth through the geometric mean,

$$M(t) = \exp\left[\frac{1}{N}\sum_{i}\ln x_{i}(t)\right] = \left[\prod_{i}x_{i}(t)\right]^{\frac{1}{N}}$$
(24)

Under this model, correlation between entities is present at all orders in the original domain and all the nodes in the network communicate instantaneously¹.

Thus, the emergence of the Gompertz curve at the macroscopic level suggests that the system is correlated, or coherent, presumably as a result of the simultaneous exposure to the same underlying stressor, but also due to the implied log-normal nature of the microscopic entities, where multiplication replaces addition as the aggregating operator [40]. We can now further appreciate Richard's parametrization as a transition from non-collaborative to collaborative growth as $\theta \to 0$. This feature of θ was also obtained by Petroni et al. [22] by interpreting the θ -logistic growth rate in (8) as non-linear resource availability dependent on the overall magnitude, Y(t), with growth at $\theta \to 0$ named "maximally coherent". Molski and Konarski [48] made a similar observation that Gompertzian growth is the coherent state in a quantum mechanical system with a time-dependent potential, an interpretation which sheds further light on the temporal nature of the postulated stress response. This quantum mechanical system has also been used to describe coherent energy states of diatomic molecules in space [49]. In the field of quantum physics, *coherence* is a well-defined mathematical property first explored by Glauber [50] in the context of electromagnetic fields. The fact that we observe the Gompertz curve in both the microscopic quantum scales and the macroscopic national scales is noteworthy, and suggests that both systems share commonalities and means of communication.

149 3.1 Unification through a modified SIR model

136

137

138

139

140

141

142

144

145

146

However, it is possible to reconcile the SIR model with the Gompertz curve. Inspired by the observation that in the linearized approximation there are only two coupled states, infected and susceptible, we augment the interaction term to higher orders. Then, we reverse the causality where the population of infected are now dependent on the population of the susceptible instead of the other way around as in (22),

$$\frac{dx_i}{dt} = \beta x_i \sum_j a_{ij} (1 - x_j) \quad \forall i.$$
 (25)

This crucial change is based on the environmental stressor hypothesis rather than a communicable disease assumption.
Furthermore, we will for the sake of simplicity assume all nodes in the network have exactly one neighbor and that there exists a unique path between all nodes,

$$\sum_{i} a_{ij} = 1 \quad \forall j. \tag{26}$$

If we let $s_i=1-x_i$ be the probability of being susceptible, the augmented and causality-reversed SIR model infection term becomes

$$\frac{dx_i}{dt} = \beta x_i \sum_j a_{ij} (s_j + s_j^2/2 + \dots) - \alpha x_i \quad \forall i,$$
(27)

Now use the Taylor series $s + s^2/2 + ... = -\ln(1 - s)$, viz.

$$\frac{dx_i}{dt} = -\beta x_i \sum_j a_{ij} \ln(1 - s_j) - \alpha = -\beta x_i \sum_j a_{ij} \ln x_j - \alpha x_i \quad \forall i,$$
(28)

¹It is illuminating to at this point compare with Gompertz' Law of Mortality, $\frac{d}{dt} \ln (1 - M(t)) = -\beta$, for t more than 25 years, which yields a naturally uncorrelated macroscopic curve $1 - M(t) = \exp \left(-\beta t\right)$ [47]. In our context, the uncorrelated feature emerges since the force of mortality is not a function of the growth itself, as we see in the text, but rather of time.

As we are interested in the aggregate macroscopic behavior, we take averages in the log domain, exploit our setup where **A** is a single mapping from one node to another, and simplify to

$$\frac{d}{dt}\ln\left[\prod_{i}x_{i}\right]^{\frac{1}{N}} = -\beta\ln\left[\prod_{i}x_{i}\right]^{\frac{1}{N}} - \alpha. \tag{29}$$

If $\alpha \to 0$, then this equation will exhibit Gompertzian growth in the geometric ensemble average of microscopic units. An interpretation of $\alpha \to 0$ could be that the limiting factor emerges purely from the growth rate without the need for a second growth-limiting parameter, in line with the previous crucial hypothesis of causality reversal stated above. One further simplification could be seen in equating the logarithm of the geometric mean with the logarithm of the median of the set of x_i to obtain (18),

$$\frac{d}{dt}\ln M(t) = -\beta \ln M(t). \tag{30}$$

4 Conclusion

167

168

169

171

172

173

174

175

176

177

In conclusion, we have shown that Gompertzian growth follows from infinite interactions between the susceptible and infected states, and that the perceived pathogen travels at infinite speed throughout the population, rejecting the possibility of a disease propagation stage through a perceived transmission network. In this vein, Richard's parameter, θ , can be related to the number of higher order interactions with the susceptible and the infected in a SIR model [20], where infinite interactions corresponds to $\theta \to 0$.

