

# Complexity of switching chaotic maps

M. Antonelli<sup>1</sup>, L. De Micco<sup>1,2</sup>, O. A. Rosso<sup>3,4</sup> and H. A. Larrondo<sup>1,2</sup>

<sup>1</sup> Facultad de Ingeniería, Universidad Nacional de Mar del Plata, Mar del Plata, Argentina.

<sup>2</sup> CONICET.

<sup>3</sup> LaCCAN/CPMAT Instituto de Computação, Universidade Federal de Alagoas, Maceió, Alagoas, Brazil.

<sup>4</sup> Laboratorio de Sistemas Complejos, Facultad de Ingeniería, Universidad de Buenos Aires, Ciudad Autónoma de Buenos Aires, Argentina.

---

## Abstract

In the last years, digital systems such as Digital Signal Processors (DSP), Field Programmable Gate Arrays (FPGA) and Application-Specific Integrated Circuits (ASIC), became the standard in all experimental sciences.

In these systems digital implementations have a custom made numerical representations, therefore finite arithmetic needs to be investigated. Binary floating- and fixed-point are the numerical representations available. Fixed-point representation is preferred over floating-point when speed, low power and/or small circuit area is necessary.

Chaotic systems implemented in finite precision will always become periodic with period  $T$ , each of them produces specific period and statistical characteristics that need to be evaluated. It has been recently shown that it is convenient to describe the statistical characteristic using both, a non causal and a causal probability distribution function (*PDF*). The corresponding entropies, must be evaluated to quantify these *PDF*'s.

In previous works an interesting study about period was carried out using as an example two well known chaotic maps: the tent map and the logistic map. After that, switching techniques between these two maps were tested. In this paper we complement that analysis by characterizing the behaviour of these maps from an statistical point of view using causal and non causal quantifiers.

---

**VERSION: May 8, 2017**

## 1 Introduction

In the last years digital systems became the standard in all experimental sciences. By using virtual instruments, new programmable electronic devices such as Digital Signal Processors (DSP) and Reconfigurable electronics such as Field Programmable Gate Arrays (FPGA) or Application-Specific Integrated Circuits (ASIC), allow experimenters to design and modify their own signal generators, measuring systems, simulation models, etc.

The effect of finite precision in these new devices needs to be investigated. Floating-point is not always available when speed, low power and/or small circuit area are required, a fixed-point solution is better in these cases. Fixed-point representation is critical if chaotic systems must be implemented, because due to roundoff errors digital implementations will always become periodic with period  $T$  and unstable orbits with a low periods may become stable destroying completely the chaotic behavior.

Grebogi and coworkers [1] studied this subject and they saw that the period  $T$  scales with roundoff  $\epsilon$  as  $T \sim \epsilon^{-d/2}$  where  $d$  is the correlation dimension of the chaotic attractor. To have a large period  $T$  is an important property of chaotic maps, in [2] Nagaraj et. al. studied the effect of switching over the average period lengths of chaotic maps in finite precision. They saw that the period  $T$  of the compound map obtained by switching between two chaotic maps is higher than the period of each map. Liu et. al. [3] studied different switching rules applied to linear systems to generate chaos. Switching issue was also addressed in [4], author considers some mathematical, physical and engineering aspects related to singular, mainly switching systems. Switching systems naturally arise in power electronics and many other areas in digital electronics. They have also interest in transport problems in deterministic ratchets [5] and it is known that synchronization of the switching procedure affects the output of the controlled system.

Stochasticity and mixing are also relevant, to characterize these properties several quantifiers were studied [6]. Among them the use of an entropy-complexity representation ( $H-C$  plane) and causal-noncausal entropy ( $H_{BP}-H_{Val}$  plane) deserves special consideration [7,8,9,6,10,11]. A fundamental issue is the criterium to select the probability distribution function (PDF) assigned to the time series. Causal and non causal options are possible. Here we consider the non-causal traditional PDF obtained by normalizing the histogram of the time series. Its statistical quantifier is the normalized entropy  $H_{Val}$  that is a measure of equiprobability among all allowed values. We also consider a causal PDF that is obtained by assigning ordering patterns to segments of trajectory of length  $D$ . This PDF was first proposed by Bandt & Pompe in a seminal paper [12]. The corresponding entropy  $H_{BP}$  was also proposed as a quantifier

by Bandt & Pompe. Amigó and coworkers proposed the number of forbidden patterns as a quantifier of chaos [13]. Essentially they reported the presence of forbidden patterns as an indicator of chaos. Recently it was shown that the name forbidden patterns is not convenient and it was replaced by missing patterns (MP) [14]. **PONER ALGO DE BPW**

In this paper we study the statistical characteristics of five maps, two well known maps: (1) the tent map (TENT) and (2) logistic map (LOG), and three additional maps generated from them: (3) SWITCH, generated by switching between TENT and LOG; (4) EVEN, generated by skipping all the elements in odd position in SWITCH time series and (5) ODD, generated by discarding all the elements in an even position in SWITCH time series. Binary floating- and fixed-point numbers are used, these specific numerical systems may be implemented in modern programmable logic devices.

The main contributions of this paper are:

- (1) the definition of different statistical quantifiers and their relationship with the properties of the series generated by the studied chaotic maps.
- (2) the study of how these quantifiers detect the evolution of stochasticity and mixing of the chaotic maps according as the numerical precision varies.
- (3) the effect on the period and the statistical properties of the time series of switching between two different maps.
- (4) the effect on the period and the statistical properties of the time series of skipping values in the switched maps.

Organization of the paper is as follows: section 2 describes the statistical quantifiers used in the paper and the relationship between their value and characteristics of the causal and non causal PDF considered; section 3 shows and discusses the results obtained for all the numerical representations. Finally section 4 deals with final remarks and future work.

## 2 Information theory quantifiers

Dynamical systems are systems that evolve in time. In practice, one may only be able to measure a scalar time series  $\mathcal{X}(t)$  which may be a function of variables  $\mathcal{V} = \{v_1, v_2, \dots, v_k\}$  describing the underlying dynamics (i.e.  $d\mathcal{V}/dt = f(\mathcal{V})$ ). Given a time series or other observational data, we try to infer properties of an unfamiliar system from the analysis of measured time record of its behavior (time series). how much information are these data revealing about the dynamics of the underlying system or processes? The information content of a system is typically evaluated via a probability distribution function (PDF)  $P$  describing the apportionment of some measurable or observable quantity,

generally a time series  $\mathcal{X}(t)$ . We can define Information Theory quantifiers as measures able to characterize relevant properties of the PDF associated with these time series, and in this way we should judiciously extract information on the dynamical system under study. These quantifiers represent metrics on the space of PDFs for data sets, allowing to compare different sets and classifying them according to the properties of underlying processes - broadly, stochastic vs. deterministic.

In our case, we are interested in chaotic dynamics. Thus we are interested in metrics which take the temporal both order of observations explicitly into account; i.e. the approach is fundamentally *causal* and *statistical* in nature. In a purely statistical approach, correlations between successive values from the time series are ignored or simply destroyed via construction of the PDF; while a causal approach focuses on the PDFs of data sequences.

The quantifiers selected are based on symbolic counting and ordinal pattern statistics. For an application of alternative quantifiers based on Symbolic Dynamics to environmental data, we refer to [?]. The metrics to be used can be broadly classified along two categories: those which quantify the *information content* of data versus those related to their *complexity*. Note that we are referring to the space of probability density functions here, not physical space. For the sake of clarity and simplicity, we only introduce Information Theory quantifiers that are defined on discrete PDFs in this section, since we are only dealing with discrete data (time series). However, all the quantifiers also have definitions for the continuous case [15,?] .

### 2.1 Shannon entropy and statistical complexity

Entropy is a basic quantity that can be regarded to as a measure of the uncertainty associated (information) to the physical process described by  $P$ . When dealing with information content, the Shannon entropy is often considered as the foundational and most natural one [15,?]. Regarded as a measure of uncertainty, is the most paradigmatic example of these information quantifiers.

Let a  $P = \{p_i; i = 1, \dots, N\}$  with  $\sum_{i=1}^N p_i = 1$ , be a discrete probability distribution, with  $N$  the number of possible states of the system under study. The Shannon's logarithmic information measure reads

$$S[P] = - \sum_{i=1}^N p_i \ln [p_i] . \quad (1)$$

If  $S[P] = S_{\min} = 0$ , we are in position to predict with complete certainty which of the possible outcomes  $i$ , whose probabilities are given by  $p_i$ , will

actually take place. Our knowledge of the underlying process described by the probability distribution is maximal in this instance. In contrast, our knowledge is minimal for a uniform distribution  $P_e = \{p_i = 1/N, \forall i = 1, \dots, N\}$  since every outcome exhibits the same probability of occurrence, and the uncertainty is maximal, i.e.,  $S[P_e] = S_{\max} = \ln N$ . These two situations are extreme cases, therefore we focus on the ‘normalized’ Shannon entropy,  $0 \leq \mathcal{H} \leq 1$ , given as

$$\mathcal{H}[P] = S[P]/S_{\max} . \quad (2)$$

Contrary to information content, there is no universally accepted definition of complexity. Here, we focus on describing the *complexity of time series* and do not refer to the complexity of the underlying *systems*. A complex system does not necessarily generate a complex output. In fact, “simple” models might generate complex data, while “complicated” systems might produce output data of low complexity [?].

An intuitive notion of a quantitative complexity attributes low values both to perfectly ordered data (i.e. with vanishing Shannon entropy) as well as to uncorrelated random data (with maximal Shannon entropy. For example, the statistical complexity of a simple oscillation or trend (ordered), but also of uncorrelated white noise (unordered) would be classified as low. Between the two cases of minimal and maximal entropy, data are more difficult to characterize and hence the complexity should be higher. We seek some functional  $\mathcal{C}[P]$  quantifying structures present in the data which deviate from these two cases. These structures relate to organization, correlational structure, memory, regularity, symmetry, patterns, and other properties [?]. One would like to assume that the degree of correlational structures would be adequately captured by some functional  $\mathcal{C}[P]$  in the same way that Shannon’s entropy  $S[P]$  [15] “captures” randomness. Clearly, the ordinal structures present in a process is not quantified by randomness measures, and consequently, measures of statistical or structural complexity are necessary for a better understanding (characterization) of the system dynamics represented by their time series [?]. A suitable measure of complexity can be defined as the product of a measure of information and a measure of disequilibrium, i.e. some kind of distance from the equiprobable distribution of the accessible states of a system. In this respect, Rosso and coworkers [16] introduced an effective *Statistical Complexity Measure* (SCM)  $\mathcal{C}$ , that is able to detect essential details of the dynamical processes underlying the dataset. Based on the seminal notion advanced by López-Ruiz *et al.* [?], this statistical complexity measure [17,16] is defined through the functional product form

$$\mathcal{C}[P] = \mathcal{Q}_J[P, P_e] \cdot \mathcal{H}[P] \quad (3)$$

of the normalized Shannon entropy  $\mathcal{H}$ , see Eq. (2), and the disequilibrium  $\mathcal{Q}_J$

defined in terms of the Jensen-Shannon divergence  $\mathcal{J}[P, P_e]$ . That is,

$$\mathcal{Q}_J[P, P_e] = Q_0 \mathcal{J}[P, P_e] = Q_0 \{S[(P + P_e)/2] - S[P]/2 - S[P_e]/2\}, \quad (4)$$

the above-mentioned Jensen-Shannon divergence and  $Q_0$ , a normalization constant such that  $0 \leq \mathcal{Q}_J \leq 1$ :

$$Q_0 = -2 \left\{ \frac{N+1}{N} \ln(N+1) - \ln(2N) + \ln N \right\}^{-1}, \quad (5)$$

are equal to the inverse of the maximum possible value of  $\mathcal{J}[P, P_e]$ . This value is obtained when one of the components of  $P$ , say  $p_m$ , is equal to one and the remaining  $p_j$  are zero.

The Jensen-Shannon divergence, which quantifies the difference between probability distributions, is especially useful to compare the symbolic composition between different sequences[?, ?, ?]. Note that the above introduced SCM depends on two different probability distributions: one associated with the system under analysis,  $P$ , and the other the uniform distribution,  $P_e$ . Furthermore, it was shown that for a given value of  $\mathcal{H}$ , the range of possible  $\mathcal{C}$  values varies between a minimum  $\mathcal{C}_{min}$  and a maximum  $\mathcal{C}_{max}$ , restricting the possible values of the SCM [18]. Thus, it is clear that important additional information related to the correlational structure between the components of the physical system is provided by evaluating the statistical complexity measure.

## 2.2 Determination of a probability distribution

The evaluation of the Information Theory derived quantifiers suppose some prior knowledge about the system; specifically for those previously introduced (Shannon entropy and statistical complexity), a probability distribution associated to the time series under analysis should be provided beforehand. The determination of the most adequate PDF is a fundamental problem because the PDF  $P$  and the sample space  $\Omega$  are inextricably linked.

Usual methodologies assign to each value of the series  $\mathcal{X}(t)$  (or non-overlapped set of consecutive values) a symbol from a finite alphabet  $\mathcal{A} = \{a_1, \dots, a_M\}$ , thus creating a *symbolic sequence* that can be regarded to as a *non causal coarse grained* description of the time series under consideration. As a consequence, order relations and the time scales of the dynamics are completely lost.

*Causal information* may be duly incorporated if information about the past dynamics of the system is included in the symbolic sequence, i.e., symbols of alphabet  $\mathcal{A}$  are assigned to a portion of the phase-space or trajectory. Bandt and Pompe (BP)[12] introduced a simple and robust symbolic methodology

that takes into account time ordering of the time series by comparing neighboring values in a time series. The causality property of the PDF allows the quantifiers (based on this PDFs) to discriminate between deterministic and stochastic systems [22]. The symbolic data are: *(i)* created by ranking the values of the series; and *(ii)* defined by reordering the embedded data in ascending order, which is tantamount to a phase space reconstruction with embedding dimension (pattern length)  $D$  and time lag  $\tau$ . In this way, it is possible to quantify the diversity of the ordering symbols (patterns) derived from a scalar time series. Note that the appropriate symbol sequence arises naturally from the time series, and no model-based assumptions are needed. The procedure is the following:

- Given a series  $\{x_t, \forall t = 0, \Delta t, \dots, N\Delta t\}$ , a sequence of vectors of length  $D$  is generated.

$$(s) \mapsto (x_{t-(d-1)\Delta t}, x_{t-(d-2)\Delta t}, \dots, x_{t-\Delta t}, x_t) \quad (6)$$

Each vector turns out to be the “history” of the value  $x_t$ . Clearly, the longer the length of the vectors  $D$ , the more information about the history would the vectors have but a higher value of  $N$  is required to have an adequate statistics.

- The permutations  $\pi = (r_0, r_1, \dots, r_{D-1})$  of  $(0, 1, \dots, D-1)$  are called “order of patterns” of time  $t$ , defined by:

$$x_{t-r_{D-1}\Delta t} \leq x_{t-r_{D-2}\Delta t} \leq \dots \leq x_{t-r_1\Delta t} \leq x_{t-r_0\Delta t} \quad (7)$$

In order to obtain a unique result it is considered  $r_i < r_{i-1}$  if  $x_{t-r_i\Delta t} = x_{t-r_{i-1}\Delta t}$ . In this way, all the  $D!$  possible permutations  $\pi$  of order  $D$ , and the PDF  $P = \{p(\pi)\}$  is defined as:

$$p(\pi) = \frac{\sharp\{s | s \leq N - D + 1; (s) \text{ has type } \pi\}}{N - D + 1} \quad (8)$$

In the last expression the  $\sharp$  symbol denotes cardinality.

Thus, an ordinal pattern probability distribution  $P = \{p(\pi_i), i = 1, \dots, D!\}$  is obtained from the time series. In this way the vector defined by Eq. (11) is converted into a unique symbol  $\pi$ . We set  $r_i < r_{i-1}$  if  $x_{s-r_i} = x_{s-r_{i-1}}$  for uniqueness. The only condition for the applicability of the BP method is a very weak stationary assumption: for  $k \leq D$ , the probability for  $x_t < x_{t+k}$  should not depend on  $t$ . Regarding the selection of the parameters, Bandt and Pompe suggested working with  $4 \leq D \leq 6$  for typical time series lengths, and specifically considered a time lag  $\tau = 1$  in their cornerstone paper.

Recently, the permutation entropy was extended to incorporate also amplitude information. Weighting the probabilities of individual patterns according to their variance alleviates potential issues regarding to ‘high noise, low signal’

patterns, because low-variance patterns that are strongly affected by noise are down-weighted in the resulting ‘weighted ordinal pattern distributions’. Hence, a potential disadvantage of ordinal pattern statistics, namely the loss of amplitude information, can be addressed by introducing weights in order to obtain a “weighted permutation entropy (WPE)”. Non-normalized weights are computed for each temporal window for the time series  $X$ , such that

$$w_j = \frac{1}{D} \sum_{k=1}^D (x_{j+k-1} - \bar{X}_j^D)^2. \quad (9)$$

In the equation above  $x_{j+k-1} - \bar{X}_j^D$  denotes the arithmetic mean of the current embedding vector of length  $D$  and its variance  $w_j$  is then used to weight the relative frequencies of each ordinal pattern  $p_j$ . Originally, this technique was proposed to discriminate patterns immersed in low noise. We take advantage of the fact that the fixed points are not computed in the WPE.

### 2.3 Information Planes

AAA

A particularly useful visualization of the quantifiers from Information Theory is their juxtaposition in two-dimensional graphs. Three such *causal information planes* can be defined: *a)* The *causality entropy–complexity plane*,  $\mathcal{H} \times \mathcal{C}$ , is based only on global characteristics of the associated time series PDF (both quantities are defined in terms of Shannon entropies); while *b)* the *causality Shannon-Fisher plane*,  $\mathcal{H} \times \mathcal{F}$ , is based on global and local characteristics of the PDF; finally, *c)* the plane spanned by  $\mathcal{C} \times \mathcal{F}$  is of course fully redundant to the other two. In the case of  $\mathcal{H} \times \mathcal{C}$  the value range is  $[0, 1] \times [\mathcal{C}_{min}(\mathcal{H}), \mathcal{C}_{max}(\mathcal{H})]$ , while in the causality plane  $\mathcal{H} \times \mathcal{F}$  the range is presumably  $[0, 1] \times [0, 1]$ ; no limit curves have been shown to exist so far.

These diagnostic tools were shown to be particularly efficient to distinguish between the deterministic chaotic and stochastic nature of a time series since the permutation quantifiers have distinct behaviours for different types of processes, see Figs. ?? and ??, respectively.

Chaotic maps have intermediate entropy  $\mathcal{H}$  and Fisher  $\mathcal{F}$  values, while their complexity  $\mathcal{C}$  reaches larger values, very close to those of the upper complexity limit [7,?]. For regular processes, entropy and complexity have small values, close to zero, while the Fisher information is close to one. Uncorrelated stochastic processes are located in the planar location associated with  $\mathcal{H}$  near one and  $\mathcal{C}$ ,  $\mathcal{F}$  near zero, respectively. It has also been found that correlated stochastic noise processes with a power spectrum proportional to  $1/f^k$ , where



$1 \leq k \leq 3$ , are characterized by intermediate permutation entropy and intermediate statistical complexity values [7], as well as intermediate to low Fisher information [?]. In both causal information planes ( $\mathcal{H} \times \mathcal{C}$ , see Fig. ?? and  $\mathcal{H} \times \mathcal{F}$ , see Fig. ??), stochastic data are clearly localized at different planar positions than deterministic chaotic ones. These two causal information planes have been profitably used to visualize and characterize different dynamical regimes when the system parameters vary [?, ?, 10, ?, 9, ?, ?, ?, ?, ?]; to study temporal dynamic evolution [?, ?, ?]; to identify periodicities in natural time series [?]; to identify deterministic dynamics contaminated with noise [19, ?]; to estimate intrinsic time scales and delayed systems [?, ?, ?]; for the characterization of pseudo-random number generators [?, ?]; to quantify the complexity of two-dimensional patterns [?], and for other biomedical and econophysics applications (see [?] and references therein).

### BBB

In statistical mechanics one is often interested in isolated systems characterized by an initial, arbitrary, and discrete probability distribution. Evolution towards equilibrium is to be described, as the overriding goal. At equilibrium, we can suppose, without loss of generality, that this state is given by the equiprobable distribution  $P_e = \{p_i = 1/N, \forall i = 1, \dots, N\}$ . The temporal evolution of the above introduced Information Theory quantifiers, Shannon entropy  $\mathcal{H}$ , statistical complexity  $\mathcal{C}$  and Fisher information measure  $\mathcal{F}$ , can be analyzed using a two-dimensional (2D) diagrams of the corresponding quantifiers versus time  $t$ . However, the second law of thermodynamics states that, for isolated systems, entropy grows monotonically with time ( $d\mathcal{H}/dt \geq 0$ ) [?]. This implies that entropy  $\mathcal{H}$  can be regarded as an arrow of time, so that an equivalent way to study the temporal evolution of these quantifiers is using the normalized entropy  $\mathcal{H}$  as substitute for the time-axis.