We further show that the observed mortality across countries can be explained by a model where the biological system is stressed by a ubiquitous and simultaneous stressor eliciting a corresponding stress response through which gradual return to pre-epidemic conditions are mediated. The stressor is modeled as a stochastic perturbation in the log-transformed domain of effects where correlation between people or microscopic entities is present at all orders. From this model, we draw parallels between the coherent behavior of the population's mortality evolution during an epidemic and the spatial coherence of quantum mechanical systems, borrowing the definition of coherence from quantum physics [48].

Thus, we see growth on a spectrum: In one extreme we find non-collaborative growth models or models with parameterized linear interaction effects, and in the other extreme we see a field of microscopic entities coherently sharing information much like quantum entangled particles. The emergence of coherent quantum phenomena at the macroscopic level suggests that no longer can the microscopic world claim a monopoly on quantum physics, especially as it relates to biology [51]. One might be surprised to find that temporal evolution of human mortality during epidemics can behave like the spatial energy distribution of quantum coherent systems.

185 Methods

86 4.1 Estimating model parameters

- 187 The time-evolution of the SIRD and the Gompertz model are obtained with a Runge-Kutta algorithm, after first fitting
- parameters using non-linear least squares according to equations (1) and (12) in the text (Fig. 2). Fitting is done using
- Python-Scipy's non-linear least squares algorithm curve_fit with specified Jacobians and a 6-day windowed average
- of observed deaths. Observations are taken from the Github repository compiled by the Center for Systems Science and
- Engineering (CSSE) at Johns Hopkins University, Baltimore, USA [28].
- The transmission rate is modeled as the product of $\eta\xi$ in (12), and its associated variance is obtained assuming a
- bi-normal relationship between the two parameters with cross-correlation ρ [52],

$$Var[\beta] = Var[\eta]Var[\xi](1 + \rho^2) + Var[\xi] * E[\eta]^2 + Var[\eta]E[\xi]^2,$$
(31)

where estimates are obtained by substituting the model parameters with those estimated through the technique outlined above.

196 4.2 Code Availability

- Scripts used to produce the plots are downloadable in the form of a Jupyter notebook using Python from
- https://nbviewer.org/urls/emf-research.fra1.digitaloceanspaces.com/gompertz/gompertz_vs_sird.ipynb

199 5 Author contributions

200 X conceived the idea and wrote the manuscript while Y advised.

201 6 Competing interests

The authors declare no competing interests.

203 Supplement

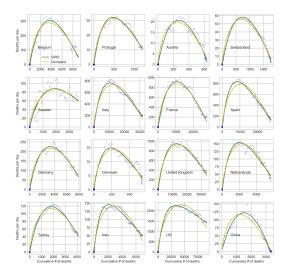


Figure 2: Cumulative diseased plotted against number of days after $Y(t)/Y_{max} > 0.005$, comparing an SIRD with a Gompertz model. Both models are fit using non-linear least squares according to equations (1) and (12) in the text. Fitting is done with Python-Scipy and a 6-day windowed average of deaths. Observations are taken from the Github repository compiled by the Center for Systems Science and Engineering (CSSE) at Johns Hopkins University, Baltimore, USA [28].

Table 1: Estimated values for the SIRD model fitted to various countries' observations as shown in Figure 1 using notation from eq. (12). Standard deviation estimates is given in the paranthesis corresponding to the last significant digit of the reported number.