Two causal information planes are defined (the term causality remembers the fact that temporal correlations between successive samples are taken into account through the Bandt and Pompe PDF recipe used to estimate both Information Theory quantifiers): *a)* The *causality entropy-complexity plane*,  $\mathcal{H} \times \mathcal{C}$ , is based only on global characteristics of the associated time series PDF (both quantities are defined in terms of Shannon entropies); while *b)* the *causality Shannon-Fisher plane*,  $\mathcal{H} \times \mathcal{F}$ , is based on global and local characteristics of the PDF. In the case of  $\mathcal{H} \times \mathcal{C}$  the variation range is  $[0, 1] \times [\mathcal{C}_{min}, \mathcal{C}_{max}]$  (with  $\mathcal{C}_{min}$  and  $\mathcal{C}_{max}$  the minimum and maximum statistical complexity values, respectively, for a given  $\mathcal{H}$  value [18]), while in the causality plane  $\mathcal{H} \times \mathcal{F}$  the range is  $[0, 1] \times [0, 1]$ .

These two diagnostic tools were shown to be particularly efficient to distinguish between the deterministic chaotic and stochastic nature of a time series since the permutation quantifiers have distinctive behaviors for different types

of motion. According to the findings obtained by Rosso *et al.* [7,?,?], chaotic maps have intermediate entropy  $\mathcal{H}$  and Fisher  $\mathcal{F}$  values, while complexity  $\mathcal{C}$  reaches larger values, very close to those of the limit. For regular processes, entropy and complexity have small values, close to zero, while the Fisher is close to one. Finally, totally uncorrelated stochastic processes are located in the planar location associated with  $\mathcal{H}$  near one and,  $\mathcal{C}$ ,  $\mathcal{F}$  near to zero, respectively. It has also been found that  $1/f^\alpha$  correlated stochastic processes with  $1 \leq \alpha \leq 3$  are characterized by intermediate permutation entropy and intermediate statistical complexity values [7], as well as, intermediate low Fisher information [7,?,?]. Moreover, note that in both causal information planes the localization of these stochastic behavior look like a separation border with respect to chaotic, which are localized above of it. In addition, these two causal information planes have been profitably used to visualization and characterization of different dynamical regimes when the system parameters vary [?,?,10,?,9,?,?,?,?]; to study time dynamic evolution [?,?,?]; identifying periodicities in natural time series [?]; identification of deterministic dynamics contaminated with noise [19,?]; estimating intrinsic time scales and delayed systems [?,?,?]; characterization of pseudo-random number generators [?,?] measure of complexity of two-dimensional patterns [?] among other biomedical and econophysics applications (see [?] and references therein).

YOOOOO

Two representation planes are considered:  $H_{BP}$  vs  $H_{hist}$  [8] and  $H_{BP}$  vs  $C_{BP}$  [7]. In the first plane a higher value in any of the entropies,  $H_{BP}$  and  $H_{hist}$ , implies an increase in the uniformity of the involved PDF. The point (1, 1) represents the ideal case with uniform histogram and uniform distribution of ordering patterns. In the second plane not the entire region  $0 < H_{BP} < 1$ ,  $0 < C_{BP} < 1$  is achievable. In fact for any PDF the pairs  $(H, C)$  of possible values fall between two extreme curves in the plane  $H$ - $C$  [23]. Fig. 1 shows two regions labeled as deterministic and stochastic. In fact transition from one region to the other are smooth and the division is a bit arbitrary. A more detailed discussion can be seen in [7]. Ideal random systems having uniform Bandt & Pompe PDF, are represented by the point (1, 0) [24] and a delta-like PDF corresponds with the point (0, 0).

We also used the number of MP as a quantifier[14]. As shown recently by Amigó *et al.* [25,26,27,28], in the case of deterministic one-dimensional maps, not all the possible ordinal patterns can be effectively materialized into orbits, which in a sense makes these patterns forbidden. Indeed, the existence of these forbidden ordinal patterns becomes a persistent fact that can be regarded as a new dynamical property. Thus, for a fixed pattern-length (embedding dimension  $D$ ) the number of forbidden patterns of a time series (unobserved patterns) is independent of the series length  $N$ . Remark that this independence does not characterize other properties of the series such as proximity and

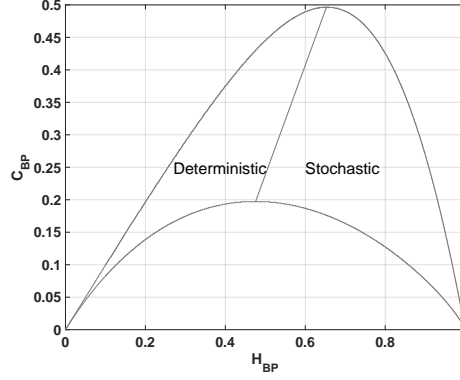


Figure 1. Entropy-Complexity plane.

correlation, which die out with time [26,28].

A full discussion about the convenience of using these quantifiers is out of the scope of this work. Nevertheless reliable bibliographic sources do exist [29,30,31,8,10,18,14].

### 3 Results

Five pseudo chaotic maps were studied, two simple maps and three combination of them. For each one floating-point representation and fixed-point numbers with  $1 \leq B \leq 53$  representation are considered, where  $B$  is the number of bits that represents the fractional part. For each one of these 54 representations 100 time series were generated using randomly chosen initial conditions within the interval  $[0, 1]$ , that is the attraction domain.

The studied maps are logistic (LOG), tent (TENT), sequential switching between TENT and LOG (SWITCH) and skipping discarding the values in the odd positions (EVEN) or the values in the even positions (ODD) respectively. **DEFINIR ACRONIMOS Y EXPLICAR MEJOR.**

Logistic map is interesting because is representative of the very large family of quadratic maps. Its expression is:

$$x_{n+1} = 4x_n(1 - x_n) \quad (10)$$

with  $x_n \in \mathcal{R}$ .

Note that to effectively work in a given representation it is necessary to change the expression of the map in order to make all the operations in the chosen representation numbers. For example, in the case of LOG the expression in

binary fixed-point numbers is:

$$x_{n+1} = 4\epsilon \text{ floor } \left\{ \frac{x_n(1 - x_n)}{\epsilon} \right\} \quad (11)$$

with  $\epsilon = 2^B$  where  $B$  is the number of bits that represents the fractional part.

The Tent map has been extensively studied in the literature because theoretically it has nice statistical properties that can be analytically obtained. For example it is easy to proof that it has a uniform histogram and consequently an ideal  $H_{val} = 1$ . The Perron-Frobenius operator and its corresponding eigenvalues and eigenfunctions may be also be analytically obtained for this map [?]. **PONER ALGUNA REFERNCIA A PERRON Y AGREGAR ALGO EN ITQS**

This map is represented with the equation:

$$x_{n+1} = \begin{cases} 2 x_n & , \text{ if } 0 \leq x_n \leq 1/2 \\ 2 (1 - x_n) & , \text{ if } 1/2 < x_n \leq 1 \end{cases} \quad (12)$$

with  $x_n \in \mathcal{R}$ .

In base-2 fractional numbers rounding, equation 19 becomes:

$$x_{n+1} = \begin{cases} 2 x_n & , \text{ if } 0 \leq x_n \leq 1/2 \\ \epsilon \text{ floor } \left\{ \frac{2 - 2 x_n}{\epsilon} \right\} & , \text{ if } 1/2 < x_n \leq 1 \end{cases} \quad (13)$$

with  $\epsilon = 2^{-B}$ .

Switching, even skipping and odd skipping procedures are shown in Fig. 2.

SWITCH map is expressed as:

$$\begin{cases} x_{n+1} = \begin{cases} 2 x_n, & , \text{ if } 0 \leq x_n \leq 1/2 \\ 2 (1 - x_n) & , \text{ if } 1/2 < x_n \leq 1 \end{cases} \\ x_{n+2} = 4 x_{n+1} (1 - x_{n+1}) \end{cases} \quad (14)$$

with  $x_n \in \mathcal{R}$  and  $n$  an even number.

Skipping is a usual randomizing technique that increases the mixing quality of a single map and correspondingly increases the value of  $H_{BP}$  and decreases  $C_{BP}$  of the time series. Skipping does not change the values of  $H_{val}$  and  $C_{val}$

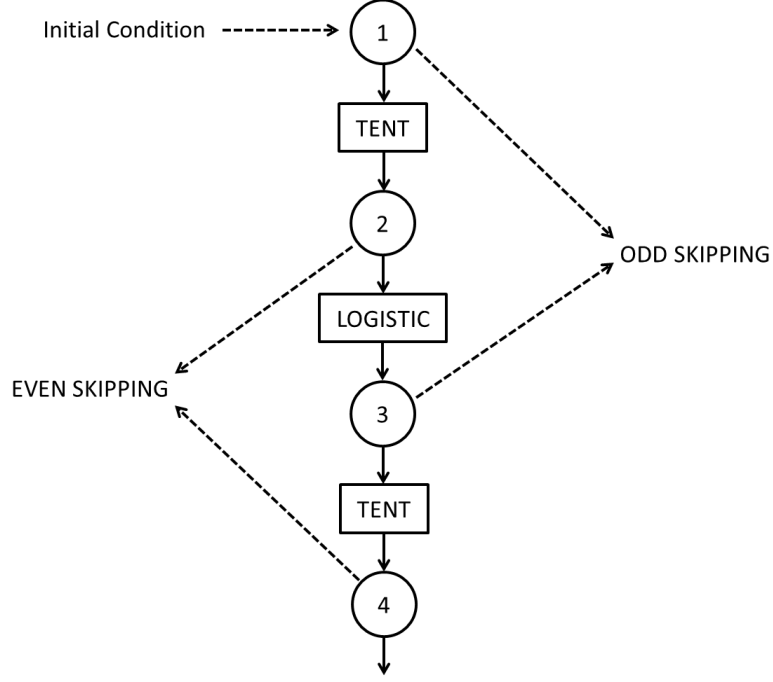


Figure 2. Sequential switching between Tent and Logistic maps. In the figure are also shown even and odd skipping strategies.

because they are evaluated using the non causal PDF (normalized histogram of values)[8].

In the case under consideration we study even and odd skipping of the sequential switching of Tent and Logistic maps:

- (1) Even skipping of the sequential switching of Tent and Logistic maps (EVEN).  
If  $\{x_n, \forall n = 1, \dots, \infty\}$  is the time series generated by 21, discard all the values in odd positions and retain the values in even positions.
- (2) Odd skipping of the sequential switching of Tent and Logistic maps. If  $\{x_n, \forall n = 1, \dots, \infty\}$  is the time series generated by 21, discard all the values in even positions and retain all the values in odd positions.

The reason for studying even and odd skipping cases is the sequential switching map SWITCH is the composition of two different maps. Even skipping may be expressed as the composition function  $\text{TENT} \circ \text{LOG}$  while odd skipping may be expressed as  $\text{LOG} \circ \text{TENT}$ . The evolution of period as function of precision was reported in [2].

Let us detail our results for each of these maps.

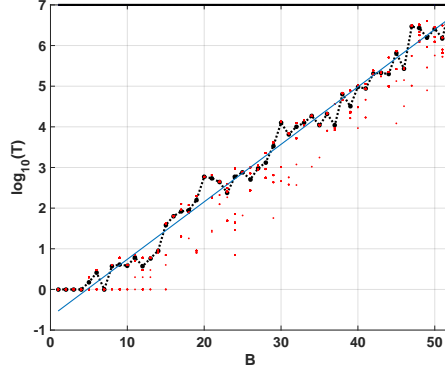


Figure 3. Period as function of precision in binary digits

Table 1

Period  $T$  as a function of  $B$  for the maps considered

map	m	b
TENT	0	0
LOG	0.139	-0.6188
SWITCH	0.1462	-0.5115
EVEN	0.1447	-0.7783
ODD	0.1444	-0.7683

### 3.1 Period $T$ as a function of $B$

Grebogi and coworkers [1] have studied how the period  $T$  is related with the precision. There they saw that the period  $T$  scales with roundoff  $\epsilon$  as  $T \sim \epsilon^{-d/2}$  where  $d$  is the correlation dimension of the chaotic attractor.

Nagaraj et als. [2] studied the case of switching between two maps. They saw that the period  $T$  of the compound map obtained by switching between two chaotic maps is higher than the period of each map and they found that a random switching improves the results. Here we have considered sequential switching to avoid the use of another random variable, because it can include its own statistical properties in the time series.

Fig. 3 shows  $T$  vs  $B$  in semi logarithmic scale. The experimental averaged points can be fitted by a straight line expressed as  $\log_2 T = mB + b$  where  $m$  is the slope and  $b$  is the  $y$ -intercept. Results for all considered maps are summarized in Table 1.

Results are compatible for those obtained in [2]. Switching between maps

increases de period  $T$  but skipping procedure decreases by almost half.

### 3.2 Quantifiers of simple maps

Here we report our results for both simple maps, LOG and TENT

#### 3.2.1 Logistic map (LOG)

Figs. 4 (a) to (f) show the statistical properties of LOG map in floating-point and fixed-point representation. All these figures show: 100 red points for each fixed-point precision ( $B$ ) and in black their average (dashed black line connecting black dots), 100 horizontal dashed blue lines that are the results of each run in floating-point and a black solid line their average. In this case, all the lines of the floating-point are overlapped.

According as  $B$  grows, statistical properties vary until they stabilize. For  $B \geq 30$  the value of  $H_{val}$  remains almost identical to the values for the floating-point representation whereas  $H_{BP}$  and  $C_{BP}$  stabilizes at  $B > 21$ . Their values are:  $\langle H_{val} \rangle = 0.9669$ ;  $\langle H_{BP} \rangle = 0.6269$ ;  $\langle C_{BP} \rangle = 0.4843$ . Note that the stable value of missing patterns  $MP = 645$  makes the optimum  $H_{BP} \leq \ln(75)/\ln(720) \simeq 0.65$ . Then,  $B = 30$  is the most convenient choice because an increase in the number of fractional digits does not improve the statistical properties.

Some conclusions can be drawn regarding BP and BPW quantifiers. For  $B = 1, 2, 3$  and 4, the averaged BP quantifiers are almost 0 while the averaged BPW quantifiers can not be calculated (see in Figs. 4 c and e the missing black dashed line). For those sequences where the initial condition where 0 all iterations result to be a sequence of zeros (the fixed point of the map), this happens when using small precisions because of the roundoffs.

When  $B$  increases,  $B = 7, 9$  and 12, the initial conditions are rounded to zero less frequently. So the generated sequences start from some value but many of them fall to zero with a short transitory. This can be seen in Figs. 4 c and e, BPW quantifiers show a high dispersion unlike BP quantifiers. This is because BPW procedure takes into account only the transient discarding fixed points, unlike BP procedure considers all the values of the sequence.

The same results are shown in double entropy planes with the precision as parameter (Fig. 5). These figures show: 100 red points for each fixed-point precision ( $B$ ) and in black their average (dashed black line connecting black dots), 100 blue dots that are the results of each run in floating-point and the black star is their average. In Fig. 5 the 100 blue points and their average are overlapped.

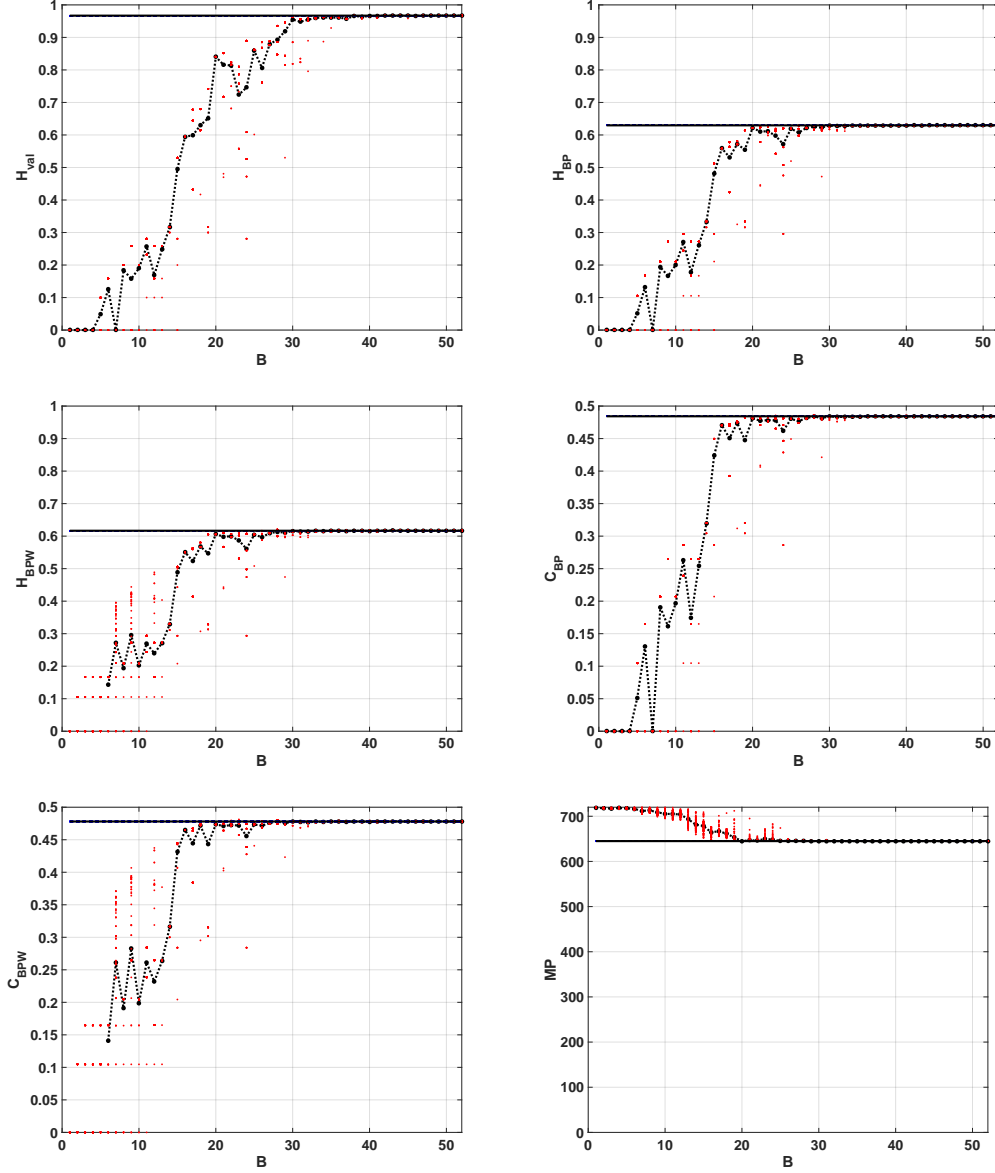


Figure 4. Statistical properties of the LOG map: (a)  $H_{val}$  vs  $B$  (b)  $H_{BP}$  vs  $B$  (c)  $C_{BP}$  vs  $B$  (d)  $MP$  vs  $B$ .

As expected, the fixed-point architecture implementation converges to the floating-point value as  $B$  increases. For both, Hbp-Hval and Hbpw-Hval, from  $B = 20$ ,  $H_{val}$  improves but  $H_{BP}$  remains constant. It can be seen that the distribution of values reaches high values ( $\langle H_{val} \rangle = 0.9669$ ) but their mixing is poor ( $\langle H_{BP} \rangle = 0.6269$ ).

In Fig. 6 we show the entropy-complexity planes. Dotted gray lines are the upper and lower margins, it is expected that a chaotic system remains near the upper margin. This results characterize a chaotic behaviour, in  $H_{BP} - C_{BP}$  plane we can see a low entropy and high complexity.