	$\eta \times 10^{-1}$	$\xi \times 10^4$	$\kappa\times 10^2$
Belgium	94(8)	3.2(2)	9.3(7)
Portugal	6.6(4)	35(2)	4.5(3)
Austria	8(1)	36(5)	11(1)
Switzerland	22(1)	13.6(7)	11.0(5)
Sweden	19(5)	8(2)	3(1)
Italy	137(3)	1.77(5)	3.71(1)
France	240(10)	1.24(7)	8.1(4)
Spain	145(2)	2.27(5)	4.75(8)
Germany	61(4)	3.9(3)	6.5(5)
Denmark	2.7(2)	96(8)	4.3(3)
United Kingdom	181(5)	1.39(5)	4.3(1)
Netherlands	29(1)	8.9(5)	4.3(2)
Turkey	34(3)	6.8(5)	6.9(6)
Iran	22.3(4)	10.4(4)	2.51(7)
US	312(5)	0.76(2)	1.83(7)
China	130(30)	3.5(5)	26(4)

Table 2: Estimated values for the Gompertz model fitted to various countries' observations, as shown in Figure 1 using notation from eq. (1).

	$\nu \times 10$	$\beta\times10^2$	$\tilde{Y}\times 10^2$
Belgium	9.3(1)	8.1(2)	10.1(9)
Portugal	6.15(3)	6.22(7)	7.0(3)
Austria	7.28(7)	8.1(2)	8.1(7)
Switzerland	8.08(6)	8.1(1)	8.9(5)
Sweden	4.9(1)	4.3(2)	6(1)
Italy	7.60(6)	5.88(7)	8.4(4)
France	10.2(1)	8.2(1)	11.0(8)
Spain	9.41(9)	7.5(1)	10.2(7)
Germany	7.56(6)	6.53(8)	8.4(4)
Denmark	5.94(7)	6.5(2)	6.8(7)
United Kingdom	8.46(5)	6.54(7)	9.3(4)
Netherlands	7.34(6)	6.52(8)	8.2(4)
Turkey	7.04(5)	6.38(8)	7.9(4)
Iran	5.87(6)	5.04(8)	6.7(4)
US	7.53(7)	5.32(7)	8.4(5)
China	9.8(1)	9.4(2)	10(1)

References

204

- [1] Harko, T., Lobo, F. S., and Mak, M. Exact analytical solutions of the susceptible-infected-recovered (sir) epidemic
 model and of the sir model with equal death and birth rates. *Applied Mathematics and Computation*, 236:184–194,
 207
- [2] Kröger, M. and Schlickeiser, R. Analytical solution of the sir-model for the temporal evolution of epidemics. part
 a: time-independent reproduction factor. *Journal of Physics A: Mathematical and Theoretical*, 53(50):505601,
 2020.
- [3] Schlickeiser, R. and Kröger, M. Analytical solution of the sir-model for the temporal evolution of epidemics: part b. semi-time case. *Journal of Physics A: Mathematical and Theoretical*, 54(17):175601, 2021.
- [4] Heng, K. and Althaus, C. L. The approximately universal shapes of epidemic curves in the susceptible-exposedinfectious-recovered (SEIR) model. *Scientific Reports*, 10(1):1–6, 2020.
- [5] Carletti, T., Fanelli, D., and Piazza, F. COVID-19: The unreasonable effectiveness of simple models. *Chaos*, *Solitons & Fractals: X*, 5:100034, 2020.
- [6] Cooper, I., Mondal, A., and Antonopoulos, C. G. A SIR model assumption for the spread of COVID-19 in different communities. *Chaos, Solitons & Fractals*, 139:110057, 2020.
- 219 [7] Postnikov, E. B. Estimation of COVID-19 dynamics "on a back-of-envelope": Does the simplest SIR model provide quantitative parameters and predictions? *Chaos, Solitons & Fractals*, 135:109841, 2020.
- [8] Muñoz-Fernández, G. A., Seoane, J. M., and Seoane-Sepúlveda, J. B. A SIR-type model describing the successive waves of COVID-19. *Chaos, Solitons & Fractals*, 144:110682, 2021.
- [9] Cooper, I., Mondal, A., Antonopoulos, C. G., and Mishra, A. Dynamical analysis of the infection status in diverse communities due to COVID-19 using a modified SIR model. *Nonlinear Dynamics*, pages 1–14, 2022.
- [10] Saikia, D., Bora, K., and Bora, M. P. COVID-19 outbreak in india: an SEIR model-based analysis. *Nonlinear Dynamics*, 104(4):4727–4751, 2021.
- [11] Gompertz, B. XXIV. on the nature of the function expressive of the law of human mortality, and on a new mode of determining the value of life contingencies. in a letter to francis baily, esq. f. r. s. &c. *Philosophical Transactions* of the Royal Society of London, 115:513–583, dec 1825. doi:10.1098/rstl.1825.0026.
- 230 [12] Bajzer, Ž., Vuk-Pavlović, S., and Huzak, M. Mathematical modeling of tumor growth kinetics. In *A survey of models for tumor-immune system dynamics*, pages 89–133. Springer, 1997.
- 232 [13] Ohnishi, A., Namekawa, Y., and Fukui, T. Universality in COVID-19 spread in view of the gompertz function.