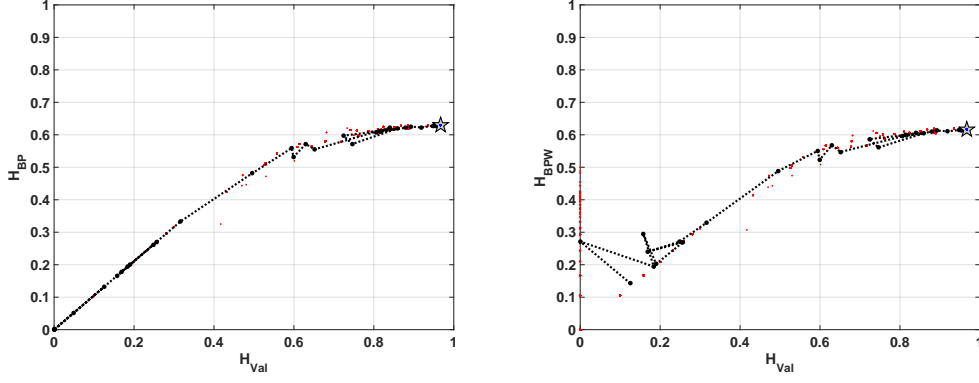


Figure 5. Evolution of statistical properties in double entropy plane of LOG map: (a)  $H_{val}$  vs  $H_{BP}$  (b)  $H_{val}$  vs  $H_{BPW}$ .

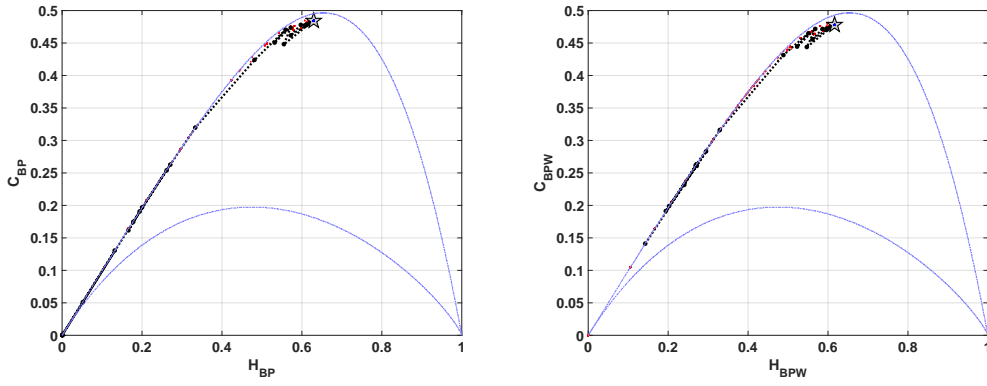


Figure 6. Evolution of statistical properties in entropy-complexity plane of LOG map: (a)  $C_{BP}$  vs  $H_{BP}$  (b)  $C_{BPW}$  vs  $H_{BPW}$ .

### 3.2.2 Tent map (*TENT*)

When this map is implemented in a computer using any numerical representation system (even floating-point!) truncation errors rapidly increase and make unstable fixed point in  $x^* = 0$  to become stable. The sequences within the attractor domain of this fixed point will have a short transitory of length between 0 and  $B$  followed by an infinite number of 0's [32,33]. This issue is easily explained in [34], the problem appears because all iterations have a left-shift operation that carries the 0's from the right side of the number to the most significant positions. Some authors [?] have proposed to add random perturbations to avoid this drawback of the Tent map. This procedure improves the statistical properties of the time series, but what really happens is that the statistical properties of the random perturbations are mixed with those of the Tent map. Here we study the Tent map “as it is” without any artifact to evaluate its real behavior, instead of theoretical statistical properties.

Figs. 7 (a) to (e) show the quantifiers for floating- and fixed-point numerical representations. Quantifiers  $H_{val}$ ,  $H_{BP}$  and  $C_{BP}$  are equal to zero for all

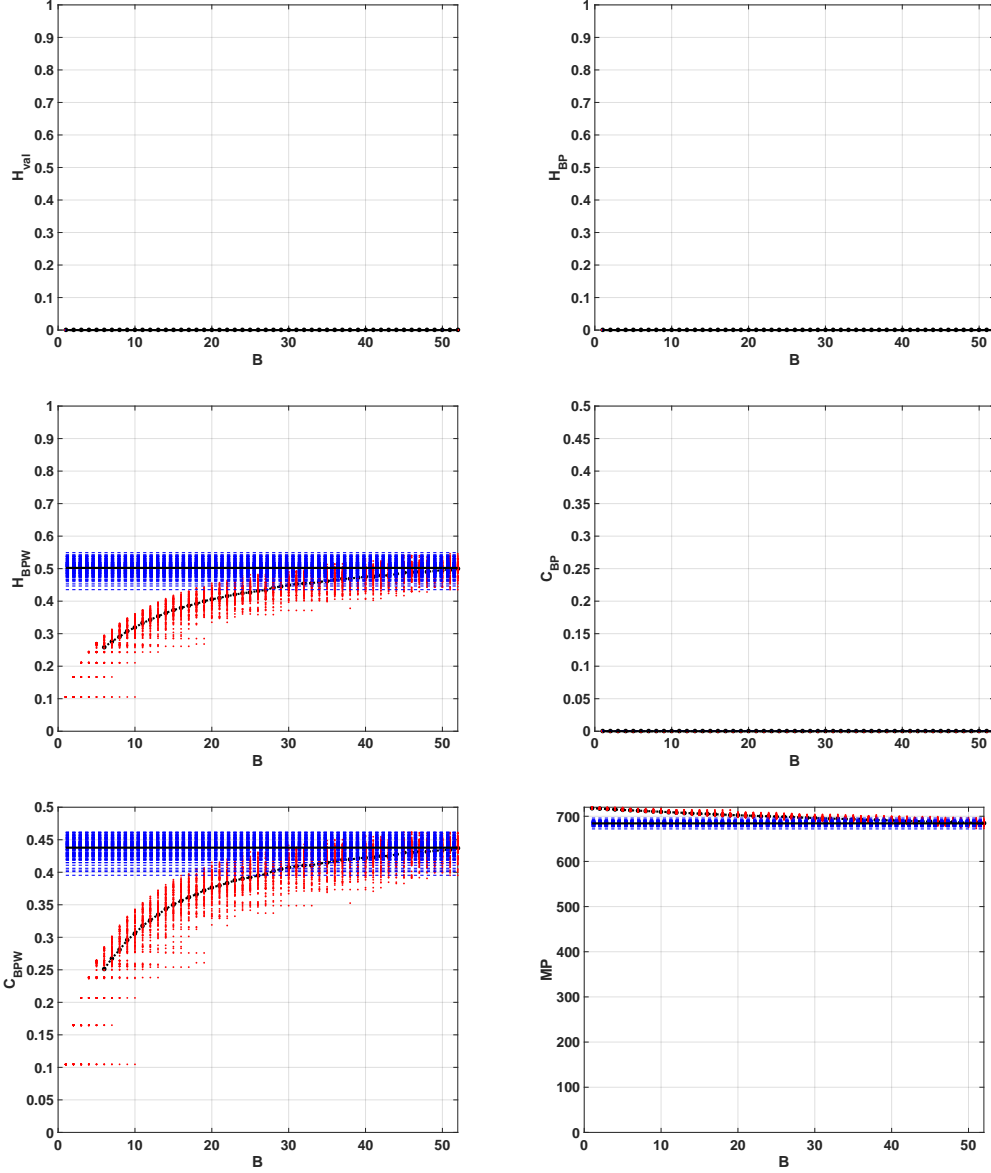


Figure 7. Statistical properties of the TENT map: (a)  $H_{val}$  vs  $B$  (b)  $H_{BP}$  vs  $B$  (c)  $C_{BP}$  vs  $B$  (d)  $MP$  vs  $B$ .

precisions, this reflects that the series quickly converge toward a fixed point for almost all sequences. In the case of  $H_{BPW}$  and  $C_{BPW}$  quantifiers they are different from zero because BPW procedure discards the elements once they reach the fixed point. The high dispersions in  $H_{BPW}$ ,  $C_{BPW}$  and  $MP$  are related to the short length of transient. These transient that converges to a fixed point has a maximum length of  $B$  iterations for fixed-point arithmetic and 54 for floating-point (double precision).

In summary in spite of using a high number of bits (with any 2-based numerical representation) to represent the digitalized TENT map it always loses the

chaotic behaviour.

### 3.3 Quantifiers of combined Maps

Here we report our results for the three combinations of the simple maps, SWITCH, EVEN and ODD.

#### 3.3.1 Sequential switching between Tent and Logistic maps (SWITCH)

Results with sequential switching are shown in Figs. 8 (a) to (f). The entropy value calculated in floating-point is  $\langle H_{val} \rangle = 0.9722$ , this value is slightly higher than the one obtained for the LOG map. For fixed-point arithmetic this value is reached in  $B = 24$ , but it stabilizes from  $B = 28$ . Regarding the ordering patterns the number of MP decreases to 586, this value is lower than the one obtained for LOG map. It means the entropy  $H_{BP}$  may increase up to  $\ln(134)/\ln(720) \simeq 0.74$ . BP and BPW quantifiers reach their maximum of  $\langle H_{BP} \rangle = 0.6546$  and  $\langle H_{BPW} \rangle = 0.6313$  at  $B = 16$ , but they stabilize from  $B = 24$ . Complexities are lower than for LOG,  $\langle C_{BP} \rangle = 0.4580$  and  $\langle C_{BPW} \rangle = 0.4578$ , these values are reached for  $B \geq 15$  but they are stable from  $B \geq 23$ . Compared with LOG, statistical properties are better with less amount of bits, for  $B \geq 24$  this map reaches optimal characteristics in the sense of random source.

Furthermore, we encountered one initial condition in floating-point long double precision with an anomalous behavior. The quantifiers based on BPW procedure cant detect an anomaly, 8 (a), (b) and (d) an horizontal blue dashed line is far from the average value but this can not be seen in the other figures 8 (c) and (e). We detected a falling to a fixed point after a long transitory, the BPW procedure discards the values corresponding with a fixed point and calculates only the transitory.

Double entropy plane  $H_{val}$  vs  $H_{BP}$  is showed in Fig. 9. The point reached in this plane for SWITCH map is similar to that reached for LOG map. The mixing is slight better in this case.

Entropy-complexity plane  $C_{BP}$  vs  $H_{BP}$  is showed in Fig. 10. If we compare with the same plane in the case of LOG (Fig. 6. a.),  $C_{BP}$  is lower for SWITCH, this fact shows a more random behaviour.

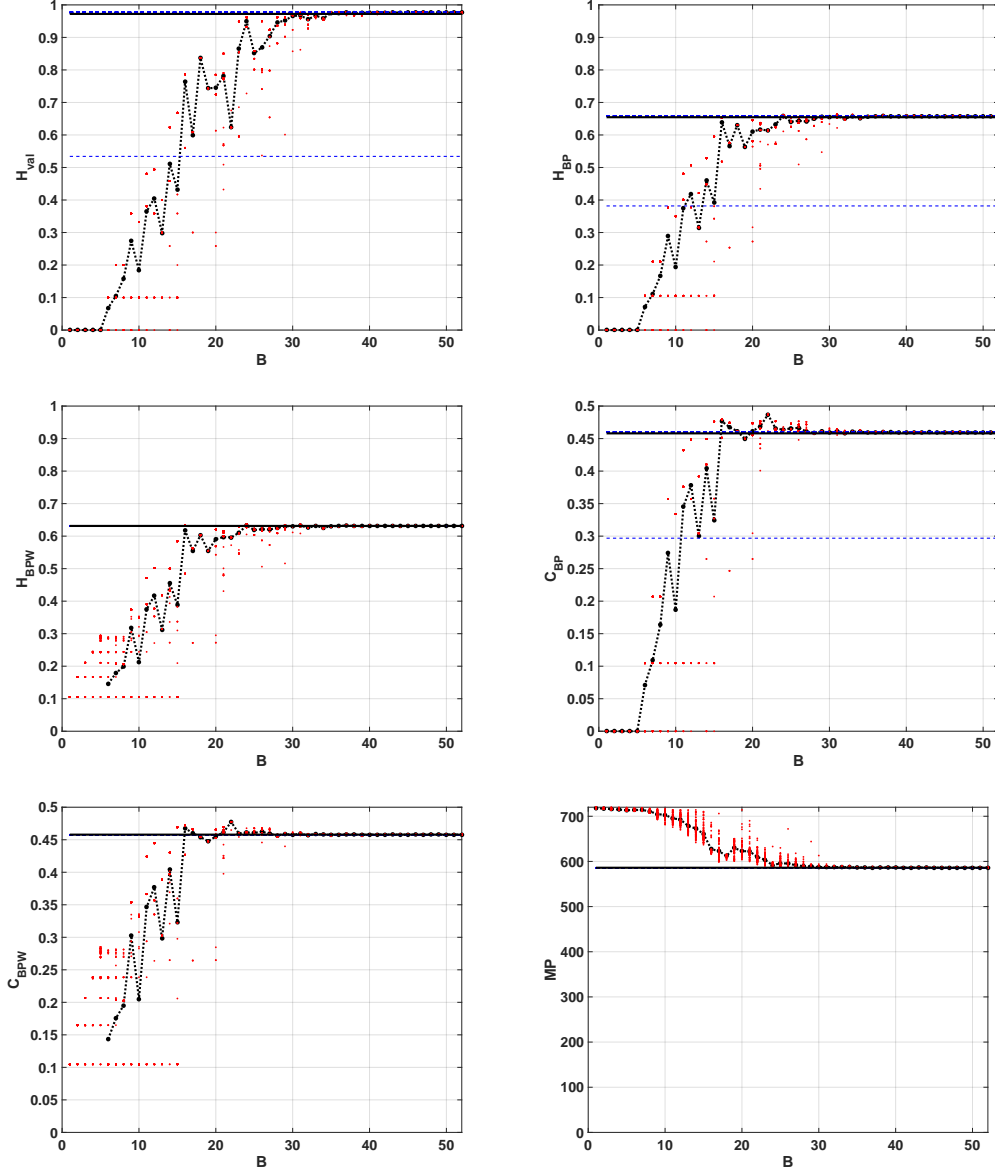


Figure 8. Statistical properties of the SWITCH map: (a)  $H_{val}$  vs  $B$  (b)  $H_{BP}$  vs  $B$  (c)  $C_{BP}$  vs  $B$  (d)  $MP$  vs  $B$ .

### 3.3.2 Skipping on sequential switching between Tent and Logistic maps (EVEN and ODD)

In figures 11.a and 12.a, quantifiers related to the normalized histogram of values slightly degrades with the skipping procedure due to finite data length. For example  $\langle H_{val} \rangle$  reduces from 0.9722 without skipping to 0.9459 for EVEN and 0.9706 for ODD. This difference between EVEN and ODD in floating point is because a high dispersion was obtained for  $H_{val}$ ,  $H_{BP}$  and  $C_{BP}$  but not for  $H_{BPW}$  or  $C_{BPW}$ .

Figs. 11.b to 11.f. and Figs. 12.b to 12.f show the results of BP and BPW

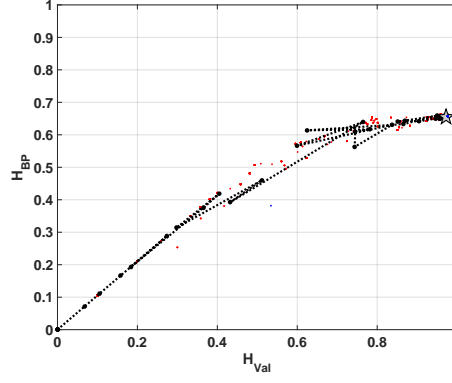


Figure 9. Evolution of statistical properties in double entropy plane of SWITCH map  $H_{val}$  vs  $H_{BP}$ .

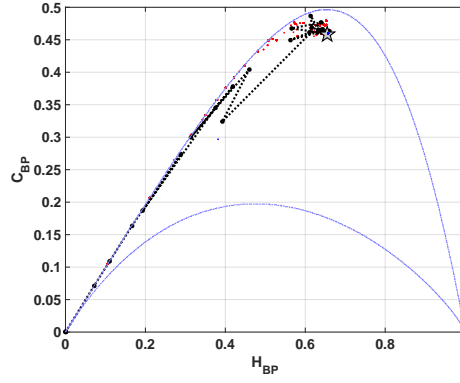


Figure 10. Evolution of statistical properties in entropy-complexity plane of SWITCH map  $C_{BP}$ .

quantifiers for EVEN and ODD respectively. Higher accuracy is required to achieve lower complexity than without using skipping. From the point of view of MP a great improvement is obtained using any of the skipping strategies but ODD is slightly better than EVEN. Missing patterns are reduced to  $MP = 118$  for EVEN and ODD, increasing the maximum allowed Bandt & Pompe entropy that reaches the mean value  $\langle H_{BP} \rangle = 0.8381$  for EVEN, and  $\langle H_{BP} \rangle = 0.9094$ . The complexity is reduced to  $\langle C_{BP} \rangle = 0.224$  for EVEN and  $\langle C_{BP} \rangle = 0.282$  for ODD. The amount of bits necessary to converge to this value is  $B > 40$  for both EVEN and ODD maps.

FALTA VER PLANOS DOBLE ENTROPÍA

FALTA VER PLANOS ENTROPÍA-COMPLEJIDAD

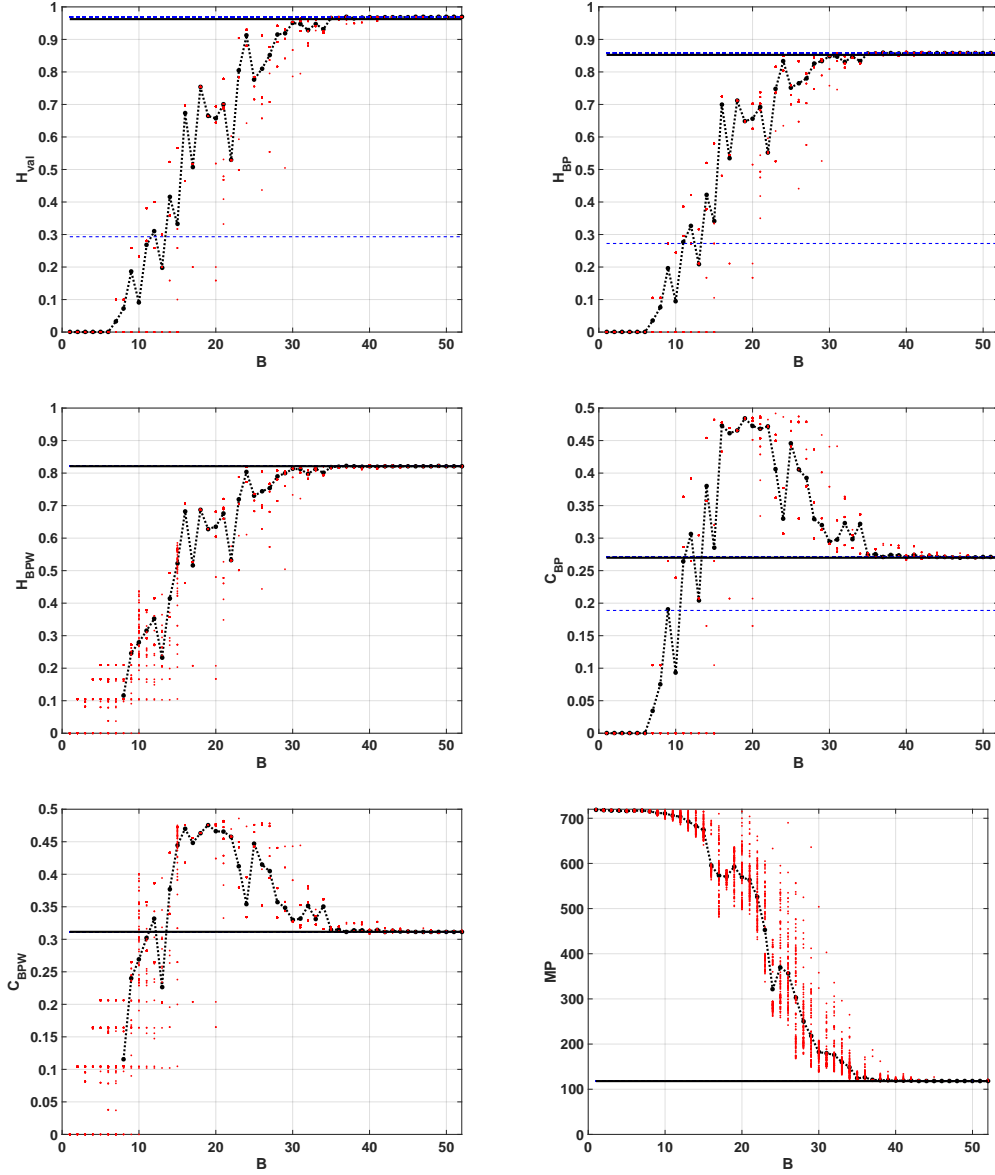


Figure 11. Statistical properties of the EVEN map: (a)  $H_{val}$  vs  $B$  (b)  $H_{BP}$  vs  $B$  (c)  $C_{BP}$  vs  $B$  (d)  $MP$  vs  $B$ .