 233 *Progress of Theoretical and Experimental Physics*, 2020(12), oct 2020. doi:10.1093/ptep/ptaa148.

- 234 [14] Rypdal, K. and Rypdal, M. A parsimonious description and cross-country analysis of COVID-19 epidemic curves. *International Journal of Environmental Research and Public Health*, 17(18):6487, sep 2020. doi:10.3390/ijerph17186487.
- [15] Català, M., Alonso, S., Alvarez-Lacalle, E., López, D., Cardona, P.-J., and Prats, C. Empirical model for short-time prediction of COVID-19 spreading. *PLOS Computational Biology*, 16(12):e1008431, dec 2020. doi:10.1371/journal.pcbi.1008431.
- 240 [16] Rodrigues, T. and Helene, O. Monte carlo approach to model COVID-19 deaths and infections using gompertz functions. *Physical Review Research*, 2(4):043381, 2020.
- ²⁴² [17] Levitt, M., Scaiewicz, A., and Zonta, F. Predicting the trajectory of any COVID19 epidemic from the best straight line. jun 2020. doi:10.1101/2020.06.26.20140814.
- ²⁴⁴ [18] Castro, M., Ares, S., Cuesta, J. A., and Manrubia, S. The turning point and end of an expanding epidemic cannot be precisely forecast. *Proceedings of the National Academy of Sciences*, 117(42):26190–26196, 2020.
- [19] Kermack, W. O. and McKendrick, A. G. A contribution to the mathematical theory of epidemics. *Proceedings of the royal society of london. Series A, Containing papers of a mathematical and physical character*, 115(772): 700–721, 1927.
- ²⁴⁹ [20] Richards, F. A flexible growth function for empirical use. *Journal of experimental Botany*, 10(2):290–301, 1959.
- 250 [21] Wu, K., Darcet, D., Wang, Q., and Sornette, D. Generalized logistic growth modeling of the covid-19 outbreak: comparing the dynamics in the 29 provinces in china and in the rest of the world. *Nonlinear dynamics*, 101(3): 1561–1581, 2020.
- Petroni, N. C., De Martino, S., and De Siena, S. Logistic and θ -logistic models in population dynamics: General analysis and exact results. *Journal of Physics A: Mathematical and Theoretical*, 53(44):445005, 2020.
- [23] Herrmann, H. A. and Schwartz, J.-M. Why COVID-19 models should incorporate the network of social interactions.
 Physical Biology, 17(6):065008, 2020.
- Vattay, G. Forecasting the outcome and estimating the epidemic model parameters from the fatality time series in COVID-19 outbreaks. *Physical Biology*, 17(6):065002, 2020.
- 259 [25] Roda, W. C., Varughese, M. B., Han, D., and Li, M. Y. Why is it difficult to accurately predict the COVID-19 epidemic? *Infectious disease modelling*, 5:271–281, 2020.
- ²⁶¹ [26] Smilkov, D., Hidalgo, C. A., and Kocarev, L. Beyond network structure: How heterogeneous susceptibility modulates the spread of epidemics. *Scientific reports*, 4(1):1–7, 2014.
- ²⁶³ [27] Mombach, J. C., Lemke, N., Bodmann, B. E., and Idiart, M. A. P. A mean-field theory of cellular growth. *EPL* (*Europhysics Letters*), 59(6):923, 2002.
- [28] Dong, E., Du, H., and Gardner, L. An interactive web-based dashboard to track COVID-19 in real time. *The Lancet infectious diseases*, 20(5):533–534, 2020.
- ²⁶⁷ [29] Hamidi, S., Sabouri, S., and Ewing, R. Does density aggravate the COVID-19 pandemic? *Journal of the American Planning Association*, 86(4):495–509, June 2020. doi:10.1080/01944363.2020.1777891.
- ²⁶⁹ [30] Hamidi, S., Ewing, R., and Sabouri, S. Longitudinal analyses of the relationship between development density and the COVID-19 morbidity and mortality rates: Early evidence from 1,165 metropolitan counties in the united states. *Health & Place*, 64:102378, July 2020. doi:10.1016/j.healthplace.2020.102378.
- 272 [31] Carozzi, F. Urban density and covid-19. SSRN Electronic Journal, 2020. doi:10.2139/ssrn.3643204.
- 273 [32] Arpino, B., Bordone, V., and Pasqualini, M. No clear association emerges between intergenerational relationships and COVID-19 fatality rates from macro-level analyses. *Proceedings of the National Academy of Sciences*, 117 (32):19116–19121, jul 2020. doi:10.1073/pnas.2008581117.
- 276 [33] Khavarian-Garmsir, A. R., Sharifi, A., and Moradpour, N. Are high-density districts more vulnerable to the COVID-19 pandemic? *Sustainable Cities and Society*, 70:102911, 2021.
- 278 [34] Barak, N., Sommer, U., and Mualam, N. Urban attributes and the spread of COVID-19: The effects of density, compliance and socio-political factors in israel. *Science of the Total Environment*, 793:148626, 2021.
- [35] De Lauro, E., De Martino, S., De Siena, S., and Giorno, V. Stochastic roots of growth phenomena. *Physica A:* Statistical Mechanics and its Applications, 401:207–213, 2014.
- 282 [36] Risken, H. The Fokker-Planck Equation. Springer, 1996.
- 283 [37] Akaike, H. Statistical predictor identification. *Annals of the institute of Statistical Mathematics*, 22(1):203–217, 1970.