## 4 Conclusions

In summary:

- Not every number base can be represented by a device with a certain base. For example, a base ten number can not be exactly represented in a conventional computer, it will always have an error inherent of the system..No todas las bases numéricas son representables con una máquina de base distinta. Por ejemplo, no se puede representar la base 10 con base 2.
- En una máquina de cálculo "a medida", como la que puede implementarse

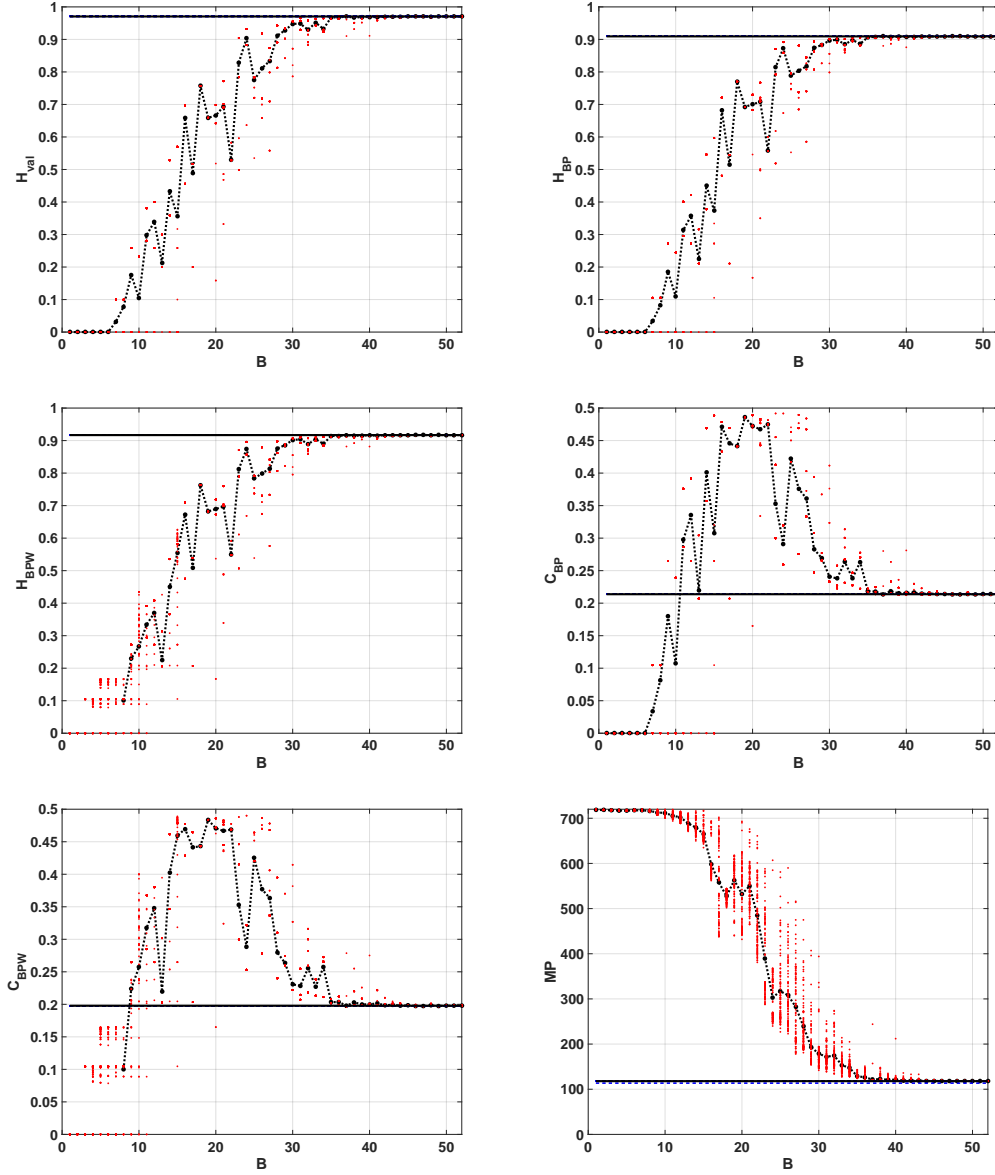


Figure 12. Statistical properties of the ODD map: (a)  $H_{val}$  vs  $B$  (b)  $H_{BP}$  vs  $B$  (c)  $C_{BP}$  vs  $B$  (d)  $MP$  vs  $B$ .

en ASICs o FPGAs existen limitaciones en el bus de datos y en la electrónica de cálculo. Si la electrónica de cálculo debe ser reducida se recomienda usar mapas que puedan ser calculados sólo con sumas y restas de la variable pseudoaleatoria.

- Los mapas que sólo tienen operaciones de shifteo en la base de la máquina de cálculo inevitablemente caerán a cero en tantas iteraciones como el largo de la mantisa de representación. Por ejemplo el tent en base 2.
- La comparación entre BP y BPW permite detectar el comportamiento del sistema. Puede detectarse si el atractor cae a un punto fijo y diferenciar si el transitorio es corto o largo, respecto de la cantidad de iteraciones del mapa.

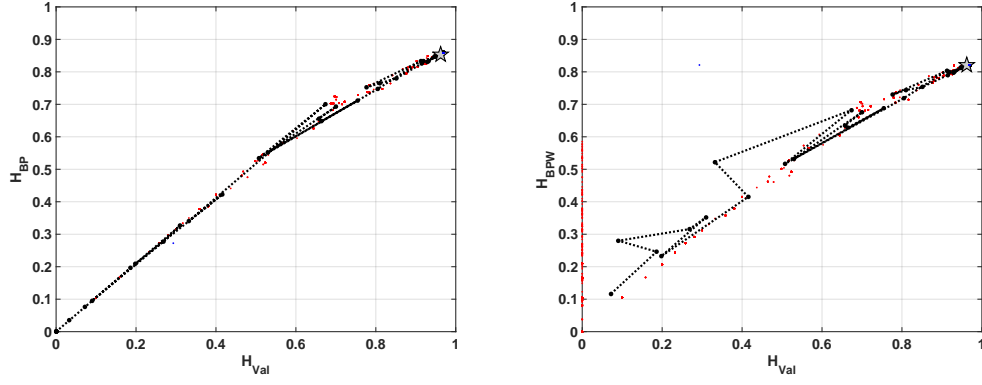


Figure 13. Evolution of statistical properties in double entropy plane of EVEN map: (a)  $H_{val}$  vs  $H_{BP}$  (b)  $H_{val}$  vs  $H_{BPW}$ .

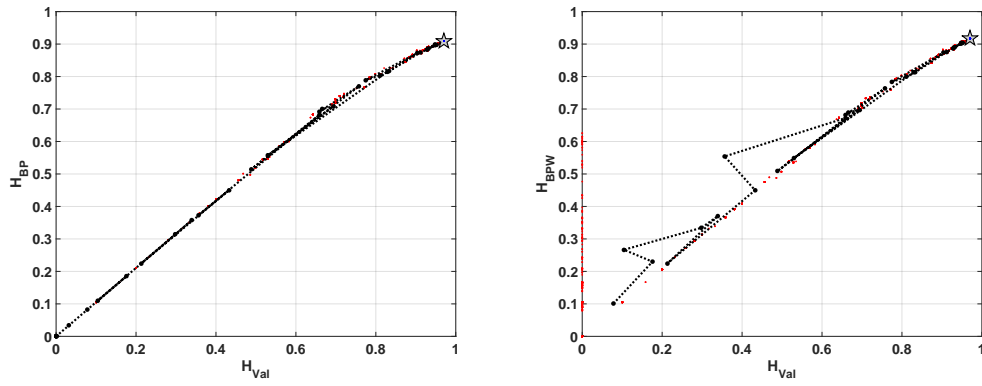


Figure 14. Evolution of statistical properties in double entropy plane of ODD map: (a)  $H_{val}$  vs  $H_{BP}$  (b)  $H_{val}$  vs  $H_{BPW}$ .

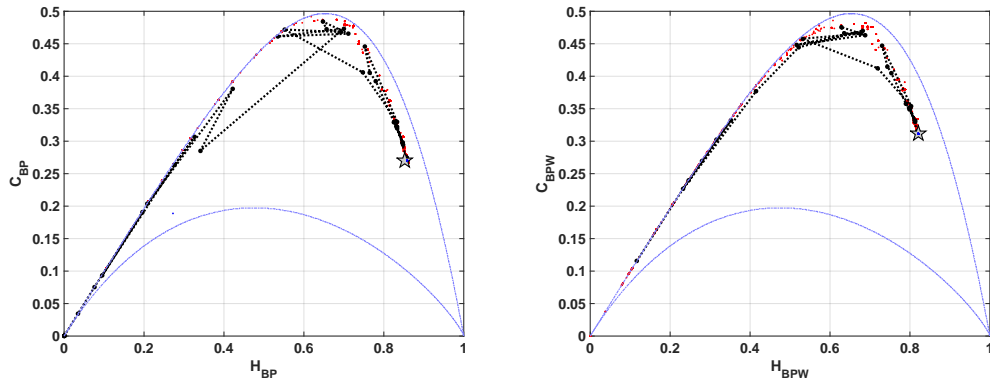


Figure 15. Evolution of statistical properties in entropy-complexity plane of EVEN map: (a)  $C_{BP}$  vs  $H_{BP}$  (b)  $C_{BPW}$  vs  $H_{BPW}$ .

- Como se menciona en el paper de referencia, el período del mapa SWITCH aumenta respecto del simple. También se nota una mejora marginal en la mezcla de la secuencia. La distribución de valores es buena en todos los casos.