- [38] Boxenbaum, H. Hypotheses on mammalian aging, toxicity, and longevity hormesis: Explication by a generalized
 gompertz function. In *Biological Effects of Low Level Exposures to Chemicals and Radiation*, pages 1–39. CRC
 Press, 2017.
- Neafsey, P. J., Boxenbaum, H., Ciraulo, D. A., and Fournier, D. J. A gompertz age-specific mortality rate model of aging, hormesis, and toxicity: Fixed-dose studies. *Drug metabolism reviews*, 19(3-4):369–401, 1988.
- ²⁹⁰ [40] Zhang, C.-L. and Popp, F.-A. Log-normal distribution of physiological parameters and the coherence of biological systems. *Medical Hypotheses*, 43(1):11–16, 1994.
- 292 [41] Skiadas, C. H. Exact solutions of stochastic differential equations: Gompertz, generalized logistic and revised exponential. *Methodology and Computing in Applied Probability*, 12(2):261–270, 2010.
- ²⁹⁴ [42] Petroni, N. C., De Martino, S., and De Siena, S. Gompertz and logistic stochastic dynamics: Advances in an ongoing quest. *arXiv preprint arXiv:2002.06409*, 2020.
- ²⁹⁶ [43] Sitaram, B. and Varma, V. Statistical mechanics of the gompertz model of interacting species. *Journal of theoretical biology*, 110(2):253–256, 1984.
- ²⁹⁸ [44] Gunasekaran, N. and Pande, L. Log normal distribution for the intrinsic abundance of species from the gompertz model. *Journal of Theoretical Biology*, 98(2):301–305, 1982.
- Chakrabarti, C. and Bhadra, S. Non equilibrium thermodynamics and stochastics of gompertzian growth. *J. Biol. Systems*, 4(2):151–157, 1996. doi:doi.org/10.1142/S0218339096000119.
- Estrada, E. and Bartesaghi, P. From networked SIS model to the gompertz function. *Applied Mathematics and Computation*, 419:126882, 2022.
- 304 [47] Shklovskii, B. A simple derivation of the gompertz law for human mortality. *Theory in Biosciences*, 123(4): 431–433, 2005.
- 306 [48] Molski, M. and Konarski, J. Coherent states of gompertzian growth. *Physical review E*, 68(2):021916, 2003.
- Morse, P. M. Diatomic molecules according to the wave mechanics. ii. vibrational levels. *Physical review*, 34(1): 57, 1929.
- [50] Glauber, R. J. Coherent and incoherent states of the radiation field. *Physical Review*, 131(6):2766, 1963.
- ³¹⁰ [51] Lambert, N., Chen, Y.-N., Cheng, Y.-C., Li, C.-M., Chen, G.-Y., and Nori, F. Quantum biology. *Nature Physics*, 9 (1):10–18, 2013.
- Nadarajah, S. and Pogány, T. K. On the distribution of the product of correlated normal random variables. *Comptes Rendus Mathematique*, 354(2):201–204, 2016.