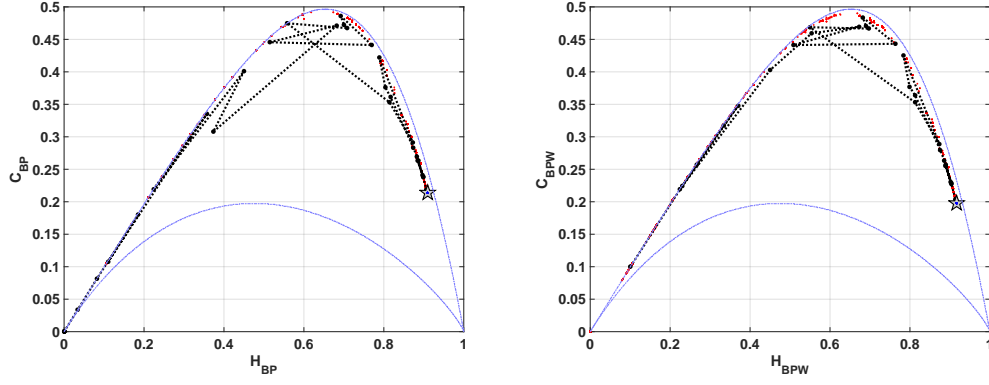


Figure 16. Evolution of statistical properties in entropy-complexity plane of ODD map: (a)  $C_{BP}^{ab}$  vs  $H_{BP}$  (b)  $C_{BPW}$  vs  $H_{BPW}$ .

- el skipping empeora el período pero mejora sustancialmente la mezcla de los valores. Esto puede verse en BP, BPW y MP.

produces a non-monotonous evolution toward the floating point result. This result is relevant because it shows that increasing the precision is not always recommended. It is specially interesting to note that some systems (TENT) with very nice statistical properties in the world of the real numbers, become “pathological” when binary numerical representations are used.

## Acknowledgment

This work was partially financed by CONICET (PIP2008), and UNMDP.

## References

- [1] C. Grebogi, E. Ott, J. A. Yorke, Roundoff-induced periodicity and the correlation dimension of chaotic attractors, *Phys. Rev. A* 38 (7) (1988) 3688–3692. doi:10.1103/PhysRevA.38.3688.
- [2] N. Nagaraj, M. C. Shastri, P. G. Vaidya, Increasing average period lengths by switching of robust chaos maps in finite precision, *Eur. Phys. J. Spec. Top.* 165 (1) (2008) 73–83. doi:10.1140/epjst/e2008-00850-4.
- [3] X. Liu, K.-L. Teo, H. Zhang, G. Chen, Switching control of linear systems for generating chaos, *Fractals An Interdiscip. J. Complex Geom. Nat.* 30 (3) (2006) 725–733. doi:10.1016/j.chaos.2005.03.020.
- [4] E. Gluskin, The nonlinear-by-switching systems (a heuristic discussion of some basic singular systems)arXiv:0801.3652.

- [5] D. G. Zarlenga, H. A. Larrondo, C. M. Arizmendi, F. Family, Complex synchronization structure of an overdamped ratchet with discontinuous periodic forcing, *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.* 80 (1) (2009) 011127. doi:10.1103/PhysRevE.80.011127.
- [6] L. De Micco, H. A. Larrondo, A. Plastino, O. A. Rosso, Quantifiers for randomness of chaotic pseudo-random number generators., *Philos. Trans. A. Math. Phys. Eng. Sci.* 367 (1901) (2009) 3281–3296. arXiv:0812.2250, doi:10.1098/rsta.2009.0075.
- [7] O. A. Rosso, H. A. Larrondo, M. T. Martin, A. Plastino, M. A. Fuentes, Distinguishing Noise from Chaos, *Phys. Rev. Lett.* 99 (15) (2007) 154102. doi:10.1103/PhysRevLett.99.154102.
- [8] L. De Micco, C. M. González, H. A. Larrondo, M. T. Martin, A. Plastino, O. A. Rosso, Randomizing nonlinear maps via symbolic dynamics, *Phys. A Stat. Mech. its Appl.* 387 (14) (2008) 3373–3383. doi:10.1016/j.physa.2008.02.037.
- [9] L. De Micco, J. G. Fernández, H. A. Larrondo, A. Plastino, O. A. Rosso, Sampling period, statistical complexity, and chaotic attractors, *Phys. A Stat. Mech. its Appl.* 391 (8) (2012) 2564–2575. doi:10.1016/j.physa.2011.12.042.
- [10] O. A. Rosso, L. De Micco, H. A. Larrondo, M. T. Martin, A. Plastino, Generalized Statistical Complexity Measure, *Int. J. Bifurc. Chaos* 20 (03) (2010) 775. doi:10.1142/S021812741002606X.
- [11] M. Antonelli, L. De Micco, H. A. Larrondo, Measuring the jitter of ring oscillators by means of information theory quantifiers, *Commun. Nonlinear Sci. Numer. Simul.* 43 (2017) 139–150. doi:10.1016/j.cnsns.2016.05.002.
- [12] C. Bandt, B. Pompe, Permutation entropy: a natural complexity measure for time series., *Phys. Rev. Lett.* 88 (17) (2002) 174102. doi:10.1103/PhysRevLett.88.174102.
- [13] J. M. Amigó, S. Zambrano, M. A. F. Sanjuán, True and false forbidden patterns in deterministic and random dynamics, *Europhys. Lett.* 79 (5) (2007) 50001. doi:10.1209/0295-5075/79/50001.
- [14] O. A. Rosso, L. C. Carpi, P. M. Saco, M. G. Ravetti, H. A. Larrondo, A. Plastino, The Amigo paradigm of forbidden/missing patterns: A detailed analysis, *Eur. Phys. J. B* 85 (12). doi:10.1140/epjb/e2012-30307-8.
- [15] C. E. Shannon, W. Weaver, The mathematical theory of communication, *Bell Syst. Tech. J.* 27 (3) (1948) 379–423. doi:10.1145/584091.584093.
- [16] P. W. Lamberti, M. T. Martin, A. Plastino, O. A. Rosso, Intensive entropic non-triviality measure, *Phys. A Stat. Mech. its Appl.* 334 (1-2) (2004) 119–131. doi:10.1016/j.physa.2003.11.005.
- [17] M. T. Martin, A. Plastino, O. A. Rosso, Statistical complexity and disequilibrium, *Phys. Lett. Sect. A Gen. At. Solid State Phys.* 311 (2-3) (2003) 126–132. doi:10.1016/S0375-9601(03)00491-2.

- [18] M. T. Martin, A. Plastino, O. A. Rosso, Generalized statistical complexity measures: Geometrical and analytical properties, *Phys. A Stat. Mech. its Appl.* 369 (2) (2006) 439–462. doi:10.1016/j.physa.2005.11.053.
- [19] O. A. Rosso, L. C. Carpi, P. M. Saco, M. Gómez Ravetti, A. Plastino, H. A. Larrondo, Causality and the entropy-complexity plane: Robustness and missing ordinal patterns, *Phys. A Stat. Mech. its Appl.* 391 (1-2) (2012) 42–55. arXiv:1105.4550, doi:10.1016/j.physa.2011.07.030.
- [20] K. Keller, H. Laufer, Symbolic Analysis of High-Dimensional Time Series, *Int. J. Bifurc. Chaos* 13 (9) (2003) 2657–2668. doi:10.1142/S0218127403008168.
- [21] K. Keller, M. Sinn, Ordinal analysis of time series, in: *Phys. A Stat. Mech. its Appl.*, Vol. 356, 2005, pp. 114–120. doi:10.1016/j.physa.2005.05.022.
- [22] O. A. Rosso, L. Zunino, D. G. Pérez, A. Figliola, H. A. Larrondo, M. Garavaglia, M. T. Martin, A. Plastino, Extracting features of Gaussian self-similar stochastic processes via the Bandt-Pompe approach, *Phys. Rev. E - Stat. Nonlinear, Soft Matter Phys.* 76 (6) (2007) 061114. doi:10.1103/PhysRevE.76.061114.
- [23] C. Anteneodo, A. R. Plastino, Some features of the López-Ruiz-Mancini-Calbet (LMC) statistical measure of complexity, *Phys. Lett. Sect. A Gen. At. Solid State Phys.* 223 (5) (1996) 348–354. doi:10.1016/S0375-9601(96)00756-6.
- [24] C. M. González, H. A. Larrondo, O. A. Rosso, Statistical complexity measure of pseudorandom bit generators, *Phys. A Stat. Mech. its Appl.* 354 (1-4) (2005) 281–300. doi:10.1016/j.physa.2005.02.054.
- [25] J. M. Amigó, L. Kocarev, J. Szczepanski, Order patterns and chaos, *Phys. Lett. Sect. A Gen. At. Solid State Phys.* 355 (1) (2006) 27–31. doi:10.1016/j.physleta.2006.01.093.
- [26] J. M. Amigó, L. Kocarev, I. Tomovski, Discrete entropy, *Phys. D Nonlinear Phenom.* 228 (1) (2007) 77–85. doi:10.1016/j.physd.2007.03.001.
- [27] J. M. Amigó, S. Zambrano, M. A. F. Sanjuán, Combinatorial detection of determinism in noisy time series, *Europhys. Lett.* 83 (6) (2008) 60005. doi:10.1209/0295-5075/83/60005.
- [28] J. Amigó, *Permutation Complexity in Dynamical Systems*, Springer Series in Synergetics, Springer Berlin Heidelberg, Berlin, Heidelberg, 2010. arXiv:arXiv:1011.1669v3, doi:10.1007/978-3-642-04084-9.
- [29] R. Wackerbauer, A. Witt, H. Atmanspacher, J. Kurths, H. Scheingraber, A comparative classification of complexity measures, *Chaos, Solitons and Fractals* 4 (1) (1994) 133–173. doi:10.1016/0960-0779(94)90023-X.
- [30] R. López-Ruiz, H. L. Mancini, X. Calbet, R. López-Ruiz, H. L. Mancini, X. Calbet, A statistical measure of complexity, *Phys. Lett. A* 209 (5-6) (1995) 321–326. doi:10.1016/0375-9601(95)00867-5.

- [31] O. A. Rosso, H. A. Larrondo, M. T. Martin, A. Plastino, M. A. Fuentes, Distinguishing noise from chaos, *Phys. Rev. Lett.* 99 (15) (2007) 154102. arXiv:1401.2139, doi:10.1103/PhysRevLett.99.154102.
- [32] M. Jessa, The period of sequences generated by tent-like maps, *IEEE Trans. Circuits Syst. I Fundam. Theory Appl.* 49 (1) (2002) 84–89. doi:10.1109/81.974880.
- [33] S. Callegari, G. Setti, P. Langlois, A CMOS tailed tent map for the generation of uniformly distributed chaotic sequences, in: *Proc. 1997 IEEE Int. Symp. Circuits Syst. Circuits Syst. Inf. Age ISCAS '97*, Vol. 2, IEEE, pp. 781–784. doi:10.1109/ISCAS.1997.621829.
- [34] S. Li, *When Chaos Meets Computers* (2004). arXiv:0405038.